

Title: Direct Deflection of Particle Dark Matter

Speakers: Asher Berlin

Series: Particle Physics

Date: October 15, 2019 - 1:00 PM

URL: <http://pirsa.org/19100077>

Abstract: Detecting light dark matter that interacts weakly with electromagnetism has recently become one of the benchmark goals of near-term and futuristic direct detection experiments. In this talk, I will discuss an alternative technique to directly detecting such particles below the GeV-scale. The approach involves distorting the local flow of dark matter with time-varying fields and measuring these distortions with shielded resonant detectors; such as LC circuits.

Direct Deflection of Particle Dark Matter

Asher Berlin

Perimeter
October 15, 2019

1908.06982 with R. D'Agnolo, S. Ellis, P. Schuster, N. Toro

Direct Detection Below an MeV

predictive cosmology

(freeze-in)

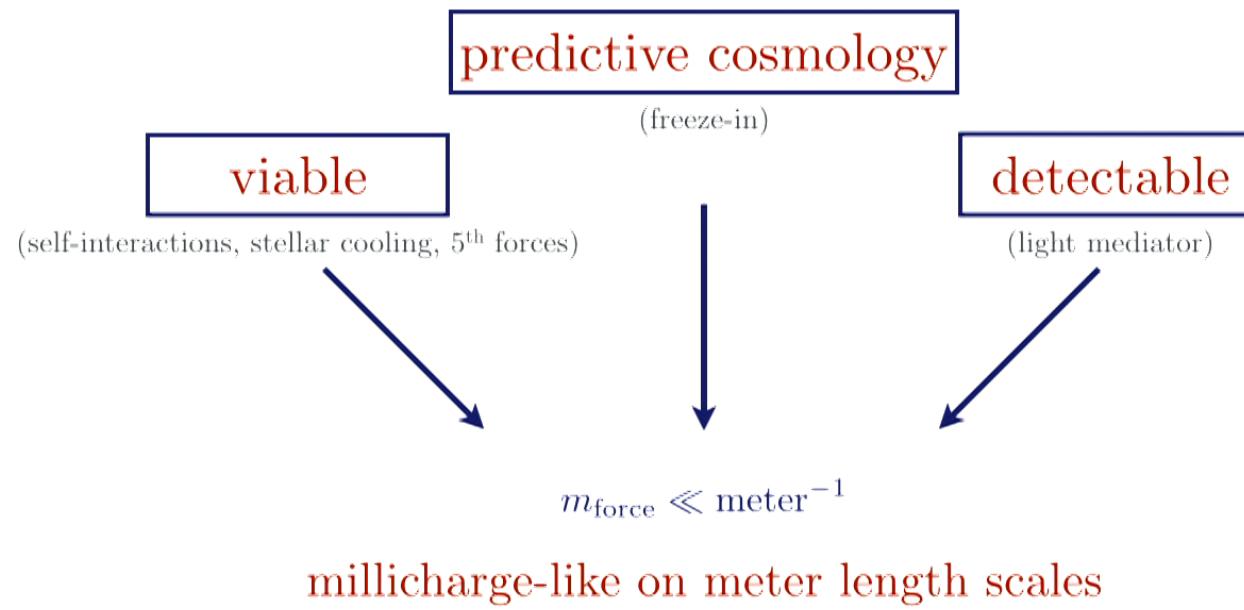
viable

(self-interactions, stellar cooling, 5th forces)

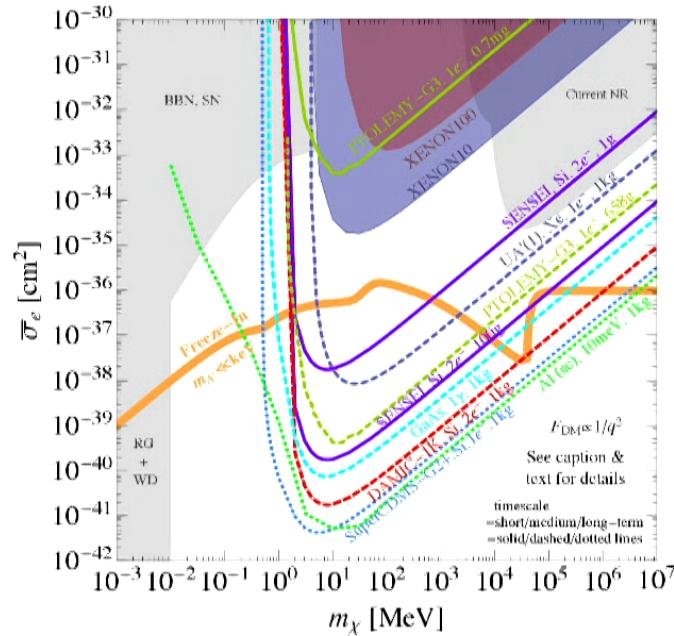
detectable

(light mediator)

Direct Detection Below an MeV



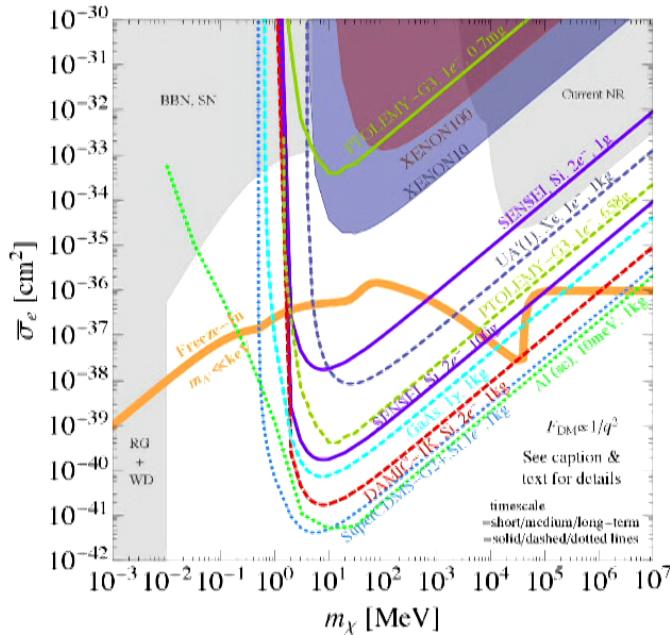
Direct Detection via Scattering



- new scattering targets
- new read-out technologies
- similar philosophy

arXiv:1707.04591

Direct Detection via Scattering



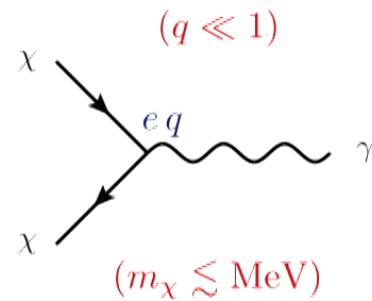
- new scattering targets
- new read-out technologies
- similar philosophy

instead, take advantage of:

small mass \rightarrow large number density, small momentum \rightarrow easier to manipulate

arXiv:1707.04591

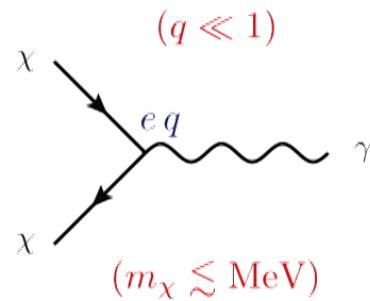
Millicharge Cosmology



χ thermalizes \implies in conflict with BBN and CMB

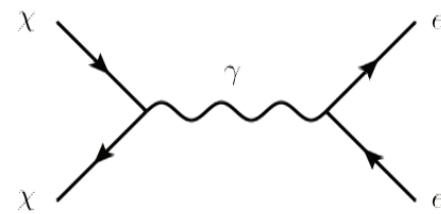
How small does q have to be?

Millicharge Cosmology



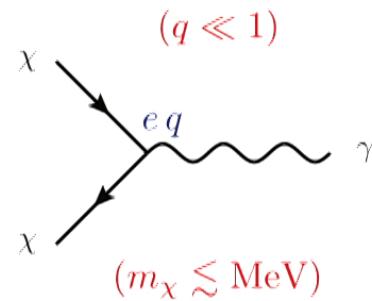
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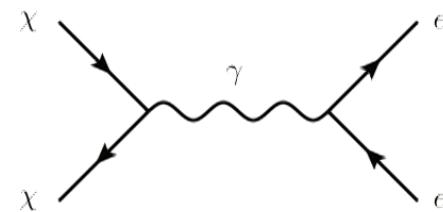
$$\Gamma \lesssim H \implies q \lesssim 10^{-9} \left(\frac{\max(m_\chi, m_e)}{\text{MeV}} \right)^{1/2}$$

Millicharge Cosmology



χ thermalizes \implies in conflict with BBN and CMB

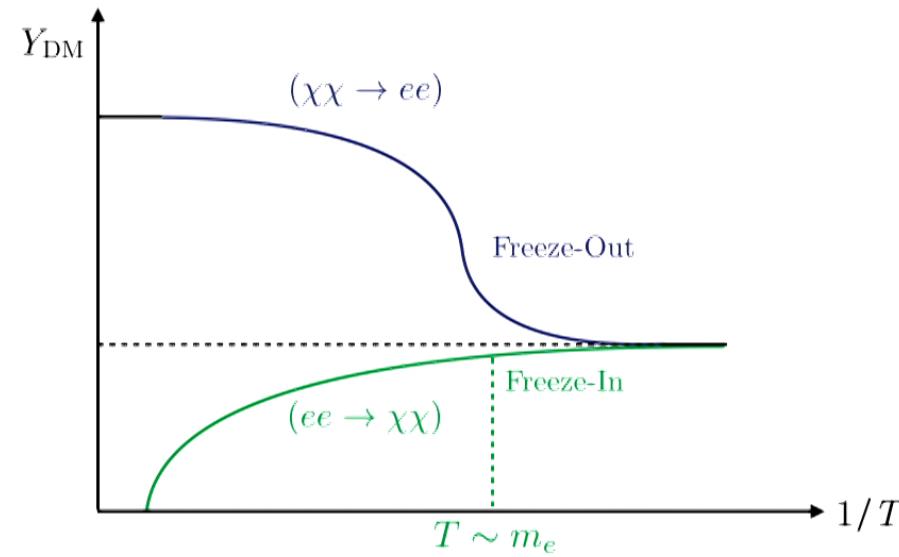
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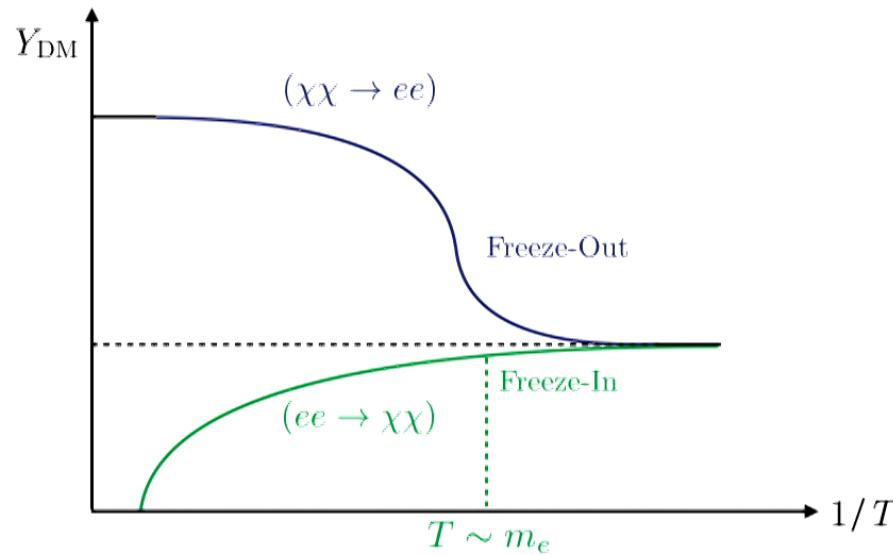
If this is the dark matter, how is it populated?

Freeze-In



arXiv:0911.1120

Freeze-In



$$\Gamma(ee \rightarrow \chi\chi) \sim \alpha_{\text{em}}^2 q^2 T, \quad n_\chi \sim n_e (\Gamma/H), \quad \rho_{\text{DM}} \sim T_{\text{eq}} T^3$$

$$\Rightarrow q \sim \frac{1}{\alpha_{\text{em}}} \left(\frac{m_e T_{\text{eq}}}{m_\chi m_{\text{pl}}} \right)^{1/2} \sim 10^{-11} \left(\frac{\text{MeV}}{m_\chi} \right)^{1/2}$$

arXiv:0911.1120

Kinetic Mixing

$$\mathcal{L} \supset -\frac{1}{4}F_{\mu\nu}^2 - \frac{1}{4}F'^2_{\mu\nu} + \frac{1}{2}m_{A'}^2 A'^2_{\mu} + \frac{\epsilon}{2} F_{\mu\nu} F'^{\mu\nu}$$

$$A_\mu \rightarrow A_\mu + \epsilon A'_\mu , \quad A'_\mu \rightarrow \frac{1}{\sqrt{1-\epsilon^2}} A'_\mu \implies \mathcal{L} \supset -\frac{1}{4}F_{\mu\nu}^2 - \frac{1}{4}F'^2_{\mu\nu} + \frac{1}{2}m_{A'}^2 A'^2_{\mu} + \mathcal{O}(\epsilon^2)$$

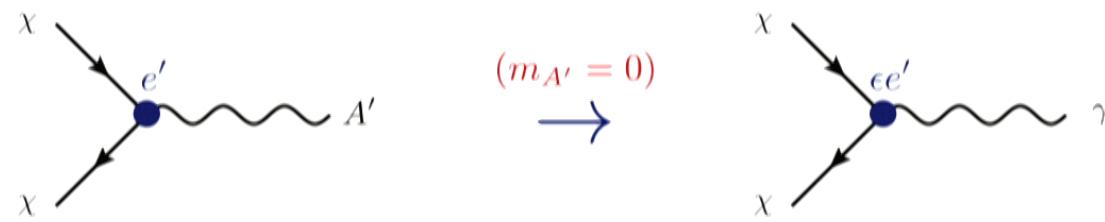
$$\mathcal{L} \supset j_\mu A^\mu + j'_\mu A'^\mu \implies \boxed{\mathcal{L} \supset j_\mu (A^\mu + \epsilon A'^\mu) + j'_\mu A'^\mu + \mathcal{O}(\epsilon^2)}$$

$$\begin{cases} A_{\text{vis}} = A + \epsilon A' & (\text{the massless photon}) \\ A_{\text{inv}} = A' - \epsilon A \end{cases}$$

$$\boxed{m_{A'} = 0 \implies \mathcal{L} \supset -\frac{1}{4}F_{\text{vis}}^2 - \frac{1}{4}F_{\text{inv}}^2 + j_\mu A_{\text{vis}}^\mu + j'_\mu (A_{\text{inv}}^\mu + \epsilon A_{\text{vis}}^\mu)}$$

Kinetic Mixing

$$\mathcal{L} \supset -\frac{1}{4}F_{\text{vis}}^2 - \frac{1}{4}F_{\text{inv}}^2 + j_\mu A_{\text{vis}}^\mu + j'_\mu (A_{\text{inv}}^\mu + \epsilon A_{\text{vis}}^\mu)$$



$$q_{\text{eff}} \sim \epsilon e' / e$$

(exact millicharge limit)

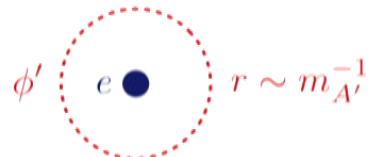
Pseudo-Millicharge

$(m_{A'} \neq 0)$

$$\mathcal{L} \supset j_\mu (A^\mu + \epsilon A'^\mu) + j'_\mu A'^\mu$$

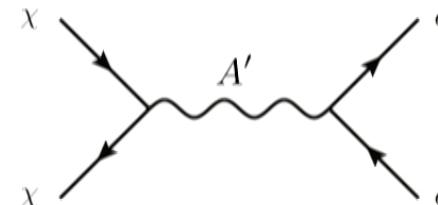
$$(\partial^2 + m_{A'}^2) A'^\mu = j'^\mu + \epsilon j^\mu$$

$$A'^\mu = (\phi', A') \implies \begin{cases} (\nabla^2 - \partial_t^2 - m_{A'}^2) \phi' = -(\rho' + \epsilon \rho) \\ (\nabla^2 - \partial_t^2 - m_{A'}^2) A' = -(j' + \epsilon j) \end{cases}$$



$$\rho' = 0, \rho = e \delta^3(x)$$

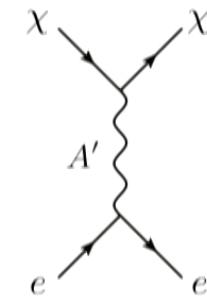
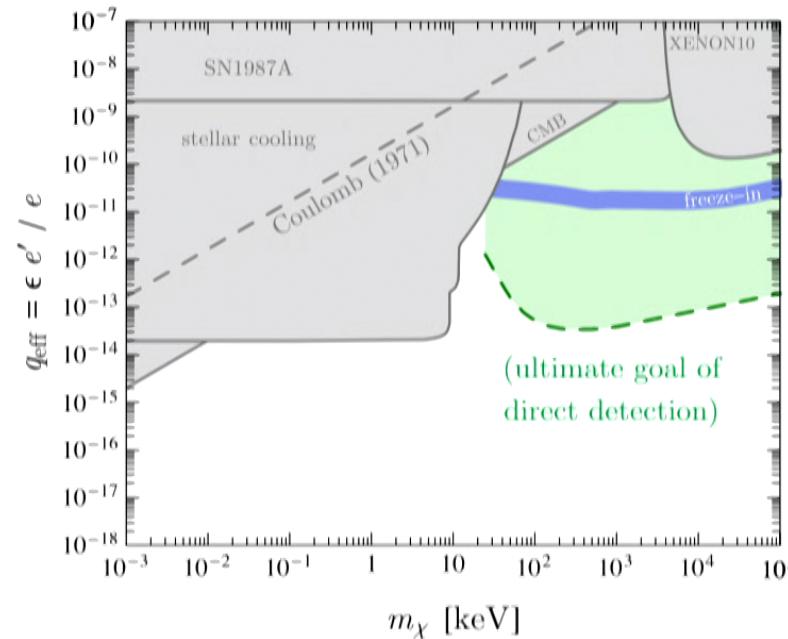
$$\implies \phi'(r) = \frac{\epsilon e}{4\pi} \frac{e^{-m_{A'} r}}{r}$$



$$m_{A'} \ll \sqrt{s} \implies \sigma \sim \frac{(e' \epsilon e)^2}{s} \sim \frac{(q_{\text{eff}} e^2)^2}{s}$$

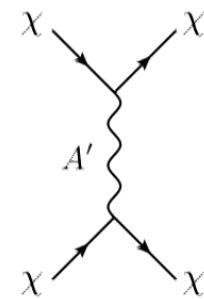
$$q_{\text{eff}} \sim \epsilon e' / e \text{ for } r \ll 1/m_{A'}$$

Parameter Space

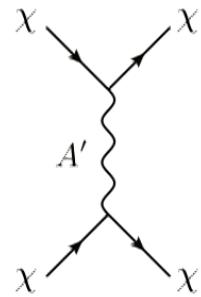


$$\sigma(\chi e \rightarrow \chi e) \propto \frac{1}{q_{\text{tr}}^4} \sim \frac{1}{(\alpha_{\text{em}} m_e)^4} \quad (m_{A'} \ll \text{keV})$$

Self-Interactions



Self-Interactions



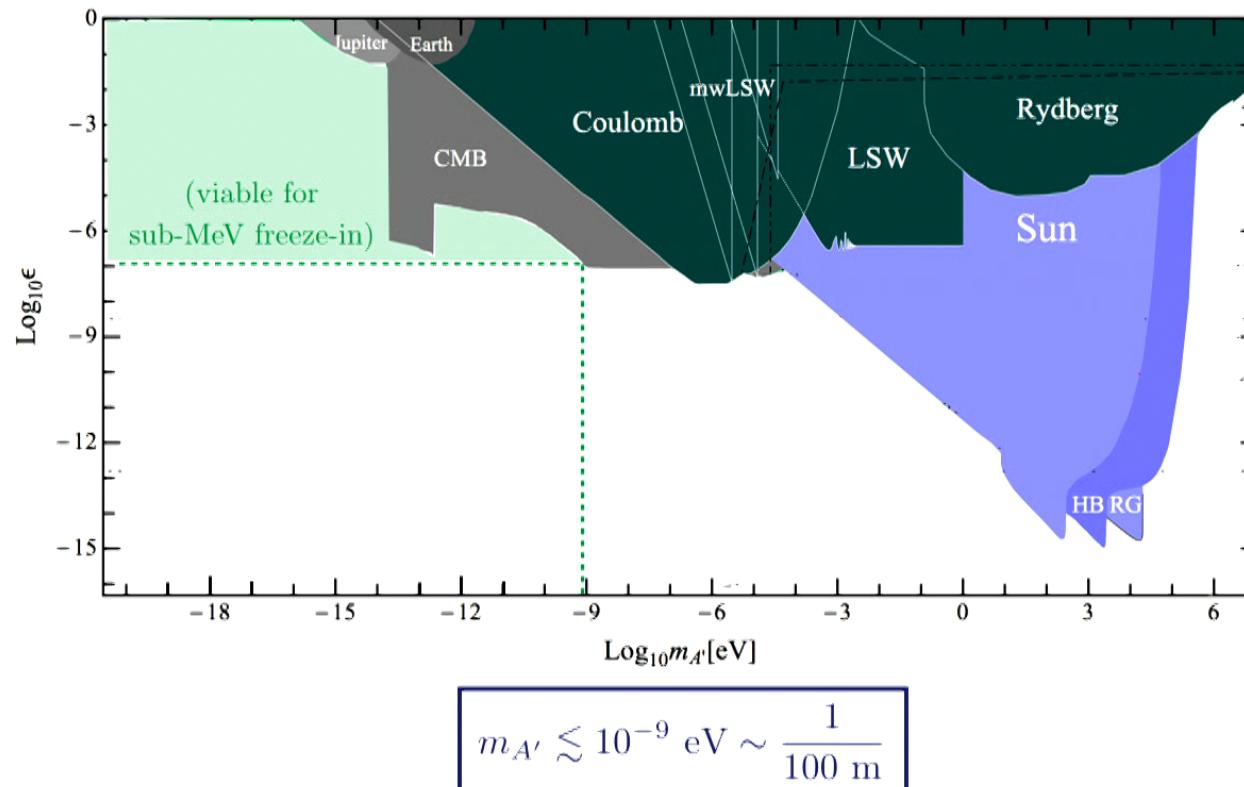
e.g., galaxy clusters $\Rightarrow \alpha' \lesssim 10^{-10} \left(\frac{m_\chi}{\text{MeV}} \right)^{3/2}$

freeze-in $\Rightarrow q_{\text{eff}} \sim \epsilon e'/e \sim 10^{-11} \left(\frac{m_\chi}{\text{MeV}} \right)^{-1/2} \Rightarrow \boxed{\alpha' \sim \frac{10^{-24}}{\epsilon^2} \left(\frac{m_\chi}{\text{MeV}} \right)^{-1}}$

$\therefore \text{SIDM + freeze-in} \Rightarrow \boxed{\epsilon \gtrsim 10^{-7} \left(\frac{m_\chi}{\text{MeV}} \right)^{-5/4}}$

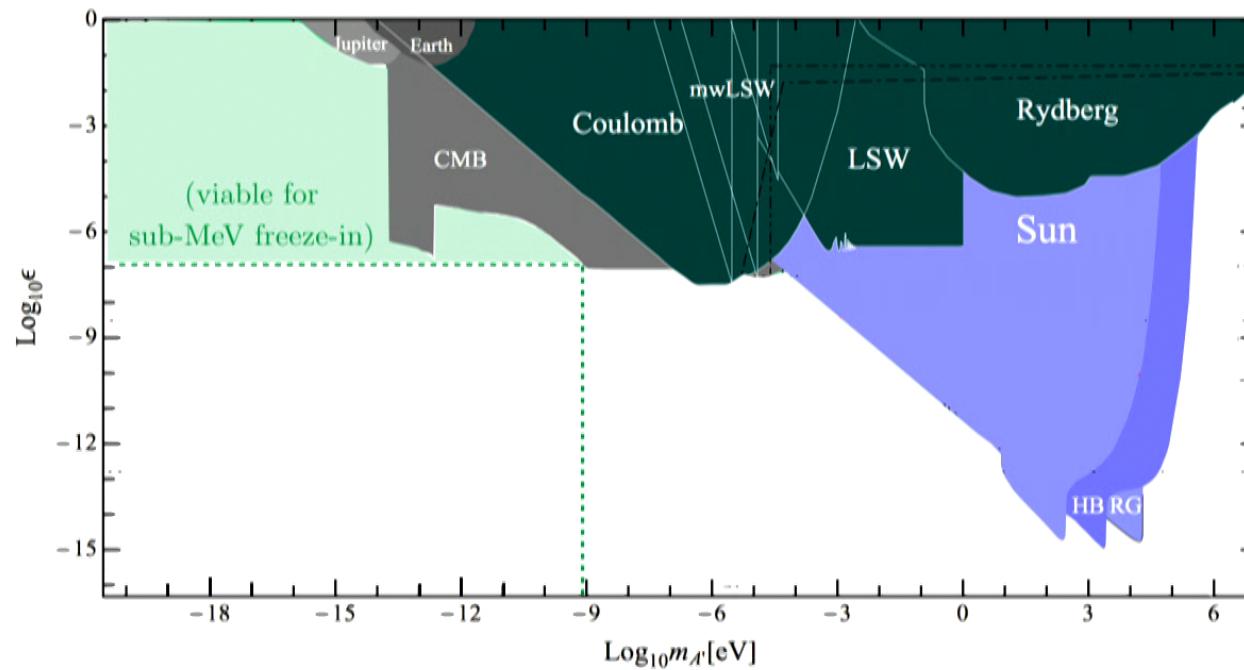
what does this imply
for the dark photon mass?

Parameter Space



arXiv:1704.05081
arXiv:1401.6077

Parameter Space

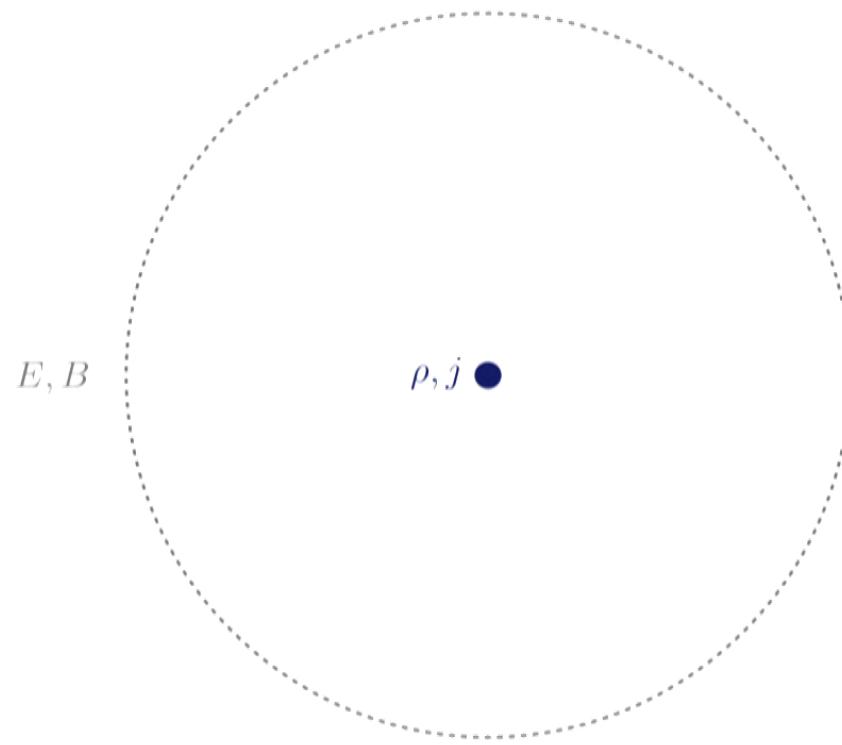


$$m_{A'} \lesssim 10^{-9} \text{ eV} \sim \frac{1}{100 \text{ m}}$$

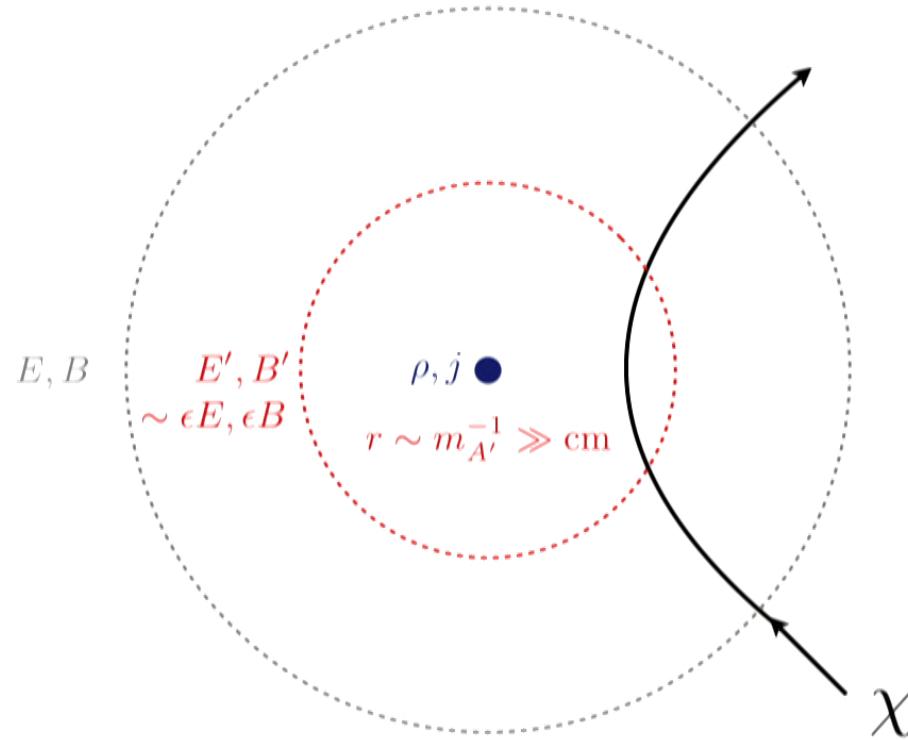
long-range forces

arXiv:1704.05081
arXiv:1401.6077

Electromagnetic Fields



Electromagnetic Fields



Active Direct Detection

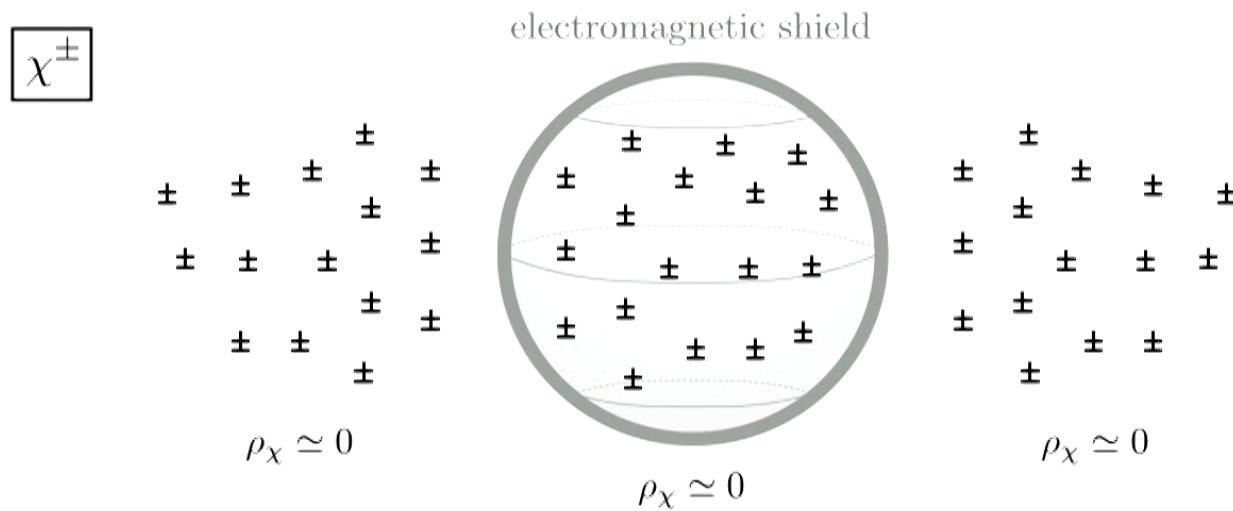
$$q_{\text{eff}} \sim \epsilon e' / e \sim 10^{-11} \left(\frac{m_\chi}{\text{MeV}} \right)^{-1/2}$$

(freeze-in)

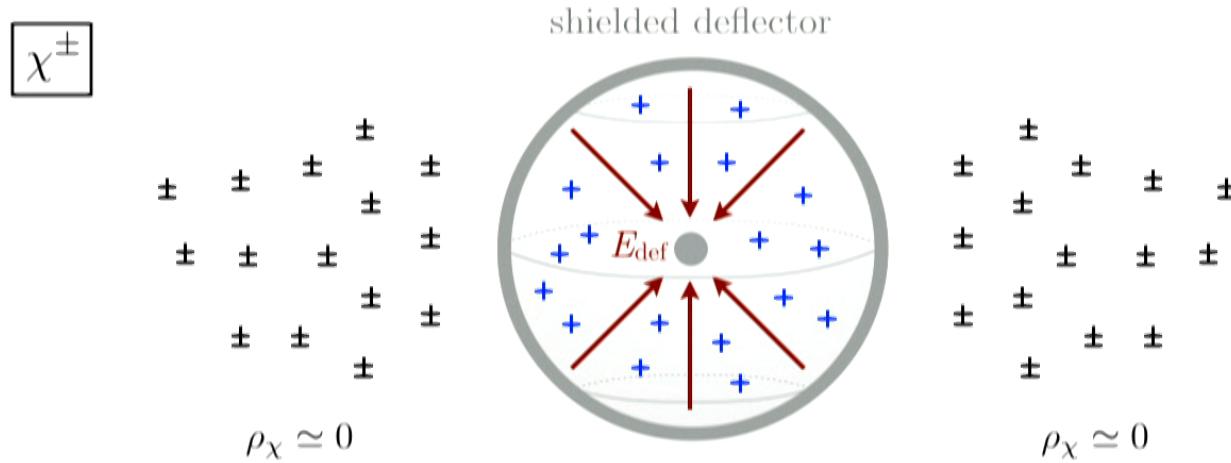
• bend it: $r_g \sim \frac{m_\chi v_\chi}{q_{\text{eff}} e B} \sim \text{meter} \times \left(\frac{m_\chi}{\text{keV}} \right)^{3/2} \left(\frac{10 \text{ T}}{B} \right)$

• stop it: $m_\chi v_\chi^2 \sim q_{\text{eff}} e \Delta V \implies \Delta V \sim \text{MV} \times \left(\frac{m_\chi}{\text{keV}} \right)^{3/2}$

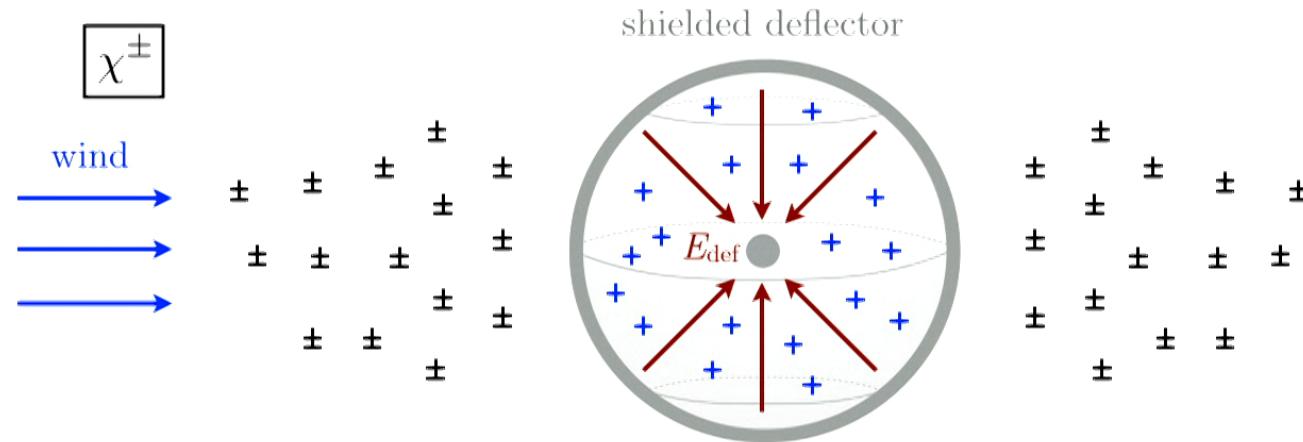
Debye Screening



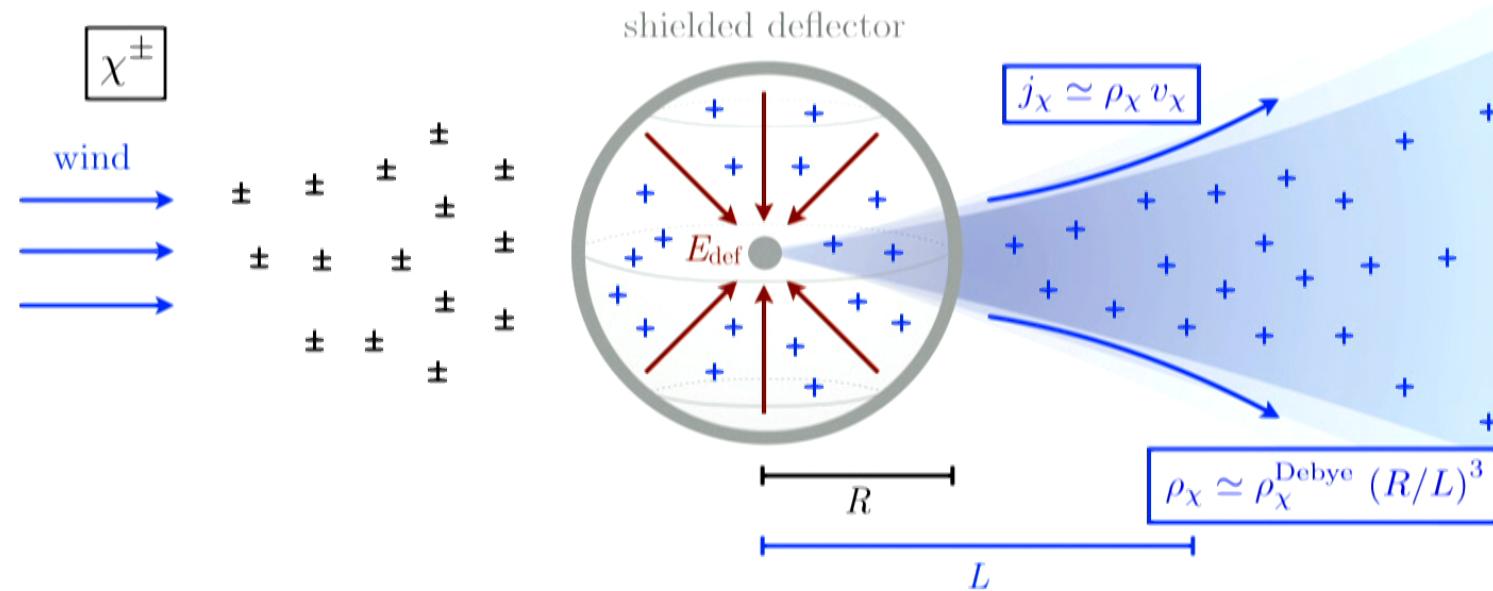
Debye Screening



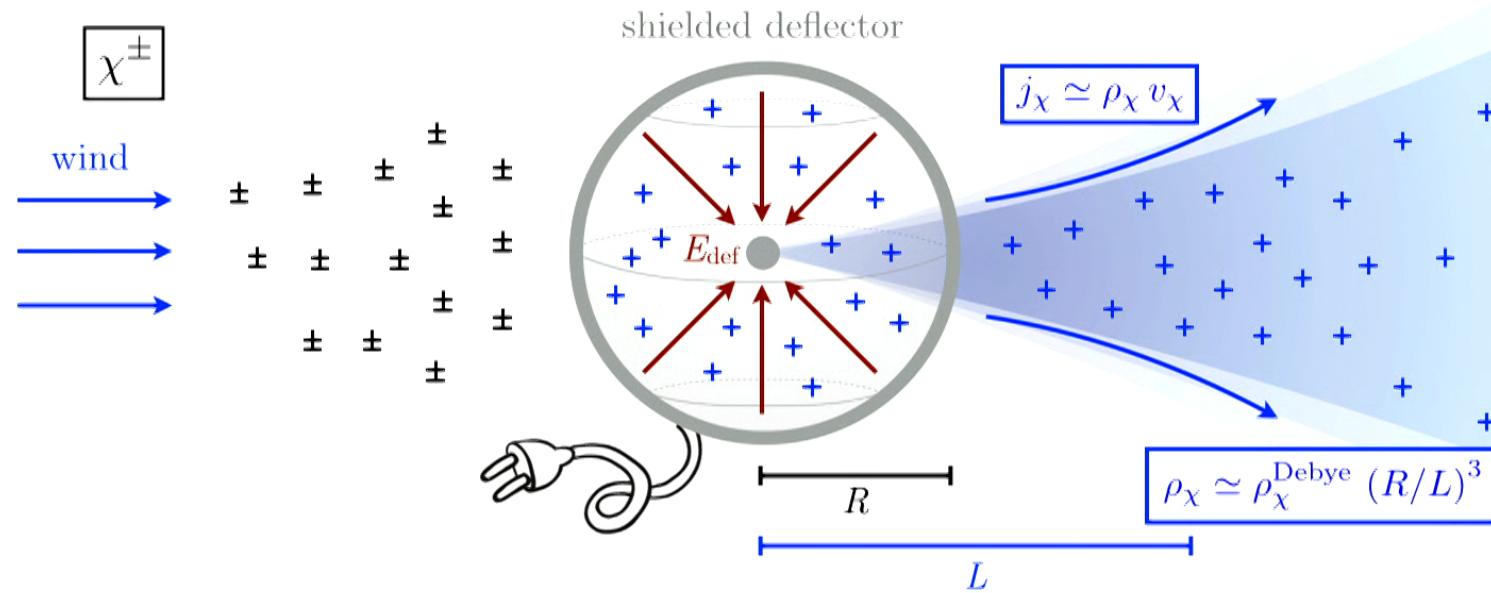
Non-Adiabatic Debye Screening



Non-Adiabatic Debye Screening



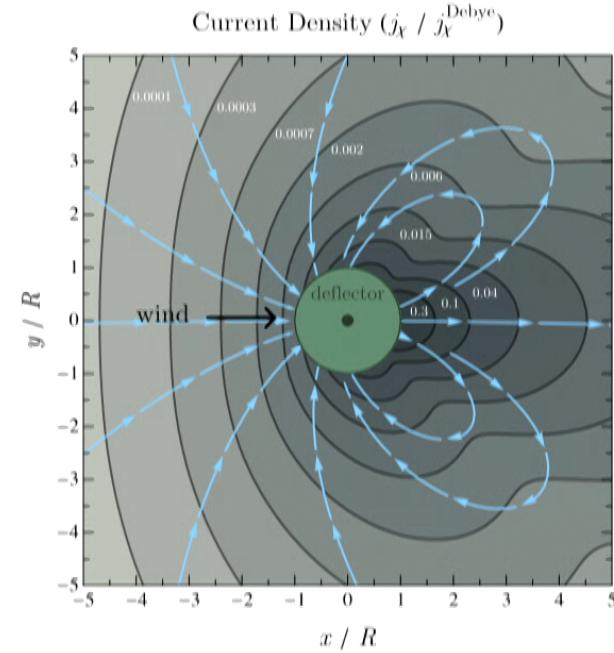
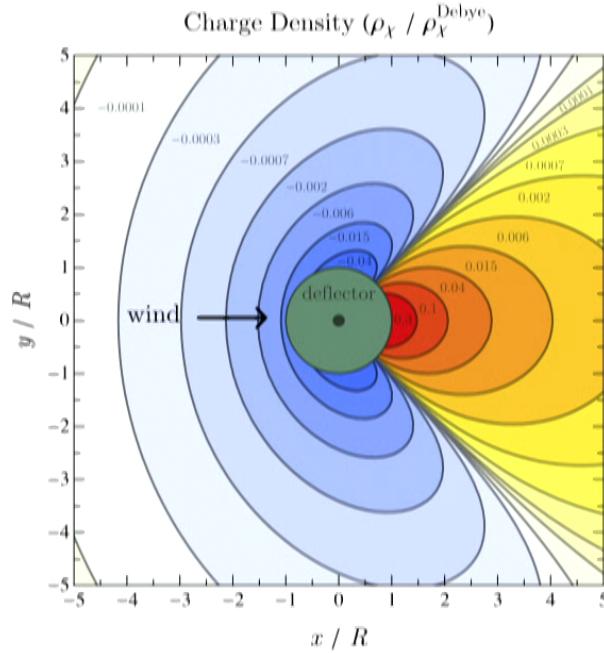
Non-Adiabatic Debye Screening



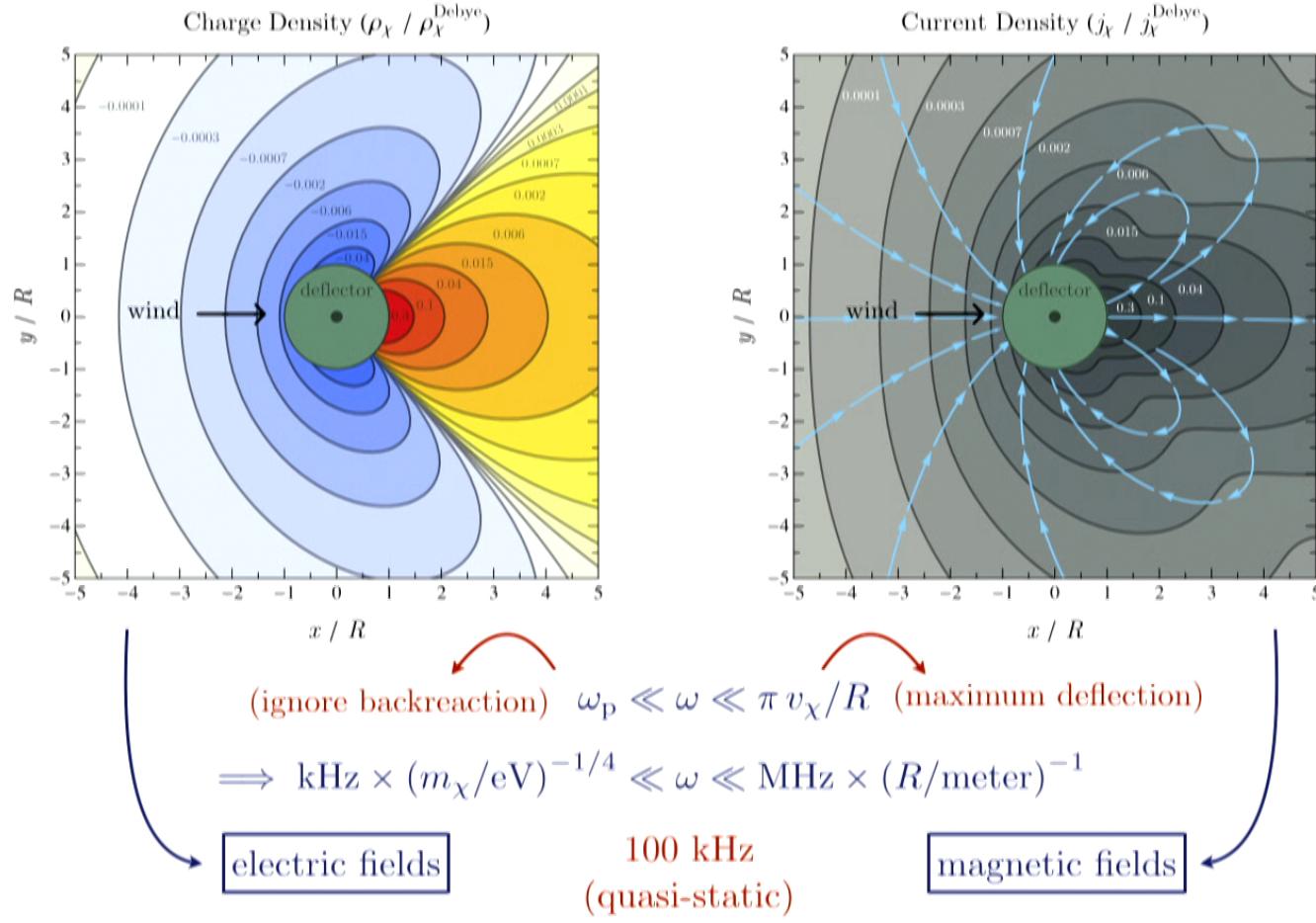
$$E_{\text{def}} \rightarrow E_{\text{def}} e^{i\omega t} \implies \rho_\chi \rightarrow \rho_\chi e^{i\omega t}, j_\chi \rightarrow j_\chi e^{i\omega t}$$

allows for resonant detection

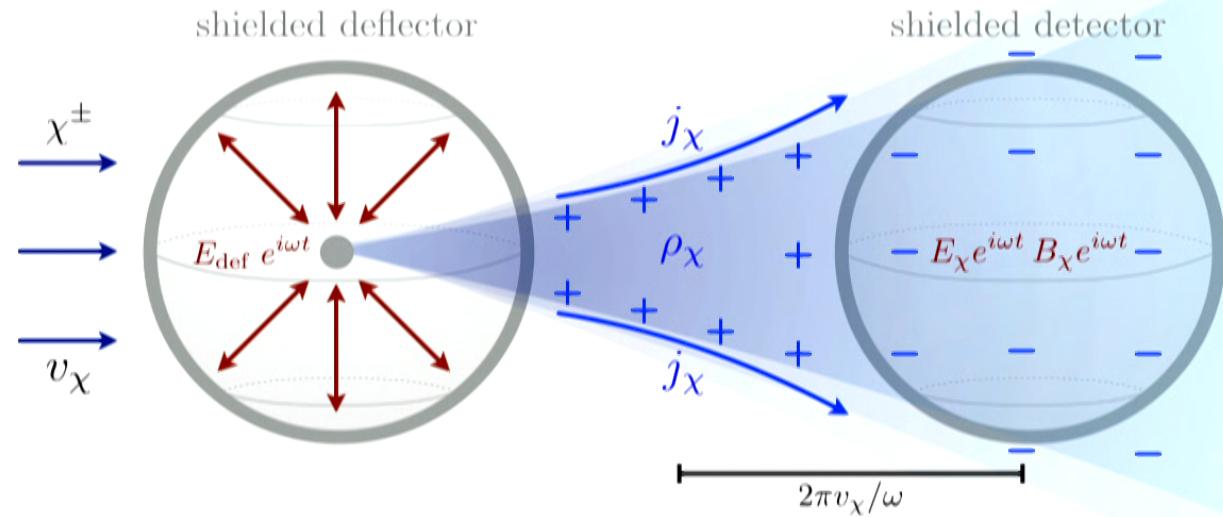
Non-Adiabatic Debye Screening



Non-Adiabatic Debye Screening



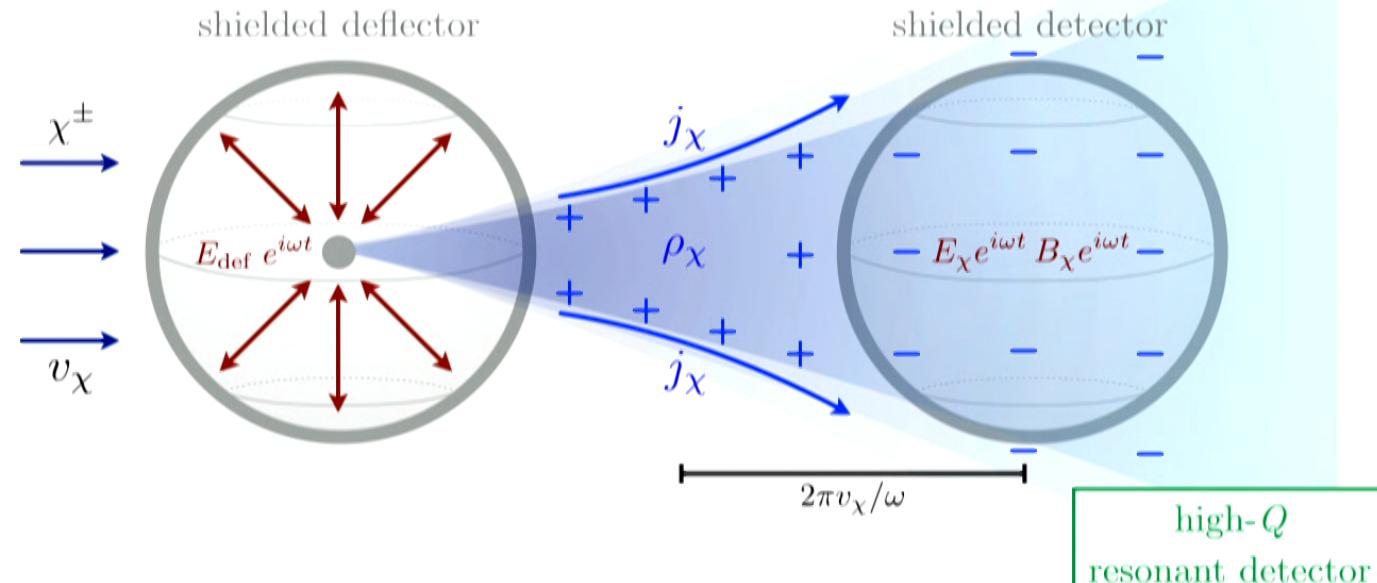
Direct Deflection



$$\text{quasi-static } (\omega \ll 1/R) \implies \begin{cases} E_\chi \sim \rho_\chi R e^{i\omega t} \\ B_\chi \sim v_\chi \rho_\chi R e^{i\omega t} \end{cases}$$

$$E_{\text{def}} \sim 10 \text{ kV/cm}, R \sim \text{meter} \implies \begin{cases} E_\chi \sim 10^{-12} \text{ kV/cm} \times (q_{\text{eff}}/10^{-10})^2 (m_\chi/\text{keV})^{-2} \\ B_\chi \sim 10^{-19} \text{ T} \times (q_{\text{eff}}/10^{-10})^2 (m_\chi/\text{keV})^{-2} \end{cases}$$

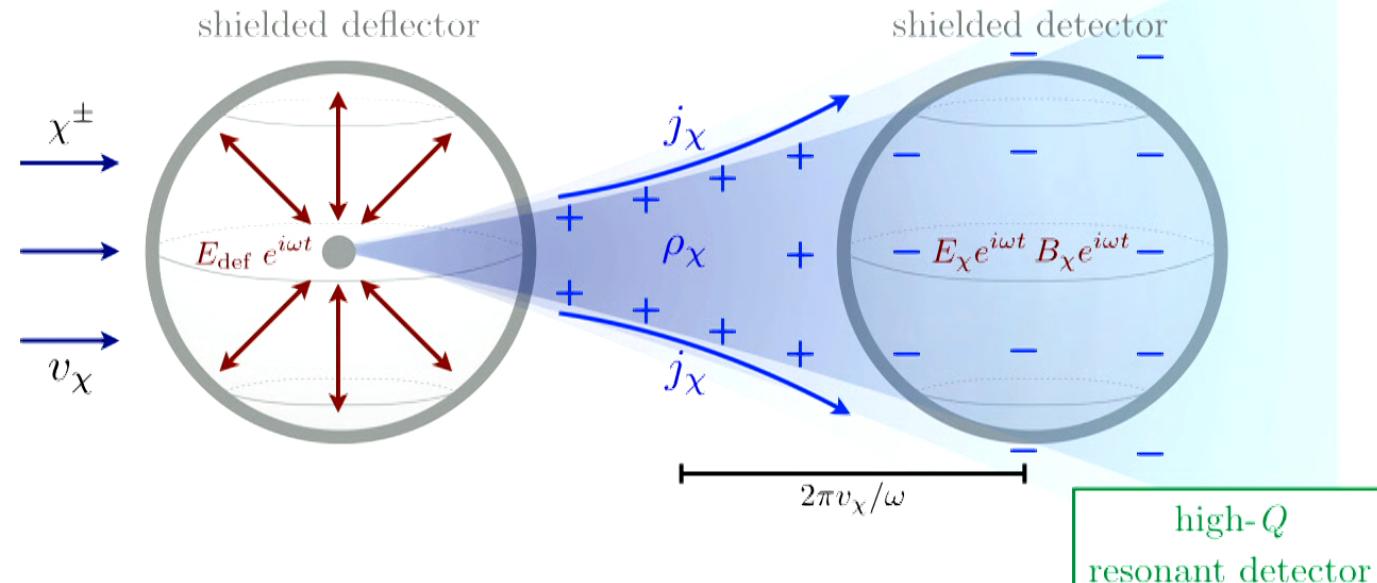
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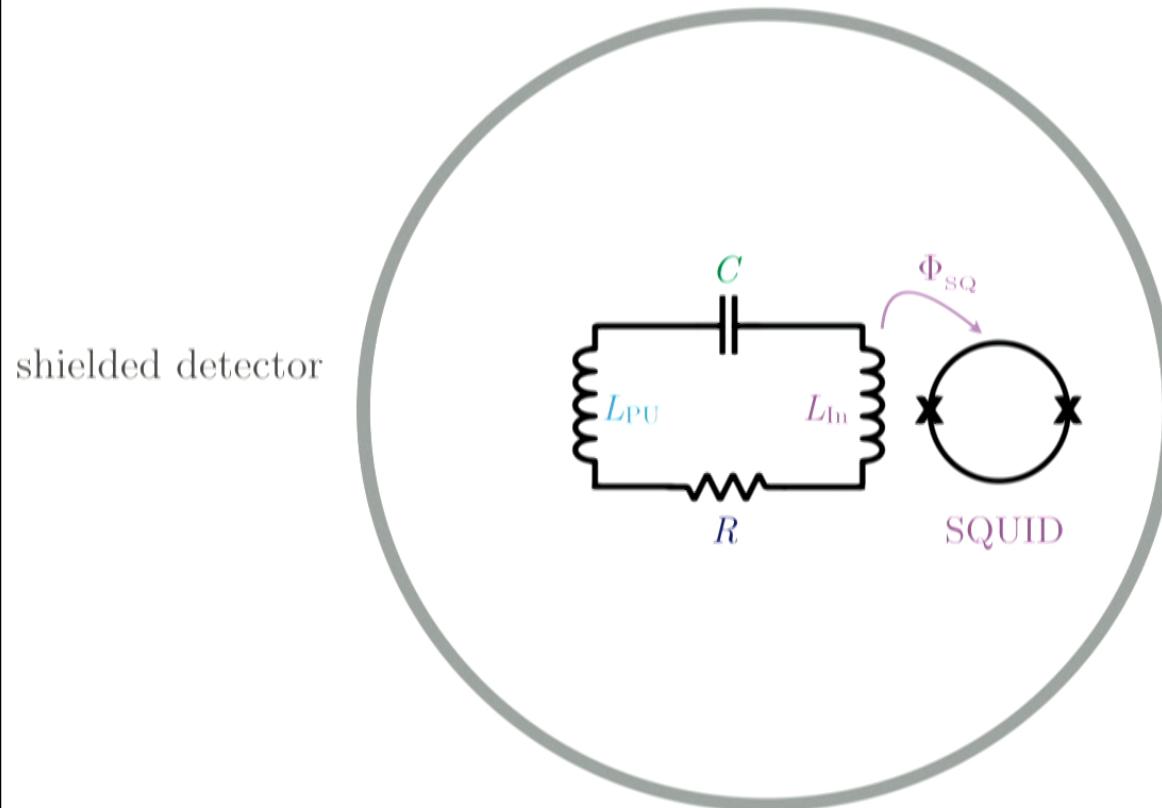
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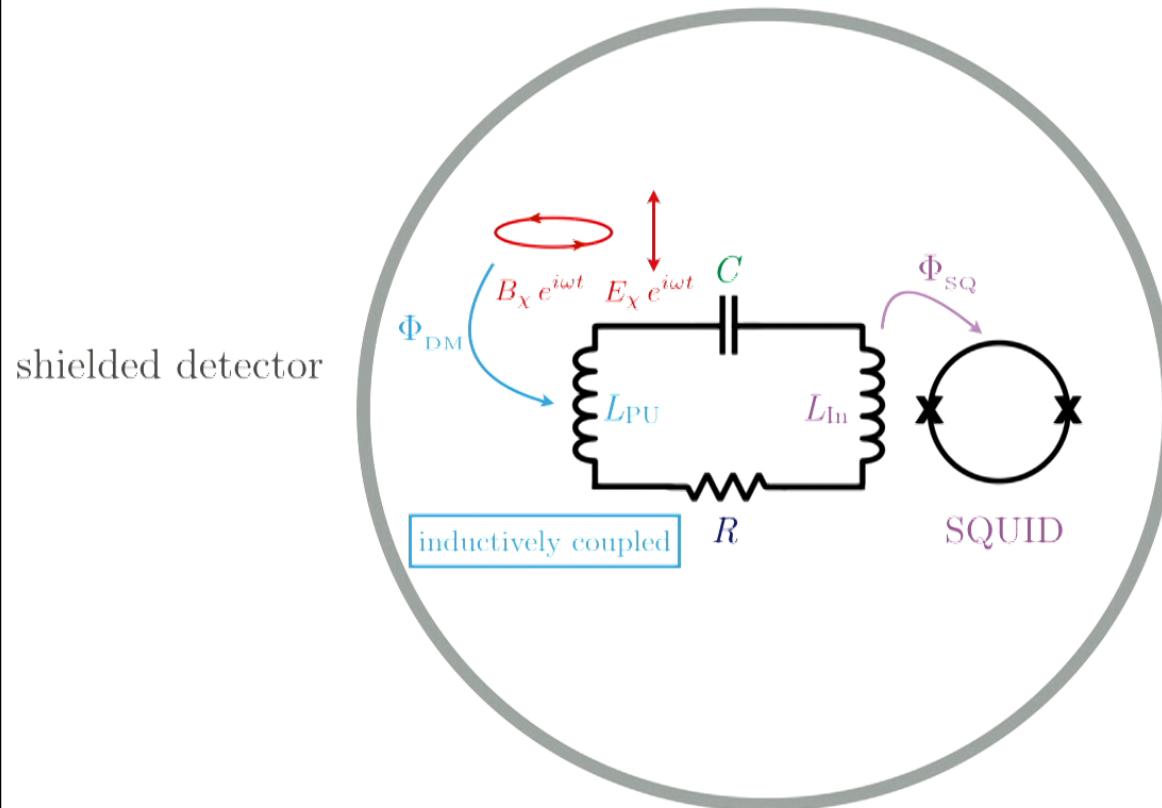
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LC Resonators

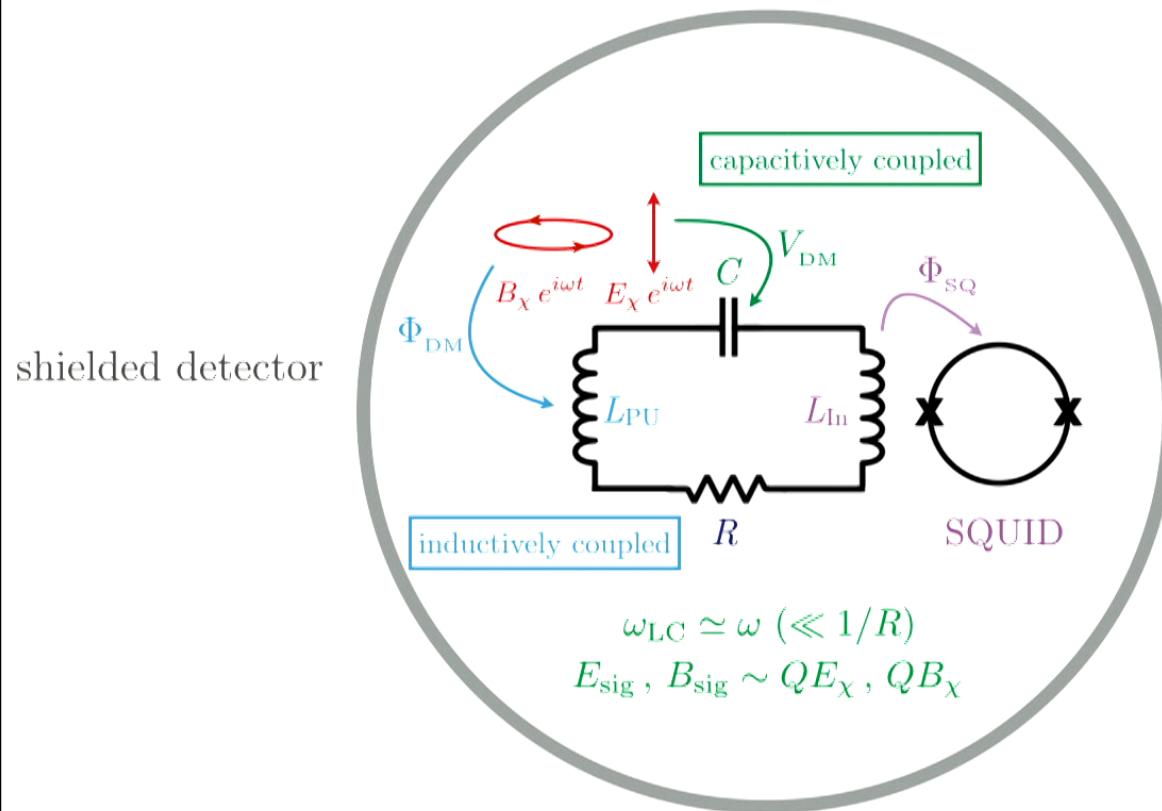


LC Resonators

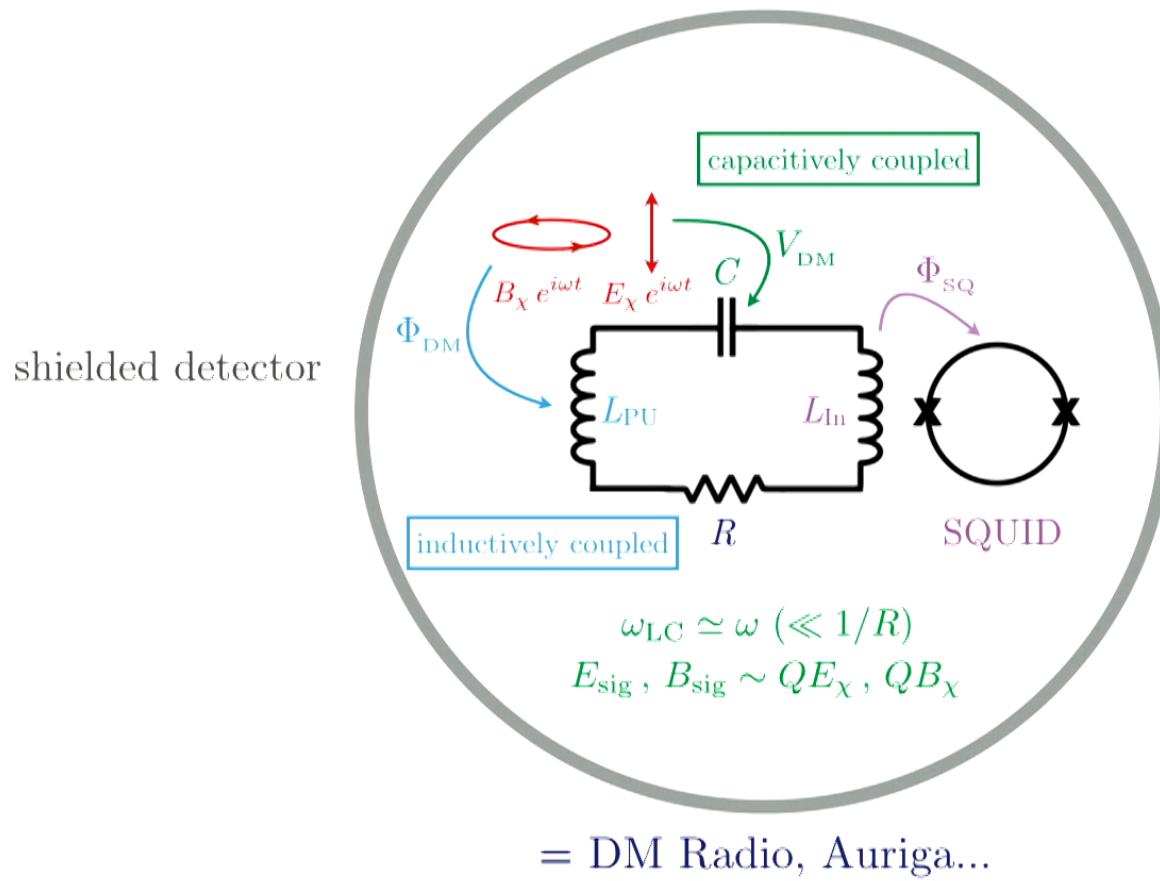


shielded detector

LC Resonators



LC Resonators



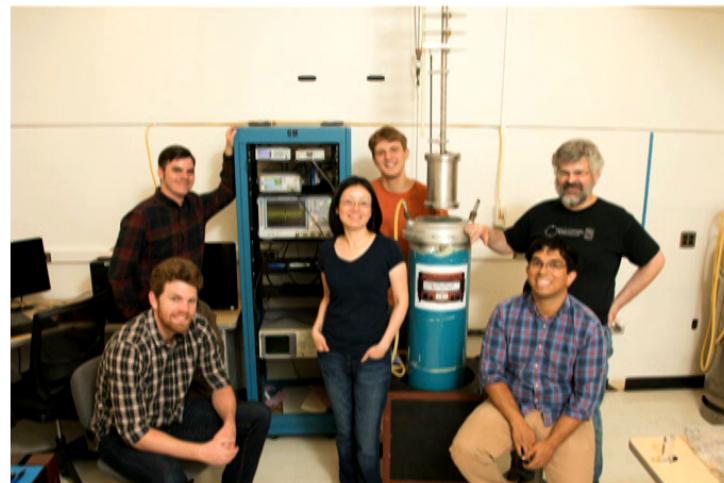
LC Resonators

Auriga
(gravity waves)



resolve thermal noise

DM Radio
(effective currents via ultralight DM)



LC Resonators

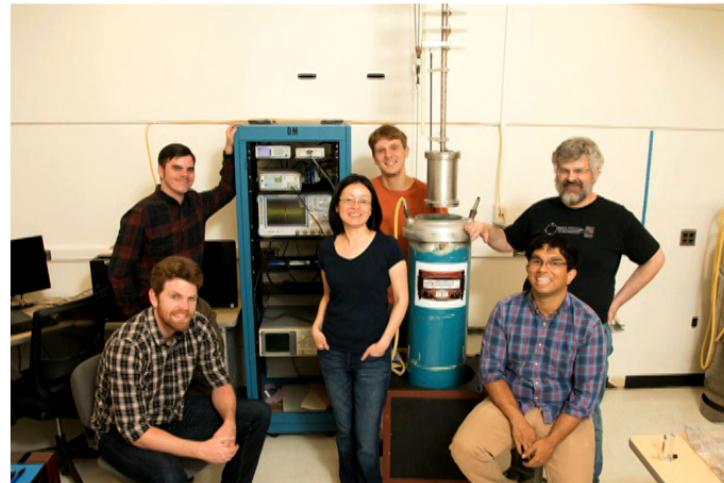
Auriga

(gravity waves)



DM Radio

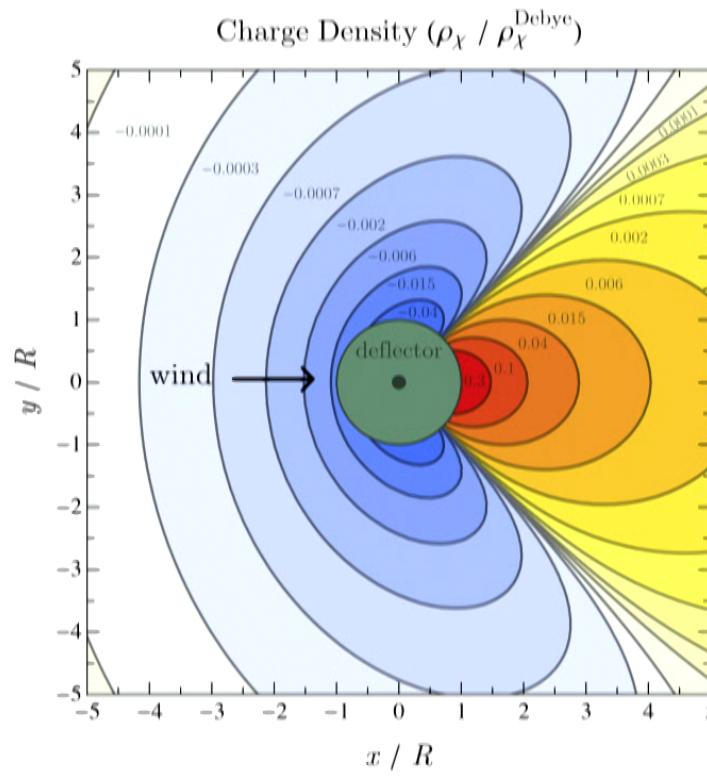
(effective currents via ultralight DM)



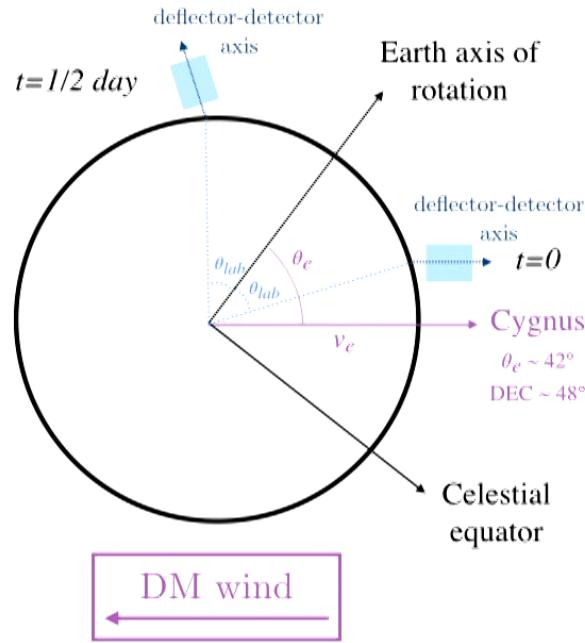
resolve thermal noise

no need to scan or operate down at kHz frequencies $\Rightarrow Q > 10^6$

Directional Dependence

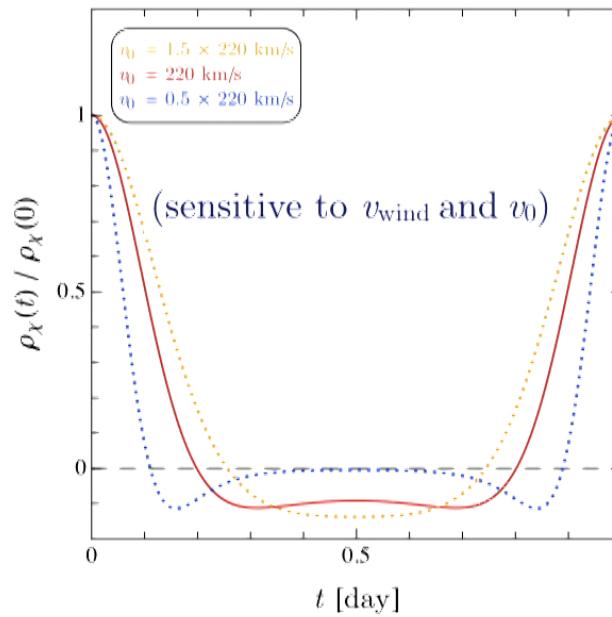
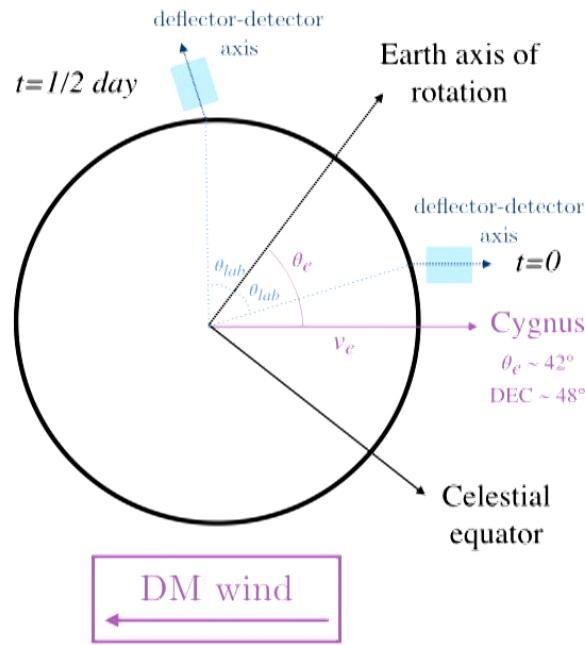


Daily Modulation



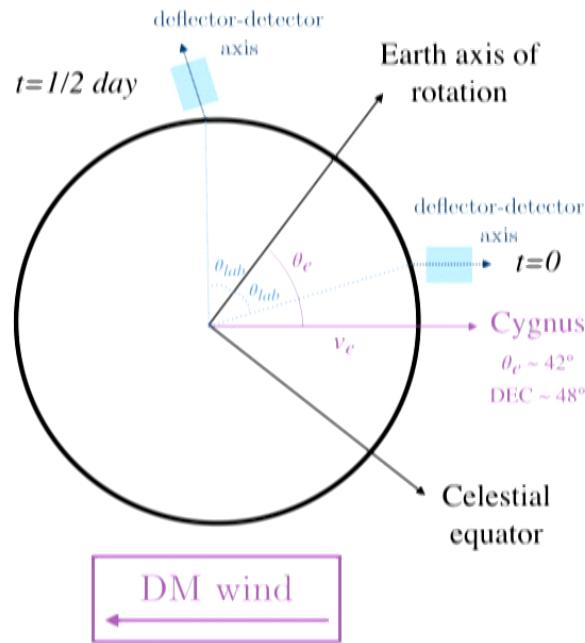
arXiv:1807.10291

Daily Modulation

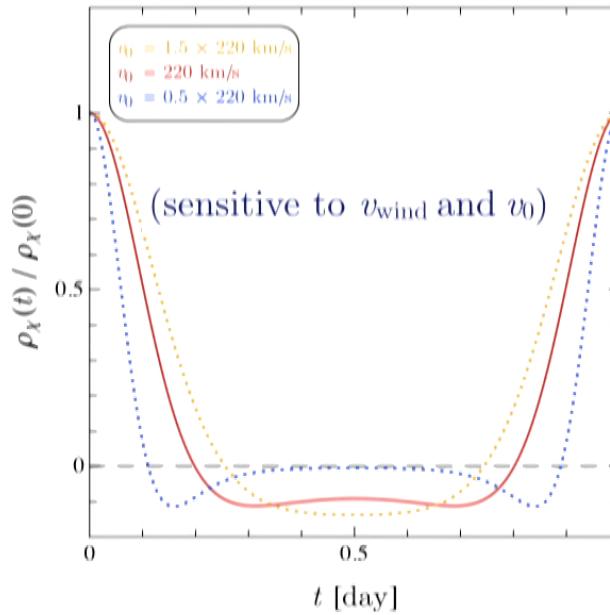


arXiv:1807.10291

Daily Modulation

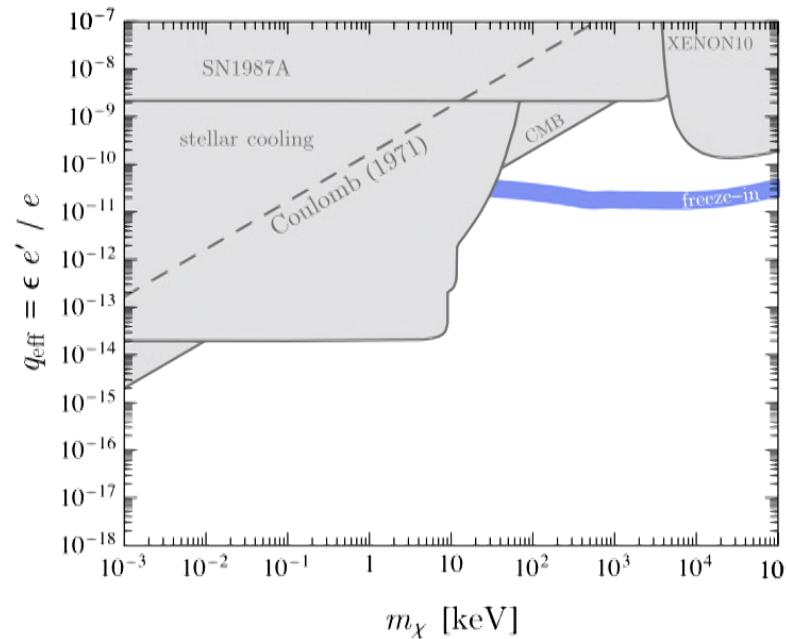


deflector: ω
signal: $\omega \pm \omega_\oplus$



arXiv:1807.10291

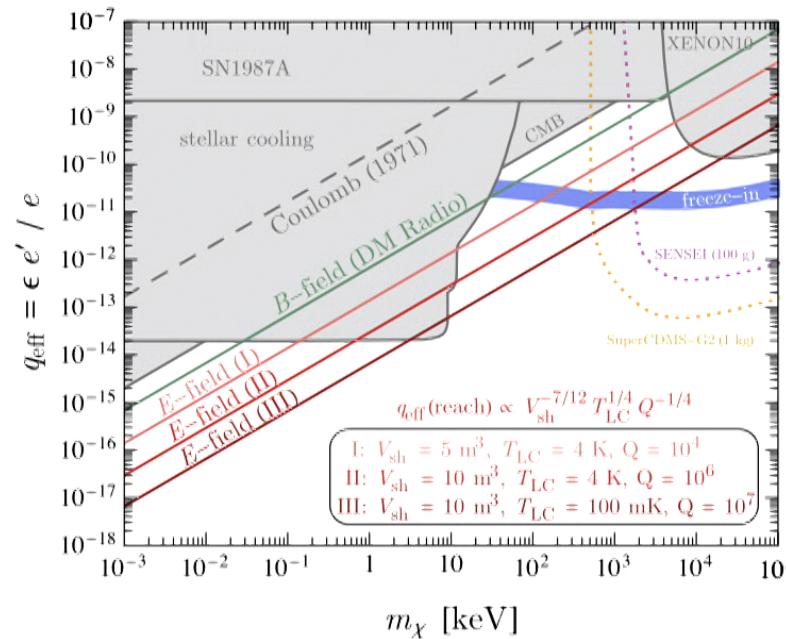
Reach Summary



$$\begin{aligned}\langle E_{\text{def}} \rangle &= 10 \text{ kV/cm} \\ \omega &= 100 \text{ kHz} \\ t_{\text{int}} &= \text{year}\end{aligned}$$

$$q_{\text{eff}}(\text{reach}) \propto m_\chi V_{\text{sh}}^{-\frac{7}{12}} \langle E_{\text{def}} \rangle^{-\frac{1}{2}} (Q \omega t_{\text{int}} / T_{\text{LC}})^{-\frac{1}{4}}$$

Reach Summary



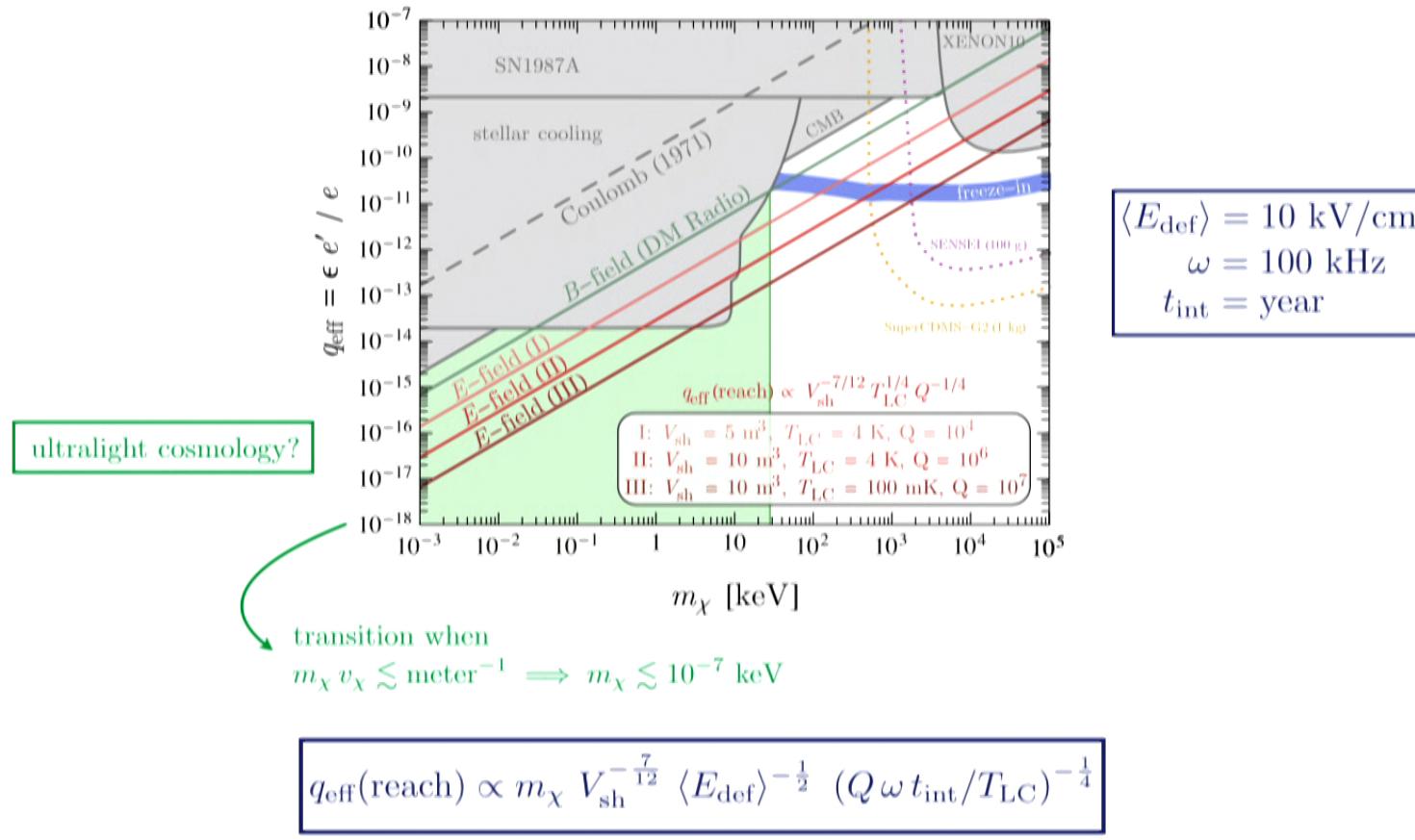
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Reach Summary



Active Direct Detection

- induced daily modulation
- electromagnetic focusing/trapping of dark matter
- optimal geometry for wind
- deflection-detection for spin-coupled forces, ...