

Title: PSI 2018/2019 - Strong Field Gravity - Lecture 9

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Collection: PSI 2018/2019 - Strong Field Gravity (East)

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URL: <http://pirsa.org/19030016>

Generalized Harmonic Formulation

$$\square x^a = H^a \quad (\text{source functions})$$

Promote H_a to independent variables

$$R_{ab} = 4\pi (2T_{ab} - g_{ab}T) \quad \text{EFEs trace-reversed}$$

$$R_{ab} = -\frac{1}{2} g^{cd} \partial_c \partial_d g_{ab} + \nabla_{(a} \Gamma_{b)} + F(g_{ab}, \partial_c g_{ab})$$

$$\Gamma_c = g^{ab} \Gamma_{abc} = g^{cd} \nabla_a \nabla^a x^d$$

$$-8\pi (2T_{as} - g_{ab}T) = \underbrace{g_{cd} \partial_c \partial_d g_{ab}}_{\substack{\text{principal} \\ \text{part}}} + \dots$$

$$2_a H_b + 2_b H_a + 2 H_d$$

$$T) = \underbrace{g^{cd} \partial_c \partial_d g_{ab}} + \partial_b g^{cd} \partial_c g_{ad} + \partial_a g^{cd} \partial_c g_{bd} +$$

principal part

$$\partial_a H_b + \partial_b H_a + 2 H_d M_{ab}^{cd} + 2 M_{db}^c M_{ca}^d$$

← some function

$$L(H_a) = 0 (**)$$

$$H_a = F_a(g_{ab})$$

$$H_a = 0$$

Har

$$+ \partial_b g^{cd} \partial_c g_{ad} + \partial_a g^{cd} \partial_c g_{bd} +$$

(*)

$$= \sum H_d \Gamma_{ab}^d + \sum \Gamma_{db}^c \Gamma_{ca}^d$$

← Some function

$$H_a = F_a(g_{ab})$$

(Simple choice)

$$H_a = 0$$

Harmonic gauge

Evolve $(g_{ab}, \partial_t g_{ab}, H_a, \partial_t H_a)$ according to $(*) + (**)$

$$C^a = \square x^a - H^a$$

Guaranteed to have $C^a = 0$ for all time?

From $*$) $R_{ab} - 4\pi(2T_{ab} - g_{ab}T) = \nabla_{(a} C_{b)}$

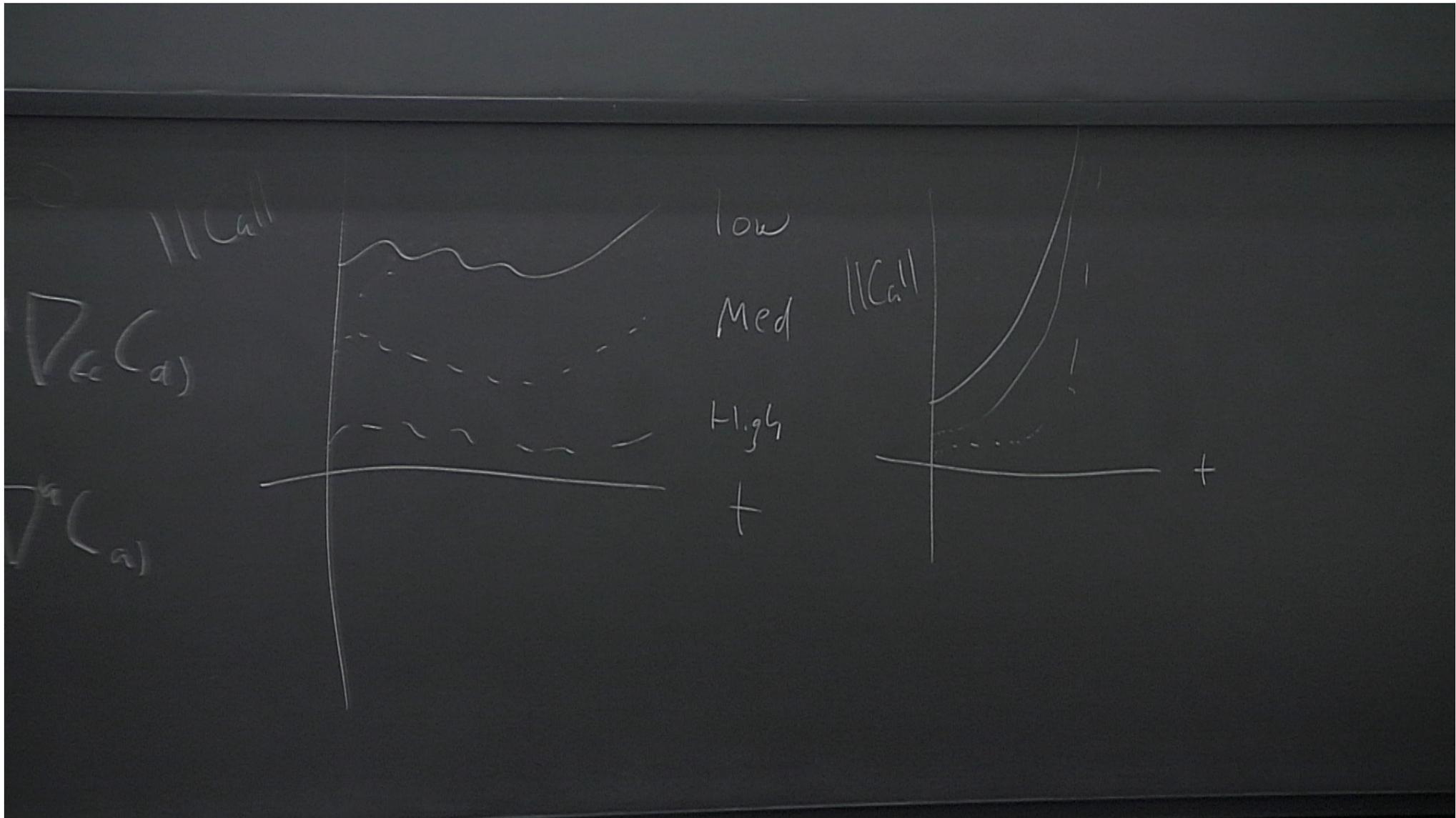
$n^a n^b \rightarrow \text{Ham}$ $n^a \gamma^d_b \rightarrow \text{Mom Cnst.}$

If $C_a = 0$ at $t=0$, $\partial_t C_a = 0$ at $t=0$

$$(R_{ab} - \frac{1}{2}Rg_{ab}) - 3\pi T_{ab} = \nabla_{(a} C_{b)} - \frac{1}{2}g_{ab}g^{cd}\nabla_{(c} C_{d)}$$

$$0 = \nabla^a \nabla_{(a} C_{b)} - \frac{1}{2}\nabla_b \nabla^a C_a$$

$$\nabla_a \nabla^a C_b = -R_b^a C_a$$



Add to (*) $K (R_a(b) - \frac{1}{2} g_{ab} n^d{}_{;d})$

$$\square C^a = -R_b^a C^b + K \nabla_b [n^b C^a]$$

$$K \sim \frac{1}{L}$$

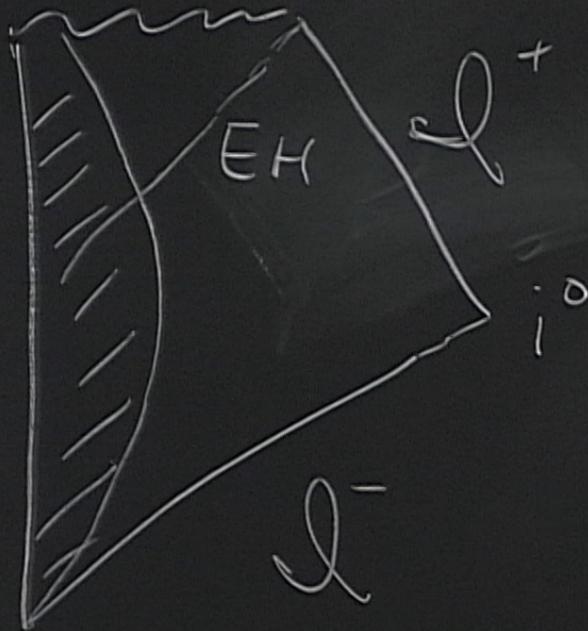
Event horizon - the boundary of the causal past

$$\text{BH is } \mathcal{B} = M - \mathcal{J}^-(\mathcal{I}^+)$$

$$\text{EH is } \partial\mathcal{B}$$

ary of the causal past of future null infinity

$\mathcal{J}^-(\mathcal{I}^+)$



Prob

- R

- IS

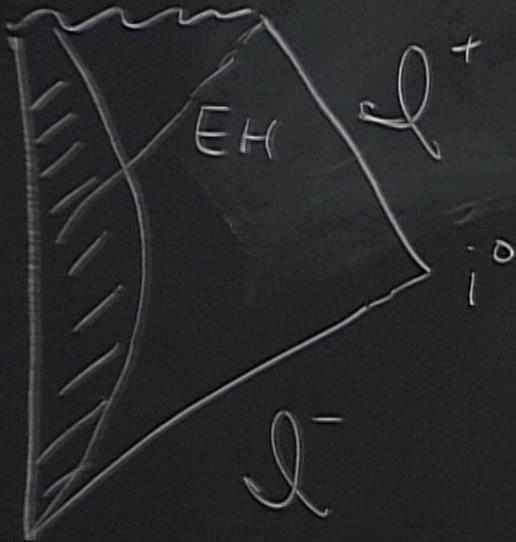
Problems with Event horizons

- Requires knowledge of whole spacetime out to \mathcal{I}^+
- Is teleological, i.e. "anticipates" future

Problems with Ev

- Requires knowledge

- Is teleological



Introduce Apparent Horizon

Timelike

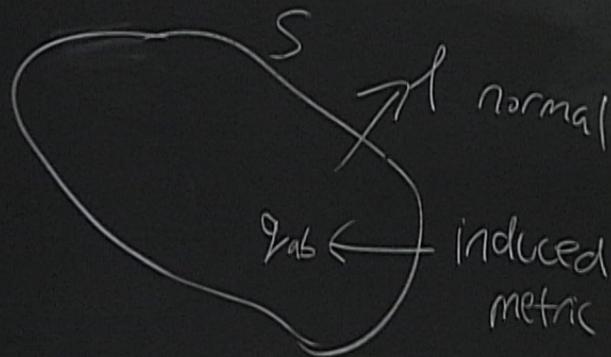
Σ_t

S is

closed 2-dim

spacelike surface

$\subset \Sigma_t$



$$\mathcal{O} = \mathcal{L}_\xi(\ln(\sqrt{g})) = g^{ab} \nabla_a b_s$$

Take l to be null

$$\begin{array}{l} \text{Outward} \\ \text{Inward} \end{array} / l_{\pm}^a = n^a \pm s^a$$

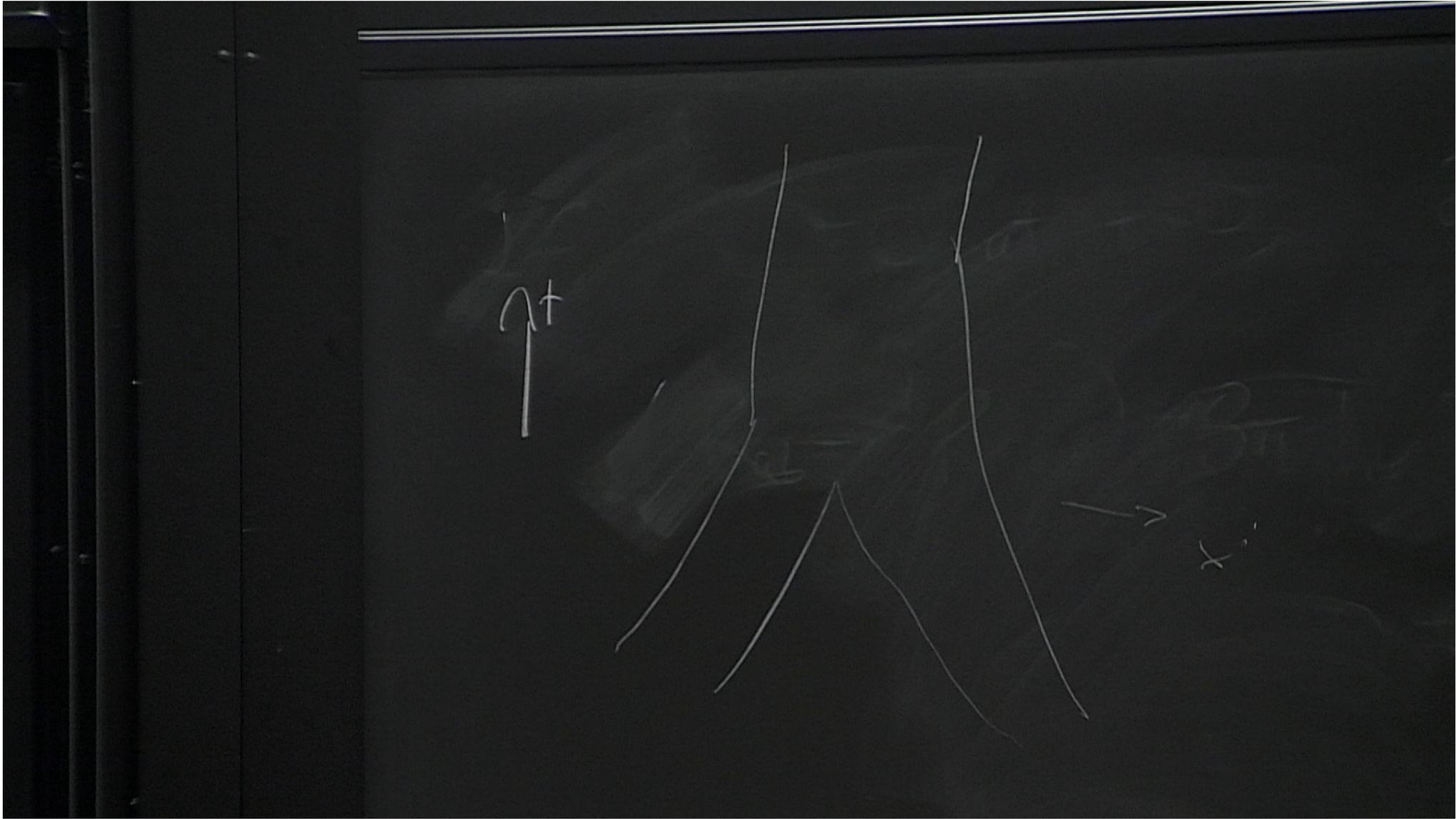
$$n^a n_a = -1, \quad s^a s_a = 1, \quad l_{\pm}^a l_{\pm}^{\pm} = 0, \quad l_a l_{\pm}^{\pm} = -2$$

Outward null expansion

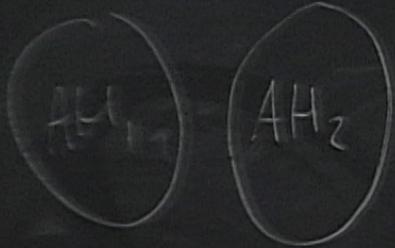
$$\Theta_{l^+} = q^{ab} \nabla_{ab} l^+ = D_i s^i + K_{ij} s^i s^j - K \text{ in } (3+1) \text{ vars}$$

If $\Theta_{l^+} \leq 0$ everywhere on S , S is outer trapped surface
 $\Theta_{l^+} = 0$ " " marginally " (MOTS)

Apparent Horizon Outermost MOTS



$$t = t_{CH} - \Delta t$$



$$t = t_{CH}$$

AH1

