Title: Entanglement, free energy and C-theorem in DCFT

Speakers: Tatsuma Nishioka

Collection: Boundaries and Defects in Quantum Field Theory

Date: August 08, 2019 - 2:45 PM

URL: http://pirsa.org/19080063

Abstract: The g-theorem is a prominent example of C-theorems in two-dimensional boundary CFT and the extensions are conjectured to hold in higher-dimensional BCFTs. On the other hand, much less is known for C-theorems in a CFT with conformal defects of higher codimensions. I will investigate the entanglement entropy across a sphere and sphere free energy as a candidate for a C-function in DCFT, and show they differ by a universal term proportional to the vev of the stress tensor. Based on this relation, I will propose to use the sphere free energy as a C-function in DCFT. This proposal unifies the previously known theorems and conjectures, and passes several checks, including a few examples in field theories and a holographic proof in simple gravity dual models of DCFTs.

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Entanglement, free energy and $\mathit{C}\text{-theorem}$ in defect CFT Outline 1 Defect conformal field theories 2 Entanglement entropy and sphere free energy in DCFT Towards a C-theorem in DCFT Tatsuma Nishioka | (University of Tokyo) | Aug 8, 2019 @ Perimeter

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Defects in quantum field theory

Defects = Non-local objects in QFTs

- Defined by boundary conditions around them
- Many examples:

1-dim: Line operators (Wilson-'t Hooft loops)

2-dim: Surface operators

Codim-1: Domain walls, interfaces and

boundaries

Codim-2: Entangling surface for entanglement

entropy

 $\int_{\mathbb{S}^2} F = 2\pi q$

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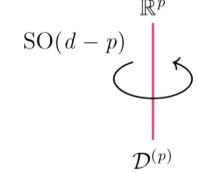
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Conformal defects

- In Euclidean CFT_d, the conformal group is SO(d+1,1)
- p-dimensional conformal defects $\mathcal{D}^{(p)}$ are either flat or spherical, preserving

SO(p+1,1): conformal symmetry on defects

SO(d-p): rotation in the transverse direction



■ Defects allow for defect local operators $\hat{\mathcal{O}}_n(\hat{x})$ \hat{x}^a : parallel coordinates on $\mathcal{D}^{(p)}$ $(a=1,\cdots,p)$

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One-point function

The residual conformal symmetry constrains one-point functions

Scalar primary:

$$\langle\,\mathcal{O}(x)\,
angle^{(\mathrm{DCFT})}=rac{a_{\mathcal{O}}}{|x_{\perp}|^{\Delta}}\;, \qquad x_{\perp}^{i}\; : \mathrm{transverse}\; \mathrm{coordinates}$$

$$(i=p+1,\cdots,d)$$

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$$(i=p+1,\cdots,d)$$

Stress tensor:

$$\langle T^{ab}(x) \rangle^{\text{(DCFT)}} = \frac{d - p - 1}{d} \frac{a_T}{|x_\perp|^d} \delta^{ab}$$

$$\langle T^{ij}(x) \rangle^{\text{(DCFT)}} = -\frac{a_T}{|x_\perp|^d} \left(\frac{p + 1}{d} \delta^{ij} - \frac{x_\perp^i x_\perp^j}{|x_\perp|^2} \right)$$

$$\langle T^{ai}(x) \rangle^{\text{(DCFT)}} = 0$$

N.B.
$$\langle \, T^{\mu\nu}(x) \, \rangle^{({
m DCFT})} = 0$$
 in BCFT $(p=d-1)$ [McAvity-Osborn 95]

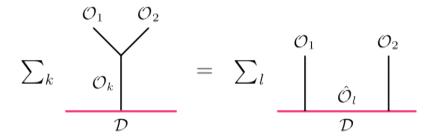
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Applications of conformal defects

Constrain bulk CFT data in defect CFT by crossing symmetry

[Liendo-Rastelli-van Rees 12, Gaiotto-Mazac-Paulos 13, Gliozzi-Liendo-Meineri-Rago 15, Billó-Gonçalves-Lauria-Meineri 16, Lemos-Liendo-Meineri-Sarkar 17, · · ·]



Understand quantum entanglement in QFT:

e.g. Mutual information as a correlator of two defects

[Cardy 13, Bianchi-Meineri-Myers-Smolkin 15, Chen-Chen-Hao-Long 17]

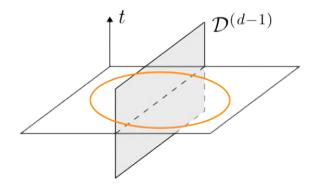
$$I(A, B) \equiv S_A + S_B - S_{A \cup B}$$
$$= \log \frac{\langle \mathcal{D}(\partial A) \mathcal{D}(\partial B) \rangle}{\langle \mathcal{D}(\partial A) \rangle \langle \mathcal{D}(\partial B) \rangle}$$

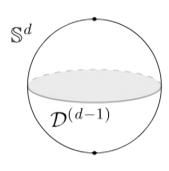
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Goal of this talk

- In DCFT we will study
 - Entanglement entropy across a sphere
 - Sphere free energy
- How do they depend on defect data?
- What is the measure of degrees of freedom associated to defects?





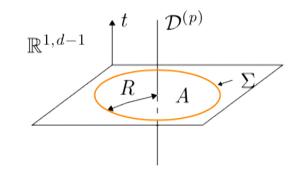
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Entanglement entropy across a sphere

- lacksquare A: a ball centered at the origin
- lacksquare $\mathcal{D}^{(p)}$: a p-dim flat defect



After a conformal transformation

(CHM map [Casini-Huerta-Myers 11, Jensen-O'Bannon 13])

The *n*-th Rényi entropy

$$S_n^{(\text{DCFT})} = \frac{1}{1-n} \log \frac{Z^{(\text{DCFT})}[\mathbb{S}_n^d]}{(Z^{(\text{DCFT})}[\mathbb{S}^d])^n}$$

 \mathbb{S}_n^d : $n\text{-fold cover of }\mathbb{S}^d$

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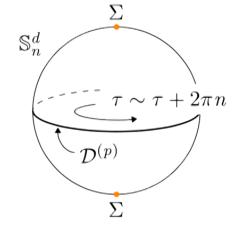
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Defect entropy

■ The excess of EE is measured by

Defect entropy

$$S_{\text{defect}} \equiv \lim_{n \to 1} \left(S_n^{(\text{DCFT})} - S_n^{(\text{CFT})} \right)$$



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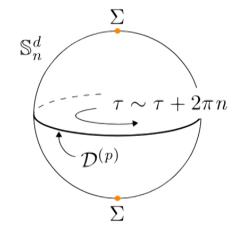
Defect entropy

■ The excess of EE is measured by

Defect entropy

$$S_{\text{defect}} \equiv \lim_{n \to 1} \left(S_n^{(\text{DCFT})} - S_n^{(\text{CFT})} \right)$$
$$= \lim_{n \to 1} \frac{1}{1 - n} \log \frac{\langle \mathcal{D}^{(p)} \rangle_n}{\langle \mathcal{D}^{(p)} \rangle^n}$$

$$\langle \mathcal{D}^{(p)} \rangle_n \equiv rac{Z^{(ext{DCFT})}[\mathbb{S}_n^d]}{Z^{(ext{CFT})}[\mathbb{S}_n^d]}$$
 $\langle \mathcal{D}^{(p)} \rangle \equiv \langle \mathcal{D}^{(p)} \rangle_1 \quad (\text{vev of } \mathcal{D}^{(p)})$



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Entanglement, free energy and C-theorem in defect CFT | Entanglement entropy and sphere free energy in DCFT

$n \to 1$ limit

Expansion around n=1 ($\delta g_{\mu\nu}=O(n-1)$)

$$\log Z^{(\text{DCFT})}[\mathbb{S}_n^d] = \log Z^{(\text{DCFT})}[\mathbb{S}^d]$$
$$-\frac{1}{2} \int_{\mathbb{S}^d} \delta g_{\mu\nu} \langle T^{\mu\nu} \rangle_{\mathbb{S}^d}^{(\text{DCFT})} + O\left((n-1)^2\right)$$

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■ In CFT [Casini-Huerta-Myers 11]

$$\langle T^{\mu\nu} \rangle_{\mathbb{S}^d}^{(CFT)} = 0$$

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■ In CFT [Casini-Huerta-Myers 11]

$$\langle T^{\mu\nu} \rangle_{\mathbb{S}^d}^{\text{(CFT)}} = 0 \qquad \Rightarrow \qquad S^{\text{(CFT)}} = \log Z^{\text{(CFT)}}$$

In DCFT

$$\langle T^{\mu\nu} \rangle_{\mathbb{S}^d}^{(DCFT)} \neq 0 \ (\propto a_T)$$

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Universal relation

Defect entropy and sphere free energy [Kobayashi-TN-Sato-Watanabe 18]

$$S_{\text{defect}} = \log \langle \mathcal{D}^{(p)} \rangle - \frac{2(d-p-1)\pi^{d/2+1}}{\sin(\pi p/2) d\Gamma(p/2+1)\Gamma((d-p)/2)} a_T$$

- Dimensional regularization is assumed
- Equality holds up to UV divergences
- lacktriangle Reproduces a known result when p=1 [Lewkowycz-Maldacena 13]
- The second term in rhs vanishes when p = d 1 (codimension-one)

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Interface entropy

■ Interface CFT, $\mathcal{D}^{(d-1)} = \mathcal{I}$

Interface entropy

$$S_{\mathcal{I}} = S^{(\text{ICFT})} - \frac{S^{(\text{CFT}_+)} + S^{(\text{CFT}_-)}}{2}$$

 \mathcal{I} CFT_{+} CFT_{-}

Universal relation:

$$S_{\mathcal{I}} = \log \langle \mathcal{I} \rangle , \qquad \langle \mathcal{I} \rangle \equiv \frac{Z^{(\text{ICFT})}[\mathbb{S}^d]}{(Z^{(\text{CFT}_+)}[\mathbb{S}^d] Z^{(\text{CFT}_-)}[\mathbb{S}^d])^{1/2}}$$

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Half-BPS Janus interfaces

■ In 2d $\mathcal{N}=(2,2)$ and 4d $\mathcal{N}=2$ SCFTs

$$Z^{(\mathrm{SCFT})}[\mathbb{S}^d](\tau,\bar{\tau}) = \left(\frac{r}{\epsilon}\right)^{-4A} \exp\left[K(\tau,\bar{\tau})/12\right]$$

- lacksquare r: radius, ϵ : UV cutoff, A: type-A central charge
- $lackbox{ } K(au, ar{ au})$: Kähler potential on a conformal manifold

[Jockers-Kumar-Lapan-Morrison-Romo 12, Gomis-Lee 12, Gerchkovitz-Gomis-Komargodski 14]

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- Half-BPS Janus interfaces [Drukker-Gaiotto-Gomis 10, Goto-Okuda 18]

$$Z^{(\mathrm{ICFT})}[\mathbb{S}^d] = Z^{(\mathrm{SCFT})}[\mathbb{S}^d](\tau_+, \bar{\tau}_-)$$

∃ Kähler ambiguity [Gomis-Ishtiaque 14]

$$K(\tau_+, \bar{\tau}_-) \to K(\tau_+, \bar{\tau}_-) + \mathcal{F}(\tau_+) + \bar{\mathcal{F}}(\bar{\tau}_-)$$

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Interface entropy as Calabi's diastasis

Using supersymmetric Rényi entropy [TN-Yaakov 13] to preserve SUSY

Interface entropy as Calabi's diastasis [Goto-Nagano-TN-Okuda, WIP]

$$S_{\mathcal{I}} = \log |\langle \mathcal{I} \rangle|$$

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Interface entropy as Calabi's diastasis

Using supersymmetric Rényi entropy [TN-Yaakov 13] to preserve SUSY

Interface entropy as Calabi's diastasis [Goto-Nagano-TN-Okuda, WIP]

$$S_{\mathcal{I}} = -\frac{1}{24} \left[K(\tau_+, \bar{\tau}_+) + K(\tau_-, \bar{\tau}_-) - K(\tau_+, \bar{\tau}_-) - K(\tau_-, \bar{\tau}_+) \right]$$

- Bulk part $(\propto A)$ cancels out
- No Kähler ambiguity
- lacksquare Agree with the known result in 2d [Bachas-BrunnerDouglas-Rastelli 13, Bachas-Plencner 16]
- Reproduces a conjectured form in 4d [Goto-Okuda 18]

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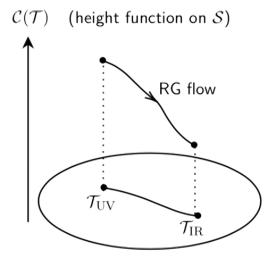
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C-theorem

C-theorem (weak)

 \exists a function $\mathcal{C}(\mathcal{T})$ on a theory space s.t.

$$\begin{split} \mathcal{T}_{\mathrm{UV}} & \xrightarrow[\mathsf{RG} \; \mathsf{flow}]{} \mathcal{T}_{\mathrm{IR}} \\ \Rightarrow & \mathcal{C}(\mathcal{T}_{\mathrm{UV}}) \geq \mathcal{C}(\mathcal{T}_{\mathrm{IR}}) \end{split}$$



$$\mathcal{S} = \mathsf{space} \; \mathsf{of} \; \mathsf{QFTs}$$

- lacksquare $\mathcal{C}(\mathcal{T})$ called a \mathcal{C} -function (pprox resource measure)
- Regarded as a measure of degrees of freedom in QFT
- Constrains the dynamics under RG if holds

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- 2d: Zamolodchikov's c-theorem [Zamolodchikov 86]
- lacksquare even $d\colon$ A-theorem $\left(\langle\,T_\mu^\mu
 angle_{\mathbb{S}^d}\propto A
 ight)$ [Cardy 88, Komargodski-Schwimmer 11]

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- 2d: Zamolodchikov's c-theorem [Zamolodchikov 86]
- lacksquare even $d\colon$ A-theorem $(\langle\,T^\mu_\mu
 angle_{\mathbb{S}^d}\propto A)$ [Cardy 88, Komargodski-Schwimmer 11]
- odd d: F-theorem [Myers-Sinha 10, Jafferis-Klebanov-Pufu-Safdi 11, Klebanov-Pufu-Safdi 11]

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- lacktriangleright Proof with entanglement entropy in $d \leq 4$ [Casini-Huerta 04, 12, Casini-Testé-Torroba 17]

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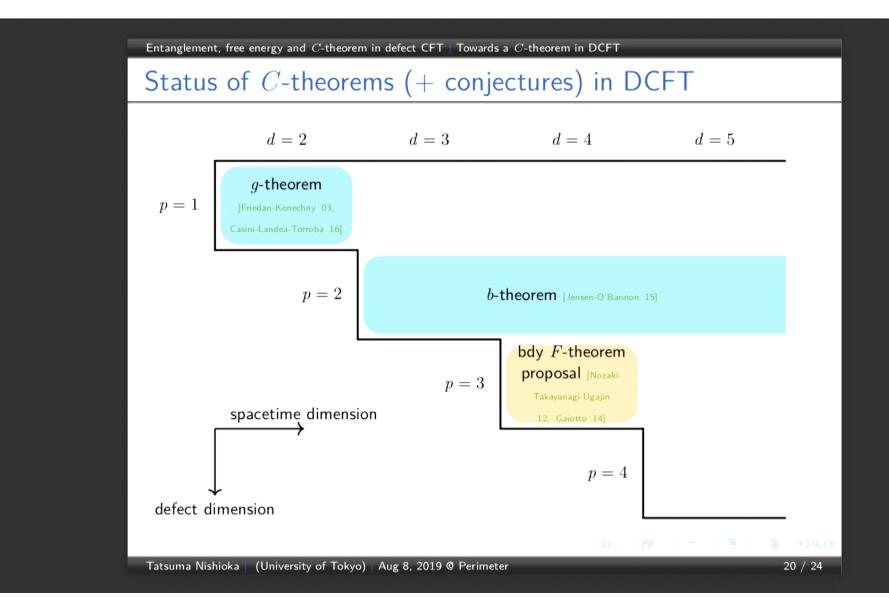
Generalized F-theorem conjecture [Giombi-Klebanov 14]

$$\tilde{F} \equiv \sin\left(\frac{\pi d}{2}\right) \log Z[\mathbb{S}^d], \qquad \tilde{F}_{\text{UV}} \ge \tilde{F}_{\text{IR}}$$

Reduces to the F- and A-theorems

$$ilde{F} = egin{cases} F & d: \text{ odd} \\ rac{\pi}{2}A & (\text{conformal anomaly}) & d: \text{ even} \end{cases}$$

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C-function in DCFT?

- Candidate of C-functions
 - entanglement entropy: holographic model (p = d 1) [Estes-Jensen-O'Bannon-Tsatis-Wrase 14]
 - sphere free energy: bdy F-thm, b-thm (p=2), Wilson loop RG flow (d=4,p=1) [Beccaria-Giombi-Tseytlin 17]
- lacktriangle These two agree when p=d-1 due to the universal relation

Are both C-functions for any d and p?

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C-theorem in DCFT: conjecture

■ Defect RG flow triggered by a relevant defect operator:

$$I = I_{\text{DCFT}} + \hat{\lambda} \int d^p \hat{x} \sqrt{\hat{g}} \,\hat{\mathcal{O}}(\hat{x})$$

Conjecture [Kobayashi-TN-Sato-Watanabe 18]

The universal part of the sphere free energy

$$\tilde{D} \equiv \sin\left(\frac{\pi p}{2}\right) \log |\langle \mathcal{D}^{(p)} \rangle|$$

does not increase along any defect RG flow

$$\tilde{D}_{\mathrm{UV}} \geq \tilde{D}_{\mathrm{IR}}$$

■ Same form as the generalized F-thm: $\tilde{F} = \sin\left(\frac{\pi d}{2}\right) \log Z[\mathbb{S}^d]$

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Checks

- Sphere free energy decreases under defect RG flows in
 - Conformal perturbation theory
 - Wilson loops (p=1) in 3d and 4d
 - Holographic models (a proof assuming null energy condition)
- Counterexamples for the monotonicity of entanglement entropy
 - Wilson loop RG flows [Kobayashi-TN-Sato-Watanabe 18], surface operators (p=2) [Jensen-O'Bannon-Robinson-Rodgers 18]
 - Holographic Wilson loops [Kumar-Silvani 16, 17] and surface operators [Rodgers 18]

$$S_{\text{defect}} = \log \langle \mathcal{D}^{(p)} \rangle - \frac{2(d-p-1)\pi^{d/2+1}}{\sin(\pi p/2) d\Gamma(p/2+1) \Gamma((d-p)/2)} a_T$$

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Summary and future work

■ Summary:

- Find the universal relation between EE and sphere free energy
- Derive the interface entropy as Calabi's diastasis
- Propose a C-theorem in DCFT

■ Future work:

- Proof in SUSY theories? (cf. F- and a-maximizations [Jafferis 10, Closset-Dumitrescu-Festuccia-Komargodskia-Seiberg 12, Intriligator-Wecht 03])
- Proof using entropic inequalities as in g-thm [Casini-Landea-Torroba 16]?
- Constrains on the dynamics of defect RG flows?
- Dependence of defect entropy on bulk marginal deformation? (cf. [Herzog-Shamir 18, Bianchi 18] for sphere free energy)

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