

Title: PSI 2019/2020 - Gravitational Physics - Lecture 3

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LECTURE 3 Lie derivative & symmetries



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Lie derivative & symmetries

Recall the defn of derivative:

$$\underbrace{\frac{dQ}{dt}}_{?}$$

$$\lim_{\delta t \rightarrow 0}$$

$$\frac{Q_{t+\delta t} - Q_t}{\delta t}$$

On a manifold these are different tangent spaces.

what does this mean?

eties

On a manifold
these are
different tgt
spaces.

Return to diffeos,
clearly form a group,
what are the generators?

$$X^\mu \rightarrow X^\mu + \delta X^\mu$$

now look at δX^μ "small"

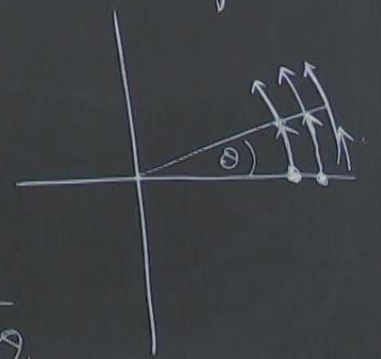
diffeos,
 in a group,
 the generators?
 $X^\mu + \delta X^\mu$
 at δX^μ 'small'

Write $\delta X^\mu = \epsilon \xi^\mu$
param cpts of vector

Vectors generate coord transfs

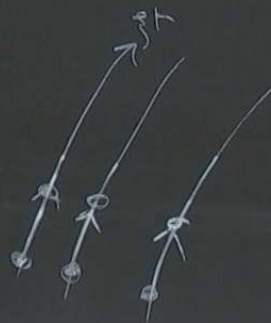
e.g. \mathbb{R}^2 : $x' = x - \epsilon y$
 $y' = y + \epsilon x$

$\xi^\mu = (-y, x) \leftrightarrow \frac{\partial}{\partial \theta}$



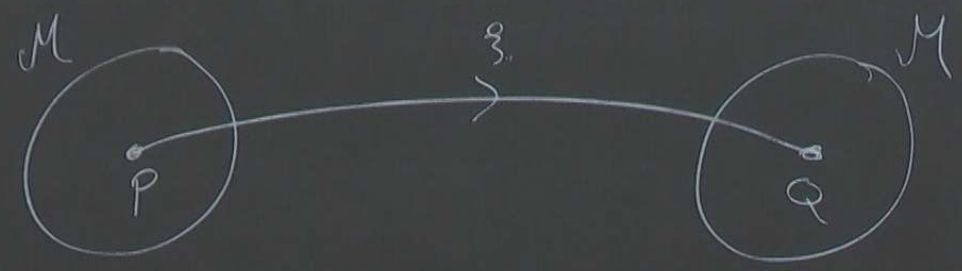
The curves to which a vector field, ξ ,
is tangent are the integral curves of ξ

$$X_{\xi}^{\mu} = X_0^{\mu} + \epsilon \xi^{\mu} + o(\epsilon^2)$$



eld, ξ ,
of ξ

What are push forward & pull back?



$$P \in X^M \longrightarrow Q = X^M + \xi \xi^M$$

under this transformation

& pull back?

$$V'^{\mu'} = \frac{\partial X'^{\mu'}}{\partial X^{\mu}} V^{\mu}$$

$$= \left(\delta_{\mu}^{\mu'} + \epsilon \Sigma_{\mu}^{\mu'} \right) V^{\mu}$$

- push forward on V^{μ} .

$$X^{\mu} + \epsilon \Sigma^{\mu}$$



under this trans

Take deriv of V along integral
curve of ξ

$$\begin{aligned} \frac{dV}{dt} &= \lim_{\epsilon \rightarrow 0} \frac{V(x')_{t+\epsilon t} - V_x}{\epsilon t} \\ &= \lim_{\epsilon \rightarrow 0} \frac{[V(x_0) + \epsilon \xi^M \partial_\nu V^M]_{x_0} - [V]_{x_0}}{\epsilon} \\ &= [\xi, V]^M \end{aligned}$$

Can generalise to arb tensors;
denote this Lie Derivative by L_{ξ}

$$(L_{\xi} \omega)_{\mu} = \xi^{\nu} \partial_{\nu} \omega_{\mu} + \xi^{\nu}{}_{,\mu} \omega_{\nu}$$

Ex. $(L_{\xi} g)_{\mu\nu} = \xi^{\sigma} \partial_{\sigma} g_{\mu\nu} + \xi^{\sigma}{}_{,\mu} g_{\sigma\nu} + \xi^{\sigma}{}_{,\nu} g_{\mu\sigma}$

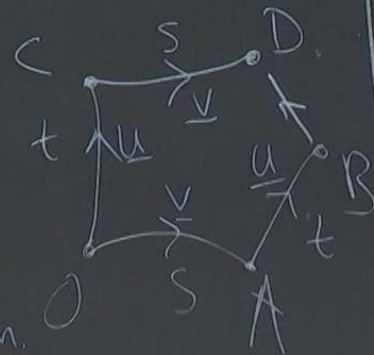
Geometry of \mathcal{L}

Lie deriv is intrinsically linked
to symmetries & coords.

Consider the commutator

$$[u, v]$$

Let A, B, C, D be as in diagram.



S, t sm
in local

S , t small, perform analysis
in local chart.

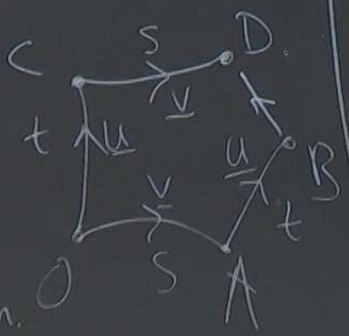
use
$$\begin{cases} X_A^M = X_0^M + S V^M + \frac{1}{2} S^2 V^\nu \partial_\nu V^M \\ U_A^M = U_0^M + S V^\nu \partial_\nu U^M \end{cases}$$

Physically linked

ords.

mutator

S in diagram.



S, t small, perform analysis "exp[sV]"
in local chart.

use
$$\begin{cases} X_A^\mu = X_0^\mu + sV^\mu + \frac{1}{2}s^2 V^\nu \partial_\nu V^\mu \\ U_A^\mu = U_0^\mu + sV^\nu \partial_\nu U^\mu \end{cases}$$

$$\begin{aligned} X_B^\mu &= X_A^\mu + tU^\mu|_A + \frac{1}{2}t^2 U^\nu \partial_\nu U^\mu|_A \\ &= X_0^\mu + sV_0^\mu + \frac{1}{2}s^2 V^\nu \partial_\nu V^\mu \\ &\quad + tU_0^\mu + stV^\nu \partial_\nu U_0^\mu + \frac{1}{2}t^2 U^\nu \partial_\nu U_0^\mu \end{aligned}$$

Similarly

$$X_D^\mu = X_0^\mu + sV^\mu + sV^\nu \partial_\nu X^\mu$$

Discrepancy

ysis "exp[sV]"

$$\frac{1}{2} s^2 V^\nu \partial_\nu V^M$$

$$t^2 U^\nu \partial_\nu U^M$$

$$U_0^M + \frac{1}{2} t^2 U^\nu \partial_\nu U_0^M$$

Similarly

$$X_D^M = X_0^M + t U_0^M + \frac{1}{2} t^2 U^\nu \partial_\nu U_0^M$$
$$+ s V_0^M + st U^\nu \partial_\nu V_0^M + \frac{1}{2} s^2 V^\nu \partial_\nu V_0^M$$

Discrepancy between D & B:

$$X_D^M - X_B^M = st (U^\nu \partial_\nu V^M - V^\nu \partial_\nu U^M)$$
$$= st [U, V]^M$$

metas $(\underline{u}_\nu - \underline{v}_\nu)$

If $[u, v] \neq 0$, can't use (s, t) as coords.

If $L_{\xi} Q = 0$, then Q is unchanged as we Lie drag along integral curves of ξ .

s coords

ged as

3.

al curves

ce.

Defn A Killing vector is a vector field along which the metric is invariant:

$$\xi^\sigma \partial_\sigma g_{\mu\nu} + \xi_{,\mu} g_{\sigma\nu} + \xi_{,\nu} g_{\mu\sigma} = 0.$$

$$L_3 g = 0 \Rightarrow g_{\mu\nu, \varphi} = 0 \quad \text{no } \varphi\text{-dep.}$$

$$L_{1,2} \frac{\sin \varphi}{\cos \theta} g_{\mu\nu, \theta} + \frac{\cos \varphi}{-\sin \theta} (g_{\theta\nu} \delta_{\mu}^{\varphi} + g_{\theta\mu} \delta_{\nu}^{\varphi})$$

$$- \frac{\csc^2 \theta \cos \varphi}{-\sin \theta} (g_{\varphi\nu} \delta_{\mu}^{\theta} + g_{\varphi\mu} \delta_{\nu}^{\theta})$$

$$- \frac{\cot \theta \sin \varphi}{\cos \theta} (g_{\varphi\nu} \delta_{\mu}^{\varphi} + g_{\varphi\mu} \delta_{\nu}^{\varphi})$$

$$\xi_3 = d\varphi.$$

φ -dep

$$\left\{ \begin{aligned} g_{\mu\nu,\theta} &= \cot\theta (g_{\varphi\nu} \delta_{\mu}^{\varphi} + g_{\varphi\mu} \delta_{\nu}^{\varphi}) \quad (\sin) \downarrow (\cos) \\ &\& \left\{ \begin{aligned} g_{\theta\nu} \delta_{\mu}^{\varphi} + g_{\theta\mu} \delta_{\nu}^{\varphi} &= \csc^2\theta (g_{\varphi\nu} \delta_{\mu}^{\theta} + g_{\theta\mu} \delta_{\nu}^{\varphi}) \\ \leftarrow g_{\mu=\varphi, \nu=\theta} &\rightarrow g_{\theta\theta} = 0 \quad \text{also, } g_{\theta\varphi} = g_{\theta\theta} g_{\varphi\varphi} = 0 \end{aligned} \right. \end{aligned} \right.$$

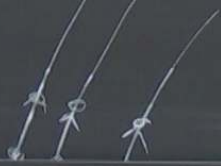
$$\mu = \nu = \theta : g_{\theta\theta,\theta} = 0$$

$$\mu = \nu = \varphi : g_{\varphi\varphi,\theta} = 2\cot\theta g_{\varphi\varphi} \rightarrow g_{\varphi\varphi} \propto \sin^2\theta.$$

$$2g_{\theta\varphi} = 0$$

$$\mu = \theta \quad \nu = \varphi.$$

$$g_{\theta\theta} = g_{\varphi\varphi} \csc^2\theta.$$



under this trans

Put together:

$$ds^2 = \underbrace{\gamma_{\mu\nu}^{(t,r)} dx^\mu dx^\nu}_{\mu, \nu = t, r} - C^2(t, r) \underbrace{[d\theta^2 + \sin^2\theta d\phi^2]}_{\text{metric on a unit } S^2}$$



under this transformation

A useful relation between \mathcal{L} and d :

$$\langle d\omega | [u, v] \rangle = u \langle \omega | v \rangle - v \langle \omega | u \rangle - \langle \omega | [u, v] \rangle$$