

Title: Modeling the pseudogap state in cuprates: quantum disordered pair density wave

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Abstract: I will briefly review the pseudogap phenomenology in high T<sub>c</sub> cuprate superconductor, especially the recent experiments, and propose a unified picture of the phenomenology under only one assumption: the fluctuating pair density wave. By quantum disordering a pair density wave, we found a state composed of insulating antinodal pairs and nodal electron pocket. We compare the theoretical predictions with ARPES results, optical conductivity, quantum oscillation and other experiments.

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# Modeling the pseudogap state in cuprates: quantum disordered pair density wave



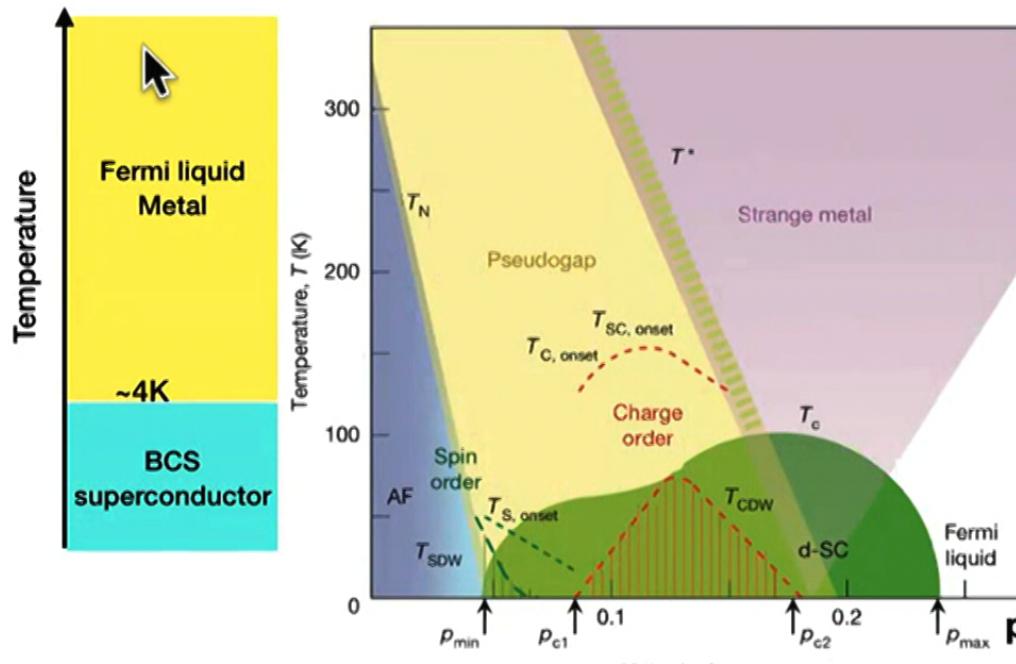
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Z. Dai, T. Senthil, P. A. Lee, arXiv: 1906.01656



**Massachusetts  
Institute of  
Technology**

## BCS superconductor      Cuprate high Tc superconductor



- **Superconductor:**  $T_c \sim 100\text{K}$
- **Strange metal:** large Fermi surface, linear  $T$  resistivity??
- **Pseudogap:** electron gap on part of the Fermi surface, Fermi arc??

Doping dependence:  
**pseudogap** is an intermediate phase between the Mott insulator at half filling, and the Fermi liquid at larger doping

Goal: use a single assumption to explain the pseudogap phenomenology at low temperature

# Outline

- **Brief review of the pseudogap**

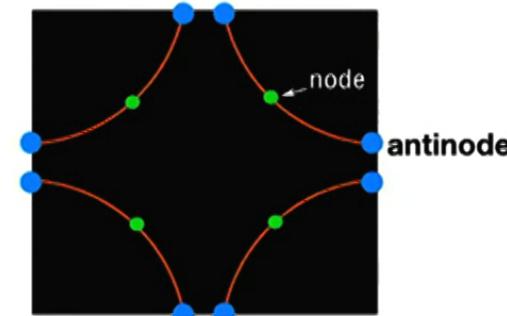
- Review of pair density wave (PDW) order
- Quantum disordered PDW: effective theory of pseudogap

*Claim: small electron pocket + insulator of electron pairs*

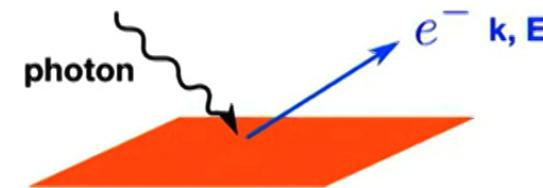
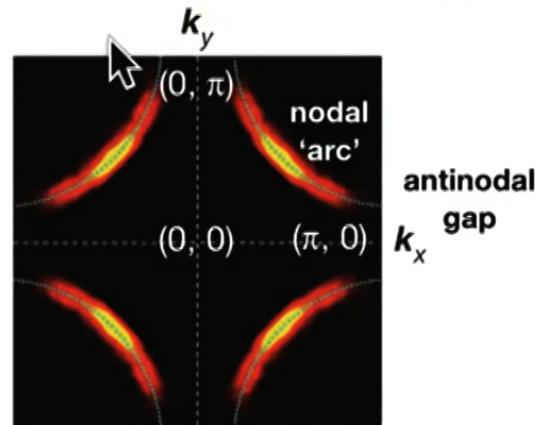
- Electron spectrum of fluctuating superconductor
- Comparison with experimental results

# Pseudogap: antinodal Gap seen in ARPES

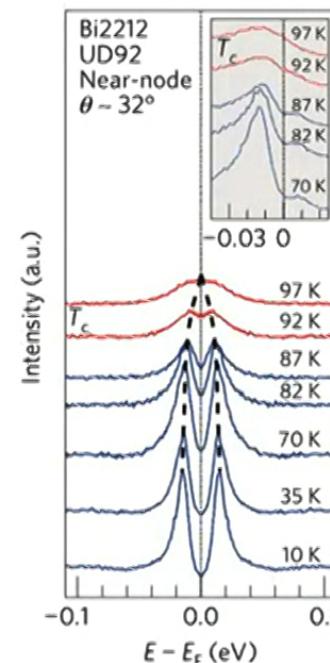
high temperature: large Fermi surface  
d wave superconductor: gapless node



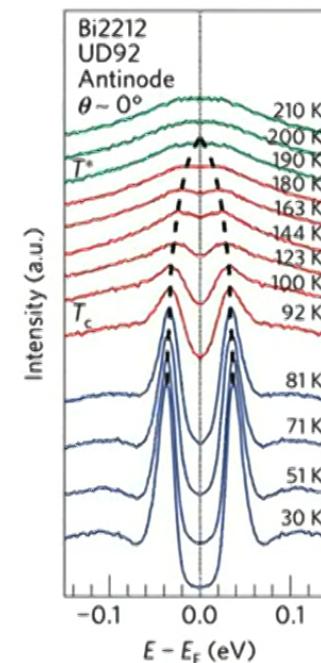
$T < T^*$  : antinodal electron gap



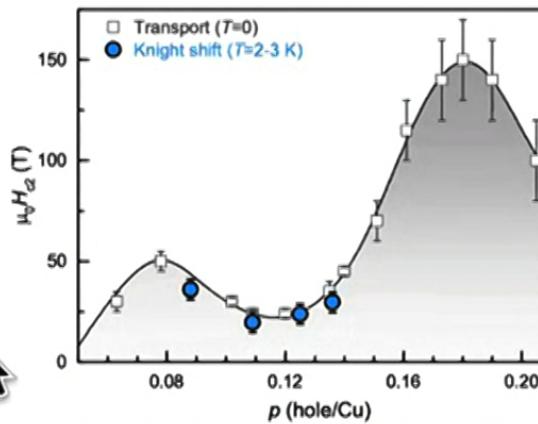
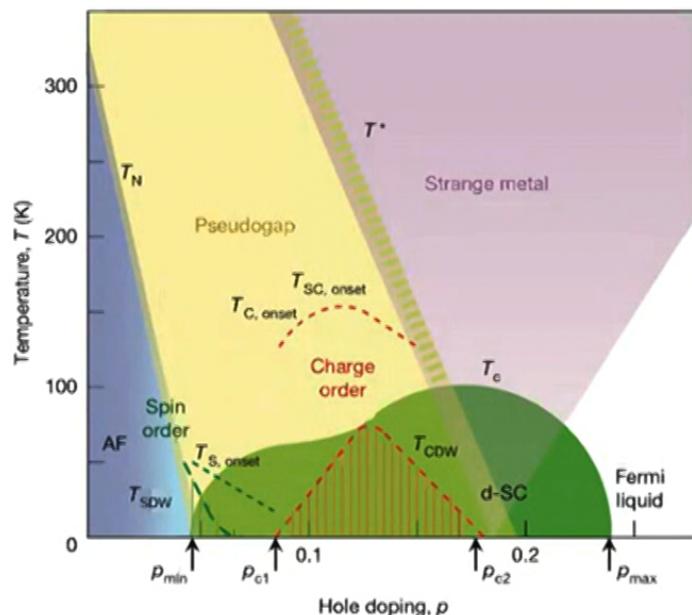
## Near-node



## Antinode



# High-field metallic pseudogap ground state



Zhou, Rui, et al. PNAS 114.50 (2017): 13148-13153.

At high field and the lowest temperature,  
quantum oscillation from a small electron  
Fermi surface, 2% of the original BZ

Fermi Liquid with small Fermi surface?

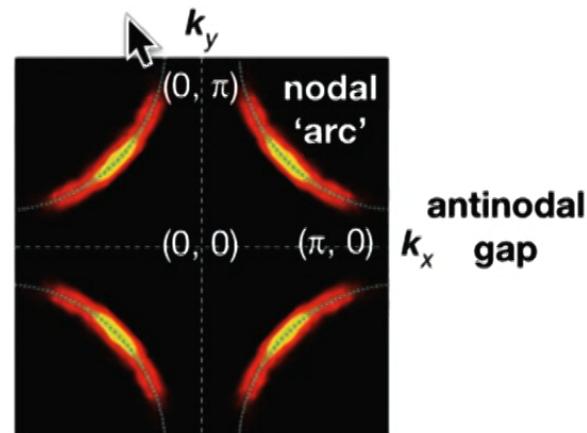
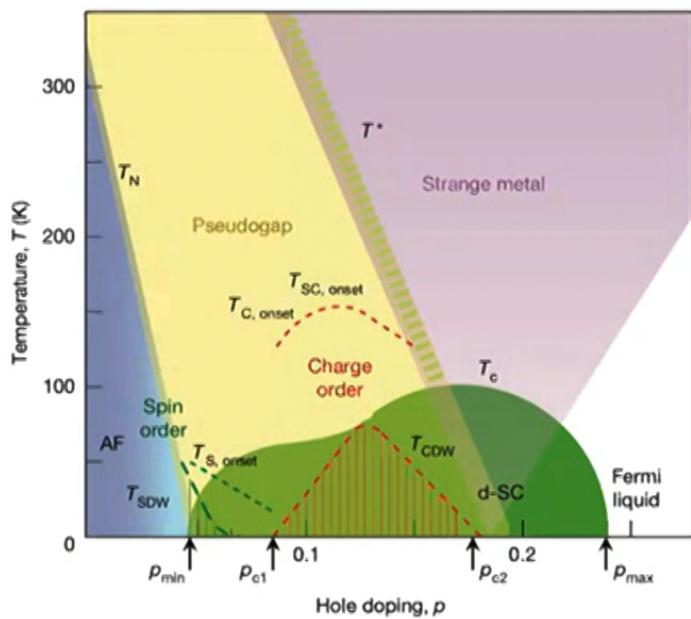
Our goal: to understand pseudogap at the  
lowest temperature,  $H > H_{c2}$ , or  $T_c \ll T \ll T^*$

Assumption: quantum fluctuating PDW

Z. Dai, T. Senthil, P. A. Lee, arXiv: 1906.01656

### **Previous theoretical approach**

- Charge density wave (CDW) alone cannot explain the pseudogap
  - Pseudogap was traditionally explained by fluctuating d wave SC



**Fluctuating SC can explain the decrease of density of states, but is difficult to explain the nodal arcs**

# Fluctuating pair density wave as a mother state

- Use the idea of fluctuating superconductivity to get a metal with pseudogap, but consider a different pairing
- PDW: electron pair condense at nonzero momentum

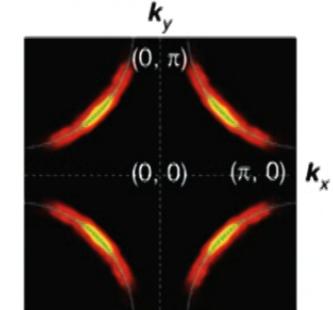
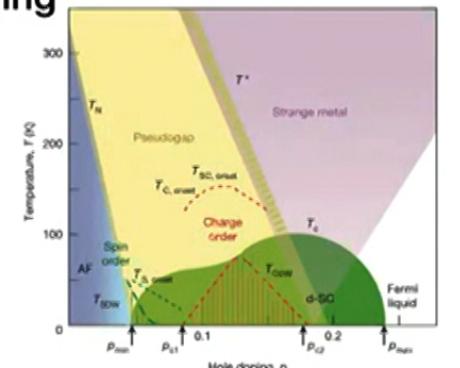
$$\Delta_{\vec{P}} = \langle c_{\vec{k},\uparrow} c_{-\vec{k}+\vec{P},\downarrow} \rangle \neq 0$$

- Bidirectional PDW, wave vector half of the CDW

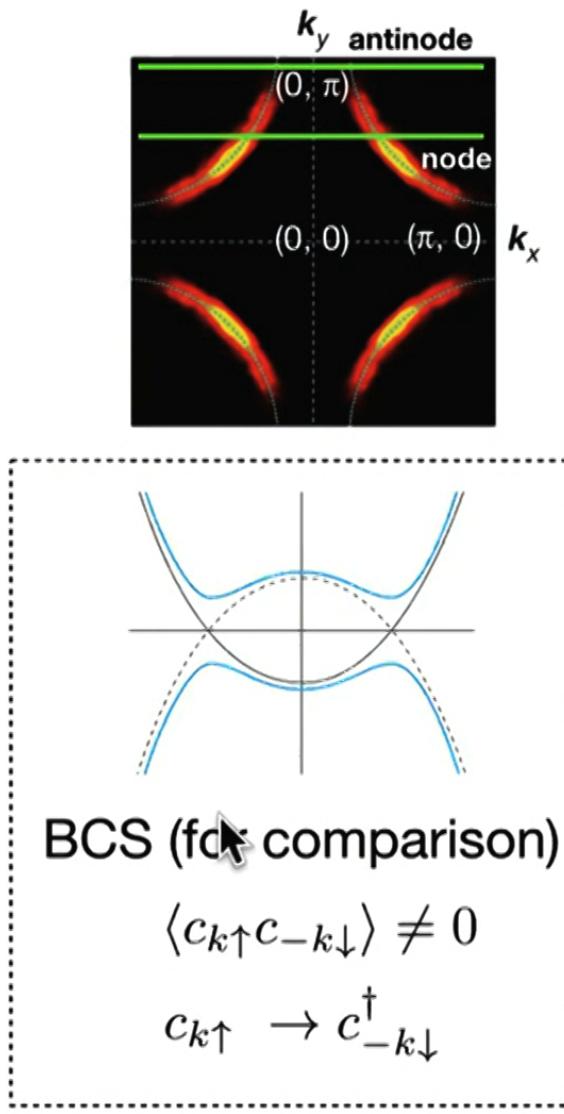
$$\vec{P} = P\hat{x}, -P\hat{x}, P\hat{y}, -P\hat{y}$$

Generate CDW at momentum  $\vec{Q}_{\text{CDW}} = 2\vec{P}$

- Pair density wave (PDW) has a number of advantages
  1. Gap out only antinodal electrons, leaves gapless ‘Fermi arc’
  2. Explains ARPES in detail
  3. Generate secondary CDW at wave vector  $\vec{Q} = 2\vec{P}$
  4. Explain the quantum oscillation at high field



P. A. Lee, PRX 4, 031017

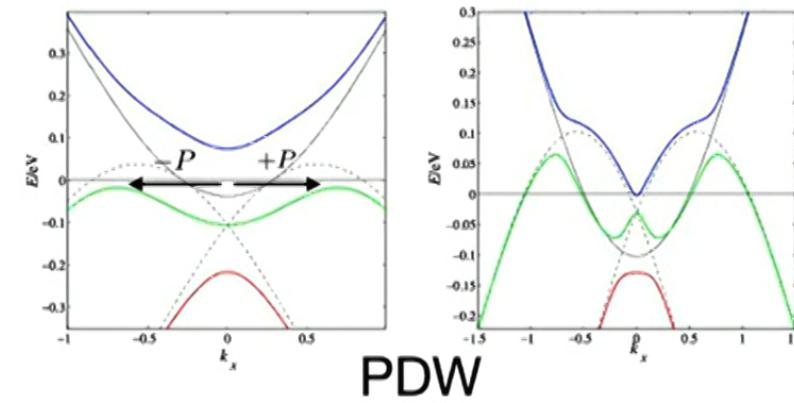


## PDW spectrum along two cuts

Mix electron and hole at different momentum

gapped antinode

gapless node



$$\langle c_{k\uparrow} c_{-k\pm P\downarrow} \rangle \neq 0$$

$$c_{k\uparrow} \rightarrow c_{-k+P\downarrow}^\dagger$$

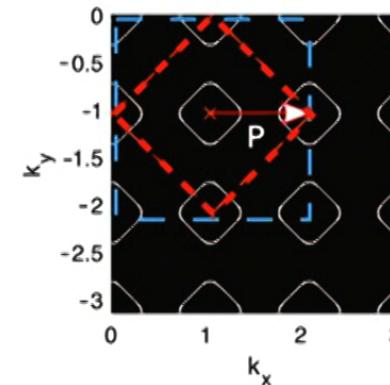
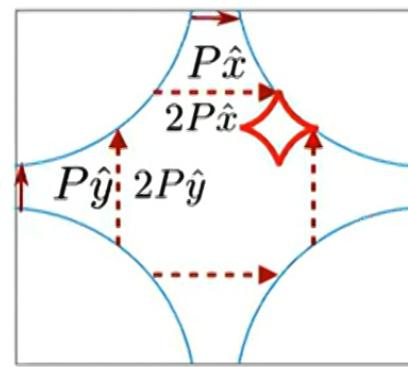
## Secondary CDW and MDW fold the B.Z.

Period-6 PDW

$$c_k \xrightarrow{\Delta_{\mathbf{P}_1}} c_{-k+\mathbf{P}_1}^\dagger \xrightarrow{\Delta_{\mathbf{P}_2}} c_{k-\mathbf{P}_1+\mathbf{P}_2}$$

PDWs at  $\pm P\hat{x}$  generate CDW at  $2P\hat{x}$

CDW connects nodal arcs into the Harrison-Sebastian pocket



Magnetization Density Wave from orbital current

PDWs at  $P\hat{x}$  and  $P\hat{y}$  generate MDW at  $P\hat{x} + P\hat{y}$

MDW sets the unit cell for the fluctuating PDW state

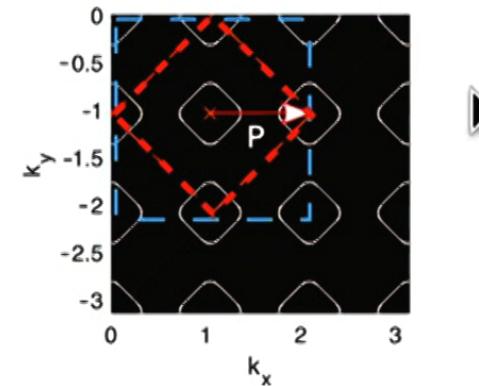
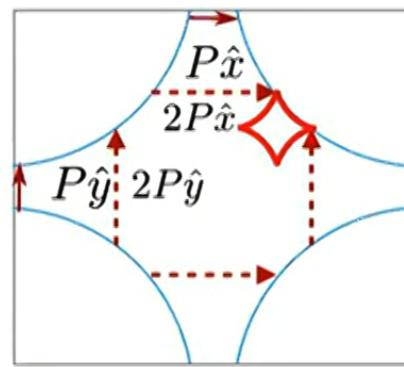
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# Outline

- Brief review of pseudogap phenomenology
- Review of pair density wave (PDW)
- Quantum disordered PDW: high field pseudogap ground state  
  
*Claim: small electron pocket + insulator of antinodal pairs*
- Electron spectrum of quantum disordered superconductor
- Comparison with experimental results

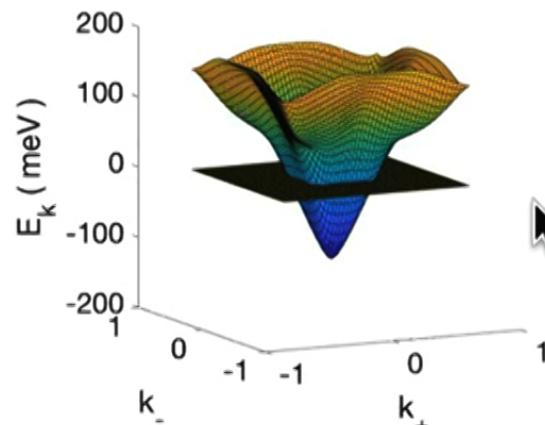
## Definition of *quantum disordered (fluctuating) order*

For an order parameter  $O$ . A *quantum disordered phase of  $O$*  refers to any phase where  $O$  does not have long range order (namely  $\langle O \rangle = 0$ ), but the phase is close to a transition into the ordered phase, thus  $\langle O(r)O(0) \rangle$  has a long correlation length relative to lattice scale.

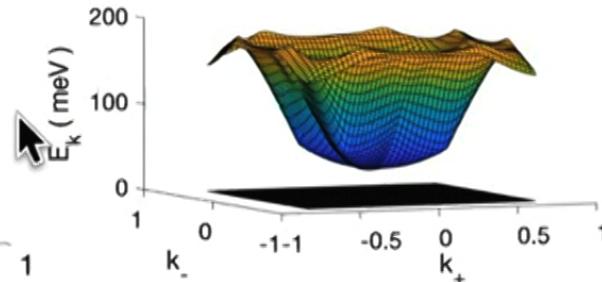
**Here  $O = PDW$**

To construct quantum disordered PDW, we want to drive the PDW state to a non-superconducting state

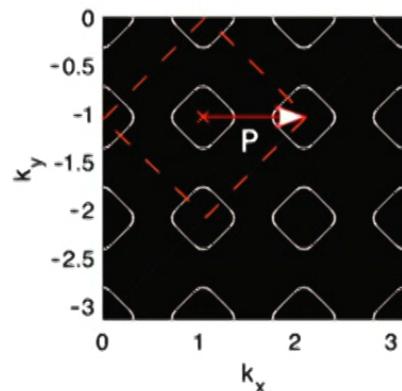
## Band structure of ordered PDW



The gapless band



The lowest gapped band

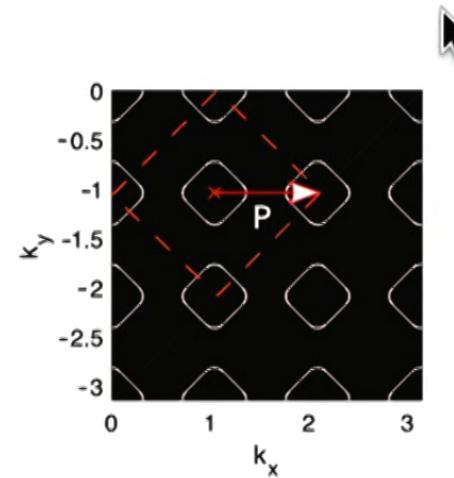
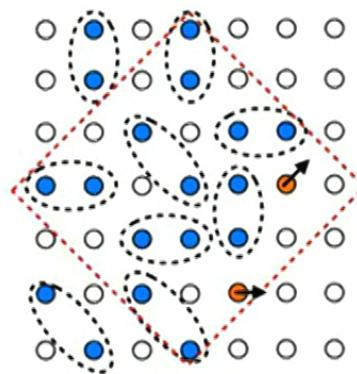


- 1 gapless band weakly coupled to PDW
- All other bands are gapped by PDW
- Treat them separately when considering the fluctuation of PDW



# The fluctuating PDW state

- Keep the secondary MDW and CDW, disorder PDW
- Treat the gapless band and gapped bands separately
- Small electron pocket + Bosonic Mott insulator



Z.Dai, T. Senthil, P. A. Lee, arXiv: 1906.01656

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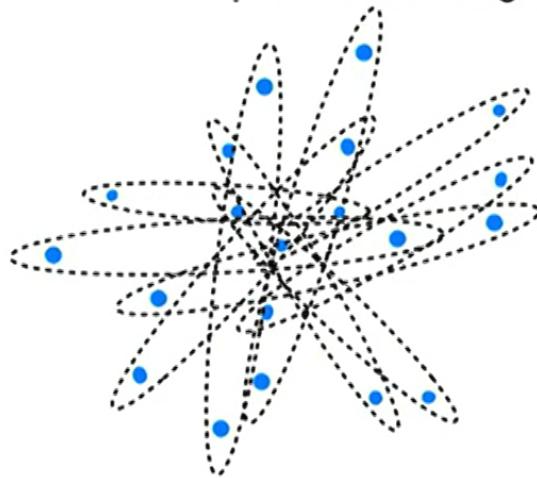
- **Electron spectrum of quantum disordered superconductor**
- Comparison with experimental results



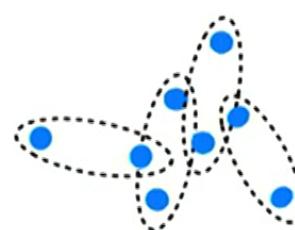
# Gapped antinodal electrons and fluctuating SC

Two possibilities when superconductivity is quantum disordered:

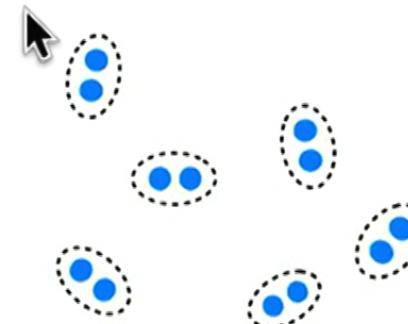
1. If the pairing is weak, the fermion gap may disappear.
2. If the pairing is strong, and the long range order is destroyed by the phase fluctuating of the pairing field, the fermion gap persists.



**BCS limit**  
pairs overlap



**Cuprate is in between**  
coherence length  $\sim 4a$



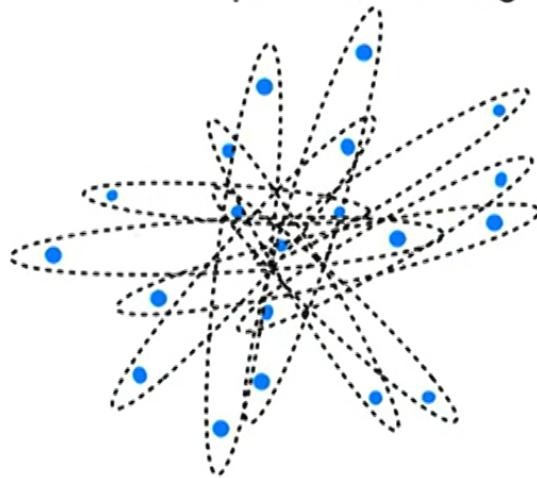
**BEC limit**  
pairing gap is just the  
binding energy



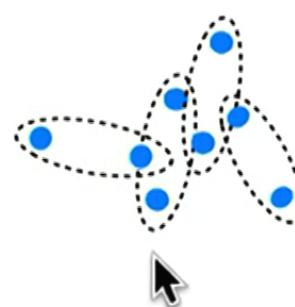
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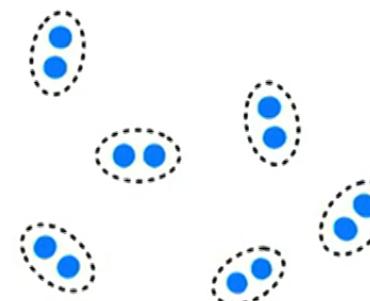
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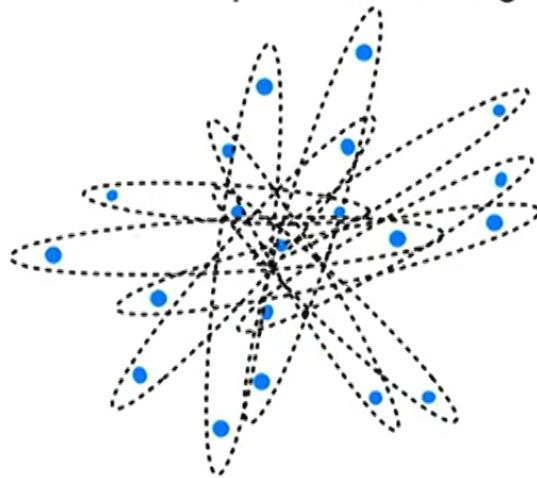
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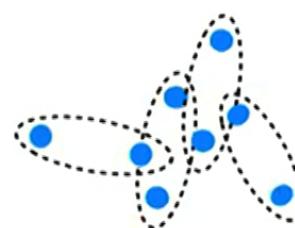
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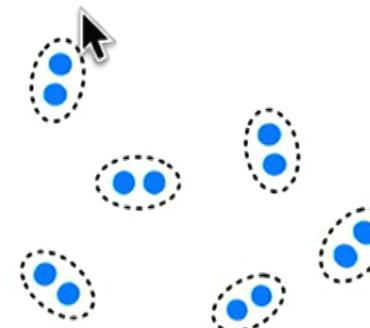
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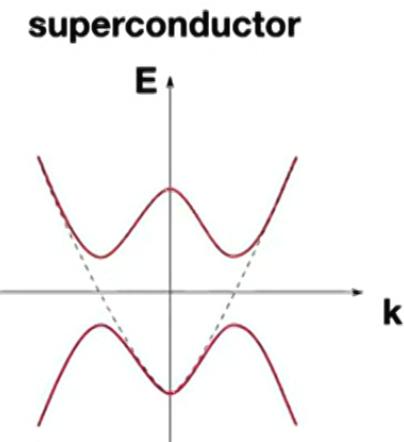
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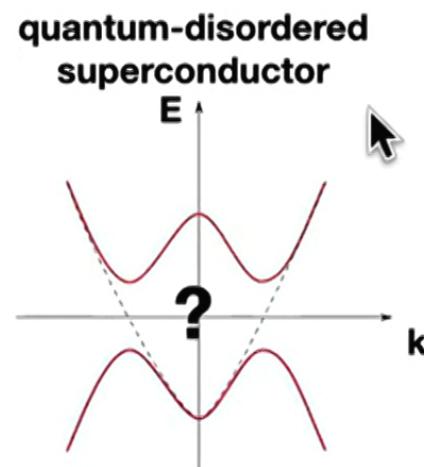
**BEC limit**  
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# Bogoliubov paradox? fluctuating s wave superconductor



Fermion spectrum in a superconductor is described by Bogoliubov band: superposition of electron and hole



How can an insulator with charge conservation have a similar band?

# Antinodal electron spectrum

Mott insulator GS: 

$$H = (c_{k\uparrow}^\dagger, c_{-k\downarrow}) \begin{pmatrix} \epsilon_k & |\Delta|e^{i\phi} \\ |\Delta|e^{-i\phi} & -\epsilon_k \end{pmatrix} \begin{pmatrix} c_{k\uparrow} \\ c_{-k\downarrow}^\dagger \end{pmatrix}$$

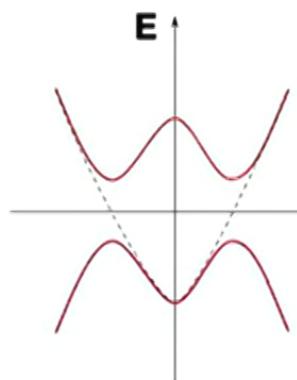
1-particle state:  
single hole 

2-particle state:  
electron + pair 

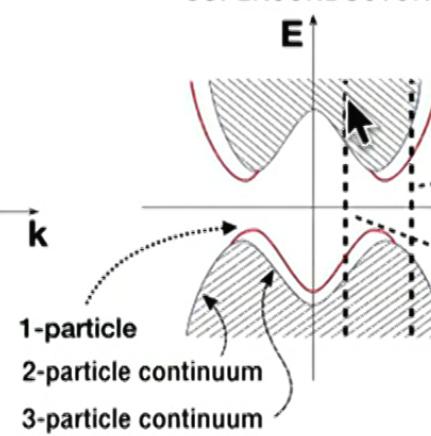
3-particle state: hole + charge 2e pair + charge -2e pair

$\psi_k = u c_{k\uparrow} + v c_{-k\downarrow}^\dagger$   
 $= |u| c_{k\uparrow} + |v| e^{i\phi} c_{-k\downarrow}^\dagger$

(a) SUPERCONDUCTOR



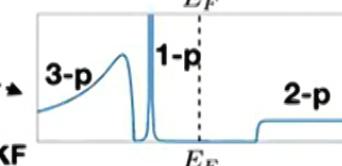
(b) QUANTUM DISORDERED SUPERCONDUCTOR



$k > k_F$



$k < k_F$

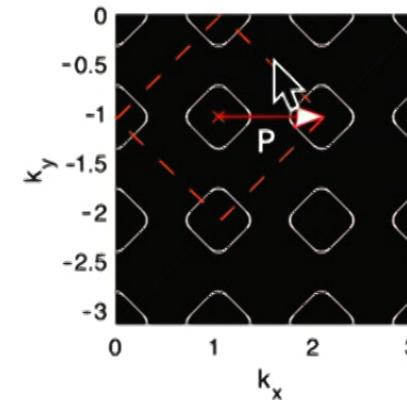
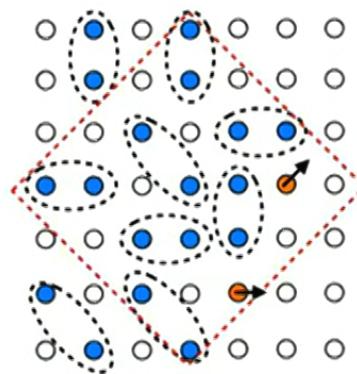


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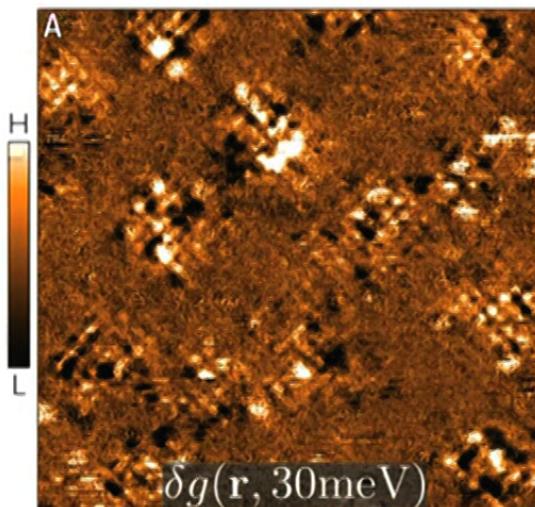
# The fluctuating PDW state

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- Treat the gapless band and gapped bands separately
- Small electron pocket + Bosonic Mott insulator



Z.Dai, T. Senthil, P. A. Lee, arXiv: 1906.01656

# Experimental observation of short range PDW



- S. D. Edkins et al., Science 364, 976 (2019)  
Y. Wang et al., PRB 97, 174510  
Z. Dai et al., PRB 97, 174511

**STM experiment in 2018, Seamus Davis group**

- static density wave in the vortex halo of the d wave SC
- **wave vector half of the previous CDW**



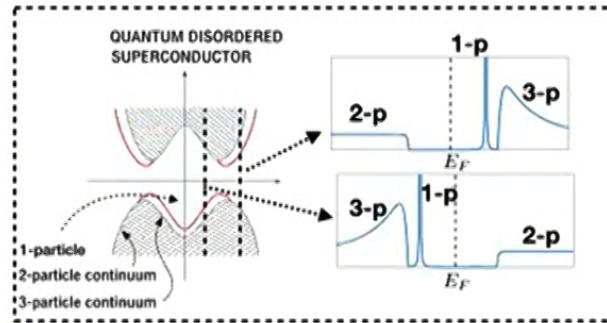
**Our interpretation:**

- short range PDW pinned by the vortex
- generate CDW at the same wave vector
- vortex halo determined by the correlation length of PDW
- small  $H_{c2}$  explained by the overlap of vortex halo

**When d wave SC is killed, PDW becomes fluctuating again**

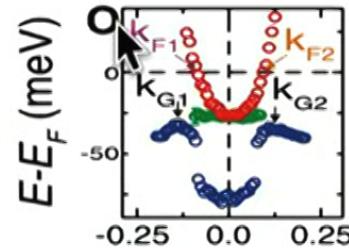
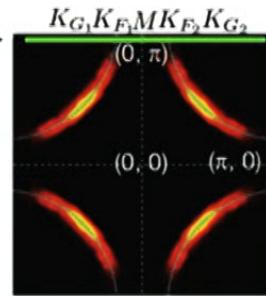
# Antinodal electron spectrum in fluctuating PDW

## Theory



## ARPES results for Bi2201

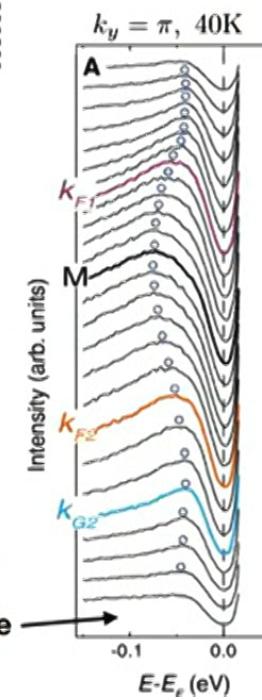
momentum along  
a cut  $k_y \sim \pi$



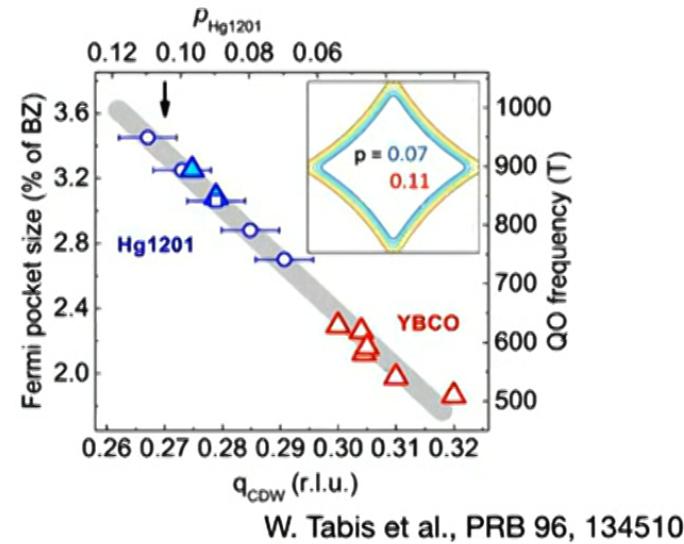
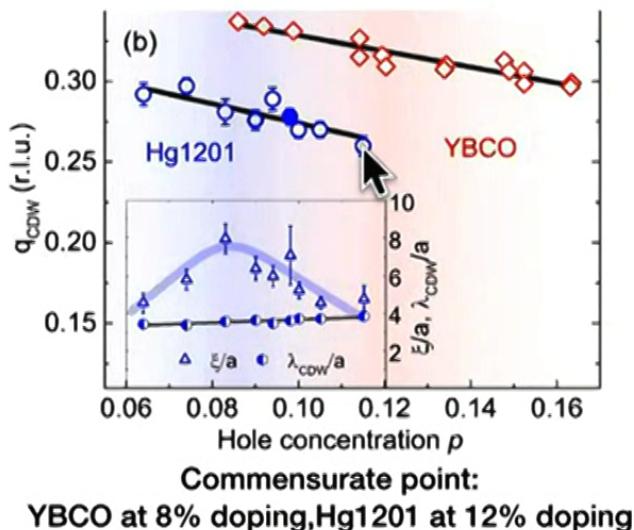
He, Rui-Hua, et al. *Science* 331.6024 (2011): 1579-1583.

Step function shows up only below  $T^*$  near antinode

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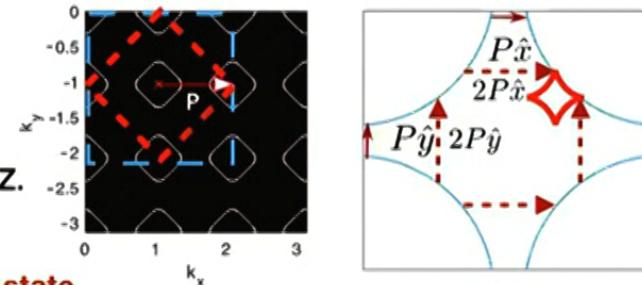


# Luttinger's theorem



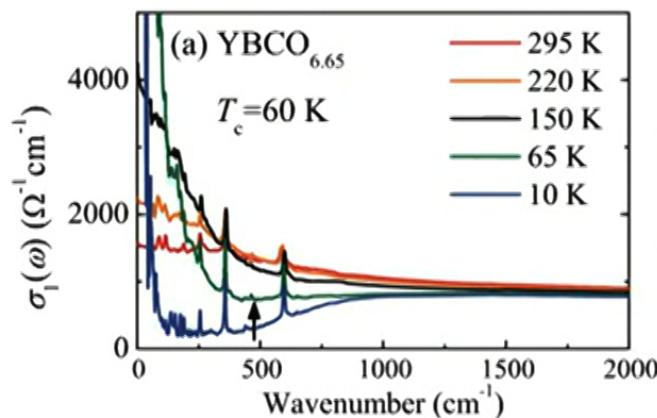
- Luttinger's theorem makes sense only in the commensurate case
- Need to pick a B.Z.
- Luttinger's theorem satisfied if we pick the MDW B.Z.
- Fluctuating PDW picture suggests the MDW B.Z.

Possibility of a non-exotic ground state

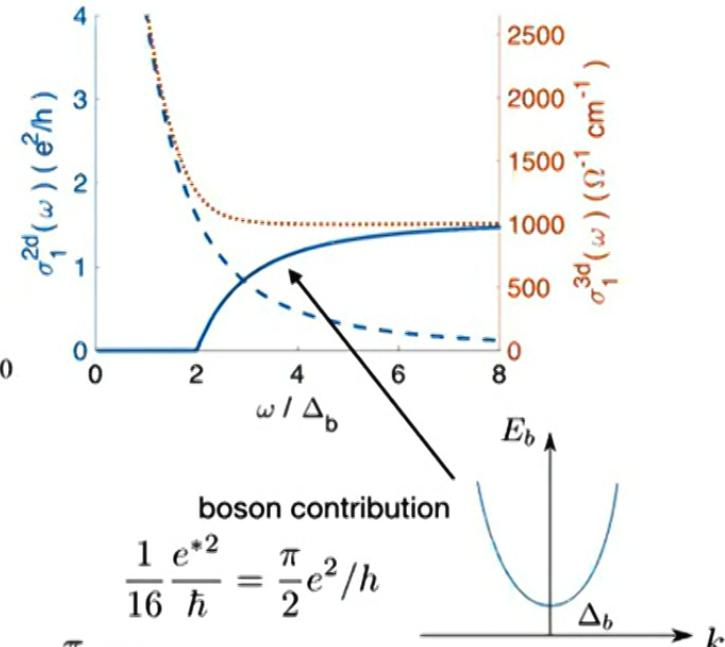


# Infrared conductivity

Experiment



Theory



**Drude peak**

$$\sigma = \frac{ne^2/m}{1/\tau - i\omega}$$

$$\text{Re}\sigma = \frac{ne^2\tau/m}{1 + \omega^2\tau^2}$$

Mid-infrared plateau almost exactly  $\frac{\pi}{2}e^2/h$  per layer!

D. N. Basov and T. Timusk, Rev. Mod. Phys. 77, 721 (2005)

Z.Dai, T. Senthil, P. A. Lee, arXiv: 1906.01656



# Conclusion

**Goal: use a single assumption to explain the pseudogap phenomenology at low temperature**

- Quantum-disordered PDW state captures the pseudogap phenomenology at low T
- Simple construction: a small-pocket Fermi liquid + bosonic Mott insulator
- Successfully explain the ARPES data in Bi2201, and the infrared conductivity plateau
- Explain CDW, small  $H_{c2}$ , quantum oscillation of a small pocket, consistent with the Luttinger's theorem
- MDW and the low-energy boson may be measured directly in the future

