

Title: PSI 2019/2020 - Gravitational Physics - Lecture 3

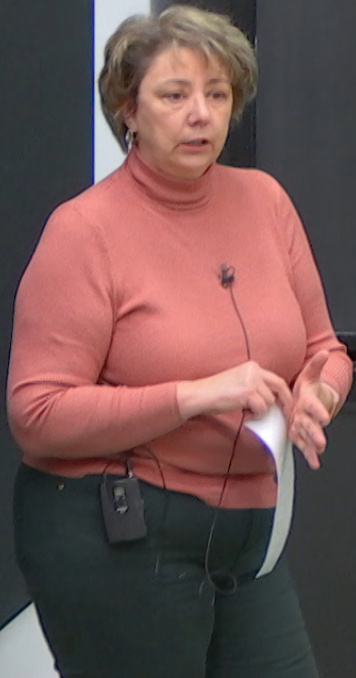
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Collection: PSI 2019/2020 - Gravitational Physics

Date: January 08, 2020 - 10:15 AM

URL: <http://pirsa.org/20010047>

LECTURE 3 Lie derivative & symmetries



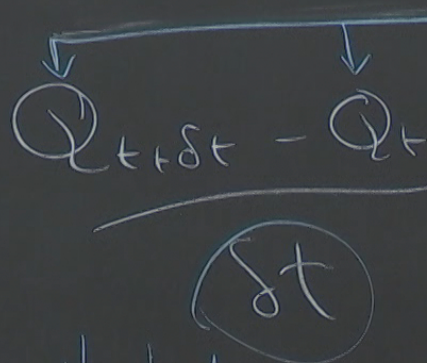


# LECTURE 3 Lie derivative & symmetries

Recall the defn of derivative:

$$\underbrace{\frac{dQ}{dt}}_{?}$$

$$\lim_{\delta t \rightarrow 0}$$

$$\frac{Q_{t+\delta t} - Q_t}{\delta t}$$


On a manifold these are different tangent spaces.

what does this mean?



eties

On a manifold  
these are  
different tgt  
spaces.

Return to diffeos,  
clearly form a group,  
what are the generators?

$$X^\mu \rightarrow X^\mu + \delta X^\mu$$

now look at  $\delta X^\mu$  "small"



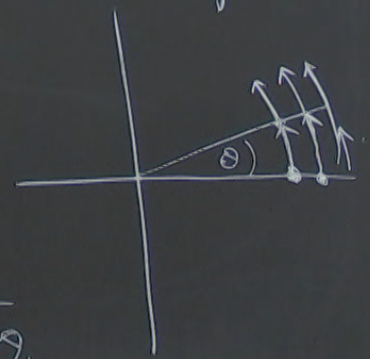
diffeos,  
 in a group,  
 the generators?  
 $X^\mu + \delta X^\mu$   
 at  $\delta X^\mu$  'small'

Write  $\delta X^\mu = \epsilon \xi^\mu$   
param      cpts of vector

Vectors generate coord transfs

e.g.  $\mathbb{R}^2$ :  $x' = x - \epsilon y$   
 $y' = y + \epsilon x$

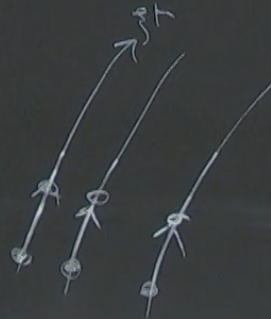
$\xi^\mu = (-y, x) \leftrightarrow \frac{\partial}{\partial \theta}$





The curves to which a vector field,  $\xi$ ,  
is tangent are the integral curves of  $\xi$

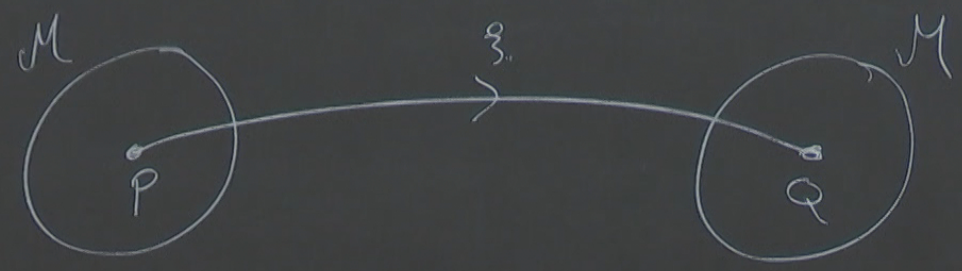
$$X_{\xi}^M = X_0^M + \epsilon \xi^M + o(\epsilon^2)$$





eld,  $\xi$ ,  
of  $\xi$

What are push forward & pull back?



$$P \in X^M \longrightarrow Q = X^M + \xi \xi^M$$

under this transformation

& pull back?

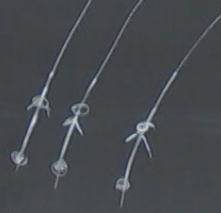
$$V'^{\mu'} = \frac{\partial X'^{\mu'}}{\partial X^{\mu}} V^{\mu}$$

$$= \left( \delta_{\mu}^{\mu'} + \epsilon \Sigma_{\mu}^{\mu'} \right) V^{\mu}$$

- push forward on  $V^{\mu}$ .

$$X^{\mu} + \epsilon \Sigma^{\mu}$$





under this trans

Take deriv of  $V$  along integral  
curve of  $\xi$

$$\begin{aligned}
 \frac{dV}{dt} &= \lim_{\epsilon \rightarrow 0} \frac{V(x')_{t+\epsilon t} - V_x}{\epsilon t} \\
 &= \lim_{\epsilon \rightarrow 0} \frac{[V(x_0) + \epsilon \xi^M \partial_\nu V^M]_{x_0} - [V(x_0) + \epsilon \xi^M \partial_\nu V^M]}{\epsilon} \\
 &= [\xi, V]^M
 \end{aligned}$$



Can generalise to arb tensors;  
denote this Lie Derivative by  $L_{\xi}$

$$(L_{\xi} \omega)_{\mu} = \xi^{\nu} \partial_{\nu} \omega_{\mu} + \xi^{\nu}{}_{,\mu} \omega_{\nu}$$

Ex.  $(L_{\xi} g)_{\mu\nu} = \xi^{\sigma} \partial_{\sigma} g_{\mu\nu} + \xi^{\sigma}{}_{,\mu} g_{\sigma\nu} + \xi^{\sigma}{}_{,\nu} g_{\mu\sigma}$



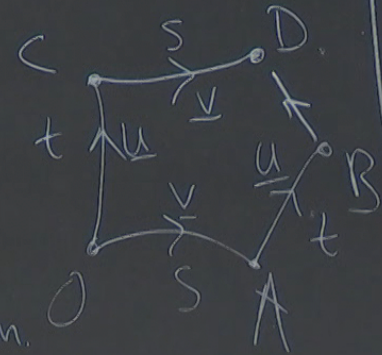
## Geometry of $\mathcal{L}$

Lie deriv is intrinsically linked  
to symmetries & coords.

Consider the commutator

$$[u, v]$$

Let  $A, B, C, D$  be as in diagram.



S, t sm  
in local



$S, t$  small, perform analysis  
in local chart.

use 
$$\begin{cases} X_A^M = X_0^M + S V^M + \frac{1}{2} S^2 V^\nu \partial_\nu V^M \\ U_A^M = U_0^M + S V^\nu \partial_\nu U^M \end{cases}$$

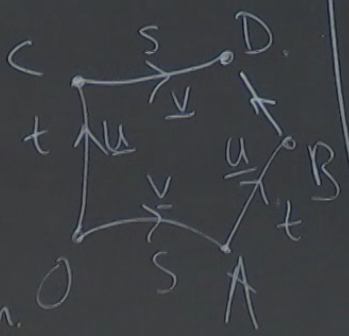


physically linked

ords.

mutator

S in diagram.



S, t small, perform analysis "exp[sV]"  
in local chart.

use 
$$\begin{cases} X_A^\mu = X_0^\mu + sV^\mu + \frac{1}{2}s^2 V^\nu \partial_\nu V^\mu \\ U_A^\mu = U_0^\mu + sV^\nu \partial_\nu U^\mu \end{cases}$$

$$\begin{aligned} X_B^\mu &= X_A^\mu + tU^\mu|_A + \frac{1}{2}t^2 U^\nu \partial_\nu U^\mu|_A \\ &= X_0^\mu + sV_0^\mu + \frac{1}{2}s^2 V^\nu \partial_\nu V^\mu \\ &\quad + tU_0^\mu + st V^\nu \partial_\nu U_0^\mu + \frac{1}{2}t^2 U^\nu \partial_\nu U_0^\mu \end{aligned}$$

Similarly

$$X_D^\mu = X_0^\mu + tU^\mu|_B + \frac{1}{2}t^2 U^\nu \partial_\nu U^\mu|_B + sV^\mu$$

Discrepancy



ysis "exp[sV]"

$$\frac{1}{2} s^2 V^\nu \partial_\nu V^M$$
$${}^\nu \partial_\nu U^M$$

$$t^2 U^\nu \partial_\nu U^M$$

$${}^\nu \partial_\nu V^M$$

$$U_0^M + \frac{1}{2} t^2 U^\nu \partial_\nu U_0^M$$

Similarly

$$X_D^M = X_0^M + t U_0^M + \frac{1}{2} t^2 U^\nu \partial_\nu U_0^M$$
$$+ s V_0^M + st U^\nu \partial_\nu V_0^M + \frac{1}{2} s^2 V^\nu \partial_\nu V_0^M$$

Discrepancy between D & B:

$$X_D^M - X_B^M = st (U^\nu \partial_\nu V^M - V^\nu \partial_\nu U^M)$$
$$= st [U, V]^M$$

metas  $(\underline{u}_\nu - \underline{v}_\nu)$



If  $[u, v] \neq 0$ , can't use  $(s, t)$  as coords.

If  $L_{\xi} Q = 0$ , then  $Q$  is unchanged as we Lie drag along integral curves of  $\xi$ .

s coords.

ged as

3.

al curves

ce.

Defn A Killing vector is a vector field along which the metric is invariant:

$$\xi^\sigma \partial_\sigma g_{\mu\nu} + \xi_{,\mu} g_{\sigma\nu} + \xi_{,\nu} g_{\mu\sigma} = 0.$$



$$L_3 g = 0 \Rightarrow g_{\mu\nu, \varphi} = 0 \quad \text{no } \varphi\text{-dep.}$$

$$L_{1,2} \quad \frac{\sin \varphi}{\cos \theta} \underline{g_{\mu\nu, \theta}} + \frac{\cos \varphi}{-\sin \theta} (g_{\theta\nu} \delta_{\mu}^{\varphi} + g_{\theta\mu} \delta_{\nu}^{\varphi})$$

$$- \frac{\csc^2 \theta \cos \varphi}{-\sin \theta} (g_{\varphi\nu} \delta_{\mu}^{\theta} + g_{\varphi\mu} \delta_{\nu}^{\theta})$$

$$- \frac{\cot \theta \sin \varphi}{\cos \theta} \underline{g_{\varphi\nu} \delta_{\mu}^{\varphi} + g_{\varphi\mu} \delta_{\nu}^{\varphi}}$$



$$\xi_3 = d\varphi.$$

$\varphi$ -dep.

$$\left\{ \begin{aligned} g_{\mu\nu,\theta} &= \cot\theta (g_{\varphi\nu} \delta_{\mu}^{\varphi} + g_{\varphi\mu} \delta_{\nu}^{\varphi}) \quad (\sin) \downarrow (\cos) \\ &\& \left\{ \begin{aligned} g_{\theta\nu} \delta_{\mu}^{\varphi} + g_{\theta\mu} \delta_{\nu}^{\varphi} &= \csc^2\theta (g_{\varphi\nu} \delta_{\mu}^{\theta} + g_{\theta\mu} \delta_{\nu}^{\varphi}) \\ \leftarrow g_{\mu=\varphi, \nu=\theta} &\rightarrow g_{\theta\theta} = 0. \text{ also, } g_{\theta\varphi} = g_{\theta\theta} g_{\varphi\varphi} = 0 \end{aligned} \right. \end{aligned} \right.$$

$$\mu = \nu = \theta : g_{\theta\theta,\theta} = 0$$

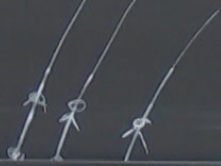
$$\mu = \nu = \varphi : g_{\varphi\varphi,\theta} = 2\cot\theta g_{\varphi\varphi} \rightarrow g_{\varphi\varphi} \propto \sin^2\theta.$$

$$2g_{\theta\varphi} = 0$$

$$\mu = \theta \quad \nu = \varphi.$$

$$g_{\theta\theta} = g_{\varphi\varphi} \csc^2\theta.$$

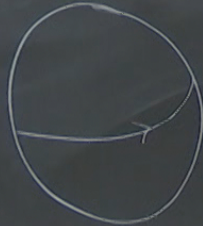




under this trans

Put together:

$$ds^2 = \underbrace{\gamma_{\mu\nu}^{(t,r)} dx^\mu dx^\nu}_{\mu, \nu = t, r} - C^2(t, r) \underbrace{[d\theta^2 + \sin^2\theta d\phi^2]}_{\text{metric on a unit } S^2}$$



under this transformation

A useful relation between  $\langle \omega | d$ :

$$\langle d\omega | [u, v] \rangle = u \langle \omega | v \rangle - v \langle \omega | u \rangle - \langle \omega | [u, v] \rangle$$