Title: Axion production and detection with superconducting RF cavities

Speakers: Vijay Narayan

Series: Particle Physics

Date: July 26, 2019 - 1:00 PM

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Abstract: We propose a novel design of a laboratory search for axions based on photon regeneration with superconducting RF cavities. Our particular setup uses a toroid as a region of confined static magnetic field, while production and detection cavities are positioned in regions of vanishing external field. This permits cavity operation at quality factors of  $Q \sim 10 \hat{A}^1 \hat{a}^{\bullet}$ . The limitations due to fundamental issues such as signal screening and back-reaction are discussed, and the optimal sensitivity is calculated. This experimental design can potentially probe axion-photon couplings beyond astrophysical limits, comparable and complementary to next generation optical experiments.

Pirsa: 19070001 Page 1/34

# New light boson?

Strong evidence for physics BSM. Where to look?

Physics in far UV can lead to low-energy observables, including new light degrees of freedom

Pirsa: 19070001 Page 2/34

## New light boson?

Strong evidence for physics BSM. Where to look?

Physics in far UV can lead to low-energy observables, including new light degrees of freedom

Axions are a generic expectation

$$\mathcal{L} \supset \frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} (\partial_{\mu} a)^2 - \frac{1}{2} m_a^2 a^2 + \frac{1}{4} g a F_{\mu\nu} \widetilde{F}^{\mu\nu}$$

Light, weakly-coupled particle. How to detect?

Pirsa: 19070001 Page 3/34

Coherent axion-photon conversion in static B<sub>o</sub>

$$P_{a\to\gamma} \sim (gB_0L)^2 \sim 10^{-19} \left(\frac{g \text{ GeV}}{10^{-11}}\right)^2 \left(\frac{B_0}{5 \text{ T}}\right)^2 \left(\frac{L}{10 \text{ m}}\right)^2$$

Coherence if  $\,qL \sim {m_a^2 L \over \omega} \lesssim 1\,$ 

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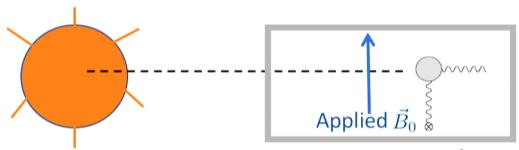
Pirsa: 19070001 Page 5/34

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"Sun shining through a Wall"



[Irastorza et al, '17]

CAST:  $g < 7.10^{-11} \text{ GeV}^{-1} \text{ for } m_a < eV$ 

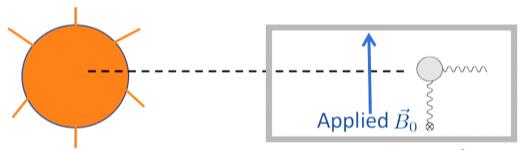
(comparable to stellar cooling bounds)

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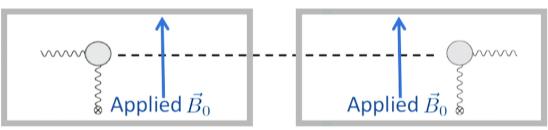
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[van Bibber et al '87, ...]

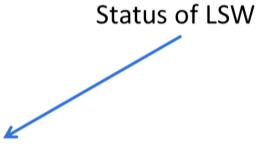
Possible at optical or radio frequencies

Key point: need large EM fields to overcome small coupling

Pirsa: 19070001 Page 8/34

# Optical vs. RF

[..., Graham et al '16]



**Optical Cavities** 

 $\omega, L^{-1}$  are independent

ALPS:  $g < 5.10^{-8} \text{ GeV}^{-1}$ for  $m_a < \text{meV}$ 

Next generation with L  $\sim$  100 m ALPS II (projected): g  $< 2.10^{-11}$  GeV<sup>-1</sup>

Pirsa: 19070001 Page 9/34

## Optical vs. RF

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**RF Cavities** 

$$\omega, L^{-1} \sim \mathcal{O}(\mathrm{GHz})$$

CROWS:  $g < 10^{-7} \text{ GeV}^{-1}$ for  $m_a < \mu eV$ 

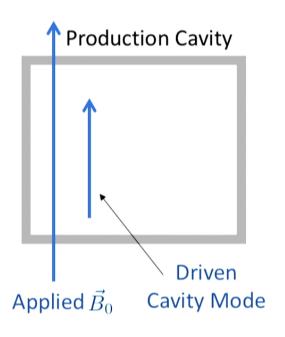
This work: possibilities with superconducting RF technology?

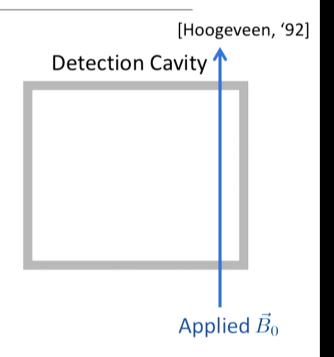
See: [Bogorad, Hook, Kahn, Soreq, '19] for very distinct proposal

Page 10/34

Status of LSW

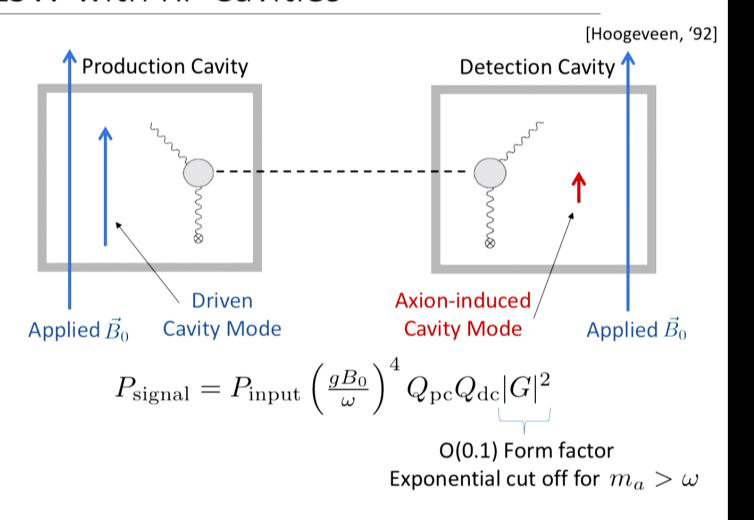
## LSW with RF Cavities





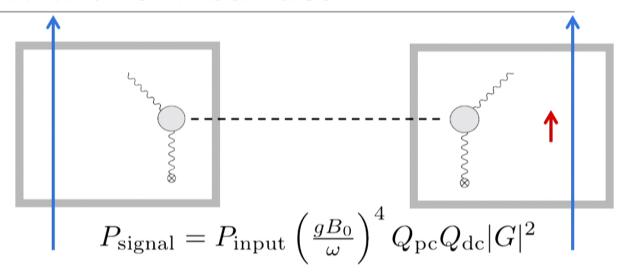
Pirsa: 19070001 Page 11/34

## LSW with RF Cavities



Pirsa: 19070001 Page 12/34

## LSW with SRF Cavities



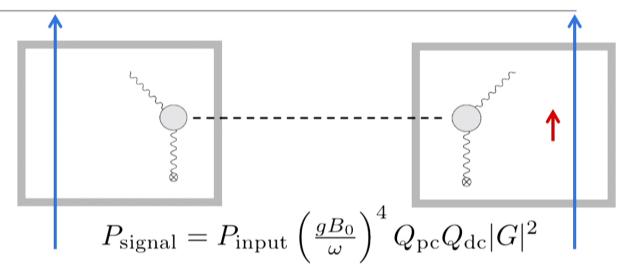
Normal conducting RF  $Q \sim 10^5 - 10^6$ 

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Superconducting RF 
$$Q \sim 10^{10}-10^{12}$$

$$Q \sim 10^{10} - 10^{12}$$

## LSW with SRF Cavities



Normal conducting RF  $Q \sim 10^5 - 10^6$ 

Superconducting RF  $Q \sim 10^{10}-10^{12}$ 

B > O(0.2 T) critical field, flux penetration degrades Q

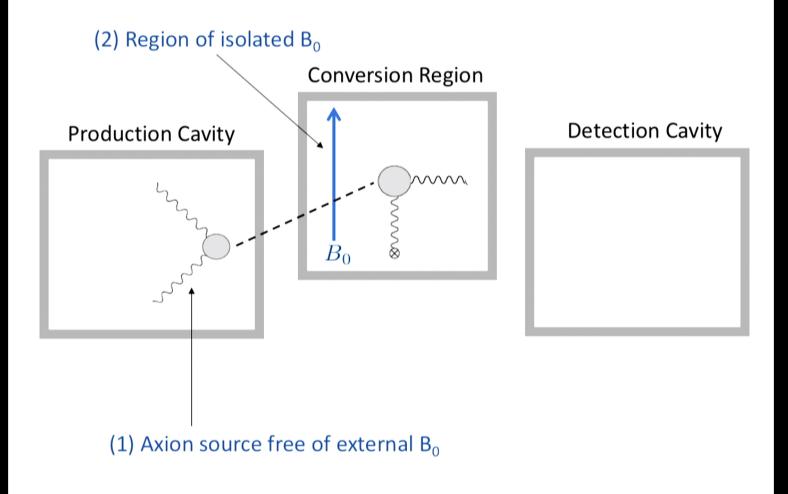
Challenge: re-design such that large B and SRF can co-exist

Pirsa: 19070001 Page 14/34

# New Design for LSW with SRF **Conversion Region Detection Cavity Production Cavity**

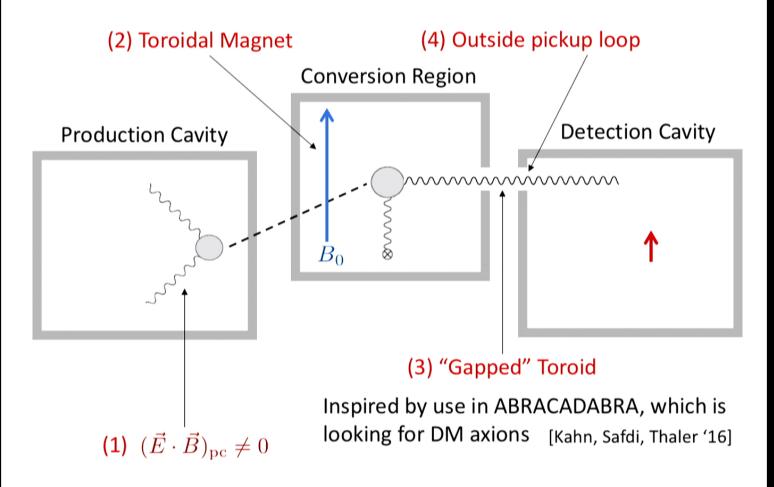
Pirsa: 19070001 Page 15/34

# New Design for LSW with SRF



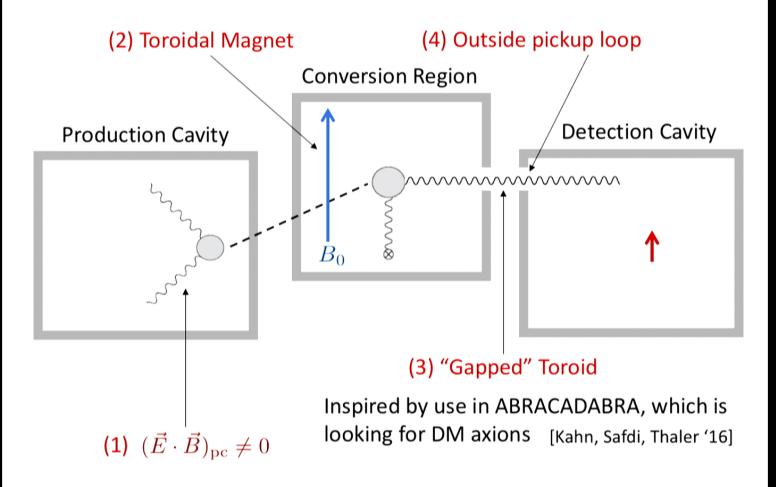
Pirsa: 19070001 Page 16/34

# New Design for LSW with SRF



Pirsa: 19070001 Page 17/34

## New Design for LSW with SRF



Pirsa: 19070001 Page 18/34

# Axion Electrodynamics

Axion EOM: 
$$(\Box + m_a^2)a(x) = -g\vec{E} \cdot \vec{B}$$

Modifies Maxwell: 
$$\vec{\nabla} \cdot \vec{E} = -g\vec{B} \cdot \vec{\nabla} a$$

$$\vec{\nabla} \times \vec{B} = \frac{\partial \vec{E}}{\partial t} - g \left( \vec{E} \times \vec{\nabla} a - \vec{B} \frac{\partial a}{\partial t} \right)$$

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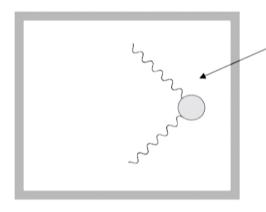
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Axion Production 
$$\implies a(x) = -ge^{i\omega t} \int_{\mathrm{pc}} d^3y \ \frac{e^{ik|\vec{x}-\vec{y}|}}{4\pi|\vec{x}-\vec{y}|} (\vec{E}\cdot\vec{B})$$

## **SRF Axion Source**

#### (1) Axion source without external B<sub>0</sub>



Cavity mode(s) such that E.B not identically vanishing

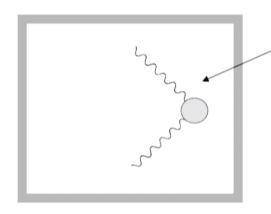
Fundamentally limited by SRF critical field (independent of Q, input power, etc.)

$$(\vec{E} \cdot \vec{B})_{\rm pc} \lesssim (0.2 \text{ T})^2$$

Pirsa: 19070001 Page 21/34

### SRF Axion Source

#### (1) Axion source without external $B_0$



Cavity mode(s) such that E.B not identically vanishing

Fundamentally limited by SRF critical field (independent of Q, input power, etc.)

$$(\vec{E} \cdot \vec{B})_{\rm pc} \lesssim (0.2 \text{ T})^2$$

Compare with normal conducting RF with external B<sub>0</sub>

$$(\vec{E} \cdot \vec{B}) \sim (0.1 \text{ T})^2 \left(\frac{P_{\text{input}}}{100 \text{ W}}\right)^{\frac{1}{2}} \left(\frac{Q_{\text{pc}}}{10^5}\right)^{\frac{1}{2}} \left(\frac{B_0}{5 \text{ T}}\right)$$

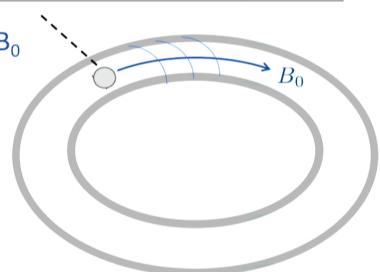
Real advantage of high-Q is on detection side!

Pirsa: 19070001 Page 22/34

# Gapped Toroid Conversion Region

(2) Confine large static B<sub>0</sub>

Generated by wrapped DC current-carrying superconducting wires



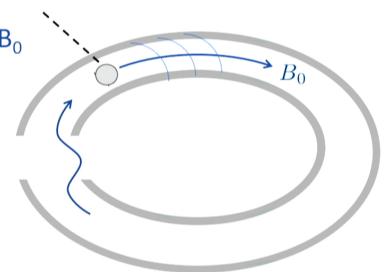
Pirsa: 19070001 Page 23/34

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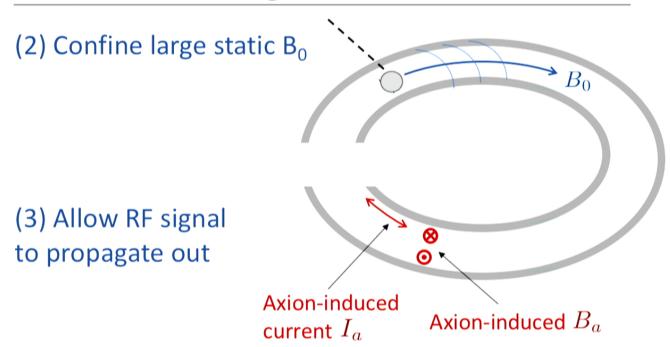
(3) Allow RF signal to propagate out



Any leakage of B<sub>0</sub> outside due to fringe effects, can be made small

Pirsa: 19070001 Page 24/34

## Conversion Region



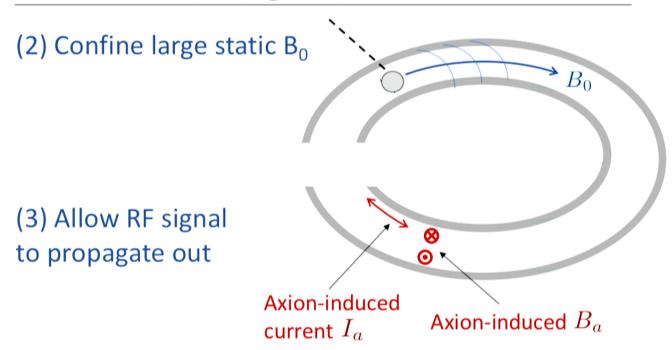
Meissner effect: B<sub>a</sub> is screened outside Sets up super-current I<sub>a</sub> on inner surface

For sufficiently large frequencies:

Ia has significant spatial gradients Signal parametrically suppressed

Pirsa: 19070001 Page 25/34

## Conversion Region



Meissner effect: B<sub>a</sub> is screened outside Sets up super-current I<sub>a</sub> on inner surface

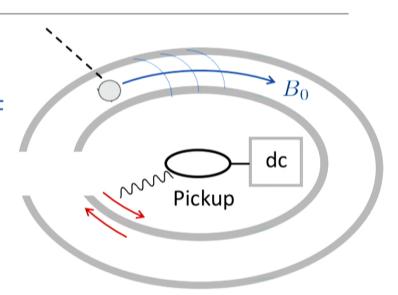
For sufficiently large frequencies:

Ia has significant spatial gradients Signal parametrically suppressed

Pirsa: 19070001 Page 26/34

## Signal Pickup

(4) Couple signal to SRF detection cavity



$$B_a \sim \frac{g^2 B_{\rm pc}^2 B_0}{\omega^2} \sim 10^{-26} \text{ T} \left(\frac{g \text{ GeV}}{10^{-11}}\right)^2 \left(\frac{B_{\rm pc}}{0.2 \text{ T}}\right)^2 \left(\frac{B_0}{5 \text{ T}}\right)$$

Amplify RF signal by Q?

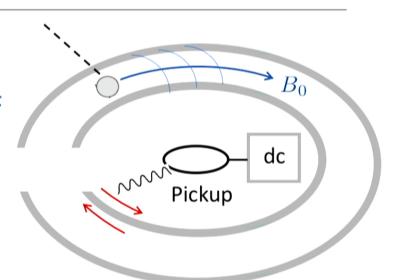
Toroid is not a perfect "source" - non-negligible back-reaction

Must account for toroid impedance as well

Pirsa: 19070001 Page 27/34

# Signal Pickup

(4) Couple signal to SRF detection cavity



Gap capacitance

Toroid inductance

Superconductor

$$C_t \sim 10^{-2} \text{ pF}\left(\frac{d}{\text{mm}}\right) \quad L_t \sim 125 \text{ nH}\left(\frac{R}{10 \text{ cm}}\right) \quad R_t \gtrsim 10^{-9} \Omega$$

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$$R_t \gtrsim 10^{-9} \ \Omega$$

Adjust pickup to impedance match  $\,L_p \sim rac{QR_t}{\omega}\,$ 

$$\Rightarrow P_{\max} \sim |I_a|^2 \frac{(L_t \omega)^2}{R_t}$$

Pirsa: 19070001 Page 29/34

Adjust pickup to impedance match  $\,L_p \sim rac{QR_t}{\omega}\,$ 

$$\Rightarrow P_{\max} \sim |I_a|^2 \frac{(L_t \omega)^2}{R_t}$$

Maintain O(GHz) natural SRF frequency  $\,\omega \sim \omega_0 \sqrt{1+rac{L}{L_p}}\,$ 

$$\Rightarrow P_{\max} \sim |I_a|^2 Q L_t \omega_0$$

Adjust pickup to impedance match  $\,L_p \sim rac{QR_t}{\omega}\,$ 

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#### Optimal signal

$$P_{
m signal} \sim |I_a|^2 (L_t \omega) \cdot \min \left\{ rac{L_t \omega}{R_t}, \; Q 
ight\}$$
Donut Cavity limited limited

#### Narrowband noise

$$P_{
m noise} = rac{T_{
m sys}}{t_{
m int}} ~~ T_{
m sys} \gtrsim \omega \sim 50 \,\, {
m mK}$$
 Quantum limited

Pirsa: 19070001 Page 32/34

## Projected Sensitivity $B_{pc} = 0.2 \text{ T}$ $B_0 = 5 \text{ T}$ $L_t = 125 \text{ nH}$ $T_{\text{sys}} = 0.1 \text{ K}$ $t_{\text{int}} = 1 \text{ year}$ (1) $R_t = 100 \text{ n}\Omega \text{ and } Q \ge 10^{10}$ Axion-photon coupling g (GeV<sup>-1</sup>) $10^{-10}$ CAST, Stellar Cooling SRF (1) **ALPS I** $10^{-11}$ $10^{-8}$ $10^{-7}$ $10^{-5}$ $10^{-6}$ $10^{-4}$ Mass $m_a$ (eV)

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  $T_{\text{sys}} = 0.1 \text{ K}$   $t_{\text{int}} = 1 \text{ year}$ 

(2) 
$$R_t = n\Omega$$
 and  $Q \ge 10^{12}$ 

