

Title: PSI 2019/2020 - Condensed Matter (Wang) - Lecture 15

Speakers: Chong Wang

Collection: PSI 2019/2020 - Condensed Matter (Wang)

Date: December 05, 2019 - 10:45 AM

URL: <http://pirsa.org/19120008>

Today: 2d topological superconductor, non-abelian anyons

Ref. Ivanov, PRL 2001.

Read & Green, PRB 2000

Kitaev, arXiv:0506438

Length & Completeness.





anyon anyons

Consider a BCS-like state in 2d spinless fermion.

$$\text{BdG: } H = \sum_k \left[ \xi(k) c_k^\dagger c_k + \frac{\Delta}{2} (e^{i\theta} c_{-k} c_k + e^{-i\theta} c_k^\dagger c_{-k}^\dagger) \right]$$

$$\xi(k) = \frac{k^2}{2m} - E_f, \quad \Delta \ll E_f, \quad e^{i\theta} = \frac{k_x + ik_y}{|k|} \quad (\text{p-ip pairing})$$

No material known to have p-ip pairing for electron.

'Fractional quantum Hall' at  $\frac{\nu_e}{3} = 5/2$  is believed to

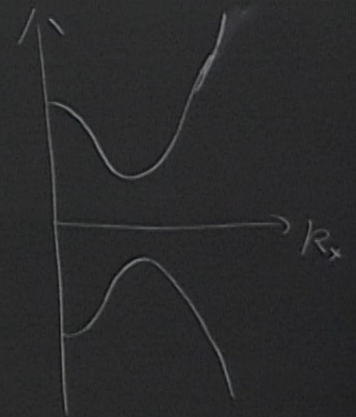
have p-ip pairing for "composite fermion".



$$H = \sum_{\substack{-\infty < k_y < \infty \\ -\infty < k_x < \infty}} (c_k^+ \quad c_{-k}^-) \begin{pmatrix} \xi(k) & \Delta e^{-i\theta} \\ \Delta e^{i\theta} & -\xi(k) \end{pmatrix} \begin{pmatrix} c_k \\ c_{-k}^+ \end{pmatrix}$$

Eigenvalue:  $\pm \sqrt{\xi^2 + \Delta^2} = \pm \bar{\epsilon}(k)$

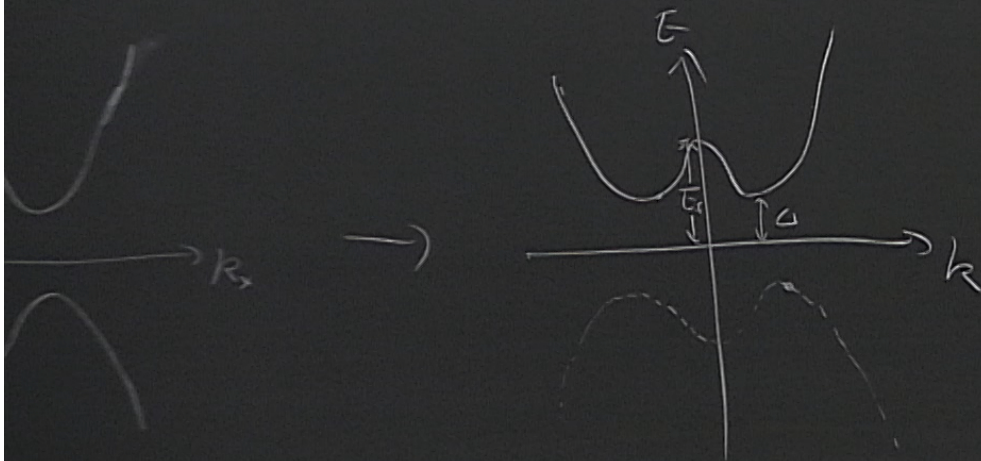
$$H = \sum_{0 < k_x} (\gamma_{k,+}^+ \quad \gamma_{k,-}^+) \begin{pmatrix} \bar{\epsilon}(k) \\ -\bar{\epsilon}(k) \end{pmatrix} \begin{pmatrix} \gamma_{k,+} \\ \gamma_{k,-} \end{pmatrix}$$





Define  $\gamma_k = \begin{cases} \gamma_{k,+} & \text{if } k > 0 \\ \gamma_{-k,-}^+ & \text{if } k < 0. \end{cases}$

$$\Rightarrow H = \sum_k E(k) \gamma_k^+ \gamma_k$$





Use a redundant formulation

Introduce a two-component  $\bar{\Psi}_k$ ,  $k \in \mathbb{R}^2$

$$H = \sum_k \bar{\Psi}_k^\dagger \begin{pmatrix} \xi(k) & \Delta e^{-i\theta} \\ \Delta e^{i\theta} & \xi(-k) \end{pmatrix} \bar{\Psi}_{-k} = \sum_k \bar{\Psi}_k^\dagger H_k \bar{\Psi}_k$$

only half of  $\bar{\Psi}_k$  is physical.

$$\bar{\Psi}_{-k} = U \bar{\Psi}_{-k}^*$$

$2 \times 2$  Unitary matrix.

Some algebra

requires  $U^\dagger H_k U = -H_{-k}^T \Rightarrow U = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$



Use a redundant formulation

Introduce a two-component  $\Psi_k, k \in \mathbb{R}^2$

$$H = \sum_k \Psi_k^\dagger \begin{pmatrix} \xi(k) & \Delta e^{-i\theta} \\ \Delta e^{i\theta} & \xi(-k) \end{pmatrix} \Psi_k = \sum_k \Psi_k^\dagger H_k \Psi_k$$

only half of  $\Psi_k$  is physical.

$$\Psi_k = \begin{pmatrix} \psi_k \\ \psi_{-k}^* \end{pmatrix} \xrightarrow{\text{Samp algebra}} \text{requires } U^\dagger H_k U = -H_{-k}^T \Rightarrow U = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

is Unitary matrix.

→ a.k.a. "particle-hole symmetry of S.C."

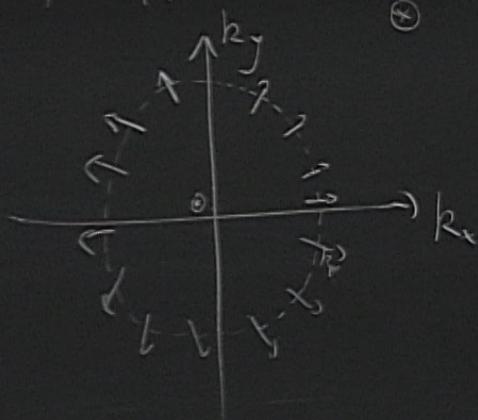
→ Redundancy, not a Symm.

⇒ A mode with  $E > 0$  → Another mode with  $-E$ .  
If a single mode at  $E = 0$ . ⇒ Stable!

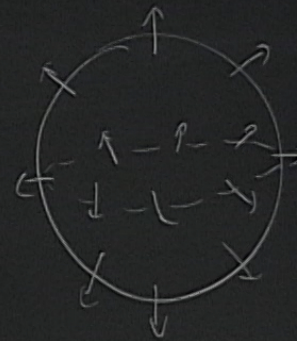


$$H_{\vec{k}} = \xi(\vec{k}) \sigma_z + \Delta \cos \theta \sigma_x + \Delta \sin \theta \sigma_y = \vec{d}(\vec{k}) \cdot \vec{\sigma}$$

Chern # of upper band = winding number of  $\vec{d}(\vec{k})$  on Bloch sphere = 1, low



identify  $\mathbb{R}^2$  as a pt

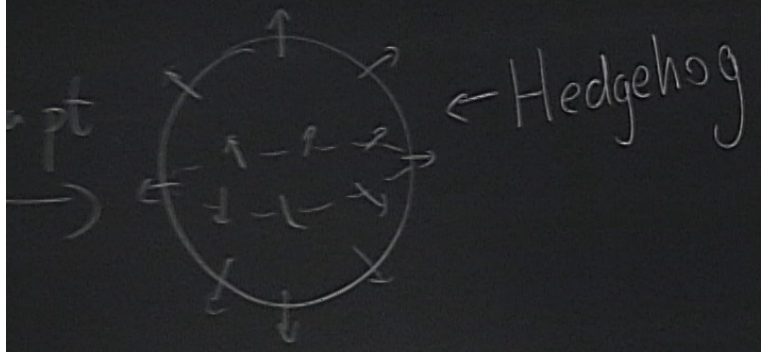


← Hedgehog



$\vec{\omega} \cdot \vec{\sigma}$   
 $\int \frac{d^3k}{(2\pi)^3}$  on Bloch sphere = 1,

lower band (unphysical) has  $C = -1$ .





$$\xi(k) = \frac{k^2}{2m} - E_f, \quad \Delta \ll E_f, \quad e^{i\theta} = \frac{k_x + ik_y}{|k|} \quad (\text{p-ip pairing})$$

No material known to have p-ip pairing for electron.

'Fractional quantum Hall' at  $\frac{n_e}{B} = 5/2$  is believed to have p-ip pairing for "composite fermion".

$\gamma_k$  if  $k_x > 0$

$\gamma_{-k}^+$  if  $k_x < 0$ .

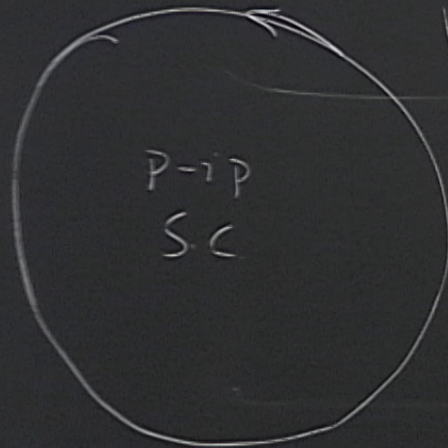
$$\sum_{\vec{k}} E(k) \gamma_{\vec{k}}^+ \gamma_{\vec{k}}$$

Dispersion  $E(k) \sim$  same as s-wave.

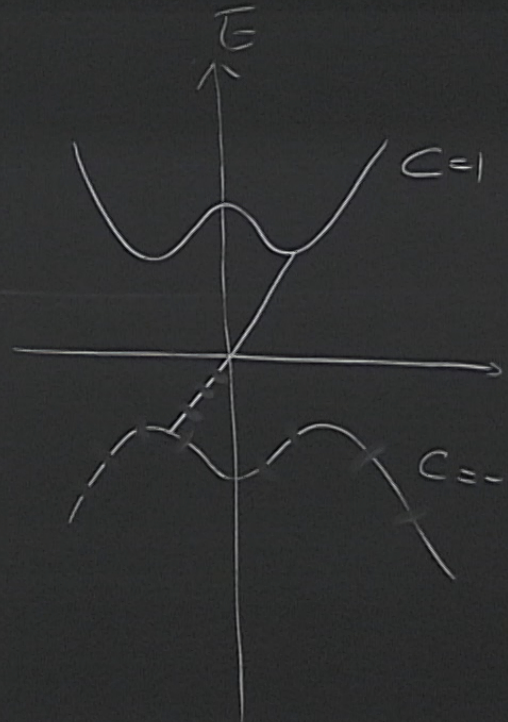
2d band  $\rightarrow$  Chern number.



Edge state:



Vacuum



Chiral fermion,



Chiral fermion, but only  $E > 0$  modes physical.

$$H = \sum_k v_x \chi_k^\dagger \chi_k, \quad \chi_{-k} = \chi_k^\dagger$$

Chiral Majorana fermion.

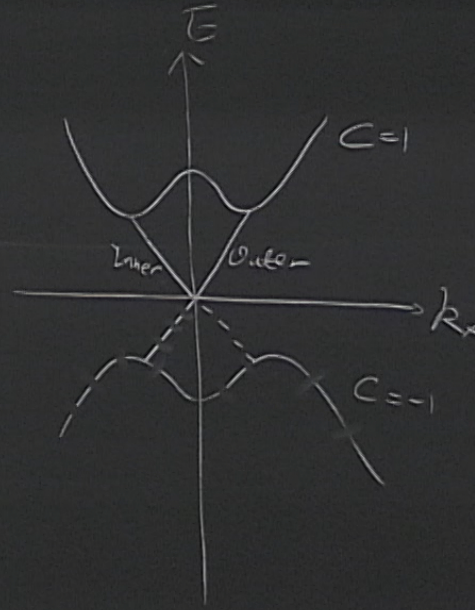
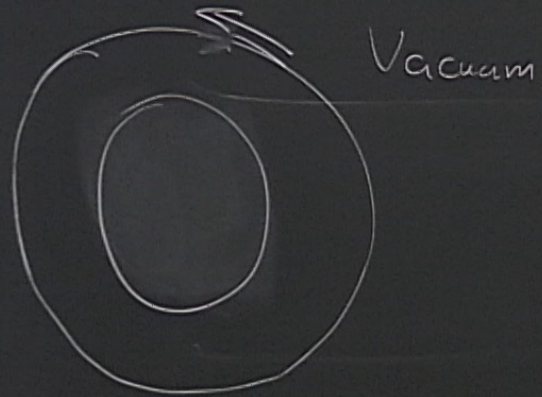
Real space:  $\chi(x) = \int_{-\infty}^{\infty} \frac{dk}{2\pi} e^{ikx} \chi(k) = \int_0^{\infty} \frac{dk}{2\pi} [e^{ikx} \chi_k + e^{-ikx} \chi_k^\dagger] : \text{Hermitian}$

$$\chi(x) = \chi^\dagger(x) \leftarrow \text{Majorana fermion (real)}$$

$$H = \chi i \partial_x \chi$$



Edge state:



Chiral fermion, but

$$H = k_x \chi_R^\dagger \chi_R$$

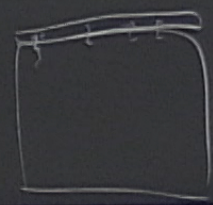
Real space:

$$\chi(x) = \chi$$

$$H = \chi i \partial_x \chi$$



$$H = \chi i \partial_x \chi$$

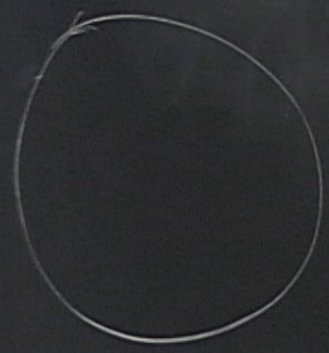


$k_x = 0$  ?? If this mode exist  $\Rightarrow$  Majorana zero mode

$\chi_0 = \chi_0^+$  should not appear alone,

Edge has anti-periodic b.c. ( $\chi(L) = -\chi(0)$ )

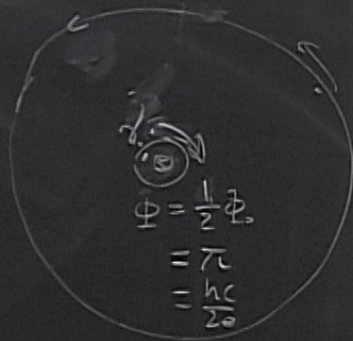
$$k_x = \frac{(2n+1)\pi}{L}$$



What



What happens with a vortex



Aharonov-Bohm phase

seen by edge fermion =  $\Phi = \pi$ .

Edge fermion has periodic BC.

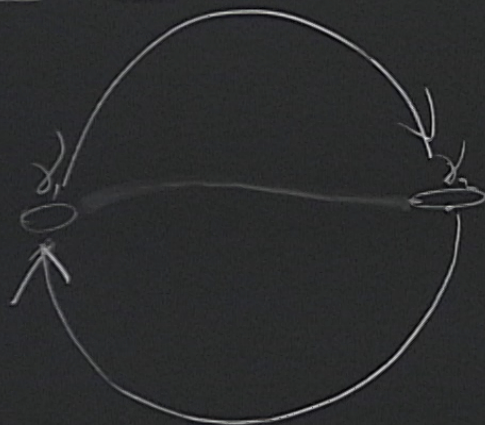
$\chi_{k=0}$ : a MZM on edge.

$\Rightarrow$  Another MZM should appear.

$\Rightarrow$  A  $\pi$ -vortex traps an unpaired MZM.

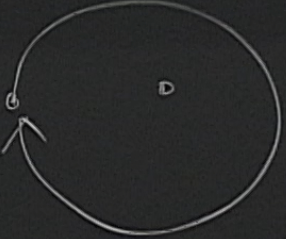


Consider 2 vertices (far away from each other)



$$2\text{-dim } \mathcal{H} = (|\Omega\rangle, \tilde{c}^\dagger |\Omega\rangle)$$
$$\hat{c} = \frac{\gamma_1 + i\gamma_2}{2}, \quad \hat{c} |\Omega\rangle = 0$$



Exchange twice  $\approx$    $\mathbb{Z}d$ : has to be trivial.  
 $\mathbb{Z}d$ : non-contractable.  
 $\Rightarrow$  Exchange doesn't have to give  $\pm 1$ .

If exchange:  $e^{i\theta} \neq \pm 1 \Rightarrow$  Abelian anyon.

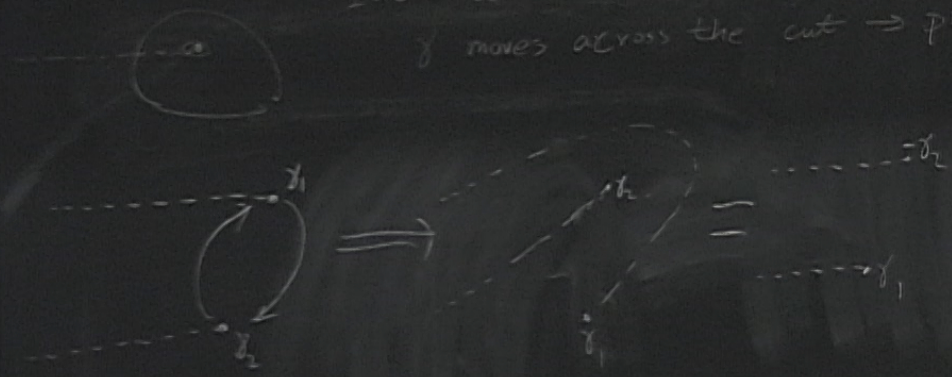
If degeneracy: Matrix  $\Rightarrow$  non-Abelian anyon.



$\Psi_k = U \Psi_{-k}^*$  algebra requires  $U^\dagger H_k U = -H_{-k} \Rightarrow U = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \Rightarrow E_k = \begin{pmatrix} - & \\ & - \end{pmatrix}$  If a single mode at

$2 \times 2$  Unitary matrix.

Introduce a branch cut associated with each vortex.  $(\text{Exchange})^2 = \begin{matrix} \gamma_1 \rightarrow -\gamma_1 \\ \gamma_2 \rightarrow -\gamma_2 \end{matrix}$   
 $\gamma$  moves across the cut  $\rightarrow$  picks up  $(-1)$ .



$E \quad \gamma_1 \rightarrow -\gamma_2$   
 $\Rightarrow \text{Exchange: } \gamma_2 \rightarrow \gamma_1$

$E = \begin{pmatrix} \gamma_1 & \gamma_2 \\ \gamma_2 & -\gamma_1 \end{pmatrix}$   
 $\gamma_1 \gamma_2 \cdot \gamma_1 = -\gamma_2$   
 $\gamma_1 \gamma_2 \cdot \gamma_2 = \gamma_1$



$$\alpha|\Omega\rangle + \beta\tilde{c}^+|\Omega\rangle \xrightarrow{e^{i\theta}} (\alpha|\Omega\rangle - i\beta\tilde{c}^+|\Omega\rangle)$$

$$c = \frac{\gamma + i\delta}{2}$$

as long as  $\alpha, \beta \neq 0$

↓ E.

$$\frac{-\gamma - i\delta}{2} = -i c^+$$



$$\alpha|\Omega\rangle + \beta\tilde{c}^+|\Omega\rangle \rightarrow e^{i\theta}(\alpha|\Omega\rangle - i\beta\tilde{c}^+|\Omega\rangle)$$

$$c = \frac{\gamma_1 - i\gamma_2}{2}$$

as long as  $\alpha, \beta \neq 0$

$$2^{\binom{N}{2}}$$

$$\downarrow \sqrt{E}$$

$$\frac{-\gamma_2 - i\gamma_1}{2} = -i c^*$$

