

Title: PSI 2018/2019 - Explorations in Quantum Gravity - Lecture 9

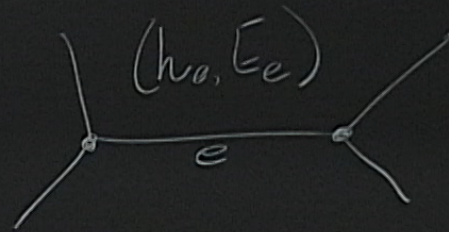
Speakers: MaÃtÃ© Dupuis

Collection: PSI 2018/2019 - Explorations in Quantum Gravity (Dupuis)

Date: March 14, 2019 - 10:15 AM

URL: <http://pirsa.org/19030076>

$(A(x), E(x)) \rightarrow$



Single Edge:

(h_{mn}, E^j)

$$\{h_{mn}, h_{m'n'}\} = 0$$

$$\{h_{mn}, E^j\} = (hT^j)_{mn}$$

$$\{E^j, E^k\} = \epsilon_{jk} E^l$$

$h_{mn} \rightarrow \hat{h}_{mn} \quad \hat{h}_{mn} |\psi\rangle$

$$[\hat{f}, \hat{g}] = i\hbar \{f, g\} \Rightarrow$$

Check $[\hat{E}^j, \hat{E}^k]$ ✓

Inner Product:

$$\langle \psi_1 | \psi_2 \rangle = \int_{SU(2)} dg \overline{\psi_1(g)} \psi_2(g)$$

$$U_{hmn} |\psi\rangle = |hmn \psi\rangle$$

$$\{\hat{f}, \hat{g}\} \Rightarrow E^j |\psi\rangle = i\hbar |L^j \psi\rangle$$

$$[\hat{E}^j, \hat{E}^k] \quad \checkmark$$

product:

$$\overline{\psi_1(g)} \psi_2(g)$$

O.N. Basis:

Peter-Weyl $\rightarrow f(g) = \sum_j d_j \sum_{mn} D_{mn}^j(g)$

$$1_{\text{group}} = \sum_{j \in \mathbb{N}/2} \sum_{m=-j}^j \sum_{n=-j}^j |jmn\rangle \langle jmn|$$

$$|jmn\rangle = \sqrt{d_j} |D_{mn}^j\rangle$$

$h |L^2 \psi\rangle$

or complete basis.

Basis:

$$f(g) = \sum_j d_j \sum_{mn} D_{mn}^j(g)$$

$$\sum_{j=-j}^j \sum_{k=-j}^j |jmn\rangle \langle jmn|$$

$$|jmn\rangle = \sqrt{d_j} |D_{mn}^j\rangle$$

TODAY: Kinematical Hilbert space.

↳ subsystem on edge \longrightarrow a graph.

↓
length op.

$$\hat{E}^0 = i\hbar L^0 = i\hbar \left. \frac{d}{dt} \right|_{t=0} R_{e^{+T}}$$

$$(R_h f)(g) = f(g^h)$$

$$|\psi\rangle = |l m \psi\rangle$$

$$\Rightarrow E |\psi\rangle = i\hbar |L' \psi\rangle$$

or complete basis.

O.N. Basis:

→ Wigner matrices

$$\text{Peter-Weyl} \rightarrow f(g) = \sum d_j \hat{f}_{mn}^j D_{mn}^j(g)$$

$$\mathbb{1}_{\text{series}} = \sum_{j \in \mathbb{N}/2} \sum_{m=-j}^j \sum_{n=-j}^j |jmn\rangle \langle jmn|$$

$$|jmn\rangle = \sqrt{d_j} |D_{mn}^j\rangle$$

$$|l, m\rangle \psi\rangle$$

$$|\psi\rangle = \int dk |L, \psi\rangle$$

$$\{|j, m, n\rangle\}$$

over complete basis

O.N. Basis

Wigner matrices

Peter-Weyl $\rightarrow f(g) = \sum_j d_j \sum_{mn} D_{mn}^j(g)$

$$f(g) = \sum_j \sum_{m,n} |j, mn\rangle \langle j, mn|$$

$$|j, mn\rangle = \sqrt{d_j} |D_{mn}^j\rangle$$

Space
edge \rightarrow a graph.

$$(R_h f)(g) = f(g^h)$$

$$(R_h D_{mn}^d)(g) = \sum_P \underbrace{D_{mp}^d(g) D_{pn}^d(g)}_{\langle g | dmp \rangle}$$

$$\Rightarrow R_h |dmn\rangle = \sum_P D_{pn}^d(g) |dmp\rangle$$

$\{|dmp\rangle, p \in \{-d, \dots, d\}\}$: subrepresentation space carries a rep D^d

$$R_h D_{mn}^d(g) = D_{mn}^d(h^{-1}g)$$

• Left multiplication $L_h D_{mn}^d(g) = D_{mn}^d(h^{-1}g) = \sum_q D_{mq}^d(h^{-1}) D_{qn}^d(g)$

$$\mathcal{H}_e = L^2(SU(2), dg) = \bigoplus_d V_d^* \otimes V_d$$

$$\begin{array}{ccc} \bullet & \xrightarrow{\quad} & \bullet \\ & \text{D}_{mn}^d(g) & \\ & \text{I}_{jmn} & \\ & \text{I}_{j} & \end{array}$$

• Left multiplication $L_h D_{mn}^d(g) = D_{mn}^d(h^{-1}g) = \sum_q D_{mq}^d(h^{-1}) D_{qn}^d(g)$

$$\mathcal{H}_e = L^2(SU(2), dg) = \bigoplus_d V_d^* \otimes V_d$$

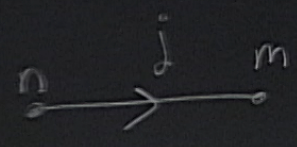
$$\begin{array}{ccc} \bullet & \xrightarrow{\quad} & \bullet \\ & \text{with } |jmn\rangle & \\ & & \text{with } j \end{array} \equiv \begin{array}{ccc} \bullet & \xrightarrow{\quad} & \bullet \\ & & \text{with } j \end{array}$$

$$= \sum_q D_{mq}^d(h^{-1}) D_{qn}^j(g) \Rightarrow \langle h | j^m n \rangle = \sum_q D_{mq}^d(h^{-1}) | j q n \rangle$$

$\{ | j q n \rangle \}_{q=-j, j}$ subrepresentation space carries a rep. $(D^j)^*$: dual contragredient rep.

$$(D^d)^*_{mn}(g) = \frac{D_{nm}^d(g^{-1})}{D_{mn}^d(g)}$$

$$= \bigoplus_j V_j^* \otimes V_j$$

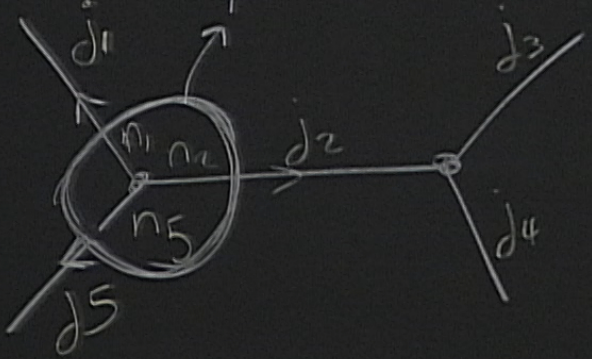


(angular momentum associated to both ends of the edge)

• Left multiplication

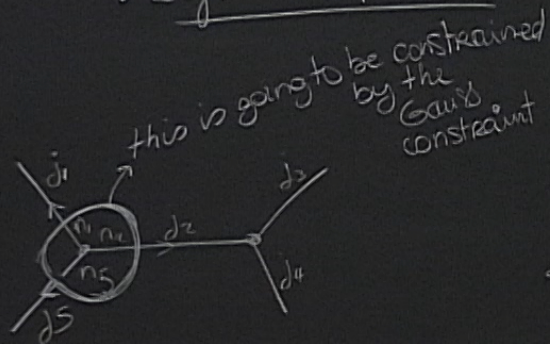
$$L_h D_{mn}^d(g) = D_{mn}^d$$

this is going to be constrained by the Gauss constraint



$$\mathcal{H}_e = L^2(SU(2))$$

• Left multiplication $L_h D_{mn}^d(g) = D_{mn}^d(h^{-1}g) = \sum_q D_{mq}^d(h^{-1}) D_{qn}^d(g) \Rightarrow L_h |jmn\rangle =$
 $|jqn\rangle$ $q=j-d$ subrup



this is going to be constrained by the Gauss constraint

$$\mathcal{H}_e = L^2(SU(2), dg) = \bigoplus_0 V_d^* \otimes V_d$$

$(D_{mn}^d(g))$
 $|jmn\rangle$

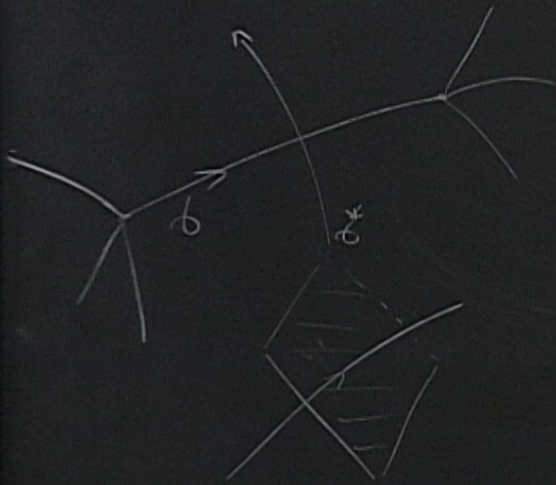
$n \rightarrow j \rightarrow m$

(angular momentum associated to both ends of the edge)

Length sp.



Length ep :



ex:

length of a curve γ^* on a 2D surface.

$$\int_{\gamma^*} \sqrt{\dot{\gamma}^a \dot{\gamma}^b g_{ab}} ds$$

ex:

length of a curve γ^* on a 2D surface.

$$\int_{\gamma^*} \sqrt{\dot{\gamma}^a \dot{\gamma}^b g_{ab}} ds$$

$\hookrightarrow g_{ab} = e_a^i e_b^j S_{ij}$

ex :

length of a curve γ^* on a 2D surface.

$$\int_{\gamma^*} \sqrt{\dot{\gamma}^{\mu a} \dot{\gamma}^{\nu b} g_{ab}} d\sigma$$

$\hookrightarrow g_{ab} = e_a^i e_b^j S_{ij}$

$$\hookrightarrow \sim \sum_{n=0}^2 \sqrt{\sum_k E_a^k E_c^k}$$

\hookrightarrow Riemann approx.

$$l_{ex}^2 := \sum_k \hat{E}_e^k \hat{E}_e^k$$

$$\begin{aligned}
 & \int_{\Sigma} \sqrt{\gamma^{xa} \gamma^{xb}} q_{ab} \, d\sigma \\
 & \quad \rightarrow q_{ab} = e_a^i e_b^j S_{ij} \\
 & \quad \rightarrow \sim \sum_{n=0}^{\infty} \sqrt{\sum_k E_{a_n}^k E_{c_n}^k} \\
 & \quad \rightarrow \text{Riemann approx.}
 \end{aligned}$$

'heuristic justification'

length of a curve γ^* on a 2D surface.

$$\int_{\gamma^*} \sqrt{\dot{\gamma}^{\mu a} \dot{\gamma}^{\mu b} g_{ab}} ds$$

$$\hookrightarrow g_{ab} = e_a^i e_b^j \delta_{ij}$$

$$\hookrightarrow \sim \sum_{n=0}^N \sqrt{\sum_k E_a^k E_{cn}^k}$$

\hookrightarrow Riemann approx

heuristic justification

$$\hat{l}_{ex}^2 := \sum_k \hat{E}_e^k \hat{E}_e^k$$

$$E^k |jmn\rangle = i\hbar \sum_p |jmp\rangle \frac{d}{dt} \Big|_{t=0} D_{pn}^j(e^{tT^k}) = i\hbar \sum_p |jmp\rangle D_{pn}^j(T^k)$$

"right multiplication"

"multiplication"

$$\hat{L}^2 |jmn\rangle = -\hbar^2 \sum_{p,q,k} D_{qp}^d(T^k) D_{pn}^d(T^k) |jmq\rangle$$
$$\hookrightarrow \sum_{p,k} D_{qp}^d(T^k) D_{pn}^d(T^k) = C_{qn}^d = D_{qn}^d(\mathbb{1}^2)$$

$$D_{pn}^d(T^k) |j m q\rangle$$

$$\sum_{p,k} D_{qp}^d(T^k) D_{pn}^d(T^k) = C_{qn}^d = D_{qn}^d(|\vec{T}|^2)$$

$$[C^d, D^d(T^k)] = 0 \quad \text{exercise: check this!}$$

$$D_{pn}^d(T^k) |j m q\rangle$$

$$\sum_{p,k} \overbrace{D_{qp}^d(T^k)} D_{pn}^d(T^k) = C_{qn}^d = D_{qn}^d(|\vec{T}|^2)$$

$$[C^d, D^d(T^k)] = 0 \quad \text{exercise: check this!}$$

- C^d is a Casimir op
- Schur's lemma: if acts on an irreducible representation vector space $V^d = \{|j m n\rangle, m \text{ fixed}, n = -d \dots d\}$

as the identity op: $C^d = \lambda \text{id}_{V^d}$

$$\sum_{j \in \mathbb{N}/2} \sum_{m=-j}^j \sum_{n=-j}^j |j m n\rangle \langle j m n|$$

$$|j m n\rangle = \sqrt{d_j} |D_{j m n}\rangle$$

- C^d is a Casimir op
- Schur's lemma: if acts on representation vector space

as the identity op:
 $C^d = \lambda \text{id}_{V_d}$

C^d : $n(d)$ Casimir (modulo a cst)

$$\lambda = -j(j+1)$$

$$l_x^2 |jmn\rangle = -\hbar^2 \sum_q |jmq\rangle C_{qn}^d = -\lambda \hbar^2 \sum_q |jmq\rangle \delta_{qn} = +\hbar^2 j(j+1) |jmn\rangle$$

l_x^2 - eigenvectors $\{|jmn\rangle\}$

- spectrum $\{l_p^2 j(j+1), j \in \mathbb{N}/2\}$ is discrete!

$$l_p = \hbar k$$

- C^d is a Casimir op
- Schur's lemma: if acts on representation vector space

as the identity op:
 $C^d = \lambda \text{id}_V$

C^d is a Casimir (modulo a cst)

$$\lambda = -j(j+1)$$

$$L^2 |jmn\rangle = -\hbar^2 \sum_q |jmq\rangle C_{qn}^d = -\lambda \hbar^2 \sum_q |jmq\rangle \delta_{qn} = +\hbar^2 j(j+1) |jmn\rangle$$

L^2 eigenvectors $\{|jmn\rangle\}$

spectrum

$\{ \hbar^2 j(j+1), j \in \mathbb{N}/2 \}$ is discrete!

$$L^3 = \hbar^2 K_{\mathfrak{so}(G,C)}$$

d_5

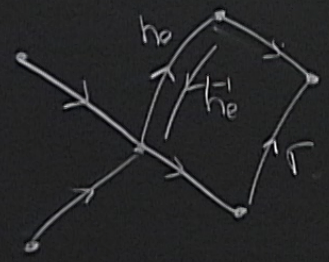
d_4

$(j m n)$

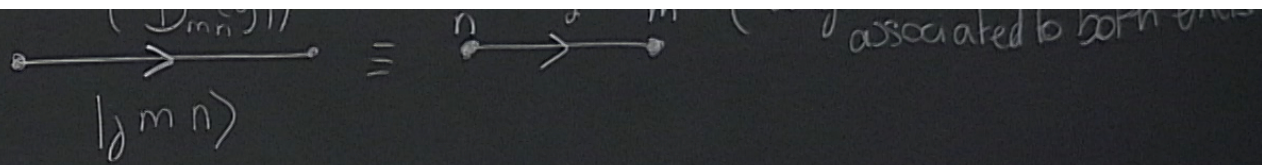
$$\mathcal{H}_2 = L^2(SU(2), dg)$$

of edges of Γ

$$\mathcal{H}_\Gamma = L^2(SU(2)^\Gamma)$$



$\int_\Gamma (h_e)$



$$\mathcal{H}_e = L^2(\text{SU}(2), dg)$$

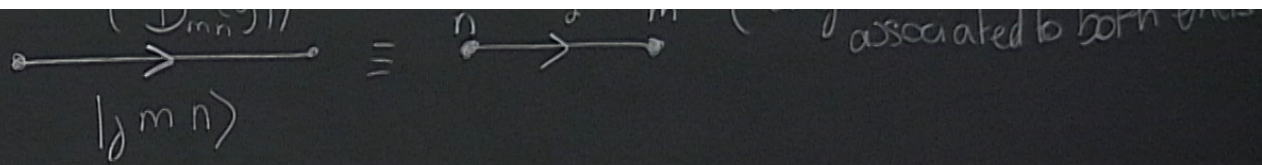
of edges of Γ

$$\mathcal{H}_\Gamma = L^2(\text{SU}(2)^{\overset{\# \text{ edges of } \Gamma}{E}}) = \bigotimes_e \mathcal{H}_e = \bigotimes_e L^2(\text{SU}(2))$$

Inner product

$$\langle \Psi_{\Gamma,1} | \Psi_{\Gamma,2} \rangle = \int_{\text{SU}(2)^E} \overline{\Psi_{\Gamma,1}(g_{e_1} \dots g_{e_E})} \Psi_{\Gamma,2}(g_{e_1} \dots g_{e_E}) dg_{e_1} \dots dg_{e_E}$$

Orthonormal basis $|\vec{j}, \vec{m}, \vec{n}\rangle_\Gamma = |j_1, m_1, n_1\rangle \otimes \dots \otimes |j_E, m_E, n_E\rangle$



$$\mathcal{H}_e = L^2(SU(2), dg)$$

of edges of Γ

$$\mathcal{H}_\Gamma = L^2(SU(2)^{\overset{\# \text{ edges of } \Gamma}{E}}) = \bigotimes_e \mathcal{H}_e = \bigotimes_e L^2(SU(2))$$

Inner product

$$\langle \Psi_{\Gamma,1} | \Psi_{\Gamma,2} \rangle = \int_{SU(2)^E} \overline{\Psi_{\Gamma,1}(g_{e_1} \dots g_{e_E})} \Psi_{\Gamma,2}(g_{e_1} \dots g_{e_E}) dg_{e_1} \dots dg_{e_E}$$

Orthonormal basis $|\vec{j}, \vec{m}, \vec{n}\rangle_\Gamma = |j_{e_1}, m_{e_1}, n_{e_1}\rangle \otimes \dots \otimes |j_{e_E}, m_{e_E}, n_{e_E}\rangle$

$(A, e) \rightarrow \text{edge}$
 (h_e, E_e)

\downarrow
 \mathcal{H}_e

\downarrow
 \mathcal{H} projective limit

\downarrow
 \mathcal{H} kinematical

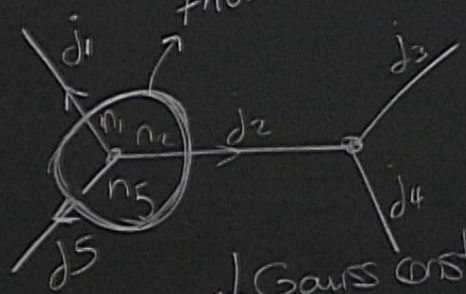
$(g_a \dots g_{e_n}) dg_a \dots dg_{e_n}$

$(g_a \dots g_{e_n})$

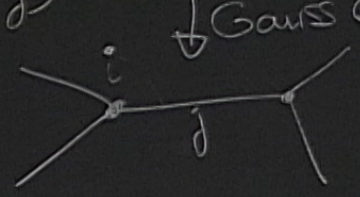
• Left multiplication

$$L_h D_{mn}^d(g) = D_{mn}^d(h)$$

this is going to be constrained by the Gauss constraint



↓ Gauss constraints



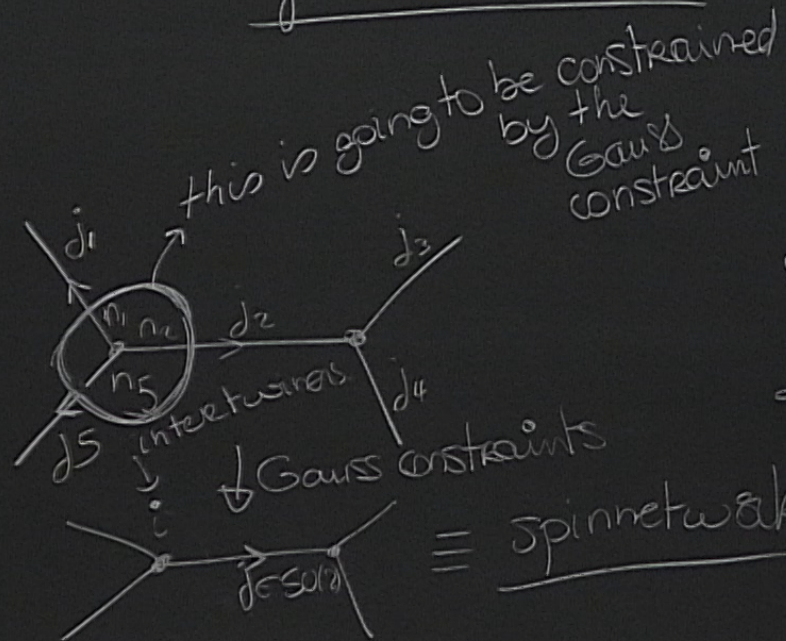
≡ spin network

$$\mathcal{H}_e = L^2(SU(2), \rho)$$

(D_m^d)
→
l, m, n

• Left multiplication

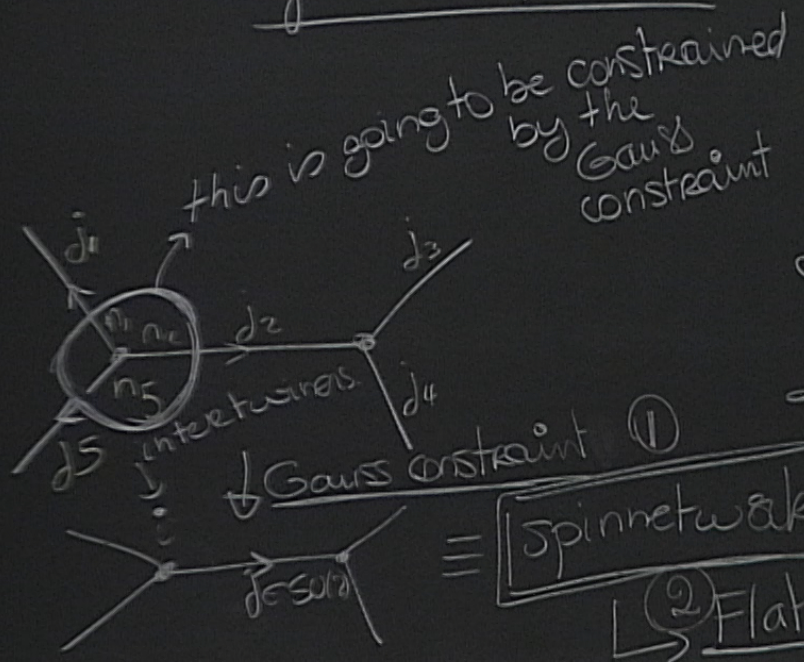
$$L_h D_{mn}^d(g) = D_{mn}^d(h^{-1}g)$$



$$\mathcal{H}_e = L^2(SU(2), \rho)$$

(D_{mn}^d)
 l, m, n

• Left multiplication $L_h D_{mn}^d(g) = D_{mn}^d(h^{-1}g) = \sum_q D_{mq}^d(h)$



$$\mathcal{H}_e = L^2(SU(2), dg) = \bigoplus_d V_d^* \otimes V_d$$

$(D_{mn}^d(g))$

$|j m n\rangle$

\rightarrow Physical Hilbert space