

Title: The Holographic Landscape of Symmetric Product Orbifolds

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Collection: Emmy Noether Workshop: The Structure of Quantum Space Time

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Abstract: I will discuss the application of Siegel paramodular forms to constructing new examples of holography. These forms are relevant to investigate the growth of coefficients in the elliptic genus of symmetric product orbifolds at large central charge. The main finding is that the landscape of symmetric product theories decomposes into two regions. In one region, the growth of the low energy states is Hagedorn, which indicates a stringy dual. In the other, the growth is much slower, and compatible with the spectrum of a supergravity theory on AdS₃. I will provide a simple diagnostic which places any symmetric product orbifold in either region. The examples I will present open a path to novel realizations of AdS₃/CFT₂.

The holographic landscape of symmetric product orbifolds

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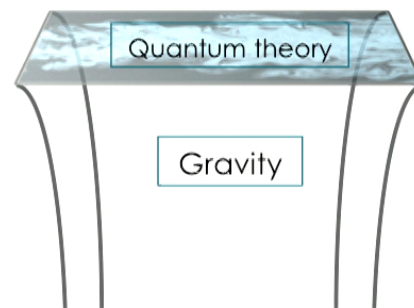
Based on
arXiv: 1611.04588 [hep-th]
arXiv: 1805.09336 [hep-th]
with [Alex Belin](#), [Joao Gomes](#) and [Christoph Keller](#)

And
arXiv: 1910.05353 [hep-th]
arXiv: 1910.05342 [hep-th]
[Alex Belin](#), [Christoph Keller](#) and [Beatrix Mühlmann](#)

Collaborators

AdS/CFT provides a non-perturbative, UV-complete definition of quantum gravity in Anti-de Sitter space.

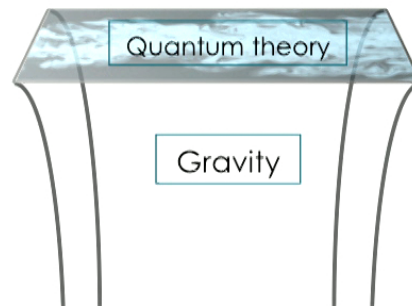
There are conditions on **CFT** such that it capture **semi-classical gravitational** features.



$\gamma^4 + 10\gamma^3 + (10\gamma^2 - 64\gamma + 108)q + \gamma^2 + 108\gamma^2 - 513\gamma + 808 + 108q^2 + \gamma^2 q^2 + \dots$

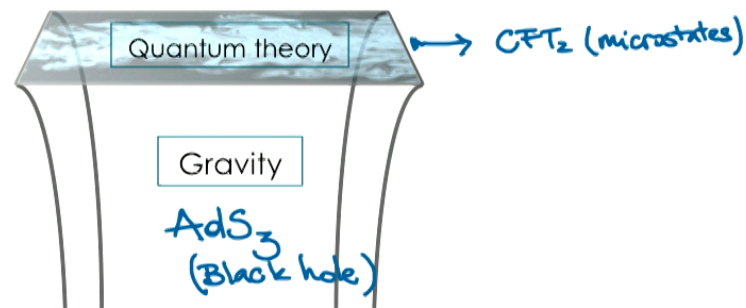
AdS/CFT provides a non-perturbative, UV-complete definition of quantum gravity in Anti-de Sitter space.

How to go about building CFTs with semi-classical gravitational features?

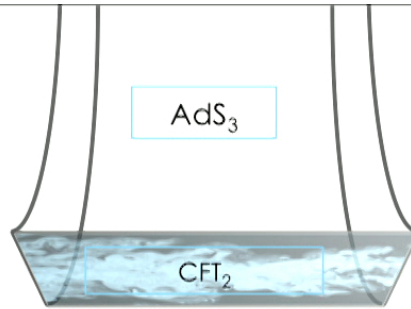


$\gamma^4 + 10\gamma^3 + (10\gamma^2 - 64\gamma^2 + 108 - 64\gamma + 10\gamma^2)q + \gamma^2 + 109\gamma^2 - 513\gamma + 808 + 808 - 513\gamma + 108\gamma^2 + \gamma^2 q^2 + \dots$

We will focus on the difficulties you encounter in AdS_3/CFT_2 .
Not universal, but it illustrates the challenges.



$y^4 + 4y^3 + y^2 - 8y^2 - y^4 + 16y + 8y^2 + y^4 + 4y^4 + y^3 - 32y^2 + y^4 + 56y + 32y^2 - y^4 + 4y^4 + y^4 + \dots$



AdS_3/CFT_2

Supergravity description

Large central charge

STRATEGY

1. Inspiration from known examples in String Theory.
2. Exploit holography.
3. Exploit **number theory**: crafting suitable counting formulas.

$y^2 + (y^2 - x^2 - a^2)z + (y^2 - x^2 - 2y^2 - a^2)z^2 + \dots$

Holographic CFTs
Conditions from gravity

The Landscape
Good, bad & promising SymProdOrb

Quantum Black Holes
Beyond area law

Outline

Holographic CFTs

Conditions from gravity

Necessary conditions on the spectrum of CFT_2
Strategy on describing the space of CFT_2

Outline



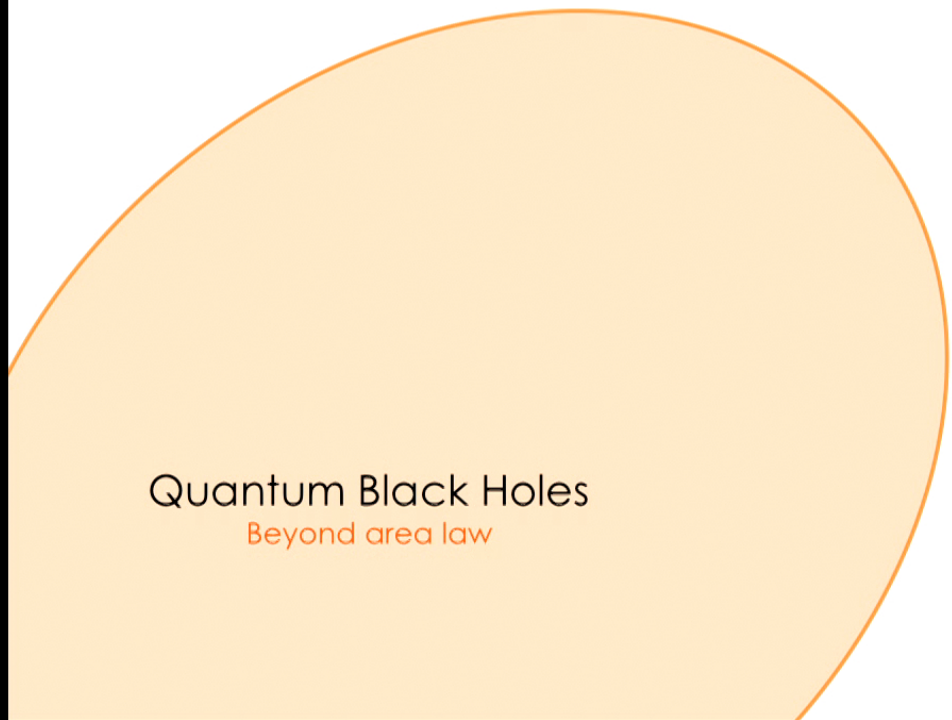
The Landscape

Good, bad & promising SymProdOrb

Outline

Implementing conditions.
Restrict the analysis to
supersymmetric states.
New results presented.

What can we learn about black hole microstates.



Quantum Black Holes
Beyond area law

Outline

Conditions from gravity

HOLOGRAPHIC CFTS

AdS₃ Gravity

The theory:

$$I_{3D} = \frac{1}{16\pi G_N} \int d^3x \sqrt{-g} \left(R + \frac{2}{\ell^2} \right) + \text{matter}$$

The spectrum:

1. Light States: Perturbative states
2. Heavy States: Black holes
3. Other stuff, e.g., multi-centered, conical defects (to be ignored today)

Universal entry in AdS₃/CFT₂: $c = \frac{3\ell}{2G_N} \gg 1$

Holographic CFT₂

We will impose two conditions

1. Black hole regime
2. Perturbative regime

$y^1+2+2y^1(2y^3+2y^2-2y^1+4-2y+2y^2+2y^3)q+(y^5-2y^4-6y^3-4y^2+5y^1+12+5y-4y^2-6y^3-2y^4+y^5)q^2+$

Heavy states: BTZ Black Hole

Background:

$$ds^2 = -\frac{(r^2 - r_+^2)(r^2 - r_-^2)}{r^2} dt^2 + \frac{\ell^2 r^2}{(r^2 - r_+^2)(r^2 - r_-^2)} dr^2 + r^2 \left(d\phi + \frac{r_+ r_-}{r^2} dt \right)^2$$

+ U(1) CS field

Thermodynamics:

$$S_{BH} = \frac{A_H}{4G_N}, \quad M = \frac{r_+^2 + r_-^2}{8G_N \ell^2}, \quad J = \frac{r_+ r_-}{4G_N \ell}$$

Huge degeneracy!

Supersymmetry:

$$\ell M = J, \quad T_H = 0 \quad + \text{quantization of U(1)}$$

Holographic CFT₂

1. Black hole regime:

$$A_H \gg G_N$$

Very massive black hole
with lots of entropy

$$S_{\text{BH}} = \ln d(c, E)$$

$$= 2\pi \sqrt{\frac{cE}{6}} + \dots$$

$$= \frac{A_H}{4G} + \dots$$

symmetry

[Strominger+Vafa; Strominger]

Holography

$y^1+2y^2+(2y^3+2y^4-2y^5+4y^6+2y^7+2y^8)q+(y^9-2y^{10}-6y^{11}-4y^{12}+5y^{13}+12+5y^{14}y^2-6y^{15}2y^{16}+y^{17}q^2+\dots$

Holographic CFT₂

2. Perturbative regime:

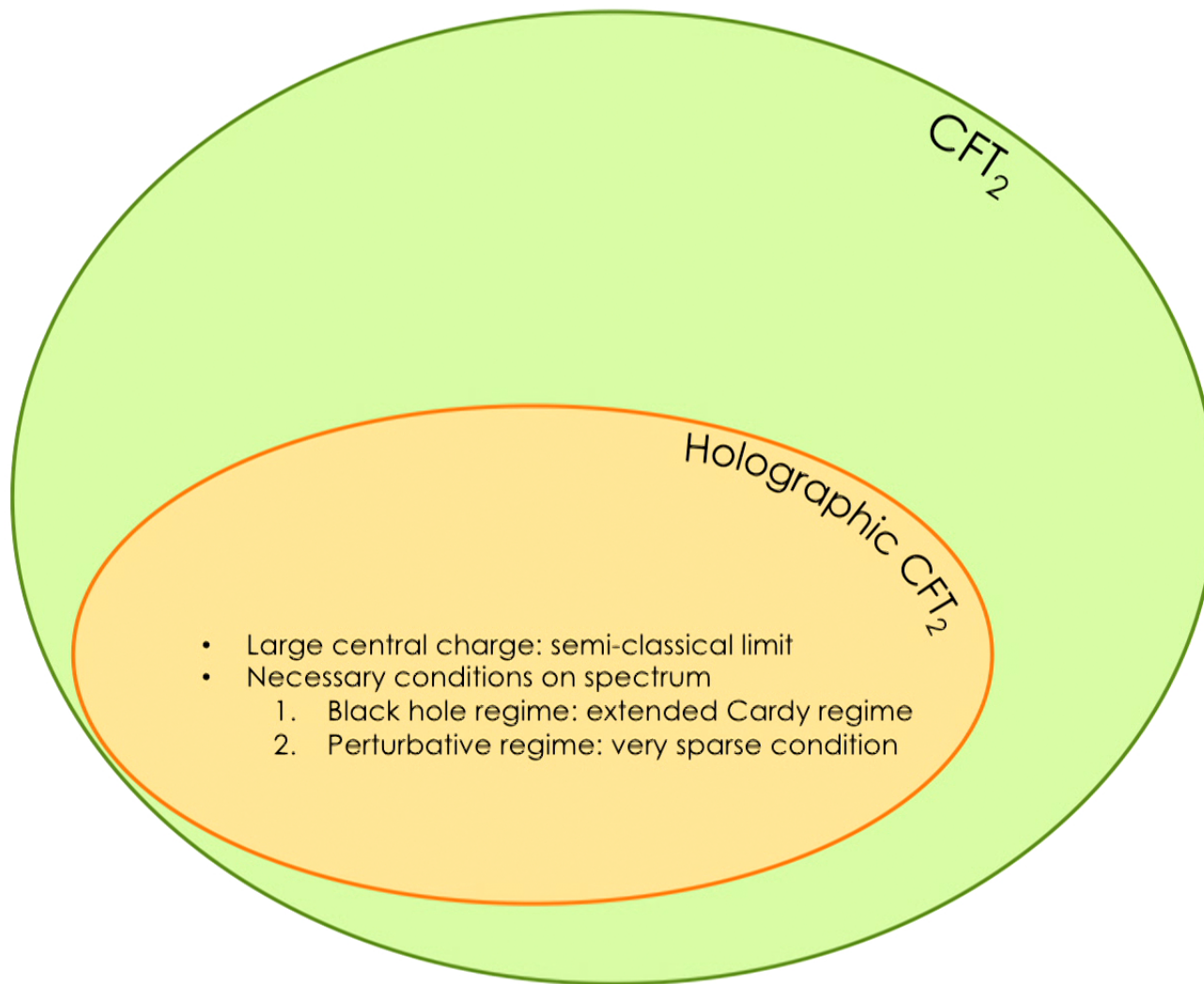
Light = Energy is $O(1)$ in Planck units.
Perturbative excitations that do not form a black hole.

- Presence of Hawking-Page transition [Keller; Hartman, Keller, Stoica]
- Extended Cardy regime for BPS BHs [Benjamin, Cheng, Kachru, Moore, Paquette; Benjamin, Kachru, Keller, Paquette]

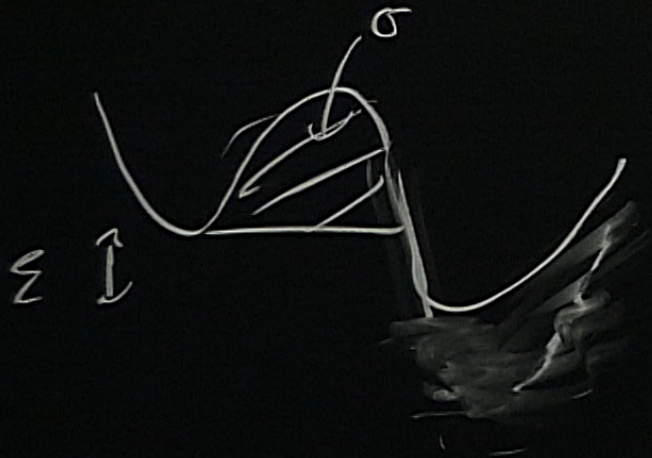
$$\ln d(E) \sim E^\alpha \quad \alpha < 1$$

Very sparse spectrum

$y^1+2y^2+(2y^3+2y^4-2y^5+4y^6+2y^7+2y^8+2y^9)q+(y^{10}-2y^{11}-6y^{12}-4y^{13}+5y^{14}+12+5y^{15}-4y^{16}-6y^{17}+2y^{18}+y^{19})q^2+\dots$



Warning: Sizes are not meaningful.



$$S = 2\pi \sqrt{\frac{E_c}{6}}$$

Cardy: $E \gg c$

EXT Cardy: $E \sim c \gg 1$

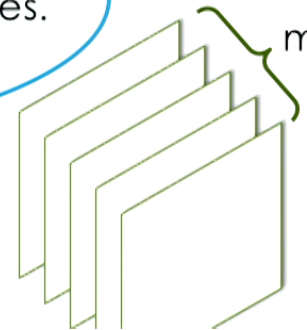
\updownarrow
 very sparse cond
 $A \gg G_N$

Strategy

We want CFTs with large central charge
We want control on the spectrum



Search within symmetric product theories.
 $\mathcal{C}_m = \mathcal{C}^{\otimes m} / S_m$



Strategy

We want CFTs with large central charge
We want control on the spectrum

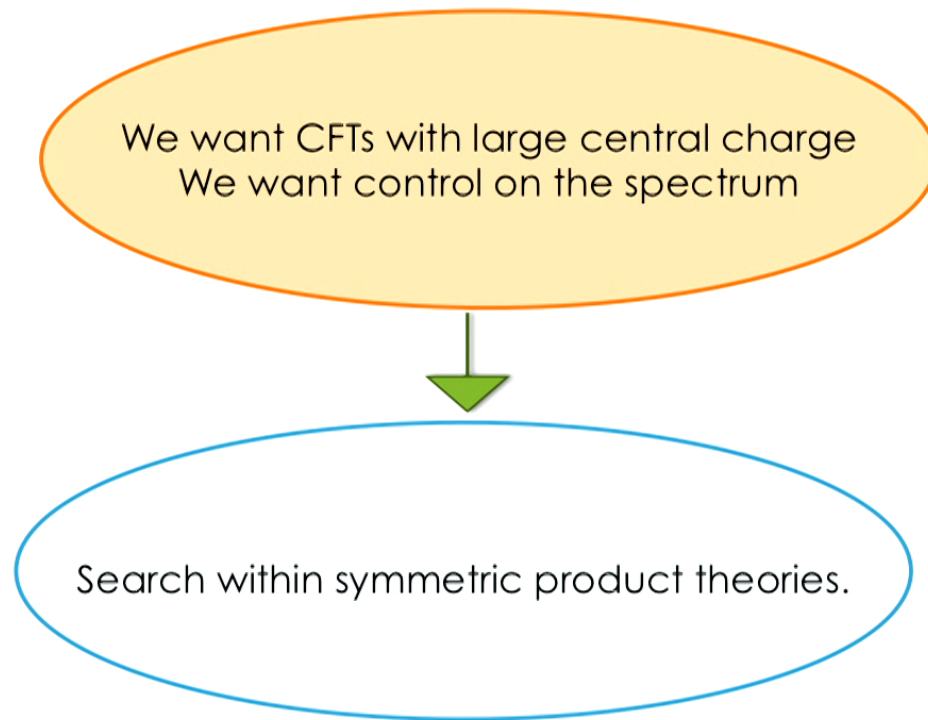


Search within symmetric product theories.

$$\mathcal{C}_m = \mathcal{C}^{\otimes m} / S_m$$

$$Z(\tau, \rho) = \sum_m Z(\tau; \mathcal{C}_m) p^m = \prod_{m>0, n \in \mathbb{Z}} \frac{1}{(1 - p^m q^n)^{d(mn)}} = \sum_{m, n} d_m(n) p^m q^n$$

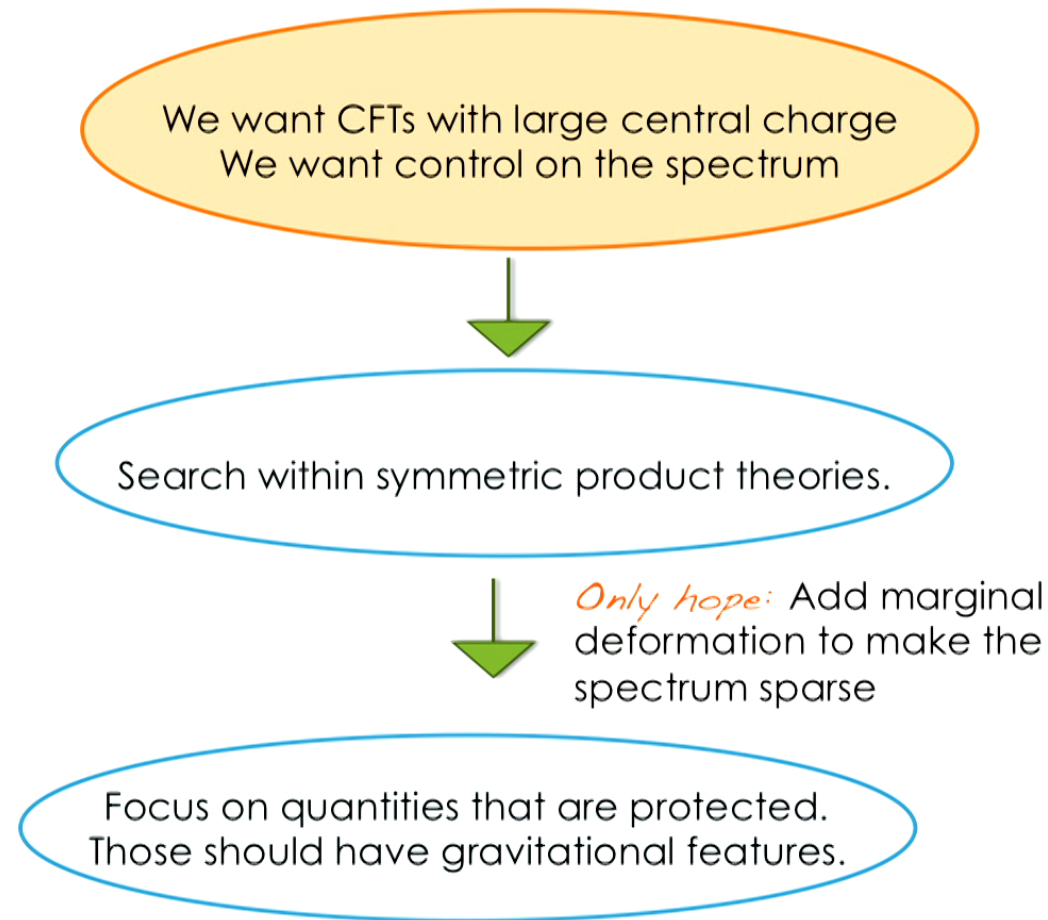
Strategy



Big Obstruction: The partition function of symmetric product CFT_2 is stringy, i.e. Hagedorn like.

$y^{l_1+l_2} (y^{l_1+l_2})^{2y_1+2y_2} (y^{l_1+l_2})^{2y_1+2y_2} (y^{l_1+l_2})^{2y_1+2y_2} \dots$

Strategy



BPS states in SCFT₂

Focus on protected quantities: the **elliptic genera**.

$$\chi(\tau, z) = \text{tr}_{RR} \left((-1)^F q^{L_0 - \frac{c}{24}} y^{J_0} \bar{q}^{\bar{L}_0 - \frac{\bar{c}}{24}} \right)$$

The elliptic genera is related to a **weak Jacobi form** of index t .

Focus on symmetric product orbifolds: easy to get **large values of $c=6r$** .

$$\mathcal{Z}(\rho, \tau, z) = \sum_r \chi(\tau, z; \text{Sym}^r(M)) e^{2\pi i \rho t r} :$$

Necessary condition on
light states in NS sector
as $r \rightarrow \infty$

$$\ln d(E) \sim E^\alpha \quad \alpha < 1$$

Note: The partition function of symmetric product CFT₂ has $\alpha=1$.
The elliptic genus can display cancellations that capture the
spectrum away from the symmetric product point.

$y^{1+10y+(10y^2-64y^3+108-64y^4+108-64y^5+10y^6)q+y^7+(108y^2-513y^3+808-513y^4+808-513y^5+108y^6+y^7)q^2+\dots}$

BPS states in SCFT₂

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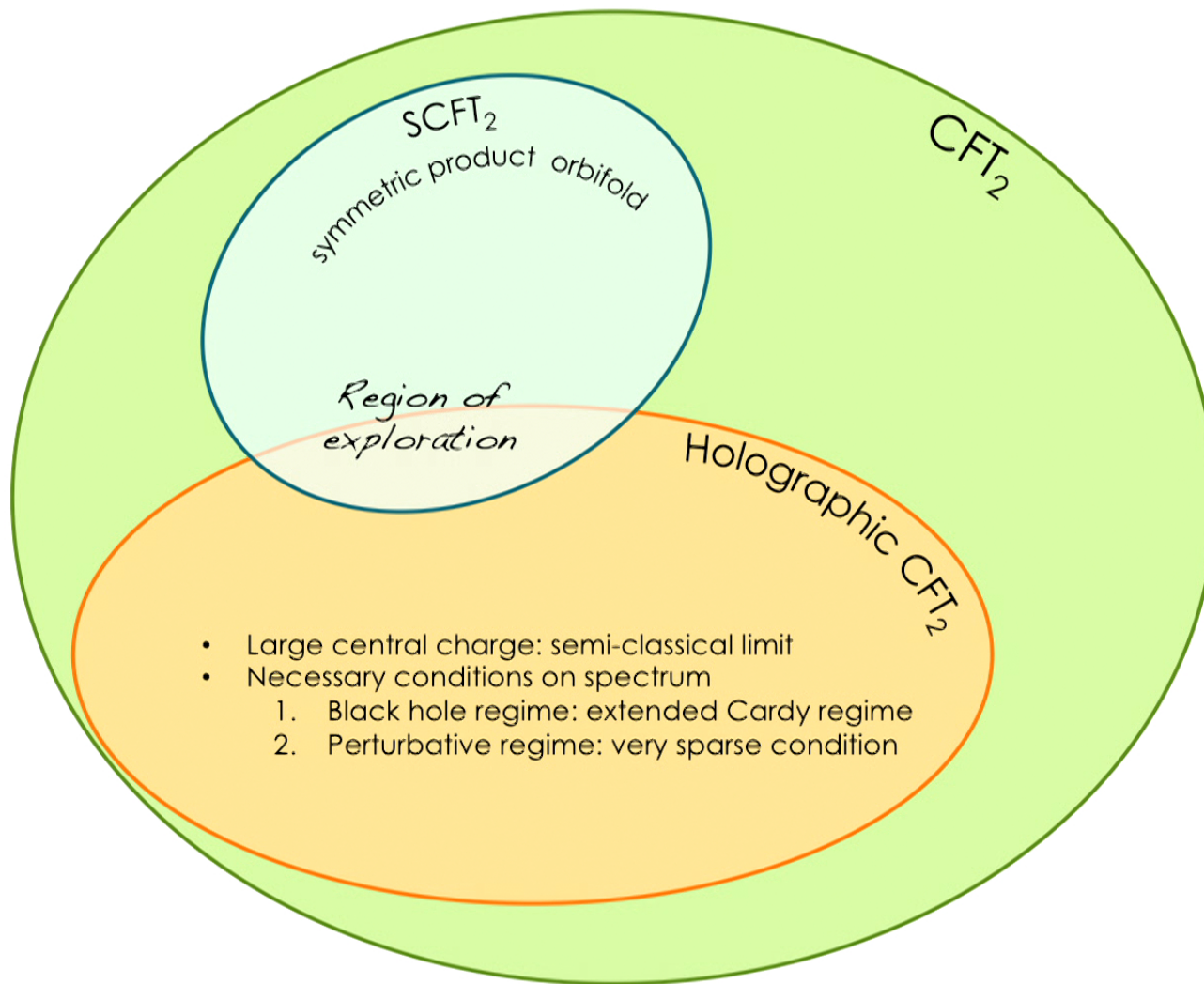
Necessary condition on
light states in NS sector
as $r \rightarrow \infty$

$$\ln d(E) \sim E^\alpha \quad \alpha < 1$$

Spoiler!

We can tell you unambiguously which wJFs are holographic, i.e $\alpha < 1$.
New examples are unveiled.

$y^{1+10y+(10y^2-64y^3+108-64y+108-64y+10y^2)q+y^5+(108y^2-513y^3+808-513y+108y^2+y^2)q^2+\dots}$



Warning: Sizes are not meaningful.

Good, bad & promising Symmetric Product Orbifolds

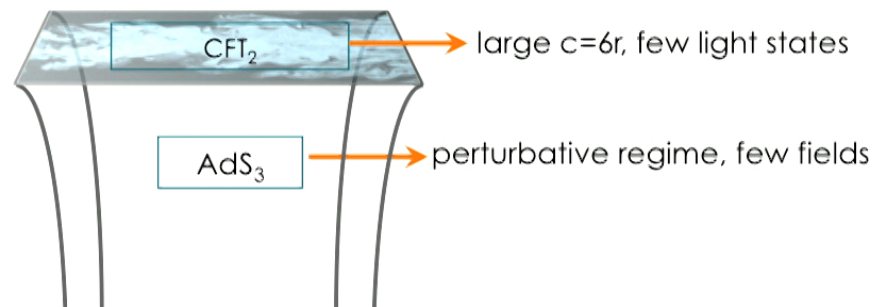
THE LANDSCAPE

The class of objects we are interested are of the form:

$$\mathcal{Z}(\rho, \tau, z) = \sum_r \chi(\tau, z; \text{Sym}^r(M)) e^{2\pi i \rho t r} = \prod_{\substack{n, l, r \in \mathbb{Z} \\ r > 0}} (1 - q^n y^l p^{tr})^{-c(nr, l)}$$

Task:

- to characterise the BPS spectrum in these theories.
- focus on the large r limit

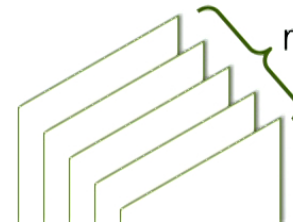


Our procedure in a few steps:

1. Select one seed theory: select a Jacobi form.
2. Perform symmetric product orbifold: increases c
3. Relate generating function to Siegel paramodular form: gives the mathematical control to extract spectrum.
4. Build a modular form that capture the light part of the CFT spectrum

$$\chi_{\text{NS},\infty} \equiv \sum_{h,l} d_{\infty}^{\text{NS}}(h,l) q^h y^l = \prod_{\substack{h \geq 0, l \in \mathbb{Z} \\ (h,l) \neq (0,0)}} \frac{1}{(1 - q^h y^l)^{f(h,l)}}$$

Generating functional for light (perturbative) states at infinite r



Promising Examples

a.k.a. SMFs that meet the two necessary conditions on the BPS spectrum

$$\ln d_{\infty}^{\text{NS}}(h) \sim h^{\alpha} \quad \alpha < 1 \quad \text{Very sparse spectrum}$$

$$\chi_{\text{NS},\infty} \equiv \sum_{h,l} d_{\infty}^{\text{NS}}(h,l) q^h y^l = \prod_{\substack{h \geq 0, l \in \mathbb{Z} \\ (h,l) \neq (0,0)}} \frac{1}{(1 - q^h y^l)^{f(h,l)}}$$

Promising Examples

We found a very **simple condition**. Given a seed Jacobi form

$$\phi_{0,t}(\tau, z) = q^0 y^{-b} + \dots$$

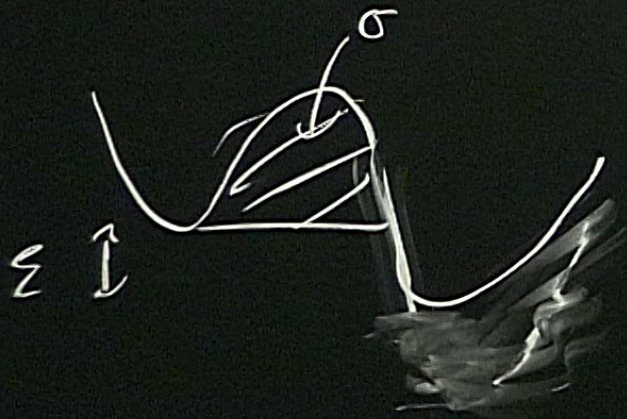
Demand **positivity** of

$$\alpha = \max_{j=0, \dots, b-1} \left(-\frac{t}{b^2} j \left(j - \frac{bl}{t} \right) - n \right)$$

Then $f(h,l)$ is sum of Kronecker functions.

Necessary condition is $b^2 \leq t$

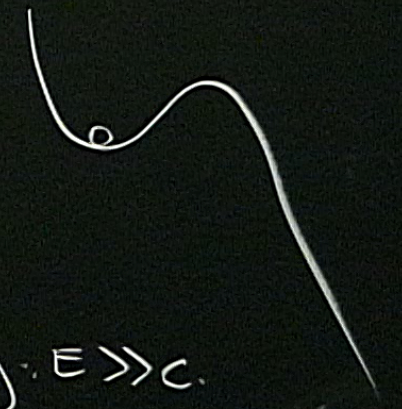
$$\chi_{\text{NS}, \infty} \equiv \sum_{h,l} d_{\infty}^{\text{NS}}(h,l) q^h y^l = \prod_{\substack{h \geq 0, l \in \mathbb{Z} \\ (h,l) \neq (0,0)}} \frac{1}{(1 - q^h y^l)^{f(h,l)}}$$



$$P$$

$$y = e^{2\pi i z}$$

$$q = e^{2\pi i z}$$



$$S = 2\pi \sqrt{\frac{E_C}{6}}$$

$$= \sqrt{E_{CL}} + \sqrt{E_{2C2}}$$

Cardy: $E \gg c$
 EXT Cardy: $E \sim c \gg 1$
 $A_H \gg G_N$
 Very sparse cond

Promising Examples

t	b	dim	t	b	dim	t	b	dim	t	b	dim
1	1	1	8	1	0	12	1	0	16	1	0
2	1	1	8	2	2	12	2	2	16	2	1
3	1	1	9	1	0	12	3	3	16	3	2
4	1	1	9	2	1	13	1	0	16	4	4
4	2	2	9	3	3	13	2	0	17	1	0
5	1	0	10	1	0	13	3	1	17	2	0
5	2	1	10	2	1	14	1	0	17	3	0
6	1	1	10	3	2	14	2	0	17	4	2
6	2	2	11	1	0	14	3	1	18	1	0
7	1	0	11	2	0	15	1	0	18	2	0
7	2	1	11	3	1	15	2	1	18	3	3
						15	3	2	18	4	3

Today

We can easily diagnose if a Symmetric Product Orbifold does the holographic trick.

$y^4 + 10y^3 + (10y^2 - 64y^2 + 108 - 64y + 10y^2)q + y^4 + 10y^3 + (10y^2 - 64y^2 + 108 - 64y + 10y^2)q^2 + \dots$

Promising Examples

t	b	dim	t	b	dim	t	b	dim	t	b	dim
1	1	1	8	1	0	12	1	0	16	1	0
2	1	1	8	2	2	12	2	2	16	2	1
3	1	1	9	1	0	12	3	3	16	3	2
4	1	1	9	2	1	13	1	0	16	4	4
4	2	2	9	3	3	13	2	0	17	1	0
5	1	0	10	1	0	13	3	1	17	2	0
5	2	1	10	2	1	14	1	0	17	3	0
6	1	1	10	3	2	14	2	0	17	4	2
6	2	2	11	1	0	14	3	1	18	1	0
7	1	0	11	2	0	15	1	0	18	2	0
7	2	1	11	3	1	15	2	1	18	3	3
						15	3	2	18	4	3

Work in Progress

Give a CFT and supergravity description.

What is the physical interpretation of the mathematical conditions?

$y^4 + 10y^3 + (10y^2 - 64y^2 + 108 - 64y + 10y^2)q + y^4 + 10y^3 + (10y^2 - 64y^2 + 108 - 64y + 10y^2)q^2 + \dots$

Good Examples

a.k.a. SMFs that meet the two necessary conditions on the BPS spectrum, **and** we know the symm prod CFT₂ and AdS₃ supergravity theory.

1. Igusa Cusp form

$$\Phi_{10}(\Omega) = \text{Exp-Lift}(2\phi_{0,1}) \quad \text{with} \quad \phi_{0,1} = \frac{1}{2}\chi(\tau, z; K3)$$

$$\frac{1}{\Phi_{10}(\Omega)} = \frac{\mathcal{Z}(\Omega)}{p\phi_{10,1}(\tau, z)} \xrightarrow{\text{light states}} \ln d(E) \sim E^{1/2}$$

- 1/4 BPS dyons in N=4 D=4 string theory [DVV]
- Quantum black hole [Sen, Dabholkar, Murthy, Gomes, ...]
- AdS₃ × S³ × K3 supergravity spectrum matches N=(4,4) SCFT [de Boer]
- CHL generalizations [David, Jatkar, Sen; Paquette, Volpato, Zimet]

Good Examples

a.k.a. SMFs that meet the two necessary conditions on the BPS spectrum, **and** we know the symm prod CFT₂ and AdS₃ supergravity theory.

2. $t=4$ paramodular form

Seed

$$\phi_{0,1}(\tau, 2z) + 4(1+a)\phi_{0,4}(\tau, z) = y^{\pm 2} + 4(1+a)y^{\pm 1} + 14 + 4a + \mathcal{O}(q)$$

Pole structure

Humbert surfaces are $H_4(2)$ and $H_1(1)$

Perturbative regime

$$\ln d(E) \sim E^{1/2}$$

Holography

$$\text{AdS}_3 \times (S^3 \times \mathbb{T}^4)/G \longleftrightarrow \text{N}=(2,2) \text{ SCFTs [Datta, Eberhardt, Gaberdiel]}$$

$y^{-10} + y^{-8} + 10y^{-2} - 64y^{-1} + 108 - 64y + 10y^2 + y^3 + 108y^2 - 513y^3 + 808 - 513y + 108y^2 + y^3 + y^4 + \dots$

Beyond area law

QUANTUM BLACK HOLES

Black holes & SMFs

The task is to extract the physics content of Fourier coefficients for large (asymptotic) values of the charges.

$$d(m, n, l) = \oint_{p=0} \frac{dp}{2\pi ip} \oint_{q=0} \frac{dq}{2\pi iq} \oint_{y=0} \frac{dy}{2\pi iy} \frac{1}{\Phi_k(\Omega)} p^{-m} q^{-n} y^{-l}$$

$$\downarrow$$
$$\Phi_k(\rho, \tau, z) = \Phi_k(t^{-1} \tau, t\rho, z)$$

Black holes & SMFs

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$$d(m, n, l) = \oint_{p=0} \frac{dp}{2\pi i p} \oint_{q=0} \frac{dq}{2\pi i q} \oint_{y=0} \frac{dy}{2\pi i y} \frac{1}{\Phi_k(\Omega)} p^{-m} q^{-n} y^{-l}$$

Asymptotic behavior of degeneracy

$$\begin{aligned} S_{\text{BH}} &= \ln d(c, E) \\ &= 2\pi \sqrt{\frac{cE}{6}} + \dots \\ &= \frac{A_{\text{H}}}{4G} + \dots \end{aligned}$$

symmetry → $d(E, c) = d(tE, t^{-1}c)$

Holography

Exchange symmetry gives us (almost) automatically an **extended Cardy regime**.

Quantum corrections

For lack of time, focus on logarithmic correction

$$S_{\text{BH}} = \ln d(c, E, J) = \frac{A_H}{4G} + \# \ln(A_H/G) + \dots$$

Scaling regime	$\tau_{1,2}^*$	A_H^2	$\ln \Lambda$
I. $E \gg 1$ $E \sim \Lambda^2, c \sim O(1), J \sim \Lambda$	Λ	Λ^2	$-(k+2)$
II. $E \sim c \gg 1$ $E \sim \Lambda, c \sim \Lambda, J \sim \Lambda$	Λ^0	Λ^2	$m_{1,1} - 2$
III. $c \gg E \gg 1$ $E \sim \Lambda, c \sim \Lambda^2, J \sim \Lambda^{3/2}$	$\Lambda^{-1/2}$	Λ^3	$m_{1,1} - 3 - \frac{k}{2}$

} Cardy regime.

} Regime relevant for BPS BHs in 4D & 5D [Sen et al]

To be explored

1. Logarithmic corrections

$$S_{\text{BH}} = \ln d(c, E, J) = \frac{A_H}{4G} + \underbrace{\#}_{\text{?}} \ln(A_H/G) + \dots$$

Universality of this correction is being tested in gravity.
Less is known microscopically.

[Charles, Larsen; AC, Godet, Larsen, Zeng]
[Mukhametzhano, Zhiboedov]

2. Mock modular forms

Decompose counting formula as quantum degeneracies of single-centered and multicentered configurations.

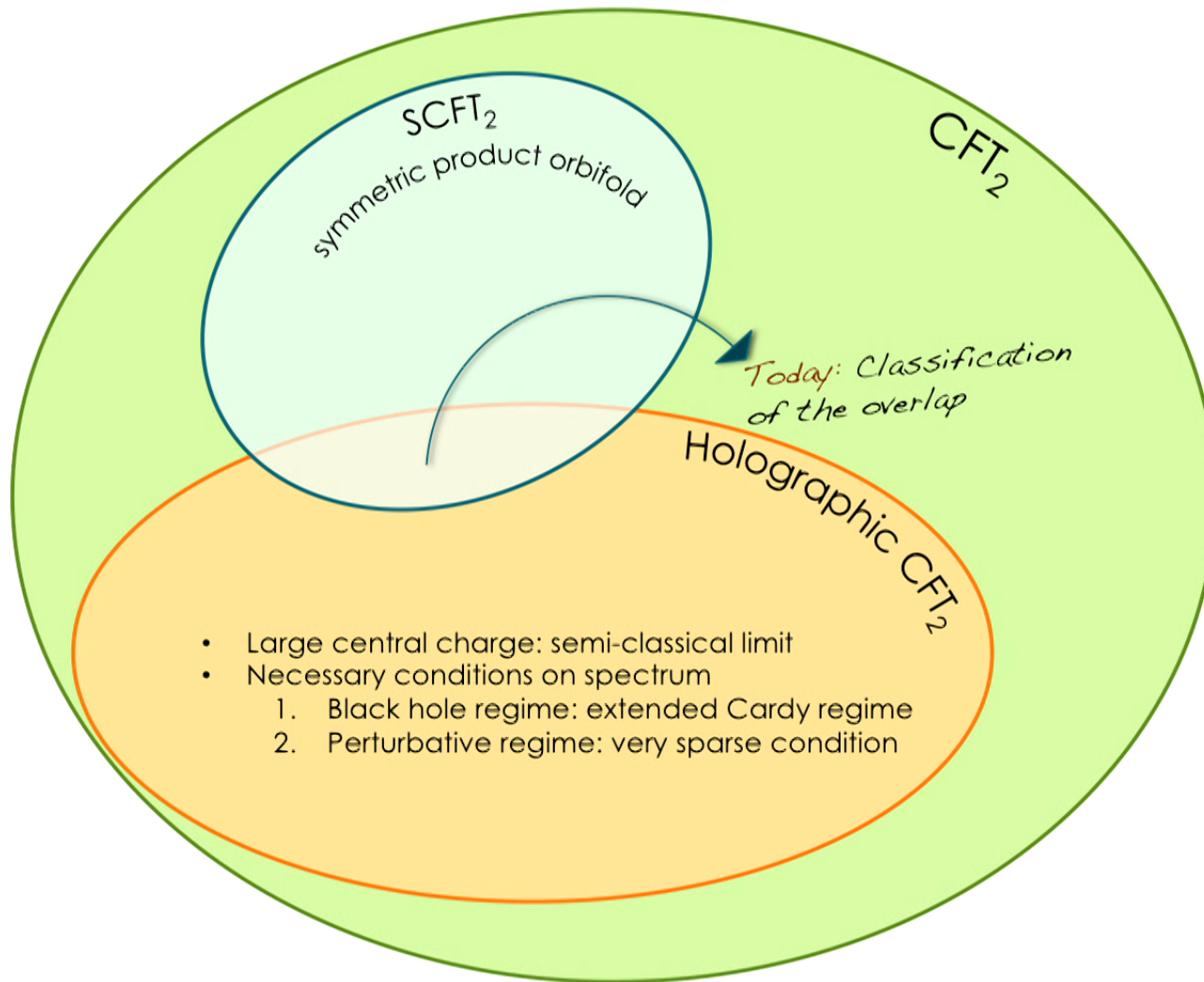
[Dabholkar, Murthy, Zagier]

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Conditions from gravity

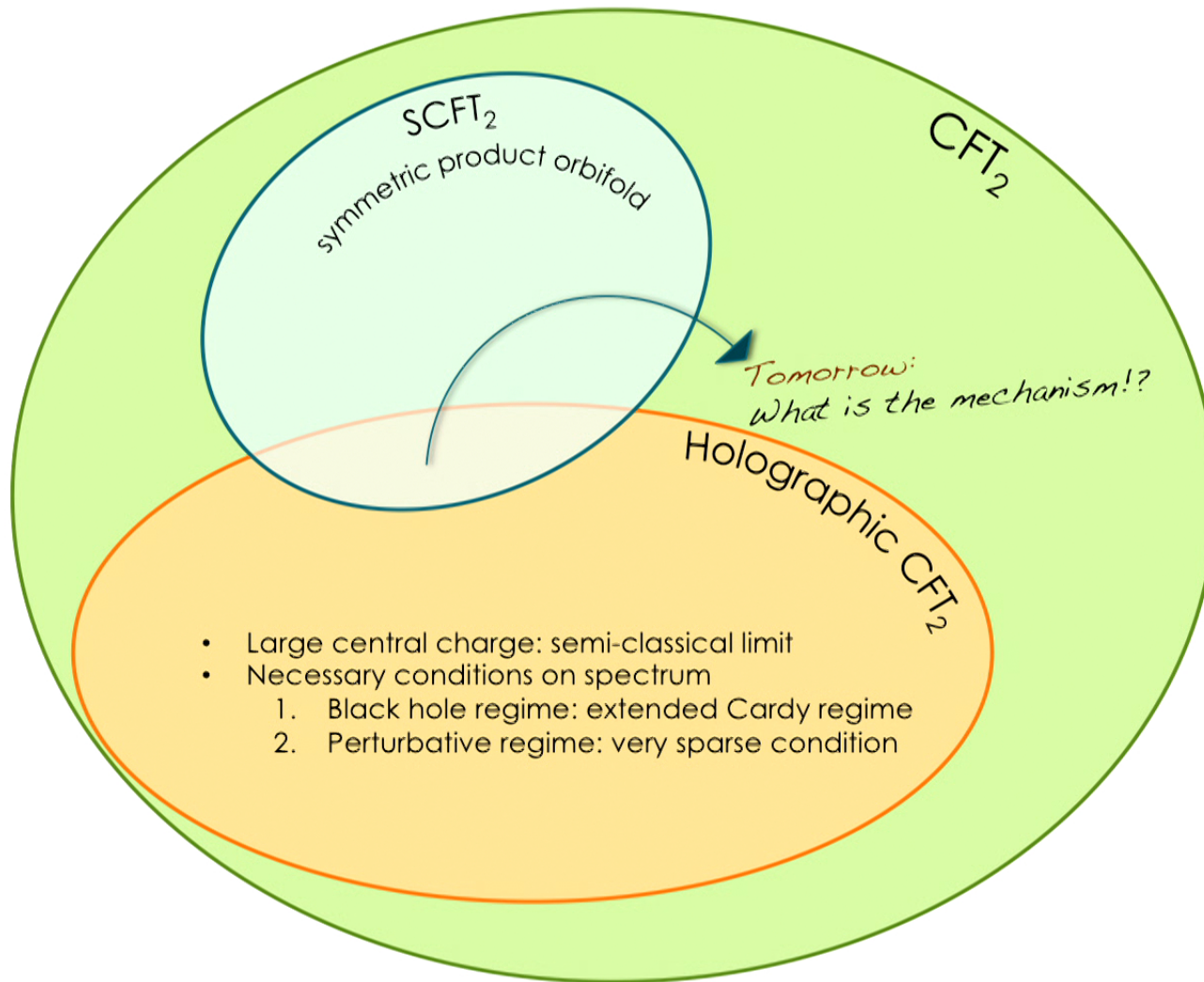
The Landscape
Good, bad & promising SMFs

Quantum Black Holes
Beyond area law

Summary



Warning: Sizes are not meaningful.



- Large central charge: semi-classical limit
- Necessary conditions on spectrum
 1. Black hole regime: extended Cardy regime
 2. Perturbative regime: very sparse condition

Warning: Sizes are not meaningful.

THANK YOU!

$$\begin{aligned} & y^{-1}+2+y-(2y^{-3}+2y^{-2}-2y^{-1}+4-2y+2y^2+2y^3)q+(y^{-5}-2y^{-4}-6y^{-3}-4y^{-2}+5y^{-1}+12+5y-4y^2-6y^3-2y^4+y^5)q^2+\dots \\ & \quad y^{-1}+1+y-(y^4+y^3-y^1-2-y+y^3+y^4)q+(-y^{-5}-2y^{-4}-2y^{-3}+3y^{-1}+4+3y-2y^3-2y^4-y^5)q^2+\dots \\ & y^{-1}+10+y+(10y^{-2}-64y^{-1}+108-64y+10y^2)q+(y^{-3}+108y^{-2}-513y^{-1}+808-513y+108y^2+y^3)q^2+\dots \\ & y^{-1}+4+y+(y^{-3}-8y^{-2}-y^{-1}+16-y+8y^2+y^3)q+(4y^{-4}-y^{-3}-32y^{-2}+y^{-1}+56+y+32y^2-y^3+4y^4)q^2+\dots \\ & \quad y^{-1}+y+(-y^5+y^{-1}+y-y^5)q+(-y^7-y^{-5}+2y^{-1}+2y-y^5-y^7)q^2+\dots \end{aligned}$$