

Title: Rethinking the origin of neutrino masses: the role of gravity

Speakers: Lena Funcke

Collection: Simplicity III

Date: September 09, 2019 - 11:00 AM

URL: <http://pirsa.org/19090068>



Rethinking the origin of neutrino masses: the role of gravity

Lena Funcke

in collaboration with Gia Dvali *et al.*

(1602.03191, 1608.08969, 1811.01991, 1905.01264)

Simplicity III, Perimeter Institute, 9 September 2019

Question: Origin of Small Neutrino Masses?

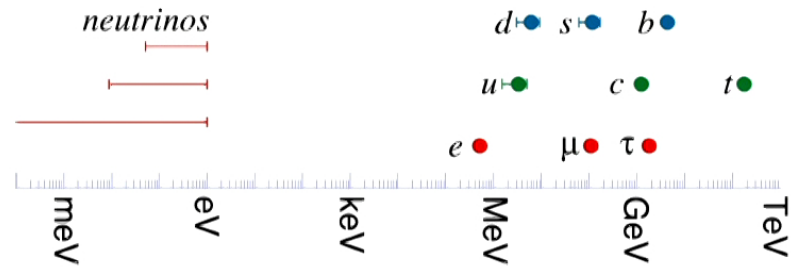


Image credits: IKEA and Murayama (2018).

Question: Origin of Small Neutrino Masses?

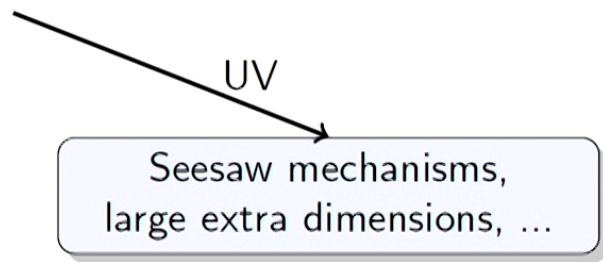
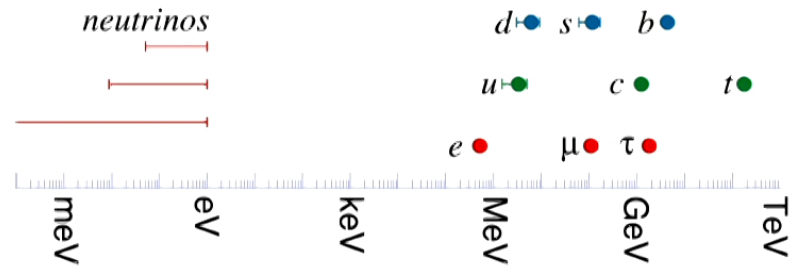


Image credits: IKEA and Murayama (2018).

Question: Origin of Small Neutrino Masses?

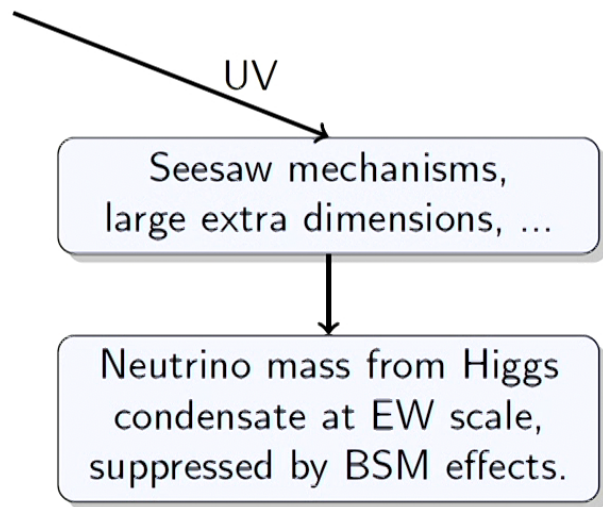
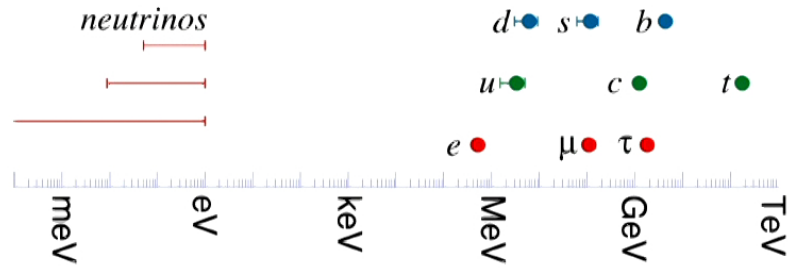


Image credits: IKEA and Murayama (2018).

$$\text{e.g. } \mathcal{L} = \frac{LHHL}{\Lambda_{UV}} \Rightarrow m_\nu = \frac{\langle H \rangle^2}{\Lambda_{UV}}$$

Question: Origin of Small Neutrino Masses?

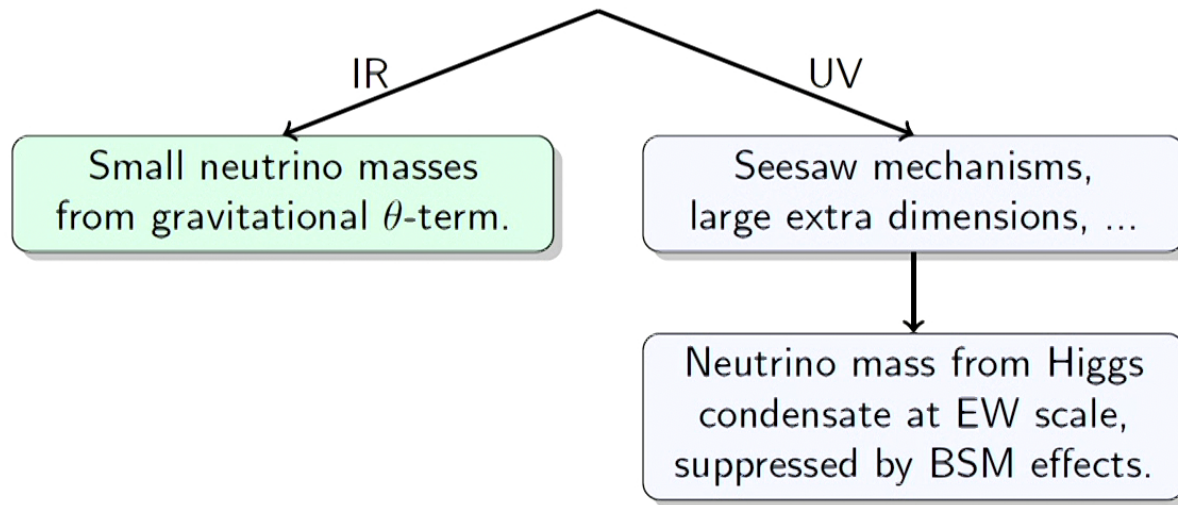
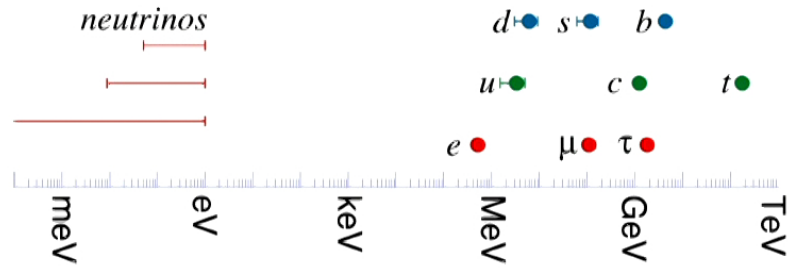


Image credits: IKEA and Murayama (2018).

$$\text{e.g. } \mathcal{L} = \frac{LHHL}{\Lambda_{UV}} \Rightarrow m_\nu = \frac{\langle H \rangle^2}{\Lambda_{UV}}$$

Question: Origin of Small Neutrino Masses?

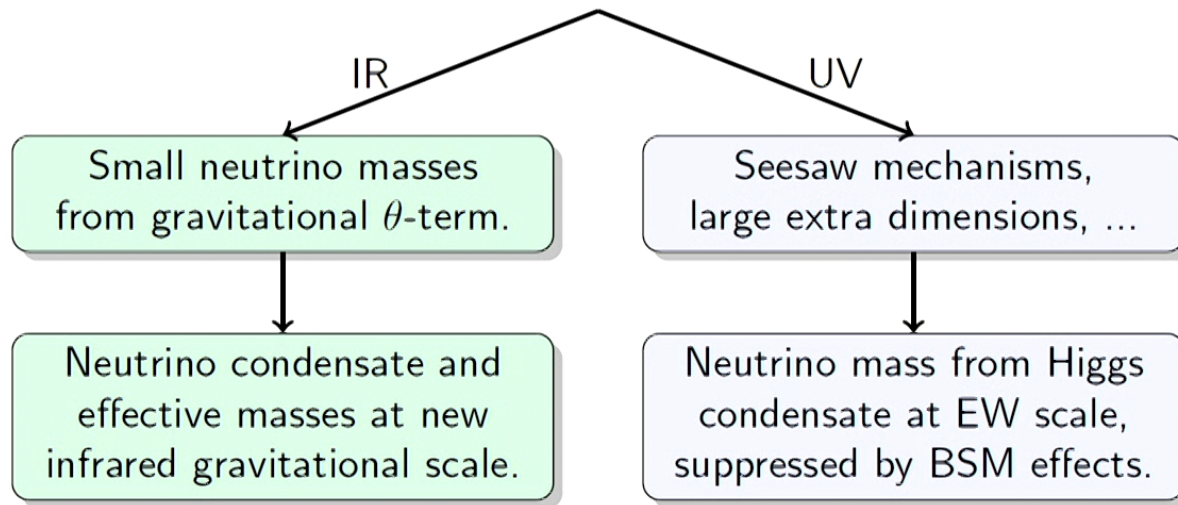
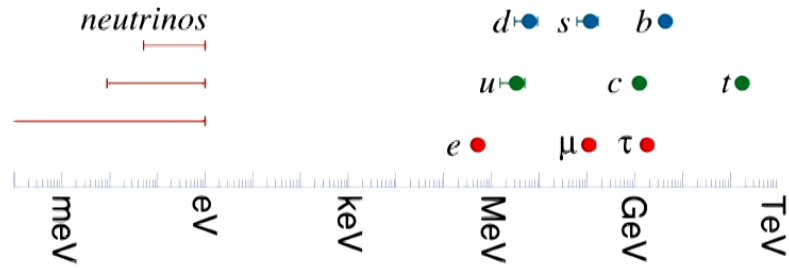


Image credits: IKEA and Murayama (2018).

$$\text{e.g. } \mathcal{L} = \frac{LHHL}{\Lambda_{UV}} \Rightarrow m_\nu = \frac{\langle H \rangle^2}{\Lambda_{UV}}$$

Analogy: Non-Perturbative QCD Vacuum



- ▶ QCD: θ -term $\mathcal{L}_{\text{QCD}} \supset \theta G\tilde{G}$ is made physical by non-perturbative effects [1].

[1] 't Hooft (1976); Witten (1979); Veneziano (1979).

Analogy: Non-Perturbative QCD Vacuum



- ▶ QCD: θ -term $\mathcal{L}_{\text{QCD}} \supset \theta G\tilde{G}$ is made physical by non-perturbative effects [1].

Quantity	QCD with $3q$
Fermion flavor symmetry	$U(3)_V \times U(3)_A \rightarrow U(3)_V$

[1] 't Hooft (1976); Witten (1979); Veneziano (1979).

Analogy: Non-Perturbative QCD Vacuum



- ▶ QCD: θ -term $\mathcal{L}_{\text{QCD}} \supset \theta G\tilde{G}$ is made physical by non-perturbative effects [1].

Quantity	QCD with $3q$
Fermion flavor symmetry	$U(3)_V \times U(3)_A \rightarrow U(3)_V$
(Pseudo)Goldstone bosons	$1(\eta') + 8(\pi, K, \eta)$

[1] 't Hooft (1976); Witten (1979); Veneziano (1979).



Analogy: Non-Perturbative QCD Vacuum



- QCD: θ -term $\mathcal{L}_{\text{QCD}} \supset \theta G \tilde{G}$ is made physical by non-perturbative effects [1].

Quantity	QCD with $3q$
Fermion flavor symmetry	$U(3)_V \times U(3)_A \rightarrow U(3)_V$
(Pseudo)Goldstone bosons	$1(\eta') + 8(\pi, K, \eta)$
Chiral $U(1)_A$ anomaly	$\partial_\mu j_5^\mu = G \tilde{G} + m_q \bar{q} \gamma_5 q$ [4]

[1] 't Hooft (1976); Witten (1979); Veneziano (1979).

[4] Adler (1969); Bell, Jackiw (1969).

$$\begin{aligned} \mathcal{L}_{\text{QCD}} &= \partial d C, & C &= A d A - \frac{3}{2} A A A, & \langle G \hat{G}, G \hat{G} \rangle_{p \rightarrow 0} &= -m_q \langle \bar{q} q \rangle \\ \mathcal{L}_G &= \partial_G d C_G, & C_G &= \Gamma d \Gamma - \frac{3}{2} \Gamma \Gamma \Gamma, & \langle R \hat{R}, R \hat{R} \rangle_{p \rightarrow 0} &= -m_b \langle \bar{u} u \rangle \end{aligned}$$

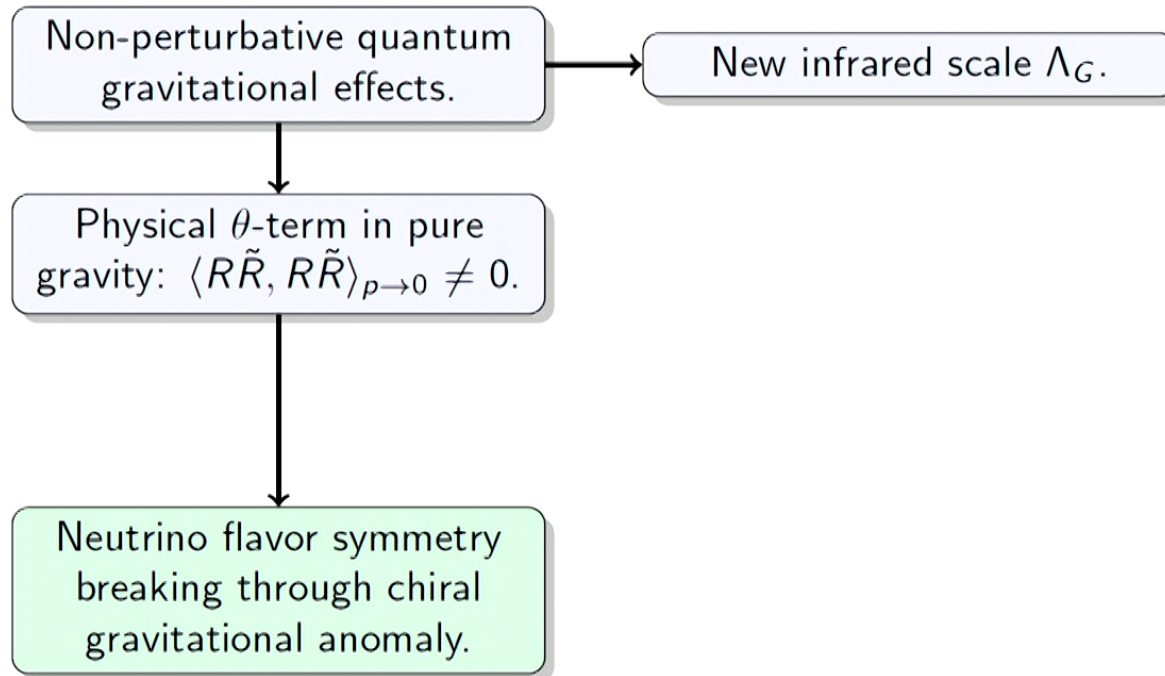
The Model: Neutrino Condensation

Non-perturbative quantum
gravitational effects.



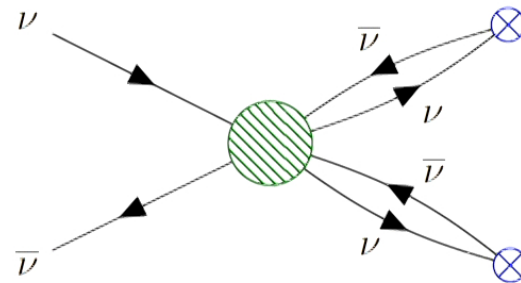
Physical θ -term in pure
gravity: $\langle R\tilde{R}, R\tilde{R} \rangle_{p \rightarrow 0} \neq 0$.

The Model: Neutrino Condensation



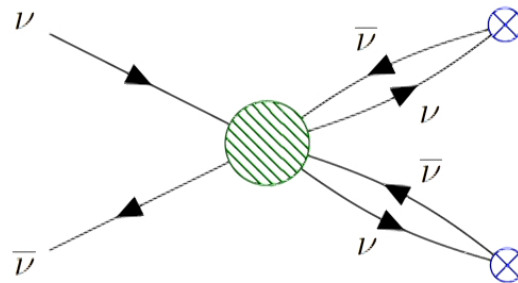
The Model: Neutrino Mass Generation

- ▶ Small effective neutrino mass generation through non-perturbative coupling to neutrino condensate.



The Model: Neutrino Mass Generation

- ▶ Small effective neutrino mass generation through non-perturbative coupling to neutrino condensate.
- ▶ Coupling analogous to 't Hooft vertex in QCD [8].



- ▶ Effective potential allows for neutrino mass hierarchy:

$$V(\hat{X}) = \sum_n \frac{1}{n} c_n \text{Tr}[(\hat{X} + \hat{X}^\dagger)^n]$$
with $\hat{X}_{\alpha L}^{\alpha R} \equiv \langle \bar{\nu}_{\alpha L} \nu_{\alpha R} \rangle$
 $\Rightarrow \partial V / \partial x_i = 0$ determines $\hat{X} = \text{diag}(x_1, x_2, x_3)$.

[8] 't Hooft (1986).

Constraints: Symmetry Breaking Scale Λ_G

Assumption: condensate $|\langle \bar{\nu}\nu \rangle| = \text{scale } \Lambda_G^3 = \text{temperature } T_{\chi\text{SB}}^3.$

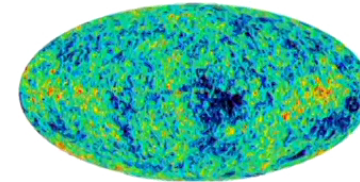
Constraints: Symmetry Breaking Scale Λ_G

Assumption: condensate $|\langle \bar{\nu}\nu \rangle| = \text{scale } \Lambda_G^3 = \text{temperature } T_{\chi\text{SB}}^3.$

$$\Lambda_G \ll m_e$$



Upper bound from
CMB constraints [9].



[9] Archidiacono, Hannestad (2014).

Image credits: NASA / WMAP Science Team [<http://map.gsfc.nasa.gov/>]

Constraints: Symmetry Breaking Scale Λ_G

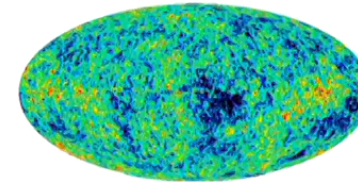
Assumption: condensate $|\langle \bar{\nu}\nu \rangle| = \text{scale } \Lambda_G^3 = \text{temperature } T_{\chi\text{SB}}^3.$

$\Lambda_G \ll m_e$



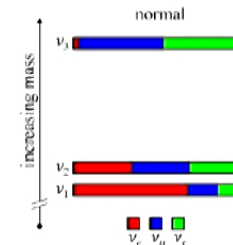
$\sim 0.3 \text{ eV}$

Upper bound from CMB constraints [9].



$\sim 4 \text{ meV}$

Lower bound from neutrino mass splitting [10].



[9] Archidiacono, Hannestad (2014). [10] Olive *et al.* (Particle Data Group) (2014).
Image credits: NASA / WMAP Science Team [<http://map.gsfc.nasa.gov/>] and Patterson (2005).

Constraints: Symmetry Breaking Scale Λ_G

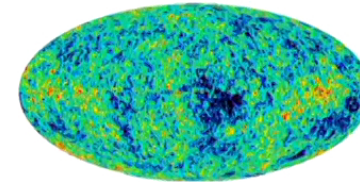
Assumption: condensate $|\langle \bar{\nu}\nu \rangle| = \text{scale } \Lambda_G^3 = \text{temperature } T_{\chi\text{SB}}^3.$

$\Lambda_G \ll m_e$



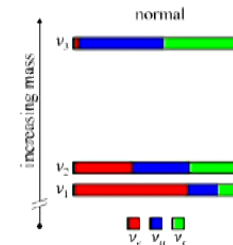
$\sim 0.3 \text{ eV}$

Upper bound from CMB constraints [9].



$\sim 4 \text{ meV}$

Lower bound from neutrino mass splitting [10].



→ Neutrino vacuum condensate $\langle \bar{\nu}\nu \rangle$ on dark energy scale.

[9] Archidiacono, Hannestad (2014). [10] Olive *et al.* (Particle Data Group) (2014).

Image credits: NASA / WMAP Science Team [<http://map.gsfc.nasa.gov/>] and Patterson (2005).

Phenomenological Implications

Weakened cosmological neutrino mass bounds.

- ▶ Relic neutrinos massless until late phase transition at $T_{\chi\text{SB}} \lesssim \Lambda_G$.

Phenomenological Implications

Weakened cosmological neutrino mass bounds.

- ▶ Relic neutrinos massless until late phase transition at $T_{\chi\text{SB}} \lesssim \Lambda_G$.
- ▶ Neutrinos decay & (partially) annihilate $\rightarrow \sum_i m_{\nu_i} \not\lesssim 0.12 \text{ eV}$ [11].

\Rightarrow Masses $m_{\nu_e} \lesssim 2.2 \text{ eV}$ [12] still allowed, measurable at



[11] Aghanim *et al.* (Planck) (2018). [12] Drexlin *et al.* (2013).

Image credit: KATRIN [<http://www.ikp.kit.edu/>].

Phenomenological Implications

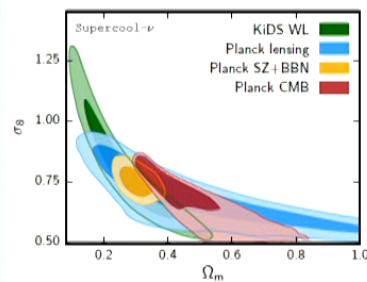
Weakened cosmological neutrino mass bounds.

- ▶ Relic neutrinos massless until late phase transition at $T_{\chi\text{SB}} \lesssim \Lambda_G$.
- ▶ Neutrinos decay & (partially) annihilate $\rightarrow \sum_i m_{\nu_i} \not\lesssim 0.12 \text{ eV}$ [11].

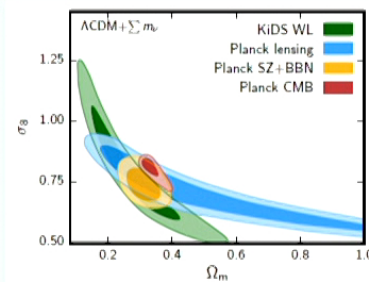
\Rightarrow Masses $m_{\nu_e} \lesssim 2.2 \text{ eV}$ [12] still allowed, measurable at



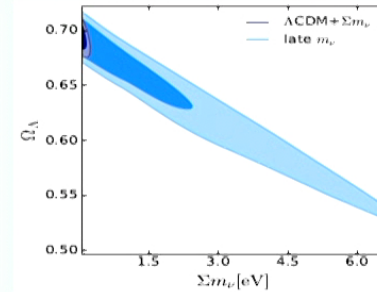
Impact on other cosmic parameters. Dark energy decay?



$\sigma_8(\Omega_m)$ for late m_ν



vs. $\sigma_8(\Omega_m)$ for ΛCDM



$\Omega_\Lambda(m_\nu)$ for both models

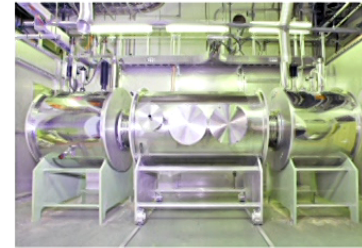
[11] Aghanim *et al.* (Planck) (2018). [12] Drexlin *et al.* (2013).

Image credit: KATRIN [<http://www.ikp.kit.edu/>]. Plots: Lorenz, LF, Calabrese, Hannestad (2018).

Phenomenological Implications

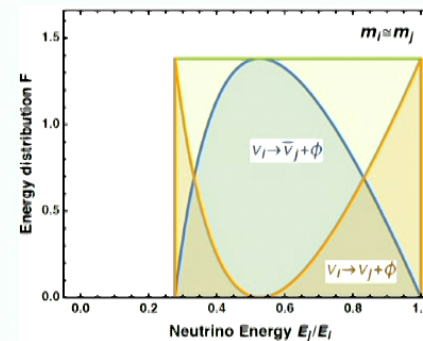
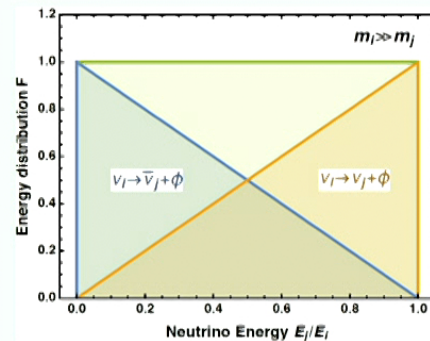
Asymmetries and sterile neutrinos:

- ▶ Asymmetric local neutrino clustering [14]?
- ▶ Resolution of cosmological tensions of hypothetical light sterile neutrinos.



Astrophysical neutrinos:

- ▶ Enhanced neutrino decays, $\nu_i \rightarrow \nu_j + \phi$ and $\nu_i \rightarrow \bar{\nu}_j + \phi$, in solar (and future IceCube) data could reveal Majorana vs. Dirac nature.

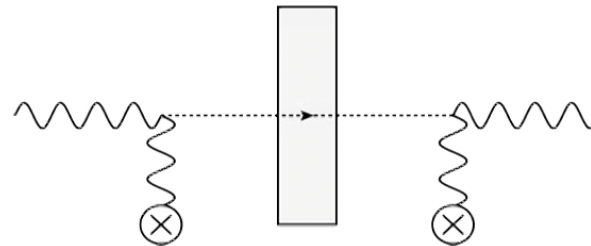
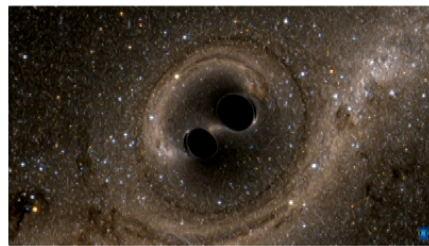


[14] with Mirzaghali, work in progress. Image credit: Betts *et al.* (2013). Plots: LF, Raffelt, Vitagliano (2019).

Phenomenological Implications

Gravity measurements:

- ▶ Different polarization intensities of gravitational waves [16].
- ▶ New attractive gravity-competing short-distance force [17].



New particle detection:

- ▶ “Shining light through walls” with axion-like bosons η_ν, ϕ_k .
- ▶ Possible deep-infrared origin of flavor-violating processes.

[16] Jackiw, Pi (2003). [17] Dvali, LF (2016b), “Domestic Axion” solution to strong CP problem.
Image credits: The SXS Project [<https://www.ligo.caltech.edu/>] and Redondo, Ringwald (2010).

Summary

Assumption: pure gravity contains physical θ -term.

Theoretical consequences:
▶ Neutrino condensation.

Summary

Assumption: pure gravity contains physical θ -term.

Theoretical consequences:

- ▶ Neutrino condensation.
- ▶ Effective small neutrino mass generation.
- ▶ Independent of Higgs or Seesaw mechanisms.
- ▶ Works for both Dirac and Majorana neutrinos.

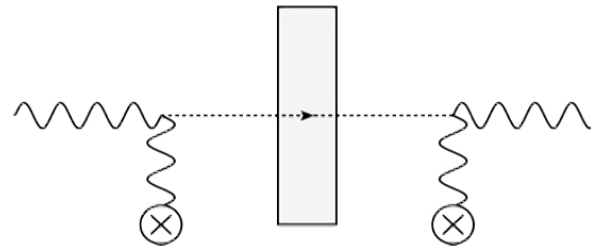
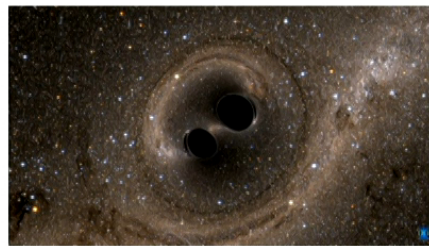
Phenomenology:

- ▶ Large neutrino masses still cosmologically allowed.
- ▶ Enhanced neutrino decays.
- ▶ Possible signatures at KATRIN, IceCube, etc.
- ▶ More details on arXiv: 1811.01991 & 1905.01264

Phenomenological Implications

Gravity measurements:

- ▶ Different polarization intensities of gravitational waves [16].
- ▶ New attractive gravity-competing short-distance force [17].



New particle detection:

- ▶ “Shining light through walls” with axion-like bosons η_ν, ϕ_k .
- ▶ Possible deep-infrared origin of flavor-violating processes.

[16] Jackiw, Pi (2003). [17] Dvali, LF (2016b), “Domestic Axion” solution to strong CP problem.
Image credits: The SXS Project [<https://www.ligo.caltech.edu/>] and Redondo, Ringwald (2010).

Summary

Assumption: pure gravity contains physical θ -term.

Theoretical consequences:

- ▶ Neutrino condensation.
- ▶ Effective small neutrino mass generation.
- ▶ Independent of Higgs or Seesaw mechanisms.
- ▶ Works for both Dirac and Majorana neutrinos.

Phenomenology:

- ▶ Large neutrino masses still cosmologically allowed.
- ▶ Enhanced neutrino decays.
- ▶ Possible signatures at KATRIN, IceCube, etc.
- ▶ More details on arXiv: 1811.01991 & 1905.01264



Thanks for listening!



Do you have any questions?