

Title: Wilson lines in AdS3 gravity

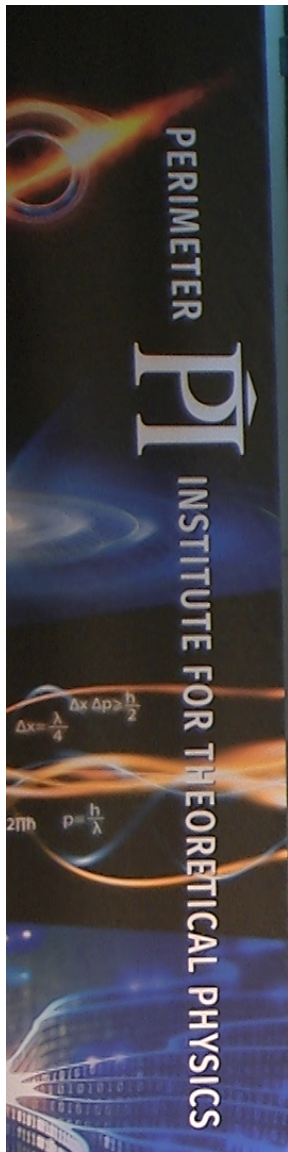
Speakers: Alejandra Castro

Collection: Boundaries and Defects in Quantum Field Theory

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Abstract: The partnership between 3D gravity and Chern-Simons theory is well-known and powerful, but many aspects of this relation are unclear. In this talk I'll discuss Wilson lines in Chern-Simons theory and their role in AdS3 gravity. Wilson lines are interesting objects to explore locality in Chern-Simons theory, and how they affect the partnership at the quantum level.



AdS₃ Gravity
 $E + \omega \Lambda < 0$

Chern-Simons
w. $SO(2,2) = SL(2) \times SL(2)$

AdS₃ Gravity vs
 EH w $\Lambda < 0$

Chern-Simons
 w. $SO(2,1) = SL(2) \times SL(2)$

Action
 e.o.m

* bndy gravitons
 (bndy conditions)

* Euc BTZ
 (thermo, regularity)

Nice:

- * Higher Spin Thys
- * Non-relativistic

Promise:

QG w/o Geometry

AdS₃ Gravity vs

EH w $\Lambda < 0$

↓
Local Properties

Geodesics
Propagators

Chern-Simons

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Local obs in CS?

Wilson lines

$$W_R(x_i, x_f) = \langle U_i | P \exp(-\oint_\gamma A) P \exp(-\oint_\gamma \bar{A}) | U_f \rangle$$

γ : curve with endpoints @ x_i, x_f

R : rep of $SO(2,2)$ $|U\rangle$ states in Rep.

→ Probe, we want fit data of a massive particle

$R \rightarrow$ Hilbert space \rightarrow highest weight reps of $SL(2)$
 \downarrow
 Commir = "mass"

$$h.w.: SL(2): \ell_0, \ell_1, \ell_{-1}$$

$$\ell_0 |h\rangle = h |h\rangle$$

$$\ell_{\pm} |h\rangle = 0$$

$$|h, k\rangle = (\ell_{-1})^k |h\rangle$$

$$C_2 = 2h(h-1) = "m^2"$$

$$\text{Say } |U\rangle = |h, k\rangle \otimes |\bar{h}, \bar{k}\rangle$$

$$W_R = \langle h, k | e^A | h, k \rangle \times \langle \bar{h}, \bar{k} | e^{\bar{A}} | \bar{h}, \bar{k} \rangle$$

AdS₃ Gravity

$$E \neq 0 \quad \Lambda < 0$$



Local Properties

Geodesics

Propagators

$$g_{\mu\nu} = \text{Tr}(A - \bar{A}^2)$$

VS

Chern-Simons

$$w. \quad SO(2,2) = \underset{\downarrow A}{SL(2)} \times \underset{\downarrow \bar{A}}{SL(2)}$$

Action

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(bndy conditions)

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$$W_R(x_i, x_f)$$

heg n-Simons

$$SO(2,1) = \underset{\downarrow \vec{A}}{SL(2)} \times \underset{\downarrow \vec{A}}{SL(2)}$$

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Wilson lines

$$W_R(x_i, x_f) = \langle U_i | P \exp(-\oint_{\gamma} A) P \exp(-\oint_{\gamma} \bar{A}) | U_f \rangle$$

$$A = L dL^{-1}$$

$$\bar{A} = R^{-1} dR \quad L, R \in SL(2)$$

$$= \langle U | G(L) \bar{G}(R^{-1}) | U \rangle$$

h.w.: s

$l_0 | h \rangle$

$l_1 | h \rangle$

$| h, k \rangle =$

$C_2 = 2h$

Say $| U \rangle = | h$

$W_R = \langle h | k$

Wilson lines

$$W_R(x_i, x_f) = \langle U | \mathcal{P} \exp(-\oint_{\gamma} A) \mathcal{P} \exp(-\oint_{\gamma} \bar{A}) | U_f \rangle$$

$$A = L dL^{-1} \quad \bar{A} = R^{-1} dR \quad L, R \in SL(2)$$

$$= \langle U | G(L) \bar{G}(R^{-1}) | U \rangle$$

$$= \langle U | LUR \rangle = \frac{e^{-\alpha(x_i, x_f)h}}{1 - e^{\alpha(x_i, x_f)h}}$$

= Bulk to bulk propagator of a massive scalar field $\bar{\Phi}$

$\alpha(x_i, x_f)$: geodesic distance on

$$\tilde{g}_{\mu\nu} = \text{Tr}(A - U \bar{A} U)^2$$

$$Q_n |U\rangle = 0 \quad Q_n = \partial_n - D_n^{\dagger}(U) \bar{Q}_n$$

$$G(L) \bar{G}(R^{-1}) |U\rangle = |LUR\rangle$$

Ex of coherent states:

1) Crosscap state

$$|U_c\rangle = \sum_{k=0}^{\infty} |h_1, k\rangle |h_1, k\rangle$$

$$(L-a - \bar{L}-a) |U_c\rangle = 0$$

2) Ishibashi state

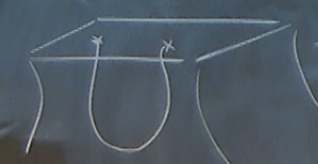
$$|U_{\text{Ish}}\rangle = \sum_k (-1)^k |h_1, k\rangle |h_1, k\rangle$$

$$(L-a - (-1)^a \bar{L}-a) |U_{\text{Ish}}\rangle = 0$$

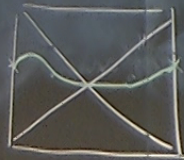
Wilson lines

$$W_R(x_i, x_f) = \langle U_i | \mathcal{P} \exp(-\oint A) \mathcal{P} \exp(-\oint \bar{A}) | U_f \rangle$$

Applications:

1)  $\rightarrow W_R(x_i, x_f) = \langle O_H, U_L, O_L, O_H \rangle$

2) Eternal BHs

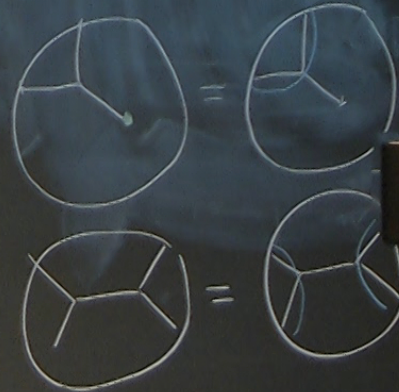


KMS cond on W_R

3) BH entropy

$$W_R(C) = e^{\frac{S_{\text{BH}}}{h} (1 - \text{loop dot})}$$

4) Junctions



5) Quantum correc. W_R

