

Title: Inversion-protected higher-order topological superconductivity in monolayer WTe₂

Speakers: Yi-Ting Hsu

Series: Condensed Matter

Date: September 10, 2019 - 3:30 PM

URL: <http://pirsa.org/19090064>

Abstract: Monolayer WTe₂, an inversion-symmetric transition metal dichalcogenide, has recently been established as a quantum spin Hall insulator and found superconducting upon gating. Here we show that generally a superconducting inversion-symmetric quantum spin Hall material whose normal state is "effectively gapped", such as gated monolayer WTe₂, can be an inversion-protected topological crystalline superconductor featuring "higher-order topology" if the superconductivity is parity-odd. We explicitly demonstrate how the bulk-boundary correspondence naturally emerges in such type of superconductors within a two-dimensional minimal model. We then study the pairing symmetry of superconducting WTe₂ with a microscopic model, and find two types of self-consistently obtained exotic pairings. First is an odd-parity pairing that possesses a nontrivial bulk symmetry indicator and hosts two Majorana Kramers pairs localizing at opposite corners. Even when the conventional pairing is energetically favored, we find that an intermediate in-plane field exceeding the Pauli limit stabilizes an equal-spin pairing aligning with the field. Our findings suggest gated monolayer WTe₂ is a playground for exotic odd-parity superconductivity, and possibly the first material realization for inversion-protected Majorana corner modes without utilizing proximity effect.

Inversion-protected higher-order topological superconductivity in monolayer WTe_2

Yi-Ting Hsu

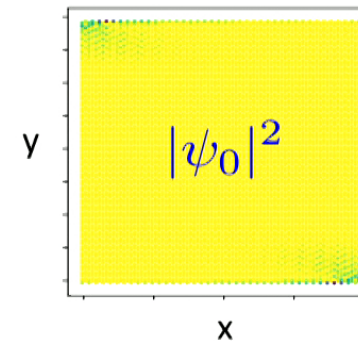
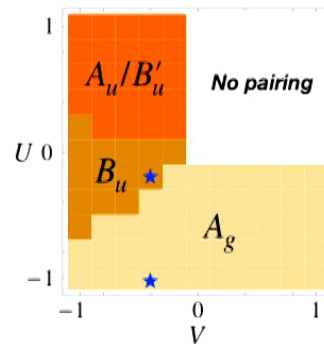
University of Maryland



YTH, Will Cole, Ruixing Zhang, Jay Sau arxiv: 1904.06361 (2019)

Perimeter Institute

09/10/19



Higher-order topological phases

❖ Signature of traditional topo phases in d dimension:

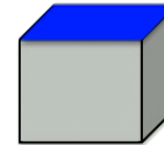
- Boundary modes in $(d-1)$ -dimension



End modes



Edge states



Surfaces

How to achieve HOT phases?

❖ Extrinsic

Protected by **surface gap in (d-1) dimension**

No bulk symmetries are involved

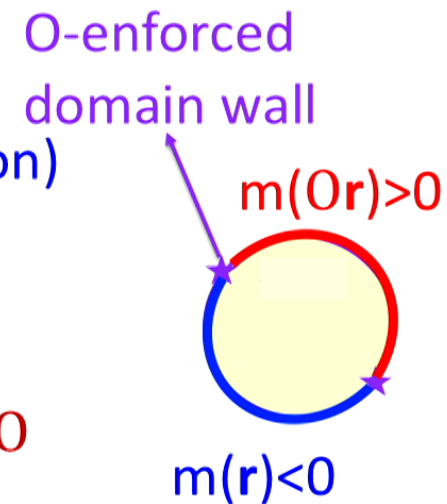
❖ Intrinsic

Protected by bulk **symmetries**

bulk gap (in d dimension)

Q: What kind of symmetries?

A: Some **non-local** bulk symmetries $\mathcal{O}$



How to achieve HOT phases?

❖ Extrinsic

Protected by surface gap in (d-1) dimension

No bulk symmetries are involved

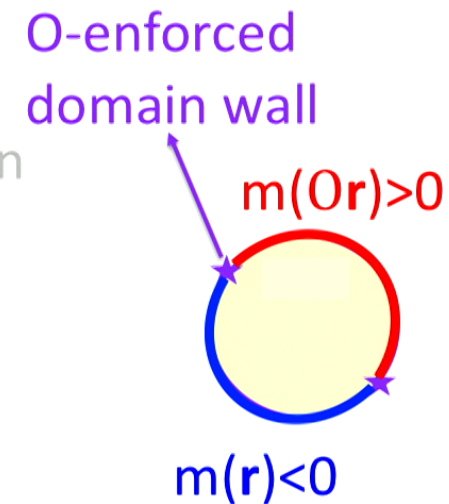
❖ Intrinsic

Protected by bulk **symmetries**

bulk gap in d dimension

Q: What kind of symmetries?

A: Some **crystalline symmetries** **0**



Crystalline symmetries vs Band topology

Q: Effects of crystalline symmetries on SPTs?

1. Allow inference of topology from symmetry eigenvalues:

E.g. Fu-Kane criteria:

For 3D time-reversal TIs with inversion symmetry,

Z_2 index is given by parity eigenvalues of the occupied bands

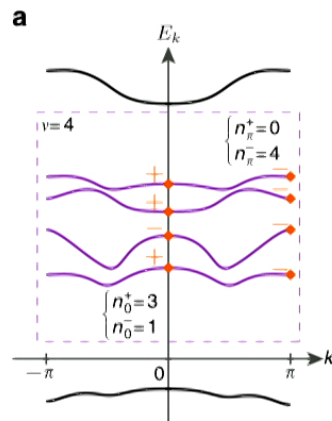
Such topological indices => **symmetry indicators**

Classify TR insulators under all space groups

Q: Effects of crystalline symmetries on SPTs?

2. Change the topological classification

Consider 3D time-reversal insulators (class AII)
+ 230 space groups (SGs):



$$X_{\text{BS}} \equiv \frac{\{\text{BS}\}}{\{\text{AI}\}}.$$
$$= X_{\text{BS}}^{(w)} \times X_{\text{BS}}^{(s)}$$

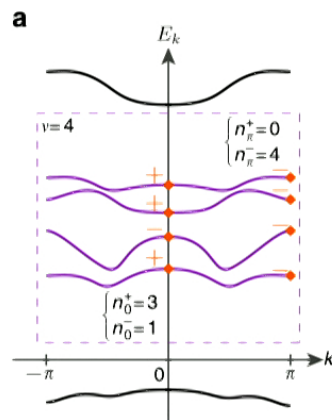
Po, Vishwanath, Watanabe, Ncomm (2017)

Classify TR insulators under all space groups

Q: Effects of crystalline symmetries on SPTs?

2. Change the topological classification

Consider 3D time-reversal insulators (class All)
+ 230 space groups (SGs):



$$X_{\text{BS}} \equiv \frac{\{\text{BS}\}}{\{\text{AI}\}}.$$
$$= X_{\text{BS}}^{(w)} \times X_{\text{BS}}^{(s)}$$

- Compatible with weak TI indices,
weak mirror Chern numbers

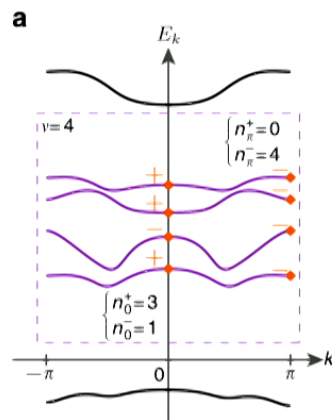
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$$= X_{\text{BS}}^{(w)} \times X_{\text{BS}}^{(s)}$$

- **NOT** compatible with known indices

- **Band topology changes** for SGs with certain crystalline symmetries

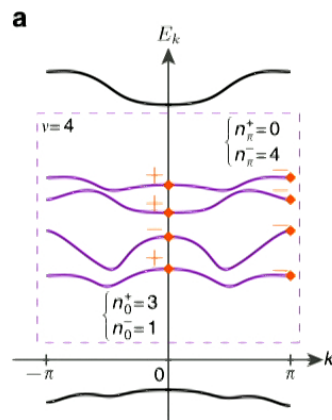
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E.g. SGs with **inversion** symmetry:

$\mathbb{Z}_4$ instead of $\mathbb{Z}_2$ (strong TI index)

Po, Vishwanath, Watanabe, Ncomm (2017)

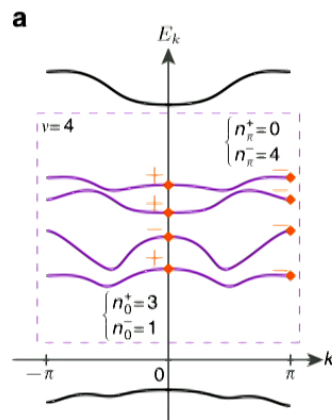
Khalaf, Po, Vishwanath, Watanabe, PRX (2018)

Classify TR insulators under all space groups

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$$X_{\text{BS}} \equiv \frac{\{\text{BS}\}}{\{\text{AI}\}}.$$
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E.g. SGs with **inversion** symmetry: Z_4

Q: How to see it's Z_4 ?

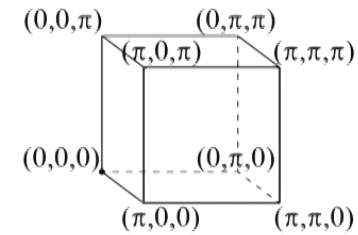
Po, Vishwanath, Watanabe, Ncomm (2017)

Khalaf, Po, Vishwanath, Watanabe, PRX (2018)

Inversion-protected TCI in 3D

❖ Z_4 Symmetry indicator:

$$\kappa_N \equiv \frac{1}{4} \sum_{K \in \text{TRIMs}} (n_K^+ - n_K^-)$$

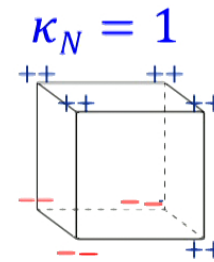
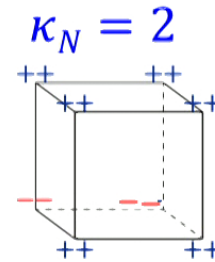
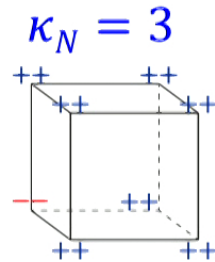
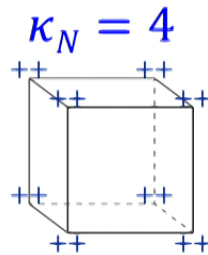
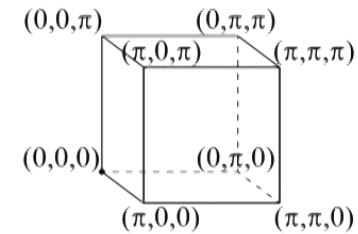


Khalaf, Po, Vishwanath, Watanabe, PRX (2018)

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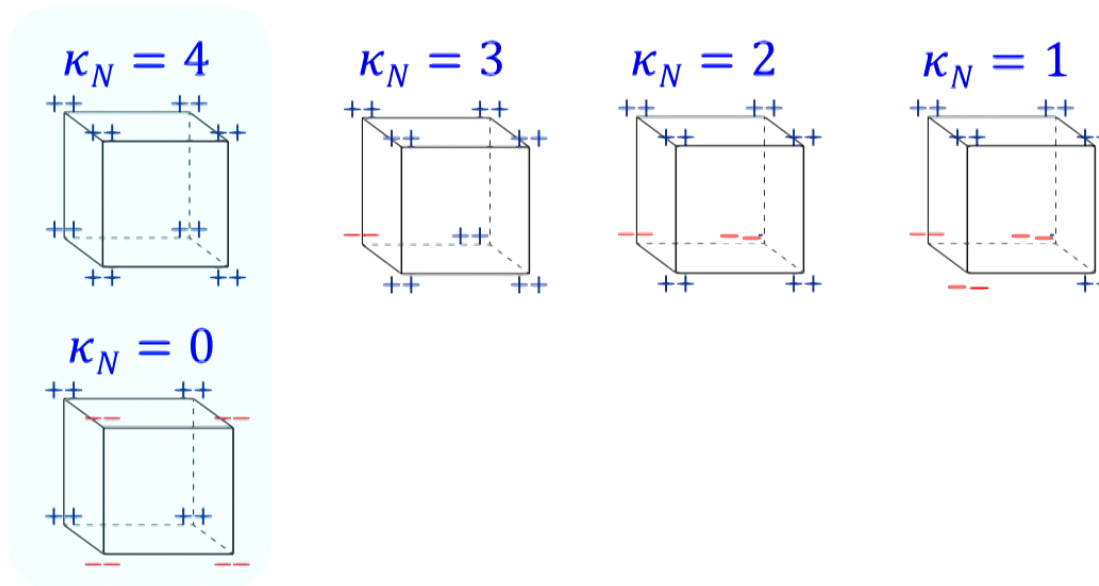
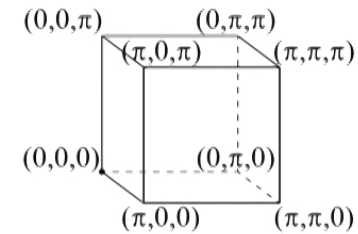


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Inversion-protected TCI in 3D

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Atomic insulators

Khalaf, Po, Vishwanath, Watanabe, PRX (2018)

Inversion-protected TCI in 3D

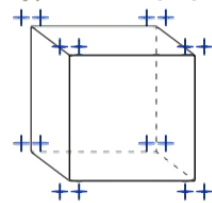
❖ Z_4 Symmetry indicator:

$$\kappa_N \equiv \frac{1}{4} \sum_{K \in \text{TRIMs}} (n_K^+ - n_K^-) \mod 4$$

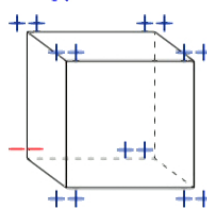
Z_2 strong index

$$\nu_N = \kappa_N \mod 2$$

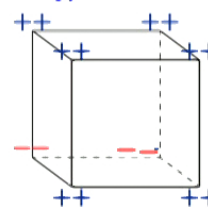
$\kappa_N = 4(0)$



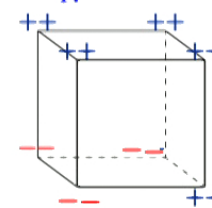
$\kappa_N = 3$



$\kappa_N = 2$



$\kappa_N = 1$



Trivial
(atomic insulator)

Surface
state

x

✓

x

✓

Khalaf, Po, Vishwanath, Watanabe, PRX (2018)

Inversion-protected TCI in 3D

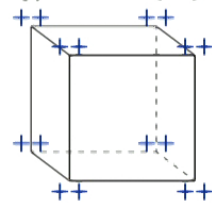
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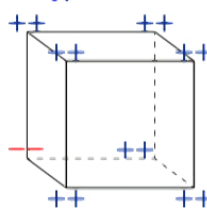
$$\nu_N = \kappa_N \mod 2$$

$$\kappa_N = 4(0)$$



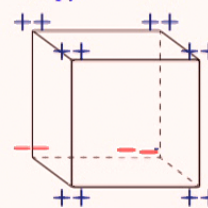
Trivial
(atomic insulator)

$$\kappa_N = 3$$



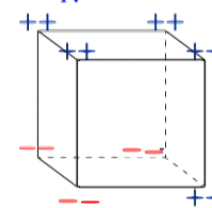
Topo

$$\kappa_N = 2$$



?

$$\kappa_N = 1$$



Topo

Surface
state

x

✓

x

✓

Khalaf, Po, Vishwanath, Watanabe, PRX (2018)

Inversion-protected TCI in 3D

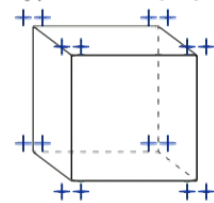
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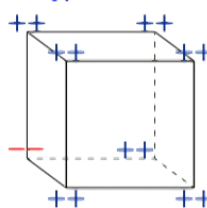
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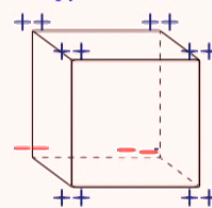
Trivial
(atomic insulator)

$$\kappa_N = 3$$



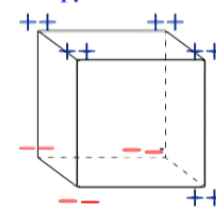
Topo

$$\kappa_N = 2$$



?
(Not atomic insulator)

$$\kappa_N = 1$$



Topo

Surface
state

x

✓

x

✓

Khalaf, Po, Vishwanath, Watanabe, PRX (2018)

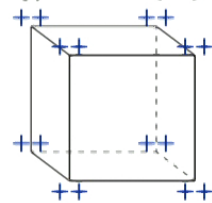
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❖ Z_4 Symmetry indicator: Khalaf, Po, Vishwanath, Watanabe, PRX (2018)

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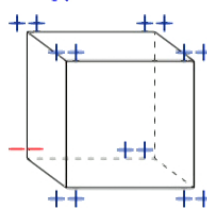
$$\nu_N = \kappa_N \bmod 2$$

$$\kappa_N = 4(0)$$



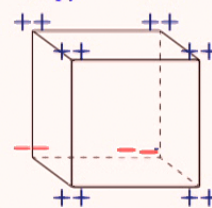
Trivial

$$\kappa_N = 3$$



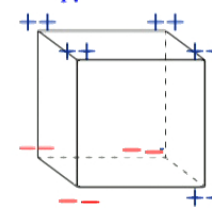
Topo

$$\kappa_N = 2$$



Higher-order
TI (?)

$$\kappa_N = 1$$



Topo

Surface
state

x

✓

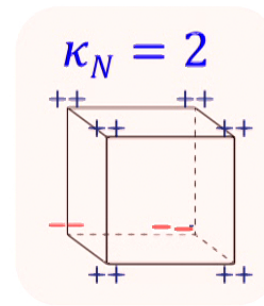
Hinge
modes (?)

✓

Inversion-protected TCI in 3D

❖ Z_4 Symmetry indicator: Khalaf, Po, Vishwanath, Watanabe, PRX (2018)

$$\kappa_N \equiv \frac{1}{4} \sum_{K \in \text{TRIMs}} (n_K^+ - n_K^-) \bmod 4 \quad \nu_N = \kappa_N \bmod 2$$



Higher-order TI (?)

Bulk-boundary correspondence
in double-TI construction

Q: Is there a **superconductor** analogue?

i.e. **Inversion**-protected TCsc with some $\kappa_{BdG}=2$

$\Rightarrow$ Intrinsic **higher-order** topo sc (HOTsc)

(Majoranas!)

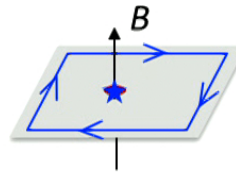
Why higher-order **superconductors**?

❖ Signature of **topo sc**:

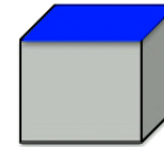
- **Boundary**: Majorana modes in **(d-1)**-dimension



End modes



Edge states
Vortex cores



Surfaces

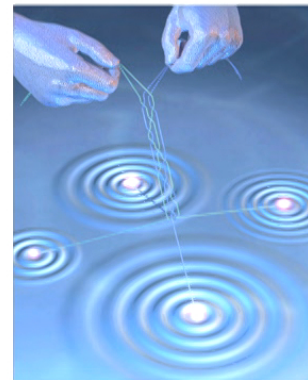
*Higher-order topo sc
(HOTsc)*



*2D case:
Majorana **corner** modes!*

Why 2D HOTsc?

- ❖ Majorana zero modes
 - **0-dimensional** zero-energy modes
 - Qubits for **topological quantum computation**



Scientific American

*How to **realize** 2D intrinsic HOTsc?*

3D inversion-protected HOTsc

- ❖ Define analogous symmetry indicator for

$$H_{\mathbf{k}}^{\text{BdG}} = \begin{pmatrix} H_{\mathbf{k}} & \Delta_{\mathbf{k}} \\ \Delta_{\mathbf{k}}^\dagger & -H_{-\mathbf{k}}^* \end{pmatrix}$$

- Define symmetry operation for H_{BdG}

$$U_{\mathbf{k}}^{\text{BdG}}(g) = \begin{pmatrix} U_{\mathbf{k}}(g) & 0 \\ 0 & \chi_g U_{-\mathbf{k}}(g)^* \end{pmatrix} \xrightarrow{\text{for normal-state } H}$$

$\chi_g = +1$: even parity pairing

$\chi_g = -1$: odd parity pairing

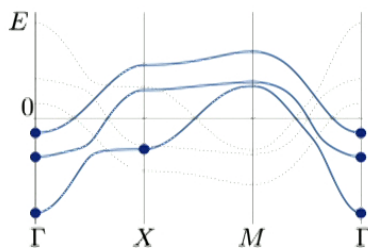
such that H_{BdG} is invariant under $U_{\mathbf{k}}^{\text{BdG}}$

Ono, Yanase, Watanabe, PR Research (2019)

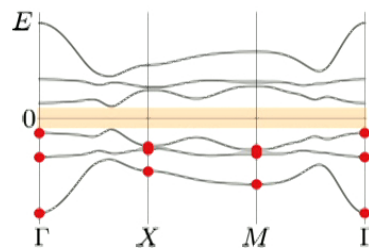
3D inversion-protected HOTsc

❖ Define **inversion** symmetry indicator for $H_{\mathbf{k}}^{\text{BdG}}$

Normal bands



BdG bands



$$\text{Inversion } P_{\text{BdG}} = \begin{pmatrix} P & 0 \\ 0 & \pm P \end{pmatrix}$$

$$\kappa_{\text{BdG}} \equiv \frac{1}{4} \sum_{\mathbf{k} \in 3\text{D TRIMs}} \sum_{\alpha=\pm 1} \alpha (n_{\mathbf{k}}^{\alpha})^{\text{BdG}} \pmod{4}$$

Metal (FS circle TRIMs)

- Even-parity pairing: κ_{BdG} is always 0
- Odd-parity pairing: κ_{BdG} is $\mathbb{Z}_4$ ($\mathbb{Z}_8$?)

Ono et al. PR research (2019)
Skurativska, Titus, Fischer (2019)

Goal: Realize inversion-protected HOTsc in 2D



1. Find a **recipe** => construct **minimal model**
2. Demonstrate **bulk- boundary correspondence**
3. Prediction for **material**

Hard: few candidates for even traditional topo sc ☹️

- Making prediction for **unconventional sc** is hard
- Challenging: **Topology + symmetry + interaction**

Beyond band structure calculation

Z_4 Symmetry indicator for 2D TCsc

- ❖ Conjecture Z_4 symmetry indicator for a 2D time-reversal superconductor (DIII) with inversion:

- $\kappa = \frac{1}{4} \sum_{\mathbf{k} \in \text{TRIM}} \sum_n \xi_{\mathbf{k}n} \mod 4$

$\xi_{\mathbf{k}n}$: parity of the occupied BdG band n

Inversion: $P_{\text{BdG}} = \begin{pmatrix} P & 0 \\ 0 & -P \end{pmatrix}$ for **odd-parity** pairing

- $\nu = \kappa \mod 2$

$\downarrow$
 Z_2 index for DIII

YTH et. al. arxiv: 1904.06361 (2019)

Z_4 Symmetry indicator for 2D TCsc

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- $$\nu = \kappa \mod 2$$

$\kappa=0$: No edge modes $\kappa=1,3$: Majorana edge modes

$\kappa=2$: No edge modes

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Inversion: $P_{\text{BdG}} = \begin{pmatrix} P & 0 \\ 0 & -P \end{pmatrix}$ for **odd-parity** pairing

$$\nu = \kappa \mod 2$$

$\kappa=0$: No edge modes $\kappa=1,3$: Majorana edge modes

$\kappa=2$: No edge modes, possible **Majorana corner modes!**

YTH et. al. arxiv: 1904.06361 (2019)



Recipe for 2D inversion-protected TCsc with $\kappa=2$

- What kind of **normal state** H_0 (**metallic**)
- What kind of **pairing** Δ

(Predictable)

Band topology $\mathbf{k} \rightarrow H_0(\mathbf{k}) \rightleftharpoons \mathbf{k} \rightarrow H_{\text{BdG}}(\mathbf{k})$
(?)

This is what we want
to understand

Recipe for $\kappa=2$ TCsc in 2D

- ❖ For 2D insulators with time-reversal & inversion:

Step1: Relate 2D TI index with
indicator for 2D Inversion-protected TCI

$$\nu_N = \kappa_N$$

Z_2 still Z_2

- For **metallic** systems (Fermi surface **circle TRIM**):

κ_N can be half integers => not Z_2 anymore

YTH et. al. arxiv: 1904.06361 (2019)

Recipe for $\kappa=2$ TCsc in 2D

- ❖ For 2D systems with time-reversal & inversion:

Step1: Relate 2D TI index with
indicator for 2D Inversion-protected TCI

$$\nu_N = \kappa_N$$

Z_2 still Z_2

- ✓ For **metallic** systems (*Fermi surface away from TRIM*):

“Effectively gapped” topological metal

YTH et. al. arxiv: 1904.06361 (2019)

Recipe for $\kappa=2$ TCsc in 2D

❖ For 2D systems with time-reversal & inversion:

Step1: Relate indices for **metallic H_0 (effectively gapped)**

$$\nu_N = \kappa_N$$

$$1 \quad 1$$

Step2: Relate symmetry indicators for H_0 and H_{BdG}

(Z_2) (Z_4)

$$\kappa = 2 \kappa_N \bmod 4 \quad \text{for weak odd-parity sc}$$

$$2 \quad 1$$

=> TCsc with $\kappa=2$!

YTH et. al. arxiv: 1904.06361 (2019)

Recipe for inversion-protected TCsc

Normal state:

Centrosymmetric quantum spin Hall material

- Gated QSH insulator
- Fermi surface away from TRIMs

Pairing:

Odd-parity sc e.g. *p*-wave

⇒ *Topo crystalline sc* with $\kappa=2$

YTH et. al. arxiv: 1904.06361 (2019)

Recipe for $\kappa=2$ TCsc in 2D

❖ For 2D systems with time-reversal & inversion:

Step1: Relate indices for **metallic H_0 (effectively gapped)**

$$\nu_N = \kappa_N$$

$$1 \quad 1$$

Step2: Relate symmetry indicators for H_0 and H_{BdG}

(Z_2) (Z_4)

$$\kappa = 2 \kappa_N \bmod 4 \quad \text{for weak odd-parity sc}$$

$$2 \quad 1$$

YTH et. al. arxiv: 1904.06361 (2019)

Recipe for inversion-protected TCsc

Normal state:

Centrosymmetric quantum spin Hall material

- Gated QSH insulator
- Fermi surface away from TRIMs

Pairing:

Odd-parity sc e.g. *p*-wave

⇒ Topo crystalline sc with $\kappa=2$

Higher-order?



Show bulk-boundary correspondence!

YTH et. al. arxiv: 1904.06361 (2019)

Bulk-boundary correspondence

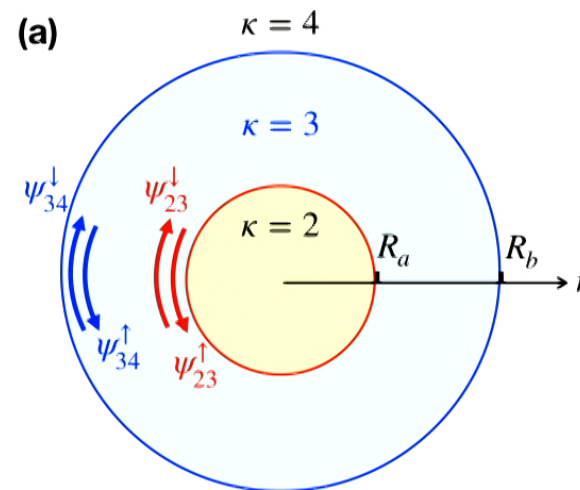
❖ What kind of boundary modes: $\kappa = 2$ | $\kappa = 4$
(vacuum)

1. Minimal model: BHZ model + odd-parity pairing
(Gated) $\kappa = 4 \rightarrow 3 \rightarrow 2$

2. Three-regime geometry

Solve for the two pairs of
helical edge modes from $\kappa = 3$

Shrink $\kappa = 3$ regime $\rightarrow 0$

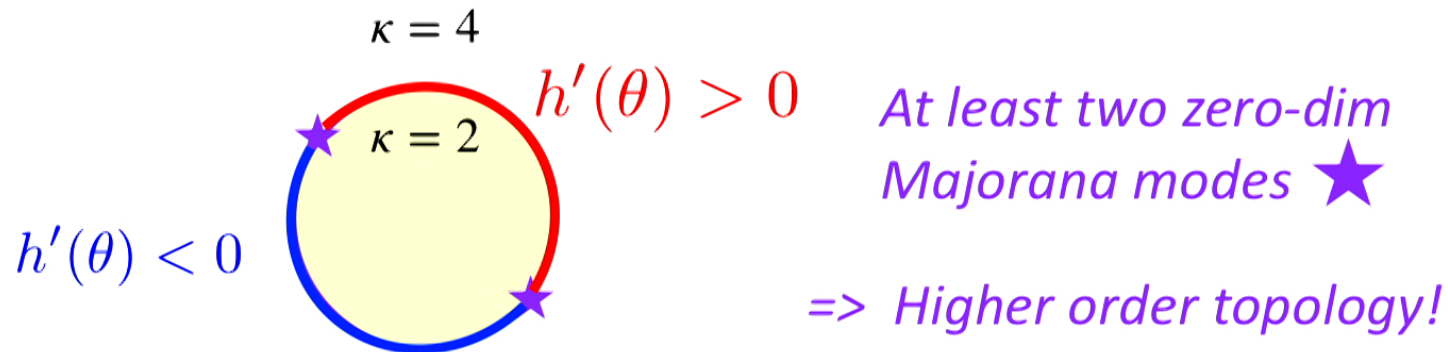


Bulk-boundary correspondence

❖ What kind of boundary modes: $\kappa = 2$ | $\kappa = 4$
(vacuum)

3. **Effective edge theory** for edge modes $\psi_{23}^{\uparrow/\downarrow}$ and $\psi_{34}^{\uparrow/\downarrow}$:

=> Lowest-order back-scattering term $h'(\theta) \propto \alpha \cos \theta + \beta \sin \theta$



$(\kappa=2)$

*Recipe for inv-protected **higher-order** sc:*

1. Centrosymmetric

2. Quantum spin Hall (FS away
from TRIMs)

3. Odd-parity pairing

Go look for the material!

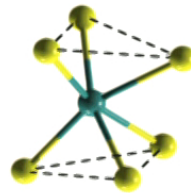
Monolayer WTe₂

❖ Van der Waal material

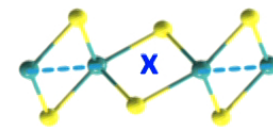
Monolayer from scotch tape

❖ **Centrosymmetric**

Unit cell



Side view



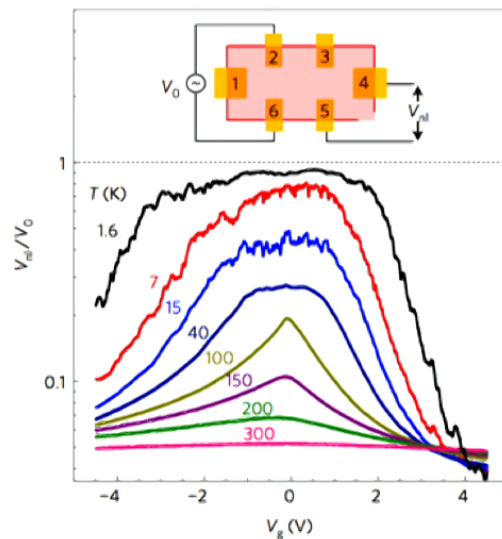
Inversion center

As-grown WTe_2 : Quantum Spin Hall Insulator

❖ Prediction: Qian et. al. Science (2014)
Zheng et. al. Adv. Mat. (2016)

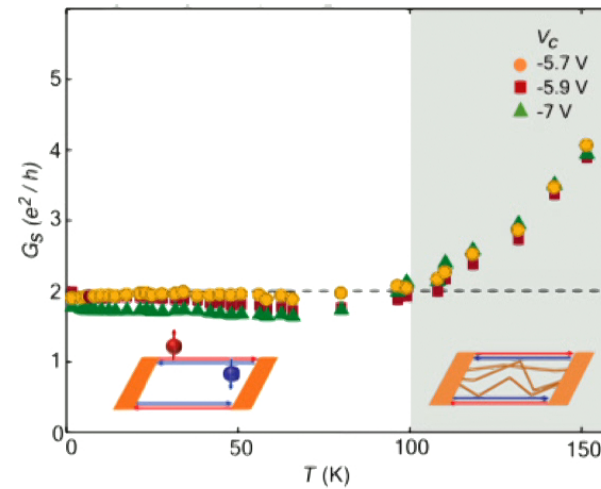
❖ Exp:

Edge conductance



Fei et. al. Nat. Phys. (2017)

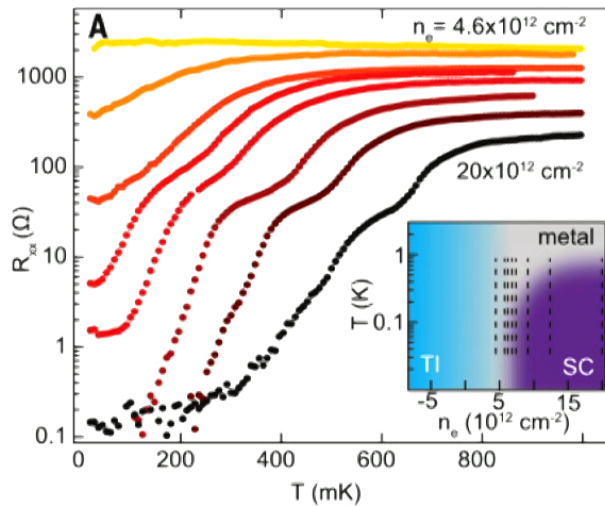
Quantized conductance



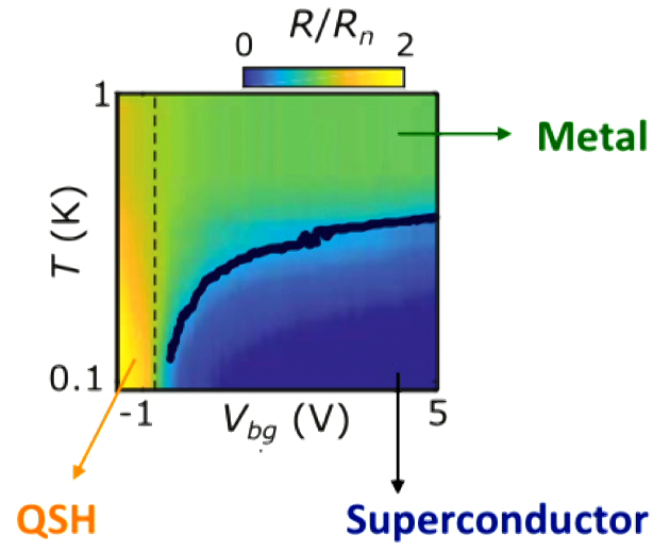
Wu et. al. Science (2018)

ARPES: Tang et. al. Nat Phys (2018)

Gated WTe₂: Superconductor



Sajadi et. al. Science (2018)



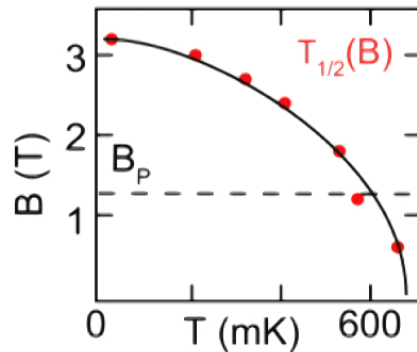
Fatemi et. al. Science (2018)

s-wave? non s-wave?

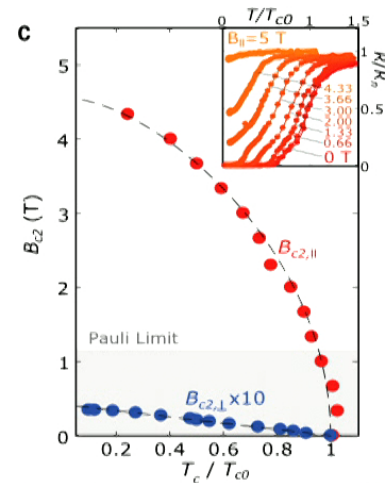
Can it be odd-parity sc?

Hints for unconventional sc

- Low carrier density (small gate voltage)
- Electronic correlation is important (PBE v.s. ARPES)
- In-plane H_{c2} higher than Pauli limit



Sajadi et. al. Science (2018)



Fatemi et. al. Science (2018)

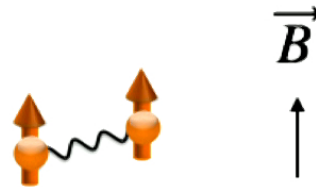
What can cause high H_{c2} ?

Possible origins for high H_{c2}

- Higher spin-orbit scattering rate

- g factor not 2

- Equal-spin pairing

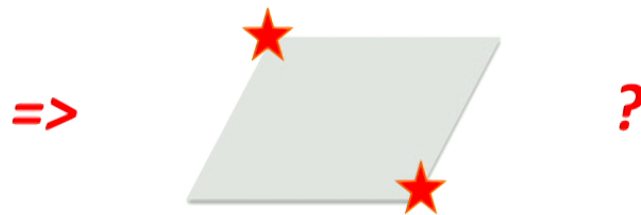


Spin-triplet: $\uparrow\uparrow, \downarrow\downarrow, \uparrow\downarrow + \downarrow\uparrow$

Parity-odd required by fermionic statistics (e.g. **p-wave**)

Gated monolayer WTe_2

- ✓ 1. Centrosymmetric
- ✓ 2. Quantum spin Hall
- ? 3. Odd-parity pairing



Gated monolayer WTe_2

(=> Inversion-protected HOTsc)



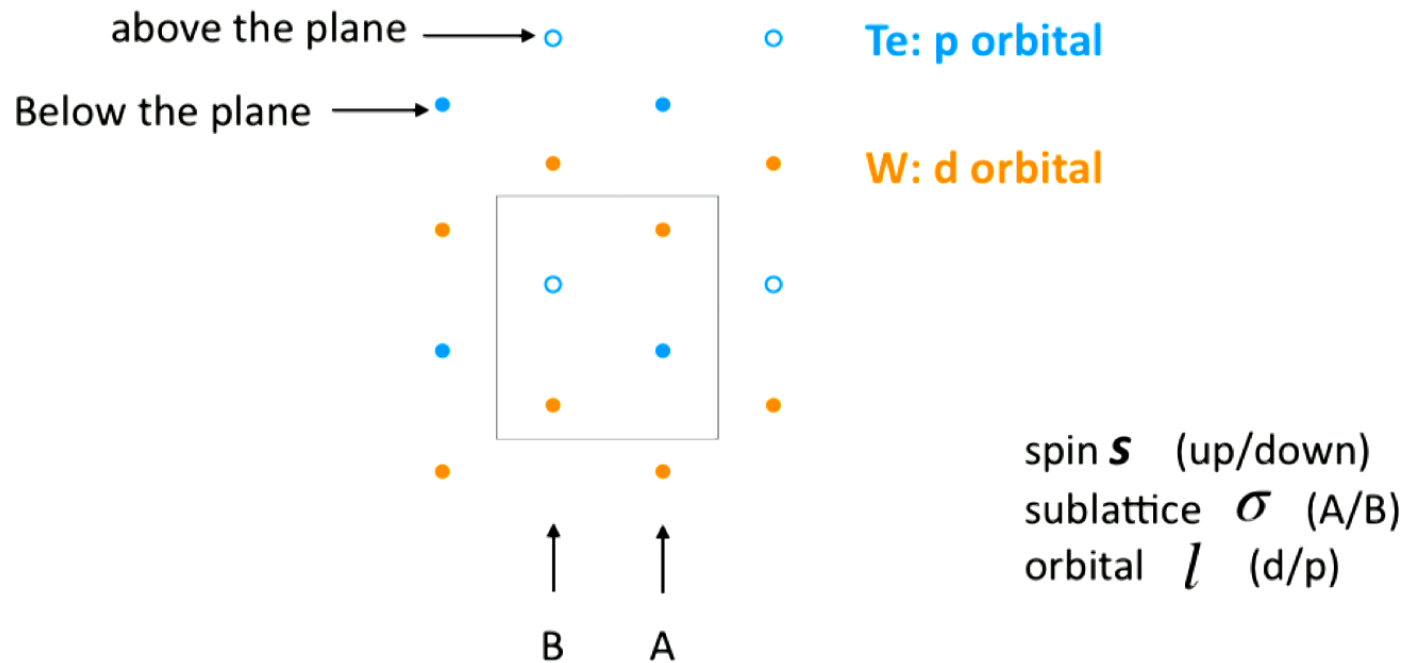
Study pairing symmetry at $H_{in-plane}=0, >0$

&

topo properties (κ , boundary modes)

Lattice symmetries

❖ 1T'- lattice of WTe_2 : (low-energy Wannier centers)



Muechler et. al. PRX (2016)

Lattice symmetries

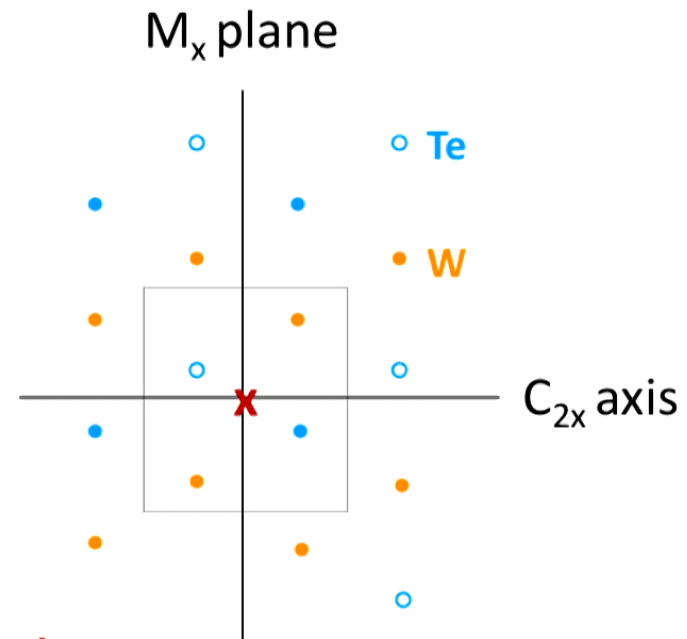
❖ Two **nonsymmorphic** symmetries: extra half translations

- Glide mirror symmetry M_x
- C_{2x} screw rotational symmetry



Inversion $P = M_x C_{2x}$

(Parity is a good quantum number)



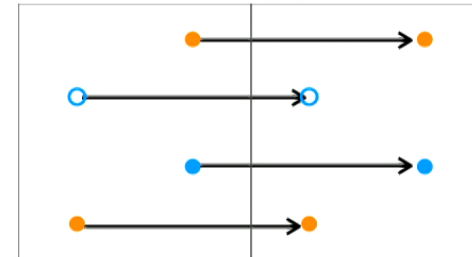
Pairing symmetries

❖ Classify possible pairing symmetries using C_{2x} and M_x

4 irreducible representations

		C_{2x}	M_x		
Parity-even	{	A_g	+	+	Spin-singlet e.g. A_g : on-site pairing
		B_g	-	-	
Parity-odd	{	A_u	+	-	Spin-triplet e.g. B_u : $\Delta_{\uparrow\downarrow+\downarrow\uparrow}$
		B_u	-	+	

Which irrep is energetically favorable?



YTH et. al. arxiv: 1904.06361 (2019)

Mean-field calculation

❖ Solve linearized gap equations

$$\Delta_{\alpha'\beta'}(\mathbf{k}') = - \sum_{\mathbf{k}''\mathbf{k}} V_{\alpha'\beta',\beta''\alpha''}(\mathbf{k}',\mathbf{k}'') \chi_{\beta''\alpha'',\alpha\beta}(\mathbf{k}'',\mathbf{k},T) \Delta_{\alpha\beta}(\mathbf{k})$$

Gap functions

Pairing interaction

Noninteracting
pairing susceptibility

α, β : spin, sublattice, orbital

$\mathbf{k}$: incoming, outgoing momenta

T : temperature

Mean-field calculation

❖ Solve linearized gap equations

$$\Delta_{\alpha'\beta'}(\mathbf{k}') = - \sum_{\mathbf{k}''\mathbf{k}} V_{\alpha'\beta',\beta''\alpha''}(\mathbf{k}',\mathbf{k}'') \chi_{\beta''\alpha'',\alpha\beta}(\mathbf{k}'',\mathbf{k},T) \Delta_{\alpha\beta}(\mathbf{k})$$

Gap functions

Pairing interaction

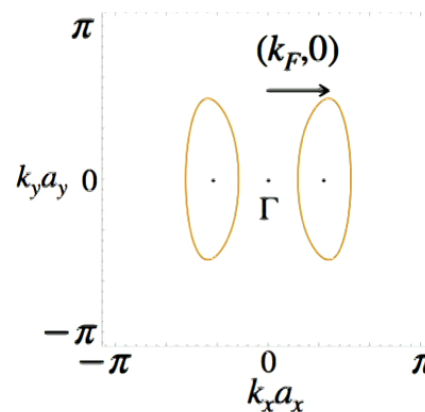
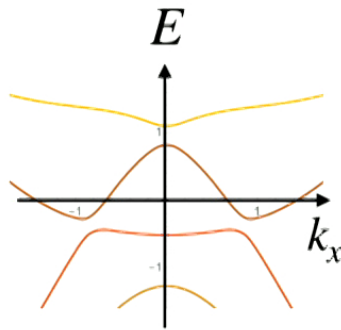
Noninteracting
pairing susceptibility

- Input: **Fermi surface, interactions V**
- Output: **Pairing symmetry** of the dominant gap function (Ag, Bu etc.)

Model: H_0

❖ Kinetic terms: **tight-binding model of Wannier orbitals from DFT**

+ finite chemical potential (gating)



Muechler et. al. PRX (2016)

Ok et. al. 1811.00551 (2018)

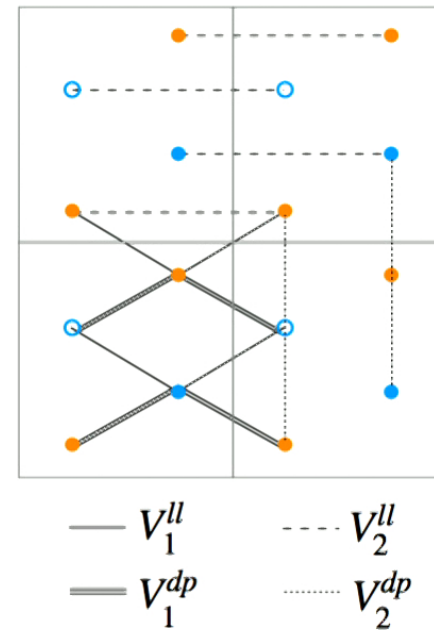
- Include both W and Te orbitals, **SOC**
- Inversion + Time-reversal => **2-fold degeneracy**
- SOC breaks $SU(2)$, but preserves S_z

Model: H_{int}

❖ Microscopic density-density interactions:

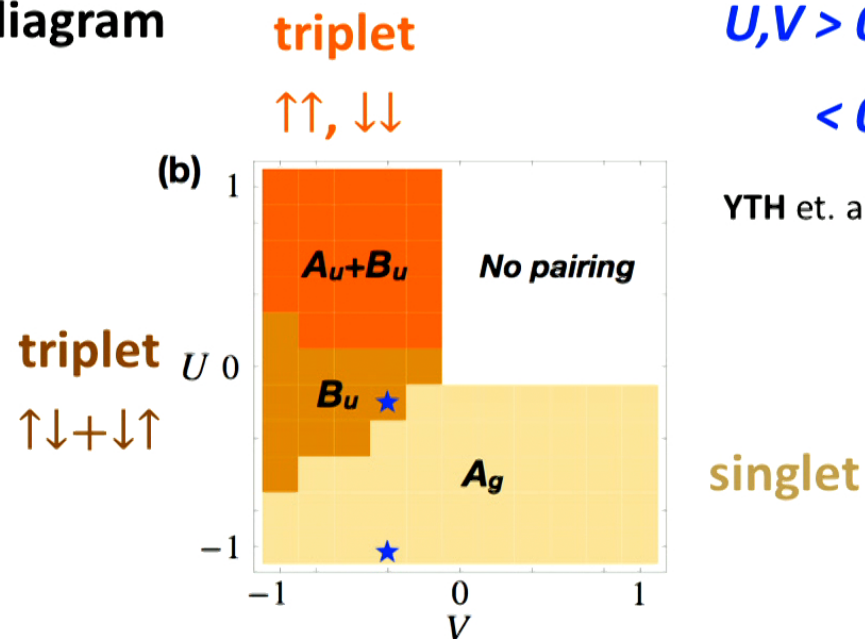
U: On-site interactions

V: Short-range interactions up to nearest neighbor unit cells



Result: pairing symmetry

❖ Phase diagram



$U, V > 0$: repulsion

< 0 : attraction

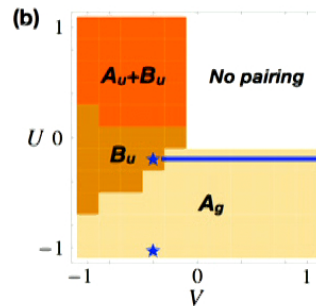
YTH et. al. arxiv: 1904.06361

Q1: Z_4 indicator $\kappa = ?$

Q2: Edge modes? Corner modes?

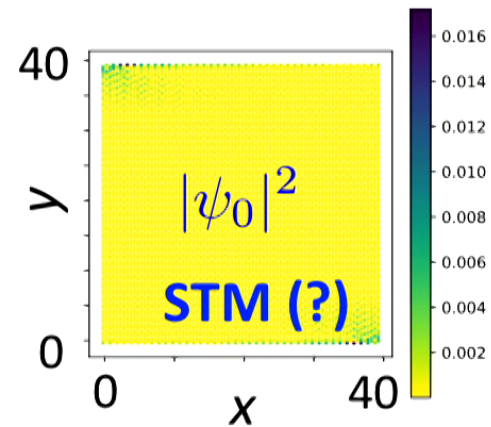
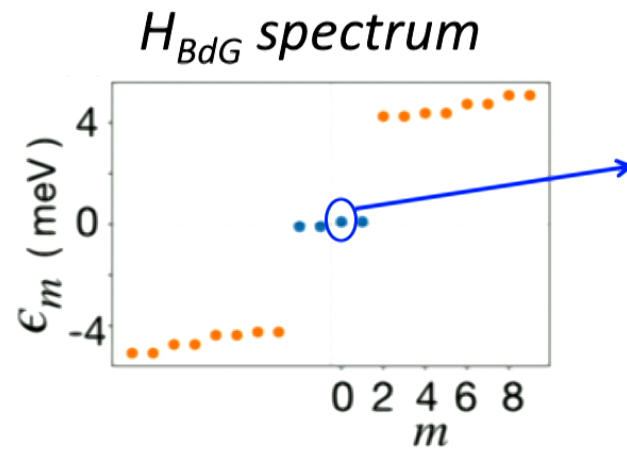
Q3: How does in-plane field affect pairing?

Topological properties



$\kappa=2$

*Inversion-protected
HOT sc !*

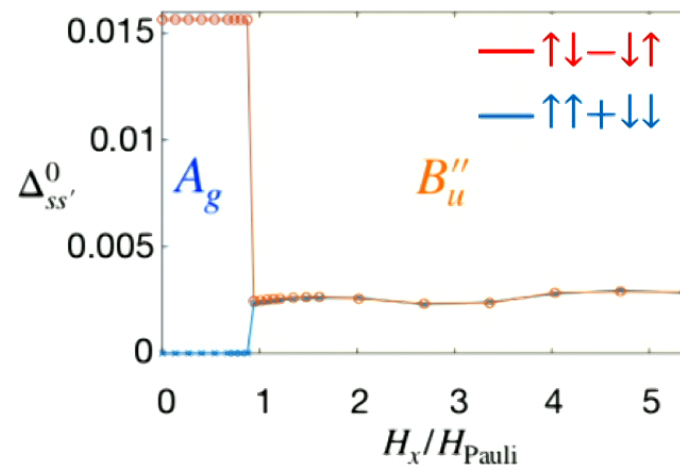
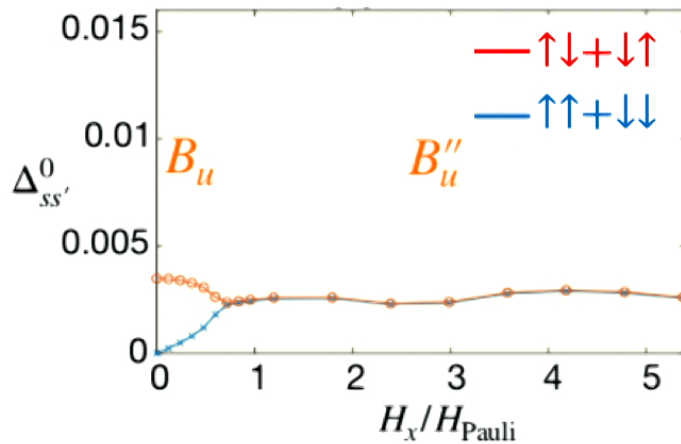
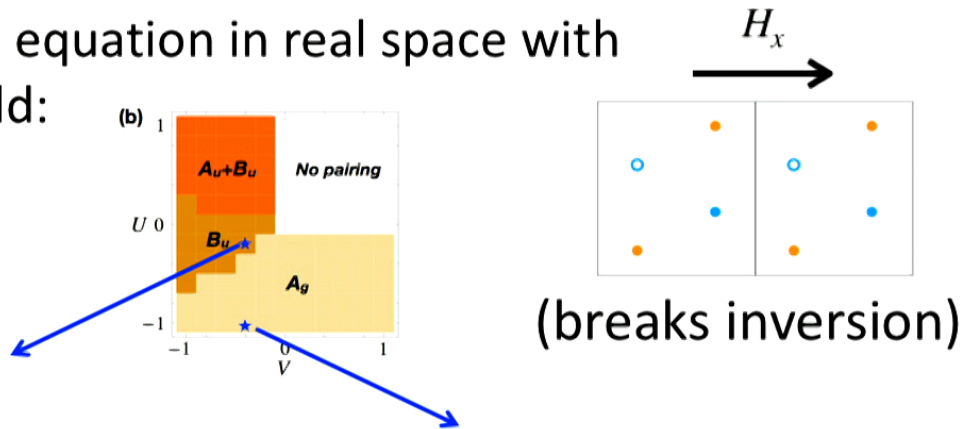


*Majorana Kramer's pairs
at corners*

YTH et. al. arxiv: 1904.06361

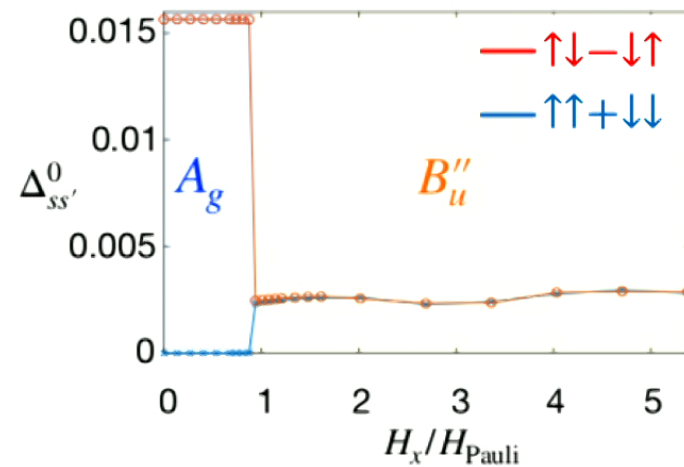
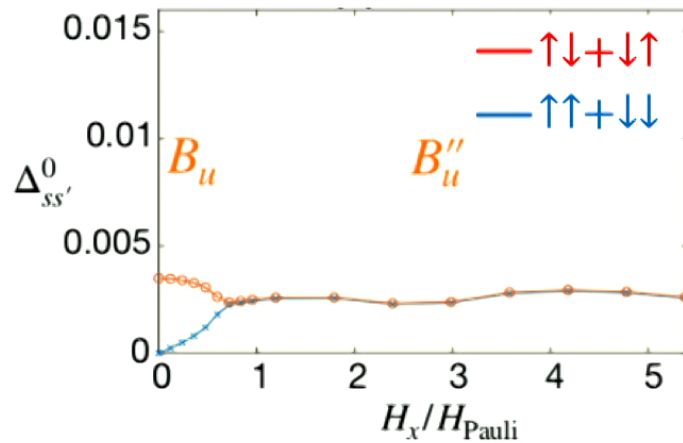
How does in-plane field affect pairing?

❖ Self-consistent gap equation in real space with in-plane magnetic field:



YTH et. al. arxiv: 1904.06361

With in-plane magnetic field H_x



- **Crossover:** $s_x=1$
Spin rotates $\uparrow\downarrow + \downarrow\uparrow \Rightarrow \rightarrow \rightarrow$
Symmetry unchanged
- **First-order transition:**
 $\text{singlet} \Rightarrow \text{triplet} \rightarrow \rightarrow$
 $C_{2x} \text{ even} \Rightarrow \text{odd}$

Bu'' : Extrinsic higher-order topo sc
(if bulk gap is not closed)

YTH et. al. arxiv: 1904.06361

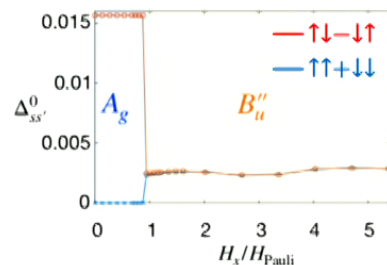
Summary: Inversion-protected HOTsc

General Recipe

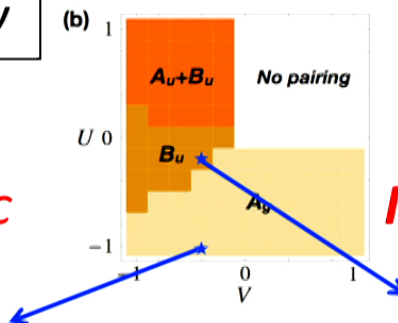
1. Centrosymmetric
2. H_0 = Quantum spin Hall
3. Odd-parity pairing

WTe_2 superconductivity

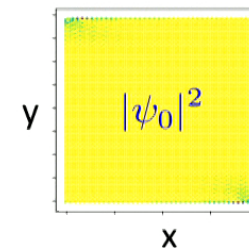
Field-induced triplet sc



(Extrinsic higher-order topo sc)



Intrinsic higher-order topological sc



YTH et. al. arxiv: 1904.06361 (2019)