

Title: PSI 2018/2019 - Explorations in Quantum Information - Lecture 6

Speakers: Eduardo Martin-Martinez

Collection: PSI 2018/2019 - Explorations in Quantum Information (Martin-Martinez)

Date: April 23, 2019 - 9:00 AM

URL: <http://pirsa.org/19040057>

KMS CONDITION. A state $\hat{\rho}$ is KMS (with inv
 translations generated by a Hamiltonian \hat{H} iff for any P
 $\hat{A}(z) = e^{i\hat{H}z} A(0) e^{-i\hat{H}z}$; $\hat{B}(z) = e^{i\hat{H}z} \hat{B}(0) e^{-i\hat{H}z}$;

i) $\langle \hat{A}(0) \hat{B}(z) \rangle_{\hat{\rho}}$, $\langle \hat{B}(z) \hat{A}(0) \rangle_{\hat{\rho}}$ are boundary values of the complex function

$$ii) \langle \hat{A}(0) \hat{B}(z+i\beta) \rangle_{\hat{\rho}} = \langle \hat{B}(z) \hat{A}(0) \rangle_{\hat{\rho}}$$

(with inverse temperature β) w.r.t. the time parameter z parametrizing any pair of (Heisenberg picture) bounded operators

the complex functions $\langle \hat{A}(0) \hat{B}(z) \rangle_{\hat{\rho}}$ and $\langle \hat{B}(z) \hat{A}(0) \rangle_{\hat{\rho}}$ holomorphic in $\begin{cases} 0 \leq \text{Im} z \leq \beta \\ -\beta \leq \text{Im} z \leq 0 \end{cases}$

today we prove: $+KMS \Rightarrow$ stationarity
 $+KMS \Leftrightarrow$ Gibbs when Gibbs is well-defined

$$\langle \hat{B}(z+i\beta) \rangle_{\hat{A}} = \langle \hat{B}(z) \rangle_{\hat{A}} \Rightarrow \langle \hat{B}(i\beta) \rangle_{\hat{A}} = \langle \hat{B}(0) \rangle_{\hat{A}} \quad (i)$$

$$\| e^{iAx} \hat{B}(iy) e^{-iAx} \| = \| \hat{B}(iy) \|$$

$$\text{trace class } \text{Tr} \left[\begin{pmatrix} \hat{A} + \hat{A} & \hat{1} \\ 0 & 0 \end{pmatrix}^{1/2} \right] < \infty$$

the strip $0 \leq \text{Im } z \leq \beta$
 $\langle \hat{B}(z) \rangle_{\hat{A}}$ is bounded on the strip, therefore it's bounded everywhere in \mathbb{C}
 $\langle \hat{B}(z) \rangle_{\hat{A}}$ has to be constant and therefore $\langle \hat{B}(z) \rangle_{\hat{A}} = \langle \hat{B}(0) \rangle_{\hat{A}}$ $z = z$
 \checkmark

Gibbs \Rightarrow KMS: Let $\hat{\rho}$ be a Gibbs state of inverse β

Gibbs \Rightarrow i) the operator $\hat{A}(z) \hat{B}(z+i\sigma) \hat{\rho} = \underbrace{\hat{A}(z)}_{B_0} \underbrace{e^{-\sigma \hat{H}}}_{TC} \underbrace{\hat{B}(z)}_{B_0} e^{-\beta \hat{H}}$

(the product of a bounded and a trace class operator is trace class. Recall

this means that in $0 < \text{Im } z < \beta$ $\langle \hat{A}(z) \hat{B}(z) \rangle_{\hat{\rho}} = \text{Tr} [\hat{A}(z) \hat{B}(z) \hat{\rho}]$

Gibbs state of inverse temp β $\hat{\rho} = \frac{1}{Z(\beta)} e^{-\beta \hat{H}}$

$$(z+i\sigma) \hat{\rho} = \underbrace{\hat{A}(0)}_{B_0} \underbrace{e^{-\sigma \hat{H}}}_{TC} \underbrace{\hat{B}(z)}_{B_0} \underbrace{e^{-(\beta-\sigma) \hat{H}}}_{TC} / Z(\beta) \quad \text{is trace-class for } 0 \leq \sigma \leq \beta$$

class operator is trace class. Reed and Simon section VI: Bounded op. Exercises 27

$$\langle \hat{A}(0) \hat{B}(z) \rangle_{\beta} = \text{Tr} [\hat{A}(0) \hat{B}(z) \hat{\rho}] = f(z) \text{ is a well-defined function of } z$$

$$(z = z+i\sigma) \partial_z f = -i \partial_{\sigma} f \Rightarrow \langle \hat{A}(0) \hat{B}(z) \rangle_{\beta} \text{ is holomorphic in the strip } \checkmark$$

$$\hat{\rho} = \frac{1}{Z(\beta)} e^{-\beta \hat{H}}$$

$\frac{e^{-\sigma \hat{H}}}{Z(\beta)}$ is trace-class for $0 \leq \sigma \leq \beta$

and ζ -mch section VI: Banded op. Exercises 27 and 29)

ζ a well-defined function of \bar{z}

$\langle \hat{\rho}(\bar{z}) \rangle_{\beta}$ is holomorphic in tk strip \checkmark

Gibbs \Rightarrow ii) $\hat{B}(z) = e^{\hat{H}z} \hat{B}(0) e^{-\hat{H}z}$ extending z to the complex plane

$$\langle \hat{A}(0) \hat{B}(z+i\beta) \rangle_{\hat{\rho}} = \text{tr} \left[\hat{A}(0) e^{-\beta \hat{H}} \hat{B}(z) e^{\beta \hat{H}} e^{-z \hat{H}} \right] / Z(\beta) = \text{Tr} \left[\hat{B}(z) \hat{A}(0) e^{-\beta \hat{H}} \right]$$

KMS \Rightarrow Gibbs: let's assume that $\hat{\rho}$ is KMS, ii) at $z=0 = i\beta$

then $[\hat{B}(0), e^{\beta \hat{H}} \hat{\rho}] = 0$ for any bounded \hat{B} so $e^{\beta \hat{H}} \hat{\rho} \propto \mathbb{1}$

to the complex plane $\hat{B}(z+i\beta) = e^{iH(z+i\beta)} \hat{B}(0) e^{-iH(z+i\beta)} =$

$$= \text{Tr} \left[\hat{B}(z) \hat{A}(0) e^{-\beta H} \right] / Z(\beta) = \text{Tr} \left[\hat{B}(z) \hat{A}(0) \hat{\rho} \right] = \langle \hat{B}(z) \hat{A}(0) \rangle_{\hat{\rho}} \quad \checkmark$$

i) at $z=0 \Rightarrow \text{Tr}(\hat{A}(0) e^{-\beta H} [\hat{B}(0), e^{\beta H} \rho]) = 0$ for any bounded op \hat{A}

$$e^{\beta H} \hat{\rho} \propto \mathbb{1} \Rightarrow \hat{\rho} = \frac{1}{Z(\beta)} e^{-\beta H} \quad \checkmark$$

$$W_{\hat{\rho}}(z, z') := \langle \hat{A}(z) \hat{A}(z') \rangle_{\hat{\rho}} \quad \text{then}$$

- KMS \Rightarrow
- ① Stationarity: $W_{\hat{\rho}}(z, z') = W_{\hat{\rho}}(0, \Delta z) \equiv W_{\hat{\rho}}(\Delta z)$
 - ② Holomorphicity of $W_{\hat{\rho}}(z)$ in the upper strip $0 \leq \text{Im} z \leq \beta$
 - ③ Complex anti-periodicity (of period β) $W_{\hat{\rho}}(\Delta z + i\beta) = -W_{\hat{\rho}}(\Delta z)$

in a QFT a state $\hat{\rho}$ is KMS (w.r.t. z of inv. temp β)

$$W_{\hat{\rho}}(\vec{x}(z), \vec{x}(z')) := \langle \hat{\phi}(t(z), \vec{x}(z)) \hat{\phi}(t(z'), \vec{x}(z')) \rangle_{\hat{\rho}} \quad \text{satisfies ①, ②, ③}$$

then

$(0, \Delta z) \equiv W_p(\Delta z)$ with $\Delta z = z' - z$
 $(2\Delta, \Delta z) \equiv W_p(\Delta z)$
 upper strip $0 \leq \text{Im } z \leq \beta$
 $W_p(\Delta z + i\beta) = W_p(-\Delta z)$

$\int_{\gamma} f(z) dz$ satisfies ①, ②, ③
 long β

$\Sigma(z)$ is the worldline generated by ∂_z