

Title: Everpresent Lambda in CosmoMC

Speakers: Nosiphiwo Zwane

Collection: Everpresent Lambda: Theory Meets Observations

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# Everpresent $\Lambda$ in CosmoMC

Nosiphiwo Zwane

University of Swaziland

November 13, 2019

work with N. Afshordi and R. Sorkin

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Nosiphiwo Zwane

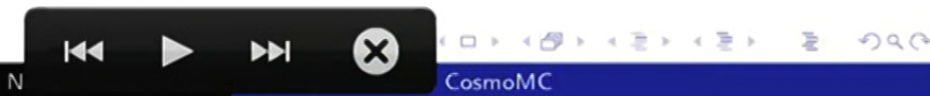
Everpresent  $\Lambda$  in CosmoMC

1:34



## Outline

- Everpresent  $\Lambda$  Models
- Everpresent  $\Lambda$  and Cosmic Microwave Background
- Baryon Acoustic Oscillations (BAO) Measurements



# Model 1

Proposed by Ahmed et al.



- The universe is assumed to be spatially homogeneous.
- In this model dark energy undergoes a random walk as described by these equations

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{1}{3}\rho_m(t) - \frac{\rho_\Lambda(t)}{3},$$
$$V(t) = \frac{4\pi}{3} \int_0^t dt' a(t')^3 \left( \int_0^{t'} dt'' \frac{1}{a(t'')} \right)^3 \text{ for } k=0,$$
$$N(t) = V(t)/\ell_p^4,$$
$$\rho_\Lambda(t + \Delta t) = \frac{S(t) + \alpha\xi\sqrt{N(t - \Delta t) - N(t)}}{V(t)}.$$

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Ahmed, Dodelson, Greene, Sorkin, Phys. Rev. D 69 (May, 2004) 103523.



## Model 2

- Simpler to simulate.
- The universe is assumed to be spatially homogeneous.
- $\Lambda$  is again assumed to be a random function of cosmic time.
- Two parameters are introduced:  $\alpha$  – controls the magnitude of fluctuations and  $\mu$  – controls the coherence time.
- $\Lambda$  is a Gaussian process with the correlation at different times being

$$\langle \hat{\Omega}_{DE}(\lambda_1) \hat{\Omega}_{DE}(\lambda_2) \rangle = \alpha^2 e^{-\frac{\mu}{2}(\lambda_1 - \lambda_2)^2}$$

where

$$\lambda = \log(a)$$

$\alpha$  = characteristic scale of the fluctuations

$\mu^{-1/2}$  = characteristic e-fold

$$\Omega_{DE} = \tanh(\hat{\Omega}_{DE})$$



## Everpresent $\Lambda$

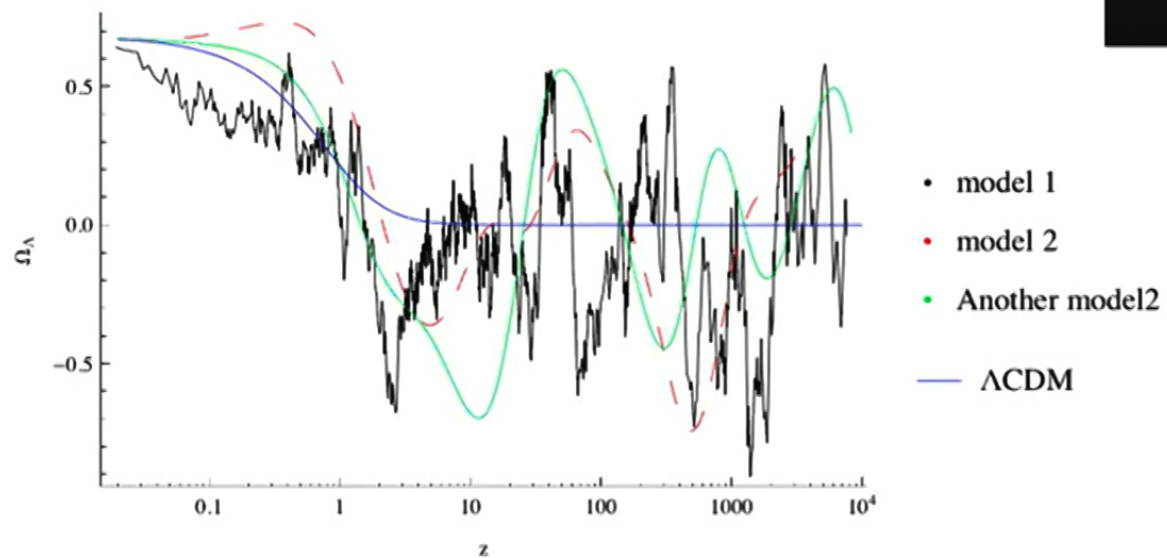
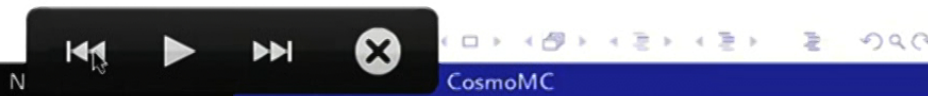


Figure :  $\Omega_{DE}$  for different Cosmological Models.



## Everpresent $\Lambda$ and the CMB



- Cosmic Microwave Background (CMB) anisotropies currently provide the most precise tests of cosmological models.
- To fit our Everpresent  $\Lambda$  model, we use CosmoMc together with CAMB.
- Model is stochastic.
- We assume the universe is flat.



## Simulations and Analysis



Table : Table for Likelihoods

	$\Lambda$ CDM	Everpresent $\Lambda$ (Model 2)
CMB: BKPLANCK	45.117	44.336
CMB: lensing	12.157	13.030
plik	1164.783	1165.278
lowTEB	10098.485	10097.864
BAO: 6DF	0.087	0.291
BAO: MGS	0.927	1.217
BAO: DR11CMASS	2.856	3.015
BAO: DR11LOWZ	1.098	1.142
BAO-Busca: DR11LyaAuto	4.265	2.636
BAO-Andreu: DR11LyaCross	4.748	4.245
Total	11332	11334



## Everpresent $\Lambda$ and the CMB

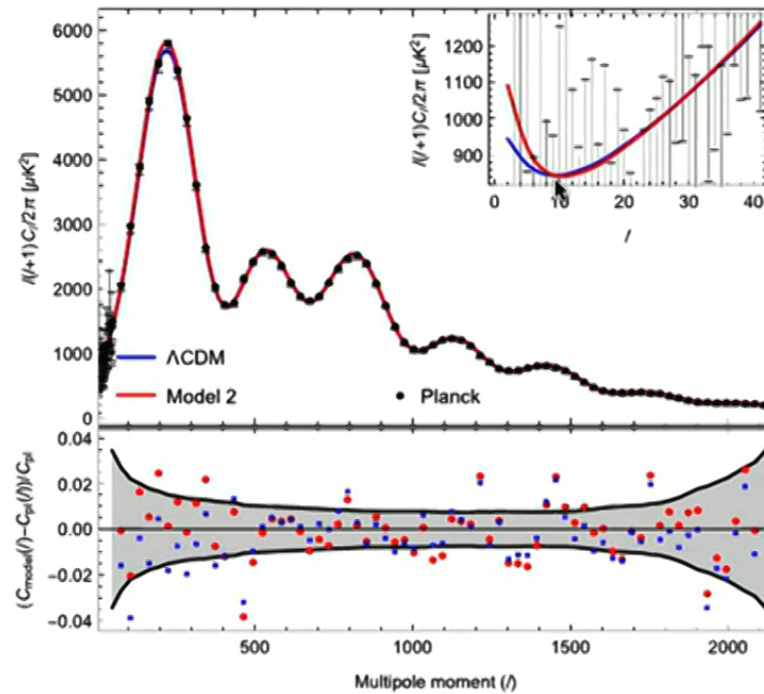


Figure : Temperature fluctuations for  $\Lambda$ CDM model ( $\chi^2 = 11334$ ), everpresent  $\Lambda$  model 2 ( $\chi^2 = 11335$ ) and Planck 2015 data.



## Everpresent $\Lambda$

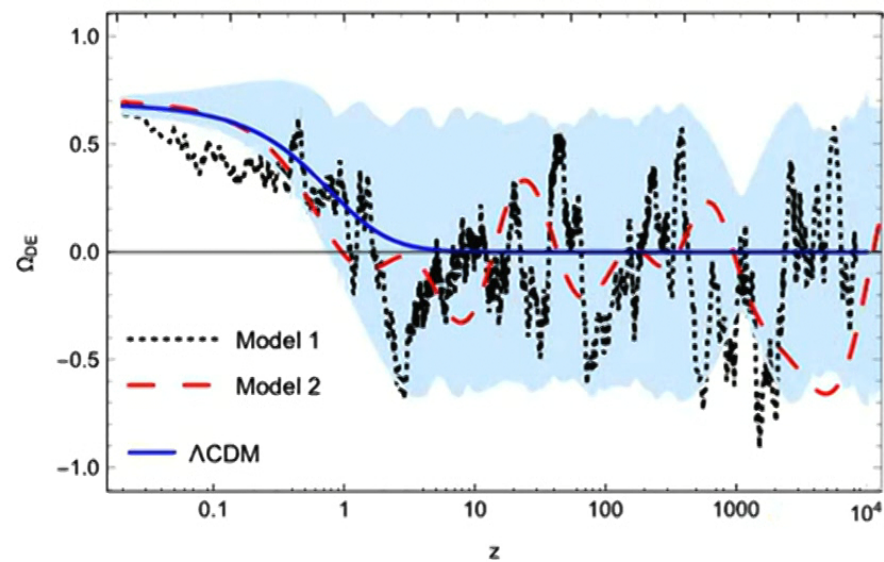


Figure :  $\Omega_{DE}$  for different Cosmological Models.

## $\chi^2$ for model



$\chi^2$  probability distribution for histories of dark energy

$$p(\chi^2(\tilde{\Omega}_{de})) = \propto D\tilde{\Omega}_{de} \exp\left(-\frac{\chi^2(\tilde{\Omega}_{de})}{2}\right) \delta(\chi^2(\tilde{\Omega}_{de}) - \chi_{red}^2(\tilde{\Omega}_{de}))$$

where

$$\chi_{red}^2 = \frac{\chi^2}{\sum_{\omega} \#\omega} \quad \text{and} \quad \chi^2 = \sum_{\omega} \frac{|\tilde{\Omega}_{de}(\omega)|^2}{\langle |\tilde{\Omega}_{de}(\omega)|^2 \rangle}$$

To check how often one can get such a good fit from everpres  $\Lambda$ , we calculate reduced  $\chi^2$  for the good fit.

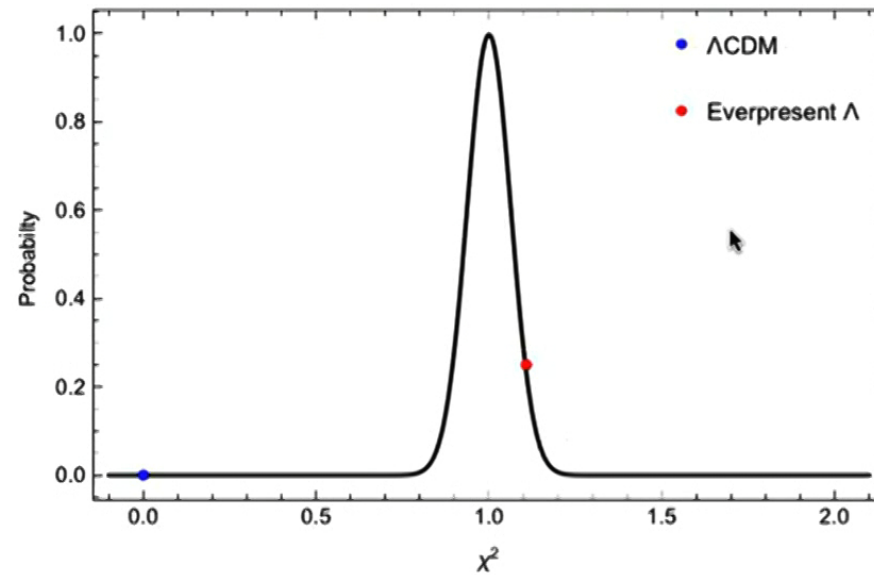


Figure : Probability of  $\chi^2$  for the models

## Baryon Acoustic Oscillations



- Baryon acoustic oscillations (BAO) matter clustering provides a standard ruler for length scales in cosmology.
- BAO can be used to measure the acceleration of the expansion of the universe (or dark energy) by comparing the sound horizon at recombination with the sound horizon,  $r_d$ , at different time.

## Baryon Acoustic Oscillations

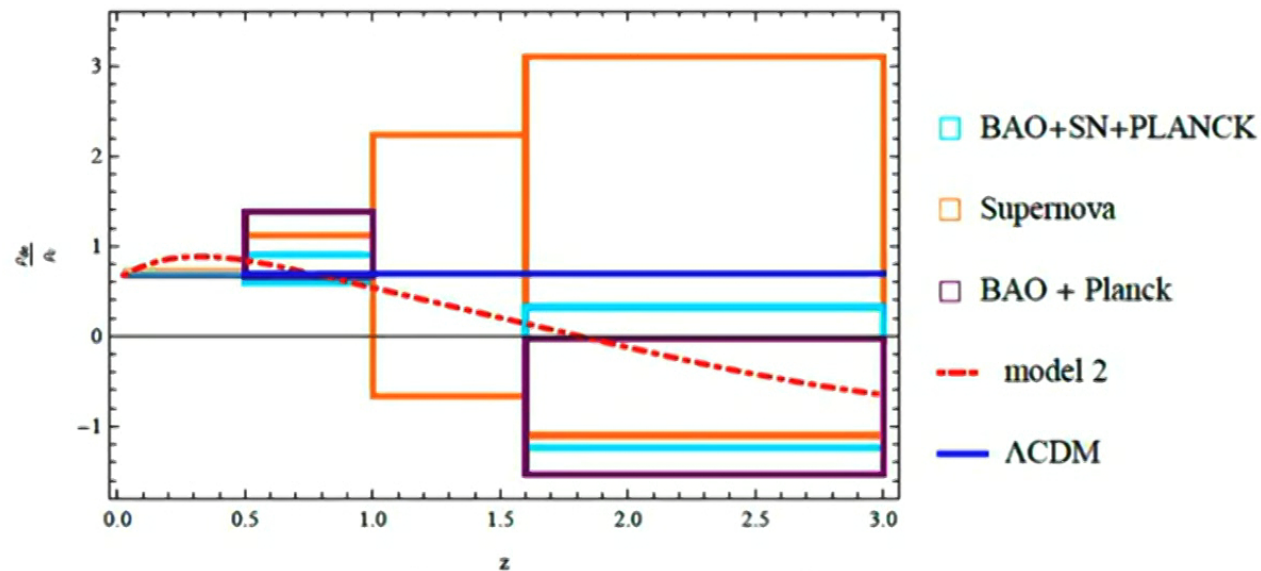


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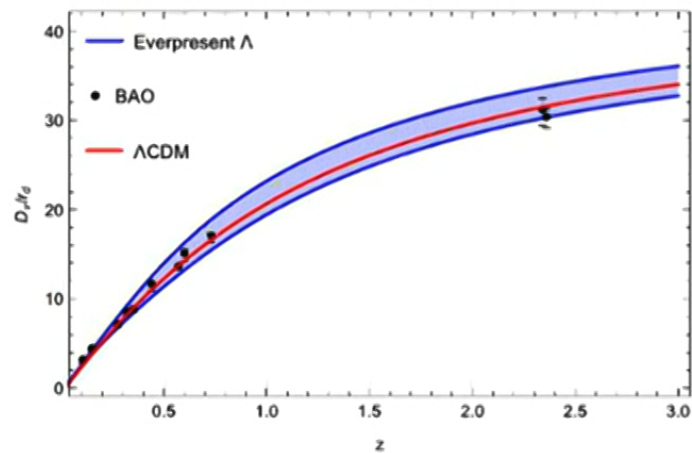
# Baryon Acoustic Oscillations (BAO)

- Baryon Oscillation Spectroscopic Survey (BOSS) 2013 results suggest that  $\Omega_{DE}(z)$  is negative at  $z = 2.34$  at  $\sim 2.5\sigma$  tension with standard  $\Lambda$ CDM.
- A number of other models that Aubourg *et al.* examined fail to fit BAO data, unless one assumes that  $\Omega_{DE} < 0$  at  $z \sim 2 - 3$ .



CosmoMC

# Baryon Acoustic Oscillations (BAO)



	$z$	Distance(Mpc)		$z$	Distance(Mpc)
6dF ( $D_v$ ) [25]	0.106	$457 \pm 27$	SDSS DR9 LRG ( $D_A$ ) [26]	0.57	$1386 \pm 45$
SDSS DR7 ( $D_v$ ) [27]	0.15	$664 \pm 25$	WiggleZ ( $D_v$ ) [28]	0.6	$2221 \pm 101$
SDSS DR7+2dF ( $D_v$ ) [29]	0.275	$1059 \pm 27$	WiggleZ ( $D_v$ ) [28]	0.73	$2516 \pm 86$
SDSS DR11 ( $D_v$ ) [30]	0.32	$1264 \pm 25$	Lya auto-corr ( $D_A$ ) [15]	2.34	$1662 \pm 96$
SDSS DR7 LRG ( $D_v$ ) [31]	0.35	$1308 \pm 25$	Lya auto-corr ( $D_H$ ) [14]	2.36	$226 \pm 8$
WiggleZ ( $D_v$ ) [28]	0.44	$1716 \pm 83$	Lya auto-corr ( $D_A$ ) [14]	2.36	$1590 \pm 60$

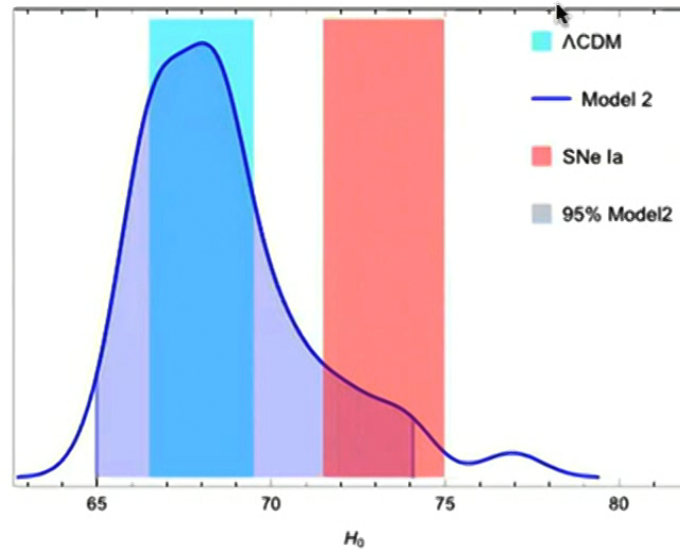
## Summary

- There are histories of dark energy that are a good fit to the CMB, these histories have  $\Omega_{de} \approx 0.7$  at  $z \approx 0$  and  $\Omega_{de} \approx 0$  at  $z \approx 1000$ .
- There are histories of dark energy that are a good fit to the BAO.
- If  $\Lambda$  really is fluctuating and “everpresent”, should become more clearly evident as observations accumulate for high redshift.



## The Value of Hubble Constant today $H_0$

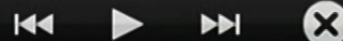
In the past few years, we have seen a divergence between the local measurements of the Hubble constant and the best fits inferred from Planck CMB and BAO observations.



- Planck Collaboration, arXiv:1502.0158.

- A. G. Riess *et al.*, arXiv:1604.0142.

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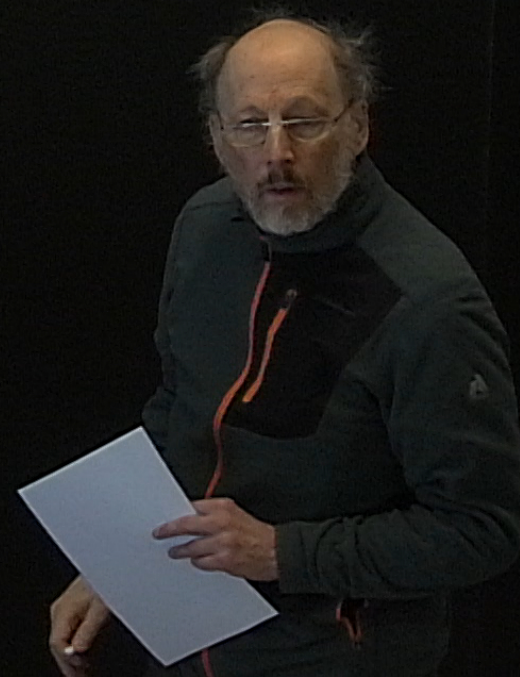
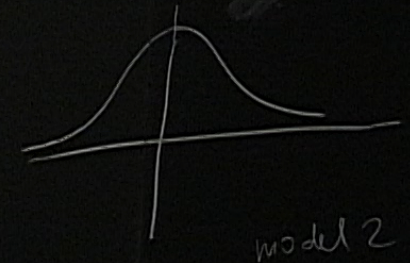
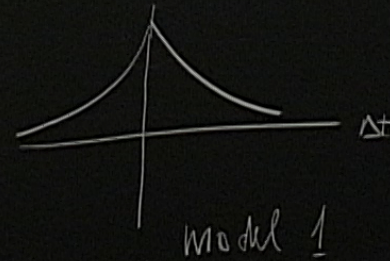
Simulations vs Observations

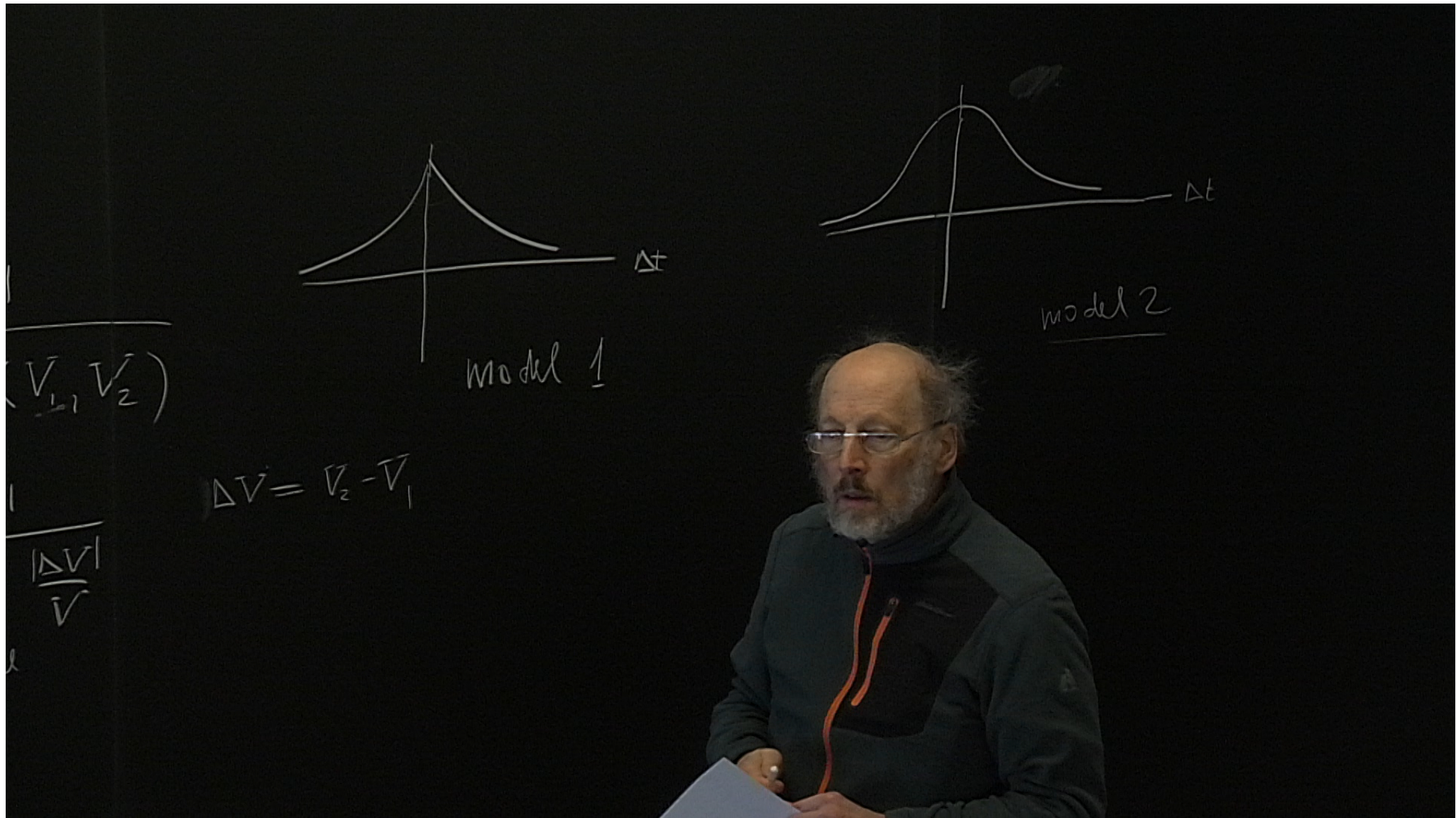
$$\langle \Lambda(t_1) \Lambda(t_2) \rangle = \frac{1}{\max(V_1, V_2)}$$

$$\frac{\langle \Lambda_1 \Lambda_2 \rangle}{\langle \Lambda^2 \rangle} \sim \frac{1}{1 + \frac{|\Delta V|}{V}}$$

small

$$\Delta V = V_2 - V_1$$



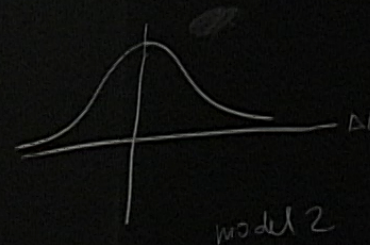
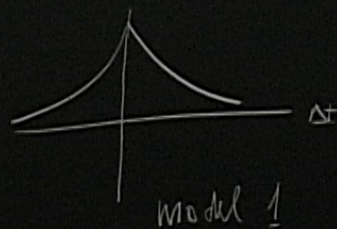


$$\langle \Lambda(t_1) \Lambda(t_2) \rangle = \frac{1}{\max(V_1, V_2)}$$

$$\frac{\langle \Lambda_1 \Lambda_2 \rangle}{\langle \Lambda_{12}^2 \rangle} \sim \frac{1}{1 + \frac{|\Delta V|}{V}}$$

small

$$|\Delta V| = V_2 - V_1$$

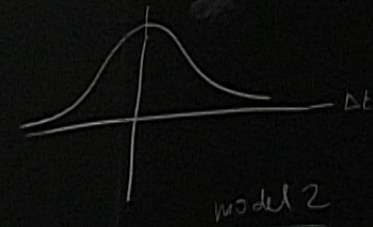
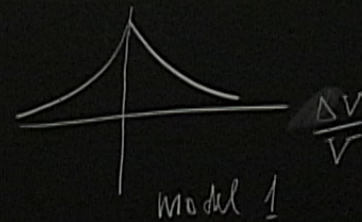


$$\langle \Lambda(t_1) \Lambda(t_2) \rangle = \frac{1}{\max(V_1, V_2)}$$

$$\frac{\langle \Lambda_1 \Lambda_2 \rangle}{\langle \Lambda_{1/2}^2 \rangle} \sim \frac{1}{1 + \frac{|\Delta V|}{V}}$$

small

$$\Delta V = V_2 - V_1$$



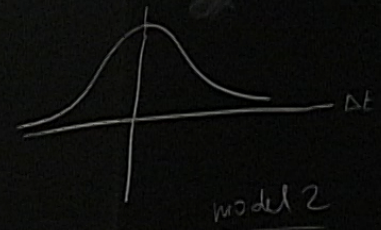
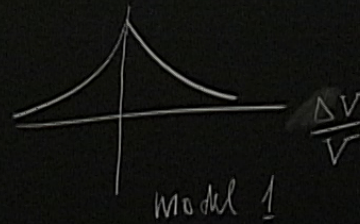
$$\Lambda(V)$$

$$\langle \Lambda(t_1) \Lambda(t_2) \rangle = \frac{1}{\max(V_1, V_2)}$$

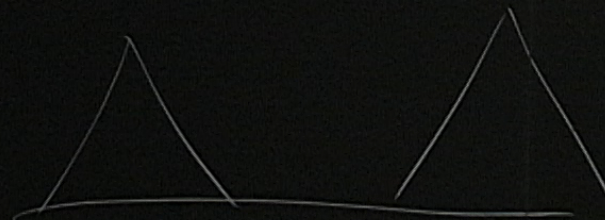
$$\frac{\langle \Lambda_1 \Lambda_2 \rangle}{\langle \Lambda_{1/2}^2 \rangle} \sim \frac{1}{1 + \frac{|\Delta V|}{V}}$$

small

$$\Delta V = V_2 - V_1$$



$$\Lambda(V)$$



$$\langle \Lambda(t_1) \Lambda(t_2) \rangle = \frac{1}{\max(V)}$$

$$\frac{\langle \Lambda_1 \Lambda_2 \rangle}{\langle \Lambda_{12}^2 \rangle} \sim \frac{1}{1 + \frac{1}{2}}$$

small