Title: Classical algorithms for quantum mean values

Speakers: David Gosset

Collection: Symmetry, Phases of Matter, and Resources in Quantum Computing

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Abstract: Consider the task of estimating the expectation value of an n-qubit tensor product observable in the output state of a shallow quantum circuit. This task is a cornerstone of variational quantum algorithms for optimization, machine learning, and the simulation of quantum many-body systems. In this talk I will describe three special cases of this problem which are "easy" for classical computers. This is joint work with Sergey Bravyi and Ramis Movassagh.

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Classical algorithms for quantum mean values

Sergey Bravyi David Gosset Ramis Movassagh

arXiv:1909.11485





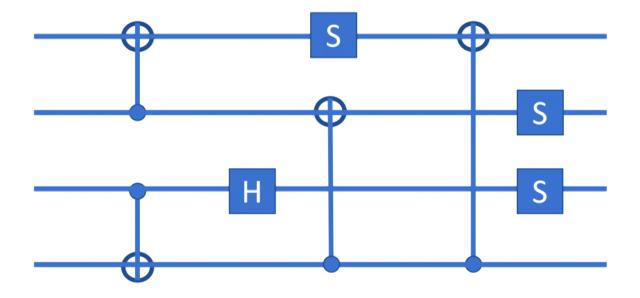




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Circuit depth

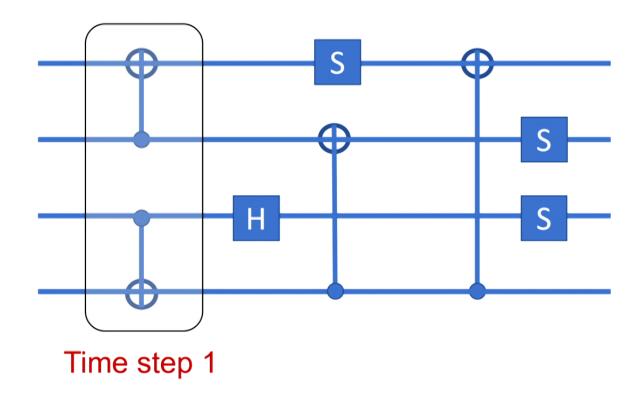
Circuit depth is the number of time steps allowing for parallel gates.



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Circuit depth

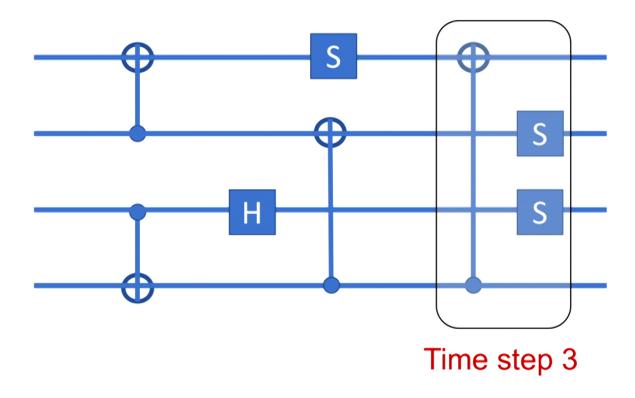
Circuit depth is the number of time steps allowing for parallel gates.



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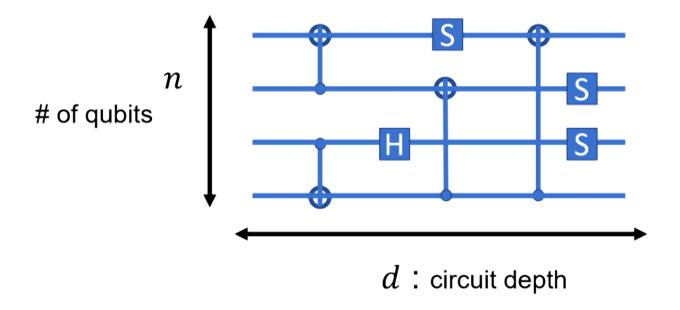
Circuit depth

Circuit depth is the number of time steps allowing for parallel gates.



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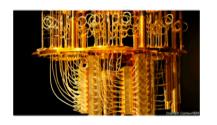
Shallow quantum circuits



We are interested in circuits with depth d = O(1).

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Why study shallow quantum circuits?



Small quantum computers: lack of error correction places limits on circuit size. So look at either few qubits (uninteresting) or low circuit depth.



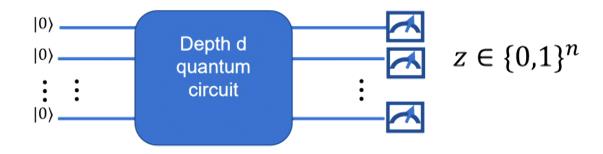
Simplicity: a restricted model of quantum computation with structure that can be exploited.



Computational power...

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What are shallow quantum circuits good for?

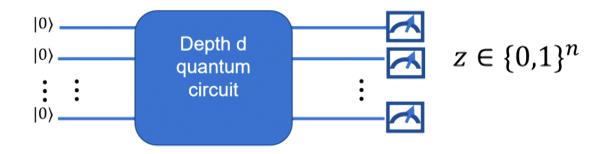


Sample from classically inaccessible probability distributions

[Terhal Divincenzo 2002] [Gao et al 17] [Bermejo-Vega et al. 17]

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What are shallow quantum circuits good for?

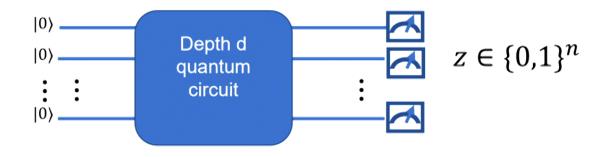


Solve certain linear algebra problems faster than classical algorithms

[Bravyi, G., Koenig 18] [Bene Watts, Kothari, Schaeffer, Tal 19] [Bravyi, G., Koenig, Tomamichel 19]

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What are shallow quantum circuits good for?



...Anything else?

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A recent family of "near-term" algorithms which has attracted great interest:

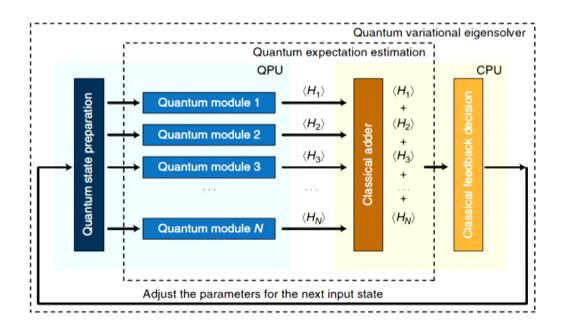


Image depicts the Variational Quantum Eigensolver paper, taken from [Peruzzo et al. 2013] ...

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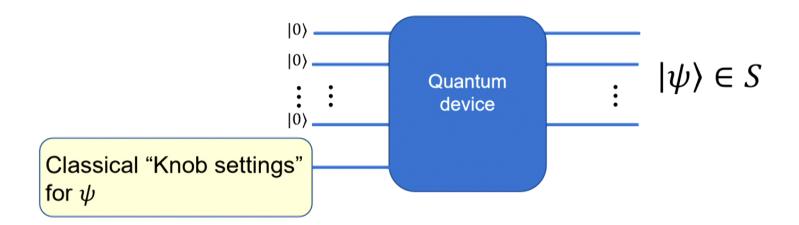
Goal: find the minimum energy of a given Hamiltonian.

$$H = \sum_{i} P_{i} \qquad E_{min} = \min_{\psi} \langle \psi | H | \psi \rangle$$

We will be interested in the case where each term is an n-qubit Pauli operator

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Mild assumption #1: your quantum device can prepare a subset S of n-qubit states ψ



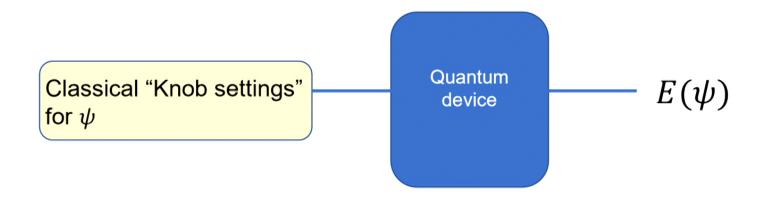
We will be interested in the case where *S* consists of states that can be prepared by constant-depth quantum circuits.

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Mild assumption #2: The device can be used to measure the energy of a given state $\psi \in S$

$$E(\psi) = \langle \psi | H | \psi \rangle = \sum \langle \psi | P_i | \psi \rangle$$

This can be achieved by computing each mean value $\langle \psi | P_i | \psi \rangle$ separately and then summing them.

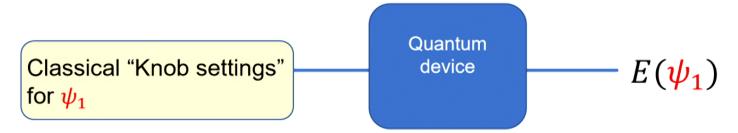


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A variational algorithm aims to compute the minimum energy **over states in** *S*

$$\min_{\psi \in S} \langle \psi | H | \psi \rangle$$

The algorithm uses the quantum device to compute energies and a classical computer to choose the knob settings:

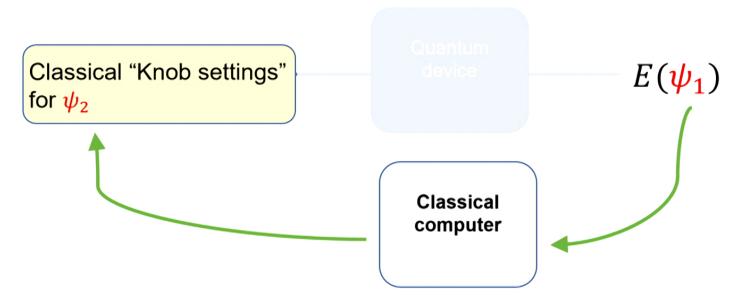


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A variational algorithm aims to compute the minimum energy **over states in** *S*

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The algorithm uses the quantum device to compute energies and a classical computer to choose the knob settings:



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What are variational quantum algorithms good for?

International Workshop on Quantum Technology and Optimization Problems QTOP 2019: Quantum Technology and Optimization Problems pp 74-85 | Cite as Variational Quantum Factoring Authors Authors and affiliations Eric Anschuetz, Jonathan Olson, Alán Aspuru-Guzik, Yudong Cao 🖂 PHYSICAL REVIEW LETTERS Quantum Machine Learning in Feature Hilbert Spaces Maria Schuld and Nathan Killoran Phys. Rev. Lett. 122, 040504 - Published 1 February 2019 nature Letter Published: 13 September 2017 Hardware-efficient variational quantum eigensolver for small molecules and quantum magnets Abhinav Kandala 🖷, Antonio Mezzacapo 🔍 Kristan Temme, Maika Takita, Markus Brink, Jerry M. Chow & Jay M. Gambetta

> nature International journal of science

Jerry M. Chow & Jay M. Gambetta

Supervised learning with quantum-

Vojtěch Havlíček, Antonio D. Córcoles ⁵⁰, Kristan Temme ⁵⁰, Aram W. Harrow, Abhinav Kandala

enhanced feature spaces

Lack of performance guarantees

Unfortunately, variational algorithms don't have performance guarantees as they are challenging to analyze:

Challenge #1: Is $\min_{\psi \in S} \langle \psi | H | \psi \rangle$ close to $\min_{\psi} \langle \psi | H | \psi \rangle$?

Challenge #2: Is the output of the algorithm close to $\min_{\psi \in S} \langle \psi | H | \psi \rangle$?

Do these algorithms really have any algorithmic speedup over classical computers...

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Do we really need a quantum computer?

The quantum computer is only used to compute mean values of observables at the output of a quantum computation

$$\langle 0^n | U^{\dagger} O U | 0^n \rangle$$

How hard is this problem? Could we use a classical computer instead?

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Let *U* be a depth d = O(1) quantum circuit.

Let *O* be a tensor product of single-qubit Hermitian operators

$$O = O_1 \otimes O_2 \otimes \cdots \otimes O_n$$

Assume $||0_i|| \le 1$

We are interested in estimating the mean value

$$\mu = \langle 0^n \big| U^{\dagger} O \ U \big| 0^n \rangle$$

$$\mu = \langle 0^n \big| U^{\dagger} O \ U \big| 0^n \rangle$$

Interesting special case:

$$0 = |x_1\rangle\langle x_1| \otimes |x_2\rangle\langle x_2| \otimes \cdots \otimes |x_n\rangle\langle x_n|$$

Then the mean value is an output probability of the quantum circuit

$$\mu = |\langle x|U|0^n\rangle|^2$$

$$O = O_1 \otimes O_2 \otimes \cdots \otimes O_n$$

$$\mu = \langle 0^n \big| U^{\dagger} O \ U \big| 0^n \rangle$$

Additive error mean value problem

Given $\epsilon = \frac{1}{poly(n)}$, compute an estimate $\tilde{\mu}$ such that

$$|\tilde{\mu} - \mu| < \epsilon$$

The additive error mean value problem can be solved efficiently on a quantum computer.

$$O = O_1 \otimes O_2 \otimes \cdots \otimes O_n$$

$$\mu = \langle 0^n \big| U^{\dagger} O \ U \big| 0^n \rangle$$

Relative error mean value problem

Given $\epsilon = \frac{1}{poly(n)}$, compute an estimate $\tilde{\mu}$ such that

$$|\tilde{\mu} - \mu| < \epsilon \mu$$

The relative error mean value problem is #P-hard.

Complexity of the mean value problem

Quantum circuit U	Observables O_j	Relative error	Additive error
Polynomial size	Pos. semidefinite	#P-hard [16]	BQP-complete
Constant depth	Close to I	P [Thm. 1]	P [Thm. 1]
Constant depth	Pos. semidefinite	#P-hard [15, 16]	BQP Subexp. classical [Thm. 4]
2D Constant depth	Hermitian	#P-hard [15, 16] Subexp. classical [17]	BPP [Thm. 5]

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Complexity of the mean value problem

Quantum circuit U	Observables O_j	Relative error	Additive error
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2D Constant depth	Hermitian	#P-hard [15, 16] Subexp. classical [17]	BPP [Thm. 5]

In the rest of the talk I will describe these 3 classical simulation algorithms...

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Case 1: Single-qubit observables are each close to the identity

Quantum circuit ${\cal U}$	Observables O_j	Relative error	Additive error
Polynomial size	Pos. semidefinite	#P-hard [16]	BQP-complete
Constant depth	Close to I	P [Thm. 1]	P [Thm. 1]
Constant depth	Pos. semidefinite	#P-hard [15, 16]	BQP Subexp. classical [Thm. 4]
2D Constant depth	Hermitian	#P-hard [15, 16] Subexp. classical [17]	BPP [Thm. 5]

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Restricted family of tensor product observables

Suppose U is a depth-d quantum circuit and consider an observable

$$O = O_1 \otimes O_2 \otimes \cdots \otimes O_n$$

where

$$\left\|O_j-I\right\|\leq \frac{0.001}{2^{5d}}$$

Closeness to identity depends only on the depth d

For 2D circuits we can replace the RHS with $O(d^{-4})$

In this part of the talk we will be interested in obtaining a (highly demanding) relative error approximation to the mean value

$$\mu = \langle 0^n \big| U^{\dagger} O \ U \big| 0^n \rangle.$$

Restricted family of tensor product observables

$$\mu = \langle 0^n \big| U^{\dagger} O \ U \big| 0^n \rangle$$

$$\mu = \langle 0^n \big| U^{\dagger} O \ U \big| 0^n \rangle \qquad O = O_1 \otimes O_2 \otimes \cdots \otimes O_n$$

$$||O_j - I|| \le \frac{0.001}{2^{5d}}$$

Example

Suppose we consider an output probability of a noisy quantum circuit

$$\mu' = \langle 0^n | \mathcal{E}^{\otimes n} (U^{\dagger} | 0) \langle 0 |^n \rangle U | 0^n \rangle.$$

$$\mathcal{E}(\rho) = (1-p)\rho + pX\rho X$$
 Flip each bit with probability p

The noisy mean value is proportional to an ideal mean value:

$$\mu' = \frac{1}{2^n}\mu$$
 with single-qubit observables $O_j = I + (1 - 2p)Z$

The above restriction is satisfied in a high noise regime $p \ge \frac{1}{2} - O(2^{-5d})$

Main result

$$\mu = \langle 0^n | U^{\dagger} O \ U | 0^n \rangle$$

$$O = O_1 \otimes O_2 \otimes \cdots \otimes O_n$$

$$||O_j - I|| \le \frac{0.001}{2^{5d}}$$

Theorem

Let $\delta \in \left(0, \frac{1}{2}\right)$ be given. There is a deterministic classical algorithm which outputs an estimate $\tilde{\mu}$ satisfying

$$|\log(\tilde{\mu}) - \log(\mu)| < \delta$$

The runtime of the algorithm is $(n\delta^{-1})^{c \cdot 2^d}$.

Solves the relative error mean value problem for this restricted set of observables Runtime can be improved for 2D geometrically local circuits

The algorithm is based on a polynomial interpolation method due to Barvinok...

Define a polynomial

$$f(\epsilon) = \langle 0^n | U^{\dagger} O(\epsilon) U | 0^n \rangle \qquad O(\epsilon) = O_1(\epsilon) \otimes O_2(\epsilon) \otimes \cdots \otimes O_n(\epsilon)$$
$$O_j(\epsilon) = (1 - \epsilon)I + \epsilon O$$

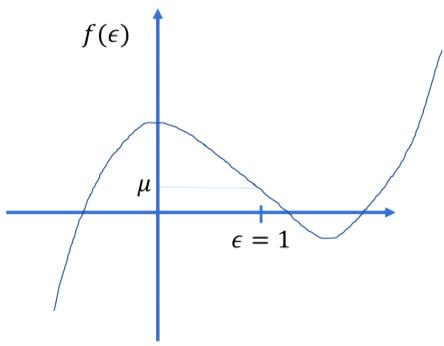
Note that f(0) = 1 and we aim to compute $\mu = f(1)$

Also note that derivatives $f^{(k)}(0)$ can be computed efficiently for small k

e.g.,
$$f^{(1)}(0) = \sum_{j=1}^{n} \langle 0^{n} | U^{\dagger}(O_{j} - I)U | 0^{n} \rangle$$

Acts nontrivially on $\leq 2^d$ qubits

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Since we know the function value and can compute derivatives at $\epsilon = 0$, it is natural to try to use a Taylor series approximation.

Barvinok: use Taylor series for the function $g(\epsilon) = \log(f(\epsilon))$ instead...

Approximate the log by its truncated Taylor series

$$g(\epsilon) = \log f(\epsilon)$$
 We want to compute $g(1)$

$$T_p(\epsilon) = g(0) + \sum_{k=1}^p \frac{\epsilon^k}{k!} g^{(k)}(0)$$

Theorem [Barvinok]

If the polynomial $f(\epsilon)$ is zero-free on the disk $|\epsilon| \leq 2$ then

$$|T_p(\epsilon) - g(\epsilon)| \le \frac{n}{(p+1)2^p}$$
 $|\epsilon| \le 1$

To achieve error δ we need only take $p = O(\log(n\delta^{-1}))$

$$g(\epsilon) = \log f(\epsilon)$$

To use Barvinok's method we need two ingredients:

1) We need to compute derivatives

$$g^{(1)}(0), ..., g^{(p)}(0)$$
 $p = O(\log(n\delta^{-1}))$

These can be computed efficiently from the derivatives $f^{(1)}(0), ..., f^{(k)}(0)$.

2) We need to show that $f(\epsilon)$ is zero-free on the disk $|\epsilon| \leq 2 \dots$

Zero-free region

$$f(\epsilon) = \langle 0^n | U^{\dagger} O(\epsilon) U | 0^n \rangle \qquad O(\epsilon) = O_1(\epsilon) \otimes O_2(\epsilon) \otimes \cdots \otimes O_n(\epsilon)$$
$$O_i(\epsilon) = (1 - \epsilon)I + \epsilon O$$

Theorem

Suppose $||O_i - I|| \le \gamma$. The polynomial f has no zeros in the disk

$$|\epsilon| \le \frac{0.001}{\gamma 2^{5d}}$$
 Depth of U

Choosing $\gamma = 0.001 \cdot 2^{-5d-1}$ suffices to make the disk radius equal to 2.

Proof sketch (zero-free region)

$$f(\epsilon) = \langle 0^n | U^{\dagger} O_1(\epsilon) \otimes \cdots O_n(\epsilon) U | 0^n \rangle$$

Write each 2×2 operator $O_j(\epsilon)$ as the upper left block of a 4×4 unitary $B_j(\epsilon)$

$$f(\epsilon) = \langle 0^{2n} \big| (U^{\dagger} \otimes I) B_1(\epsilon) \otimes \cdots B_n(\epsilon) (U \otimes I) \big| 0^{2n} \rangle$$

Define
$$V_j(\epsilon) = (U^\dagger \otimes I)B_j(\epsilon)(U \otimes I)$$
 The $V_j(\epsilon)$ each act on 2^{d+1} qubits

Then
$$f(\epsilon) = \langle 0^{2n} | V_1(\epsilon) V_2(\epsilon) \dots V_n(\epsilon) | 0^{2n} \rangle$$
 A constant depth circuit Each gate is close to identity

Proof sketch (zero-free region)

$$f(\epsilon) = \langle 0^{2n} | V(\epsilon) | 0^{2n} \rangle \qquad V(\epsilon) = V_1(\epsilon) V_2(\epsilon) \dots V_n(\epsilon) \qquad \begin{array}{l} \text{A constant-depth circuit} \\ \text{Each gate is close to identity} \end{array}$$

Now consider a probability distribution over 2n-bit strings defined by

$$p_{\epsilon}(z) = \left| \left\langle z | V(\epsilon) | 0^{2n} \right\rangle \right|^2$$

Our goal is to show that $p_{\epsilon}(0^{2n}) > 0$ for all ϵ in the disk...

Proof sketch (zero-free region)

$$p_{\epsilon}(z) = \left| \left\langle z | V(\epsilon) | 0^{2n} \right\rangle \right|^2$$

Let E_j be the event that the jth bit is 1. We show that each event $\{E_j\}_{1 \le j \le 2n}$ occurs with a small probability $q = O(2^d \gamma |\epsilon|)$ and is independent of most $D = O(2^{4d})$ of the others...

$$\Pr_{p_{\epsilon}}[E_{j}] = \langle 0^{2n} | V(\epsilon)^{\dagger} | 1 \rangle \langle 1 |_{j} V(\epsilon) | 0^{2n} \rangle = \langle 0^{2n} | A_{j} | 0^{2n} \rangle$$
All gates in $V(\epsilon)$ except $O(2^{d})$ of them can be cancelled here.

Supported on $O(2^{2d})$ qubits

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Proof sketch (zero-free region)

$$p_{\epsilon}(z) = \left| \left\langle z | V(\epsilon) | 0^{2n} \right\rangle \right|^{2}$$

Let E_j be the event that the jth bit is 1. We show that each event $\{E_j\}_{1 \le j \le 2n}$ occurs with a small probability $q = O(2^d \gamma |\epsilon|)$ and is independent of most $D = O(2^{4d})$ of the others...

The Lovasz Local Lemma then implies $p_{\epsilon}(0^{2n}) > 0$ as long as

$$\exp(1) \cdot q \cdot D < 1$$
 $|\epsilon| < O(1) \cdot 2^{-5d} \gamma^{-1}$

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Can the bound on zero-free radius be improved?

In the worst case it is possible for the zero-free radius to be exponentially small in the depth:

$$|\psi\rangle = \frac{1}{\sqrt{2}} \left(|0\rangle^{2^d} + |1\rangle^{2^d} \right) = U|0\rangle^{2^d}$$
Depth d

$$O_{j}(\epsilon) = I + \epsilon Z_{j}$$
 $f(\epsilon) = \langle \psi | O_{1}(\epsilon) \otimes \cdots \otimes O_{2^{d}}(\epsilon) | \psi \rangle$
$$= \frac{1}{2} \Big((1 + \epsilon)^{2^{d}} + (1 - \epsilon)^{2^{d}} \Big)$$
 Has a root at
$$\epsilon_{0} \approx \frac{i\pi}{2^{d+1}}$$

For random circuits the zero free radius is typically much larger...

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Zero-free radius for random unitaries

Consider observables diagonal in the Z-basis:

$$O_j(\epsilon) = I + \epsilon Z_j$$

$$O_j(\epsilon) = I + \epsilon Z_j \qquad f(\epsilon) = \langle 0^n | U^{\dagger} O_1(\epsilon) \otimes \cdots \otimes O_n(\epsilon) U | 0^n \rangle$$

Theorem

Let *U* be a random quantum circuit drawn from a unitary 2-design The polynomial f has no zeros in a disk

$$|\epsilon| \le 1 - O(n^{-1}\log(n))$$

Proof idea

Write
$$f(\epsilon) = 1 + \sum_{k=1}^{n} c_k \epsilon^k$$
 2-design property gives $\mathbb{E}[|c_k|^2] \le \frac{1}{2^n} \binom{n}{k}$

Use to show that w.h.p for ϵ in the disk we have $|f(\epsilon) - 1| \le 1/2$

Case 2: Positive semidefinite observables

Quantum circuit U	Observables O_j	Relative error	Additive error
Polynomial size	Pos. semidefinite	#P-hard [16]	BQP-complete
Constant depth	Close to I	P [Thm. 1]	P [Thm. 1]
Constant depth	Pos. semidefinite	#P-hard [15, 16]	BQP Subexp. classical [Thm. 4]
2D Constant depth	Hermitian	#P-hard [15, 16] Subexp. classical [17]	BPP [Thm. 5]

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Subexponential time classical algorithm

$$O = O_1 \otimes O_2 \otimes \cdots \otimes O_n$$

$$||O_j||=1$$

Theorem

Let $\delta \in \left(0, \frac{1}{2}\right)$ be given. There is a deterministic classical algorithm which outputs an estimate $\tilde{\mu}$ satisfying

$$\left|\tilde{\mu} - |\langle 0^n | U^{\dagger} O U | 0^n \rangle|\right| < \delta$$

The runtime of the algorithm is $e^{\tilde{O}(4^d \sqrt{n \cdot \log(\delta^{-1})})}$.

In general, the algorithm estimates the absolute value of the mean.

Solves the additive error MVP for pos. semidefinite observables.

Case 3: 2D shallow circuits

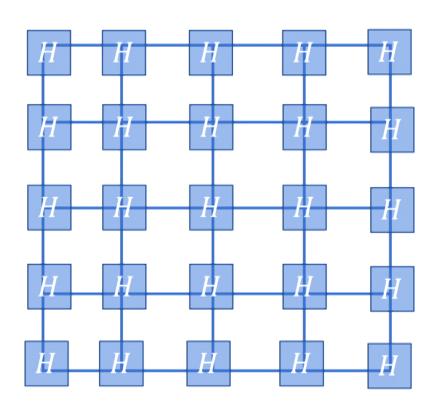
Quantum circuit U	Observables O_j	Relative error	Additive error
Polynomial size	Pos. semidefinite	#P-hard [16]	BQP-complete
Constant depth	Close to I	P [Thm. 1]	P [Thm. 1]
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2D Constant depth	Hermitian	#P-hard [15, 16] Subexp. classical [17]	BPP [Thm. 5]

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Suppose the qubits are located at the vertices of a 2D grid, and U is a depth d quantum circuit where each gate acts between nearest-neighbors.

Example:

$$U = \left(\prod_{(i,j)\in E} CZ_{ij}\right) H^{\otimes n} |0^n\rangle$$

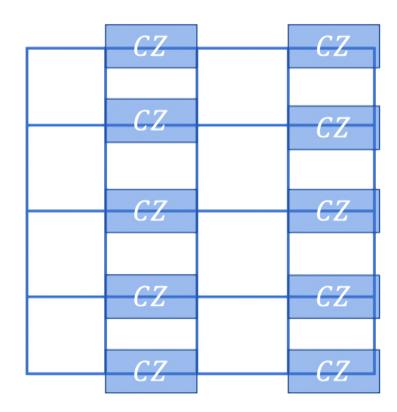


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Suppose the qubits are located at the vertices of a 2D grid, and U is a depth d quantum circuit where each gate acts between nearest-neighbors.

Example:

$$U = \left(\prod_{(i,j)\in E} CZ_{ij}\right) H^{\otimes n} |0^n\rangle$$

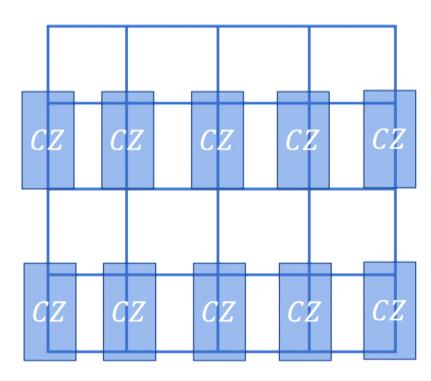


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Suppose the qubits are located at the vertices of a 2D grid, and U is a depth d quantum circuit where each gate acts between nearest-neighbors.

Example:

$$U = \left(\prod_{(i,j)\in E} CZ_{ij}\right) H^{\otimes n} |0^n\rangle$$



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Suppose the qubits are located at the vertices of a 2D grid, and U is a depth d quantum circuit where each gate acts between nearest-neighbors.

$$O = O_1 \otimes O_2 \otimes \cdots \otimes O_n \qquad \qquad \mu = \langle 0^n | U^{\dagger} O \ U | 0^n \rangle$$

Theorem

Let $\delta \in \left(0, \frac{1}{2}\right)$ be given. There is a randomized classical algorithm which, with probability at least 2/3, outputs an estimate $\tilde{\mu}$ satisfying

$$|\mu - \tilde{\mu}| \le \delta$$

The runtime is $O(n\delta^{-2}2^{O(d^2)})$. Linear time!

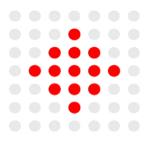
Algorithm is based on an MPS representation and Monte Carlo method...

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Express mean value as amplitude of a 2D constant depth circuit with commuting gates

$$\mu = \langle 0^n | U^{\dagger} O_1 \otimes O_2 \otimes \dots \otimes O_n U | 0^n \rangle$$
$$= \langle 0^n | Q_n Q_{n-1} \dots Q_1 | 0^n \rangle \qquad Q_n = U^{\dagger} O_j \mathsf{U}$$

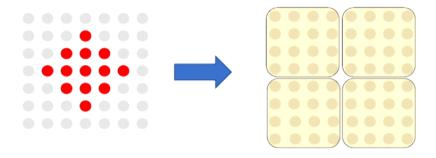
Each gate Q_i is supported on a $2d \times 2d$ square region centred at qubit j



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$$\mu = \langle 0^n | Q_n Q_{n-1} \dots Q_1 | 0^n \rangle$$

Coarse-grain: group the qubits into supersites of size $2d \times 2d$

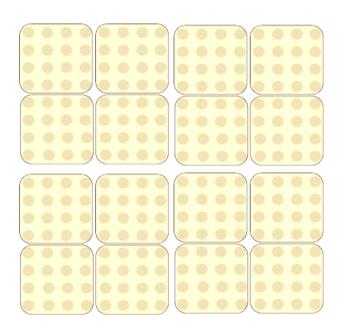


Each gate now acts nontrivially on 1 plaquette consisting of 4 supersites

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Express mean value as inner product between two Matrix Product states

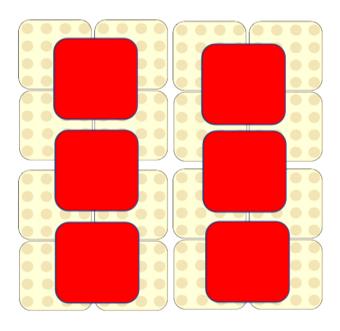
$$\mu = \langle 0^n | Q_n Q_{n-1} \dots Q_1 | 0^n \rangle$$
$$= \langle \Phi_{even} | \Phi_{odd} \rangle$$



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Express mean value as inner product between two Matrix Product states

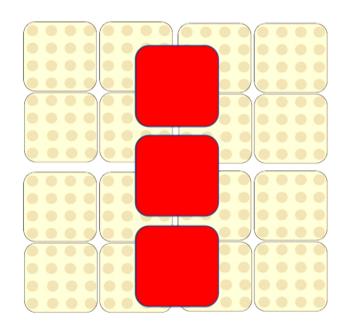
$$\mu = \langle 0^n | Q_n Q_{n-1} \dots Q_1 | 0^n \rangle$$
$$= \langle \Phi_{even} | \Phi_{odd} \rangle$$



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Inner product between MPS can be estimated in polynomial time using a Monte Carlo method [Van den Nest 2009]

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Open problems

Big question: what is the complexity of the additive-error mean value problem for constant-depth circuits?

Can the subexponential-time algorithm be generalized to the case of observables which may not be positive semidefinite?

Can the 2D algorithm be generalized to higher dimensional lattices?

Other applications of the zero-free region for quantum circuits?

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