

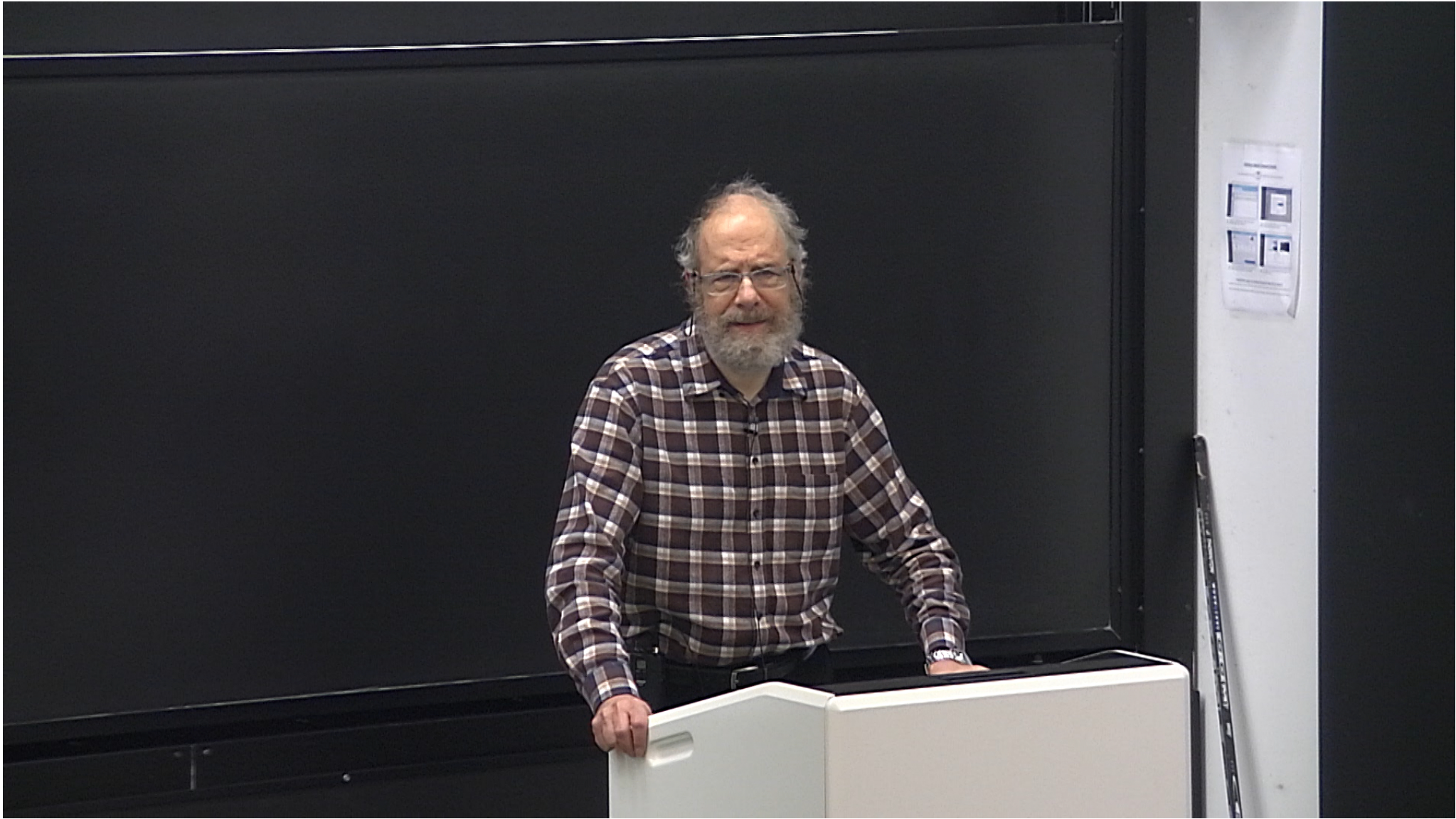
Title: PSI 2019/2020 - Quantum Field Theory II (David) - Lecture 1

Speakers: Francois David

Collection: PSI 2019/2020 - Quantum Field Theory II (David)

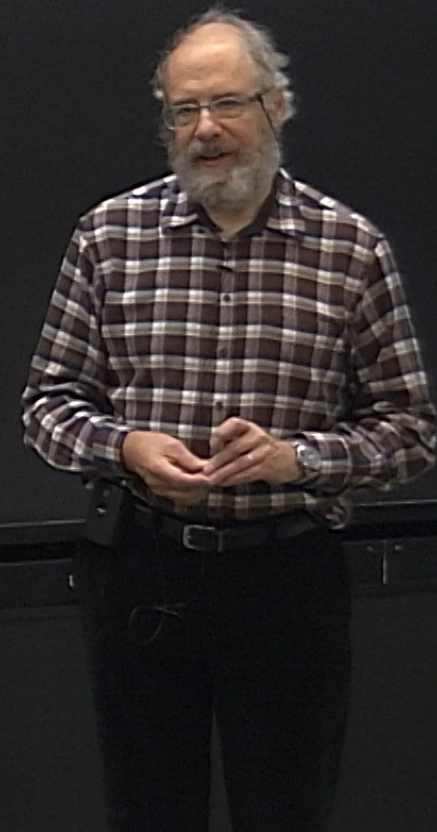
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- QFT I canonical quantization
- Quantum Theory Path Integral Quantization



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- Quantum Theory Path Integral Quantization

* QFT Functional Integral Quantization

* Renormalization Theory ← specific to QFT
 ↑ Stat. Mech & Cond. Matter ← K. Wilson

- QFT I canonical quantization

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* QFT Functional Integral Quantization

* Renormalization Theory ← specific to QFT

K. Wilson
Wilsonian Ren. Group

* Gauge Theories Non-abelian gauge theories → Higgs Mechanism

↑ Stat. Mech & Cond. Matter ←

Non relativistic QM.

Classical Theory

1 particle mass m in a potential $V(q)$: eg. harmonic oscillator

position $q(t)$ $t = \text{time}$ 1-dimensional

Velocity $\dot{q}(t) = \frac{dq(t)}{dt}$, momentum $p = m \dot{q}$

$$V(q) = \frac{m}{2} \omega^2 q^2$$

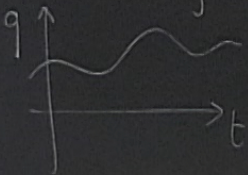
Hamiltonian

$$H(q, p) = \frac{p^2}{2m} + V(q)$$

$$\dot{q} = \frac{\delta H}{\delta p}$$

Lagrangian

Action \leftarrow trajectory



$$S[q] = \int dt \mathcal{L}(q, \dot{q}) \quad \mathcal{L}(q, \dot{q}) = \frac{m}{2} \dot{q}^2 - V(q)$$

$$S = \int dt [p \dot{q} - H(q, p)] \leftarrow \text{Legendre transformation}$$

q) : e.g. harmonic oscillator

$$V(q) = \frac{m}{2} \omega^2 q^2$$

Quantum $p \rightarrow \hat{p}$; $q \rightarrow \hat{q}$

Hamiltonian $\hat{H} = \frac{\hat{p}^2}{2m} + V(\hat{q})$

commutation relation $[\hat{q}, \hat{p}] = i\hbar$

Evolution operator

$$|\Psi_0\rangle \xrightarrow{t} |\Psi_t\rangle = \hat{U}(t) |\Psi_0\rangle,$$

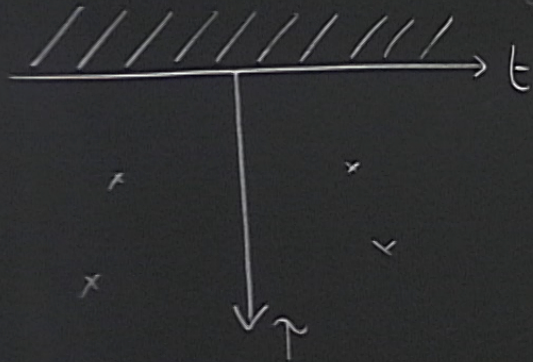
$t=0$ quantum states $t=et$

$$\hat{U}(t) = \exp\left(\frac{t}{i\hbar} \hat{H}\right)$$

\Leftrightarrow Schrödinger equation
(Schrödinger picture)

t functional of $t \rightarrow q(t)$ Lagrangian $S = \int dt [p \cdot \dot{q} - H(q, p)]$

"imaginary time"

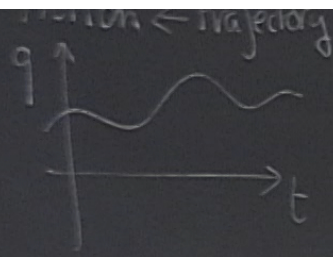


$$t = -i\tau \quad \tau = t_E \text{ "Euclidean" time}$$

$$\hat{U}(-i\tau) = \exp\left(-\frac{\tau}{\hbar} \hat{H}\right) = \hat{U}_E(\tau)$$

unitary operator

Hermitian, positive



action ← trajectory

$$S[q] = \int dt \mathcal{L}(q, \dot{q}) \quad \mathcal{L}(q, \dot{q}) = \frac{m}{2} \dot{q}^2 - V(q)$$

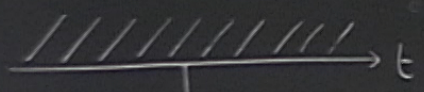
↑
functional of $t \rightarrow q(t)$ ↑ Lagrangian

$$S = \int dt [p \cdot \dot{q} - H(q, p)] \leftarrow \text{Legendre transformation}$$

$t=0$
quantum state
Physics

"imaginary time"

$$t = -i\tau \quad \tau = t_E \text{ "Euclidean" time} \quad \tau \geq 0$$



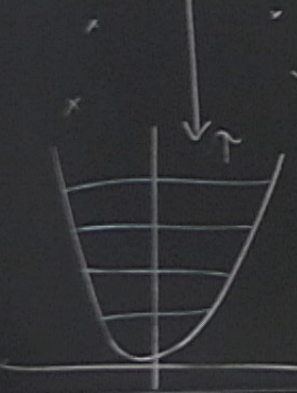
$$\hat{U}(-i\tau) = \exp\left(-\frac{\tau}{\hbar} \hat{H}\right) = \hat{U}_E(\tau) = \sum_{i=0}^{\infty} \exp\left(-\frac{\tau}{\hbar} E_i\right) |i\rangle\langle i|$$

↑
unitary operator

Hermitian, positive definite

↑
smaller and smaller
as long as $\tau \geq 0$

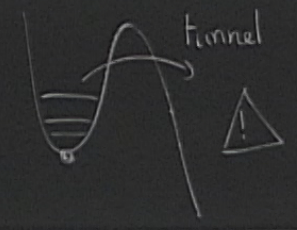
\hat{H} spectrum bounded from below, and infinite



energy e. states $|n\rangle \quad n=0,1,\dots$

$$E_n \quad E_0 < E_1 < E_2 < \dots$$

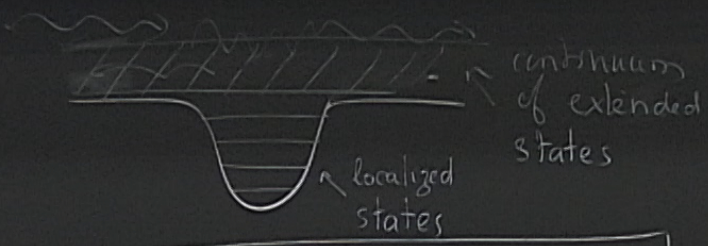
$$\hat{H} = \sum_i E_i |i\rangle\langle i|$$



$\hat{U}_E(\tau)$ Bounded operator
(operator algebra theory)
Von Neumann, Banach, etc...

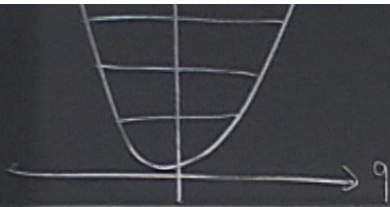
(q, p) ← Legendre transformation
 Physical (real) time
 quantum states
 $t=0$ time t
 (Schrödinger picture)

$\exp(-\frac{\tau}{\hbar} E_i) |i\rangle\langle i|$
 smaller and smaller
 as long as $\tau \geq 0$
 bounded from below, and infinite



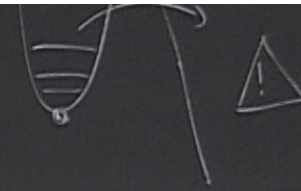
OK too
 $U(t)$ unbounded if $\text{Im}(t) > 0$
 bounded if $\text{Im}(t) \leq 0$

$\hat{U}_E(\tau)$ Bounded operator
 (operator algebra theory)
 von Neumann, Borchs, etc.



$$E_n \quad E_0 < E_1 < E_2 < \dots$$

$$\hat{H} = \sum_i E_i |i\rangle\langle i|$$

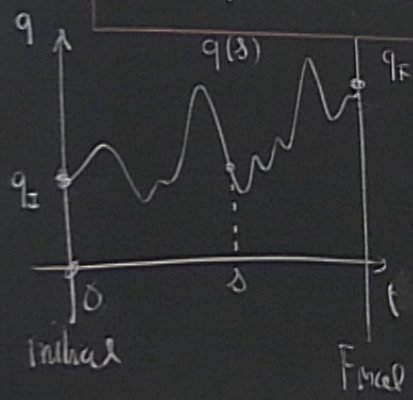


(operator
von Neuman)

Path integrals (Maité)
probability amplitude

Feynman

$$\langle q_F | U(t) | q_I \rangle = \int D[q(s)] \exp\left(\frac{i}{\hbar} S[q]\right)$$



↑
Trajectory
(history)

↑
measure on
trajectories

$q(0) = q_I$
 $q(t) = q_F$

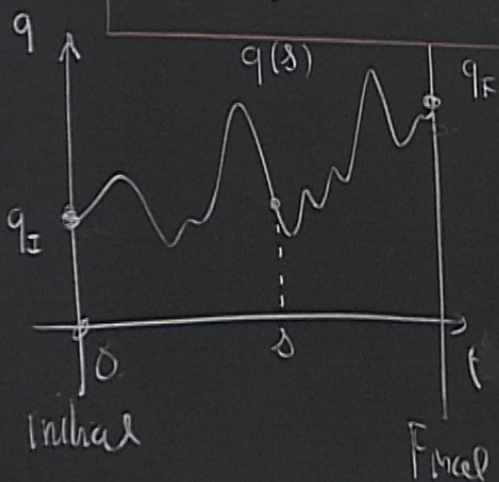


$$H = \sum_i E_i |i\rangle\langle i|$$

Path integrals (Maître)
probability amplitude

Feynman

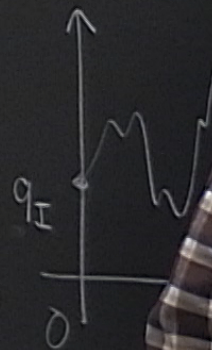
$$\langle q_F | U(t) | q_I \rangle = \int \mathcal{D}[q(s)] \exp\left(\frac{i}{\hbar} S[q]\right)$$



↑
trajectories
(history)
 $q(0) = q_I$
 $q(t) = q_F$

↑
measure on
trajectories

↑
phase



(operator algebra theory)
 von Neumann, Banach, etc...

Imaginary time

positive weight

$$\langle q_F | U_E(\tau) | q_I \rangle = \int D[q(\sigma)] \exp\left(-\frac{1}{\hbar} S_E[q]\right)$$

\uparrow Euclidean time \uparrow Euclidean Action

$$\left(\frac{1}{\hbar} S[q]\right)$$

phase quantum

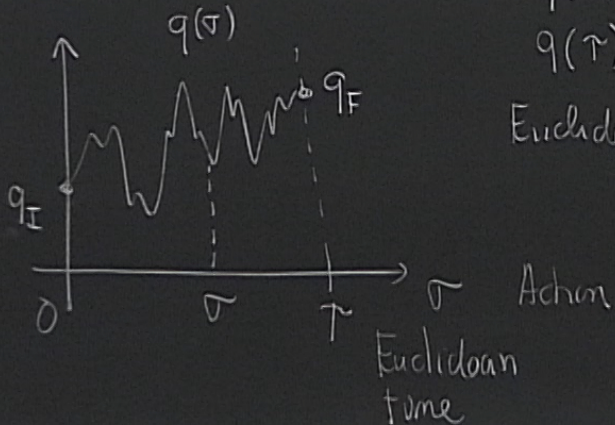
$$q(0) = q_I$$

$$q(\tau) = q_F$$

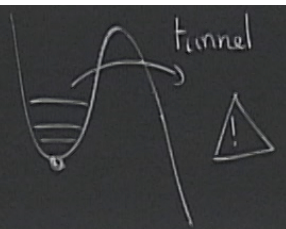
Euclidean action

$$S_E[q] = \int_0^\tau d\sigma \left[\frac{m}{2} \dot{q}^2 + V(q) \right]$$

$$S[q] = \int_0^t d\sigma \left[\frac{m}{2} \dot{q}^2 - V(q) \right]$$



$n=0,1,2,\dots$
 $E_n \quad E_0 < E_1 < E_2 < \dots$
 $\hat{H} = \sum_i E_i |i\rangle\langle i|$



$\hat{U}_E(\tau)$ Bounded operator
 (operator algebra theory)
 von Neumann, Banach, etc...

bounded if $\text{Im}(t) \leq 0$

Maité) Feynman

$$|i\rangle = \int D[q(\sigma)] \exp\left(\frac{i}{\hbar} S[q]\right)$$

Trajectories (history) →
 measure on trajectories →
 phase quantum →
 Interferences ←

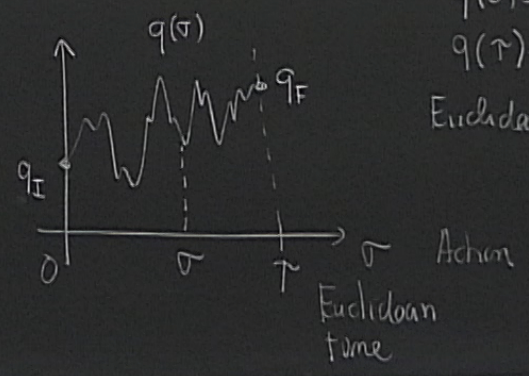
$q(0) = q_I$
 $q(t) = q_F$

Imaginary time

$$\langle q_F | U_E(\tau) | q_I \rangle = \int D[q(\sigma)] \exp\left(-\frac{1}{\hbar} S_E[q]\right)$$

Euclidean time (under $D[q(\sigma)]$)
 Euclidean Action (under $S_E[q]$)

$q(0) = q_I$
 $q(\tau) = q_F$
 Euclidean action



$$S_E[q] = \int_0^\tau d\sigma \left[\frac{m}{2} \dot{q}^2 + V(q) \right]$$

$$S[q] = \int_0^t d\sigma \left[\frac{m}{2} \dot{q}^2 - V(q) \right]$$

$\rho_E(\tau)$ Bounded operator
 operator algebra theory
 von Neumann, Borchs, etc...

bounded by $\rho_H(\tau) \leq 0$

positive weight \rightarrow probability theory on trajectories

$$\langle \dots \rangle = \int D[q(\sigma)] \exp\left(-\frac{1}{\hbar} S_E[q]\right)$$

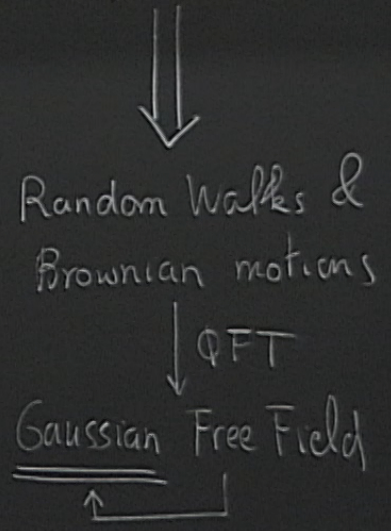
\uparrow Euclidean time \nwarrow Euclidean Action

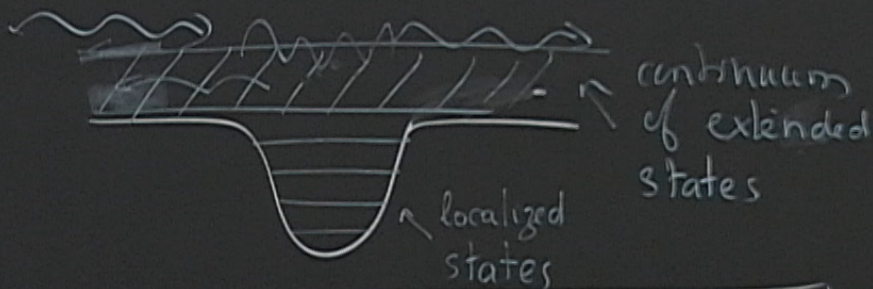
$q(0) = q_E$
 $q(\tau) = q_F$
 Euclidean action

$$S_E[q] = \int_0^\tau d\sigma \left[\frac{m}{2} \dot{q}^2 + V(q) \right]$$

$$S[q] = \int_0^\tau d\sigma \left[\frac{m}{2} \dot{q}^2 - V(q) \right]$$

$\rightarrow \sigma$ Action
 Euclidean time



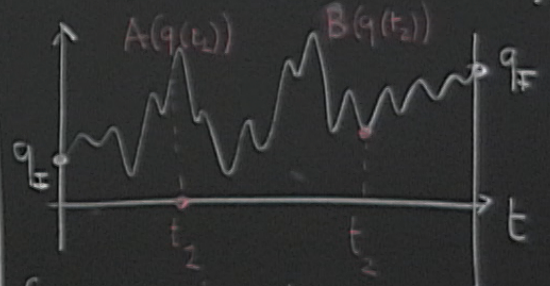


$|i\rangle\langle i|$
 smaller and smaller
 long as $\tau \geq 0$
 low, and infinite

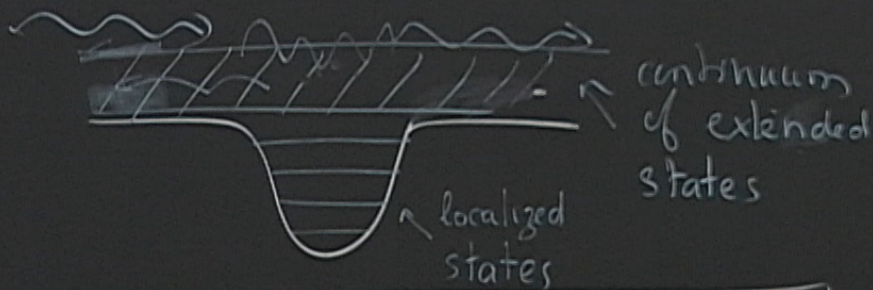
OK too
 $U(t)$ unbounded if $\text{Im}(t) > 0$
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Bounded operator
 operator algebra theory
 von Neumann, Banach, etc...

Operators (Observables)
 function $A(q) \rightarrow \hat{A} = A(\hat{Q})$



$$\int \mathcal{D}[q] \exp\left(\frac{i}{\hbar} S[q]\right) A(q(t_1)) \cdot B(q(t_2))$$

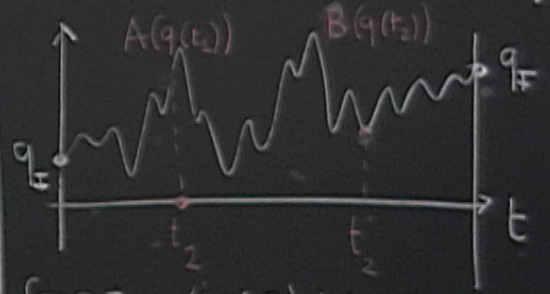


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Operators (Observables)
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$$\int \mathcal{D}[q] \exp\left(\frac{i}{\hbar} S[q]\right) A(q(t_1)) \cdot B(q(t_2))$$

$$\langle q_{\text{F}} | \hat{U}(t_1) \hat{B} \hat{U}(t_2) \hat{A} \hat{U}(t_1) | q_{\text{I}} \rangle$$

Matrix element in Schrödinger representation

Heisenberg Picture

$$|q_I\rangle = |q_I; 0\rangle$$

$$|q_F\rangle = |q_F; t\rangle$$

$$\hat{U}(-t_1)\hat{A}\hat{U}(t_1) = \hat{A}(t_1)$$

$$\hat{U}(-t_2)\hat{B}\hat{U}(t_2) = \hat{B}(t_2)$$

$$\Rightarrow \langle q_F, t_1 | \hat{B}(t_2)\hat{A}(t_1) | q_I, 0 \rangle$$

Hersenberg Picture

$$|q_I\rangle = |q_I; 0\rangle$$

$$\hat{U}(t)|q_F\rangle = |q_F; t\rangle \Rightarrow \langle q_F, t| = \langle q_F| \hat{U}(t)$$

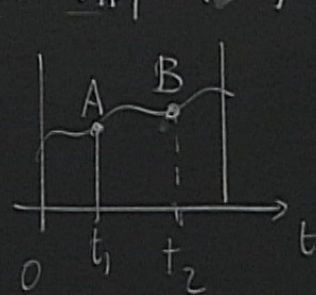
$$\hat{U}(t_1) \hat{A} \hat{U}(t_1) = \hat{A}(t_1)$$

$$\hat{U}(t_2) \hat{B} \hat{U}(t_2) = \hat{B}(t_2)$$

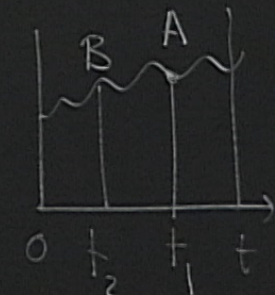
$$\Rightarrow \langle q_F, t_1| \hat{B}(t_2) \hat{A}(t_1) |q_I, 0\rangle$$

$$\hat{U}(t_2 - t_1) = \hat{U}(t_2) \hat{U}(-t_1)$$

$$\hat{U}(t - t_2) = \hat{U}(t) \hat{U}(-t_2)$$



$$0 < t_1 < t_2 < t$$



$$0 < t_2 < t_1 < t$$

Hersenberg Picture

$$|q_I\rangle = |q_I; 0\rangle$$

$$\hat{U}(t)|q_F\rangle = |q_F; t\rangle \Rightarrow \langle q_F, t| = \langle q_F| \hat{U}(t)$$

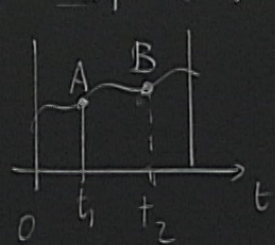
$$\hat{U}(t_1)\hat{A}\hat{U}(t_1) = \hat{A}(t_1)$$

$$\hat{U}(t_2)\hat{B}\hat{U}(t_2) = \hat{B}(t_2)$$

$$\Rightarrow \langle q_F, t | \prod [\hat{B}(t_2)\hat{A}(t_1)] | q_I, 0 \rangle$$

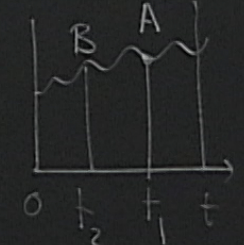
$$\hat{U}(t_2 - t_1) = \hat{U}(t_2) \cdot \hat{U}(-t_1)$$

$$\hat{U}(t - t_2) = \hat{U}(t) \cdot \hat{U}(-t_2)$$



$$0 < t_1 < t_2 < t$$

$$\hat{B}(t_2)\hat{A}(t_1)$$



$$0 < t_2 < t_1 < t$$

$$\hat{A}(t_1)\hat{B}(t_2)$$

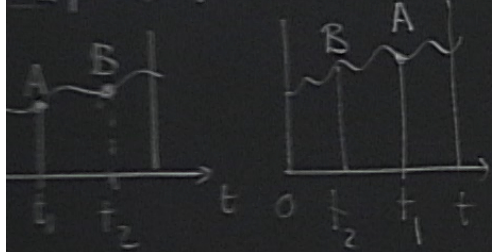
Path Integral

Time ordered product of \hat{A} and \hat{B}

$$\langle t_2 - t_1 \rangle = \hat{U}(t_2) \cdot \hat{U}(-t_1)$$

$$\langle t - t_0 \rangle = \hat{U}(t) \cdot \hat{U}(-t_0)$$

$$\langle a_{\pm} | \hat{U}(t) | b_{\pm} \rangle$$



$$\langle t_1 \rangle < \langle t_2 \rangle < t$$

$$0 < t_2 < t_1 < t$$

$$\langle t_1 \rangle \hat{A}(t_1)$$

$$\hat{A}(t_1) \hat{B}(t_2)$$

If \hat{A} and \hat{B} and \hat{H} do not commute
 Commutation relations are "hidden" in P.Int
 Path Integral

Ex: $\hat{Q} \leftrightarrow q(t) \quad \hat{P} = m \frac{d}{dt} \hat{Q}(t) \rightarrow m \dot{q}(t)$

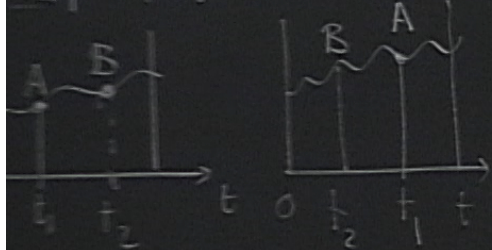
⇓

Time ordered product
 of \hat{A} and \hat{B}

$$\hat{U}(t_2 - t_1) = \hat{U}(t_2) \cdot \hat{U}(-t_1)$$

$$\hat{U}(t - t_0) = \hat{U}(t) \cdot \hat{U}(-t_0)$$

$$\langle a_2 | \hat{U}(t_2 - t_1)$$



$$t_1 < t_2 < t$$

$$0 < t_2 < t_1 < t$$

$$\hat{A}(t_1)$$

$$\hat{A}(t_1) \hat{B}(t_2)$$

If \hat{A} and \hat{B} and \hat{H} do not commute

Commutation relations are "hidden" in P.Int

Path Integral

Ex: $\hat{Q} \leftrightarrow q(t) \quad \hat{P} = m \frac{d}{dt} \hat{Q}(t) \rightarrow m \dot{q}(t)$

Recover $[\hat{Q}, \hat{P}] = i\hbar$
 ↳ quantum term

Time ordered product
 of \hat{A} and \hat{B}