

Title: PSI 2018/2019 - Strong Field Gravity - Lecture 5

Speakers: William East

Collection: PSI 2018/2019 - Strong Field Gravity (East)

Date: March 08, 2019 - 9:00 AM

URL: <http://pirsa.org/19030012>

Generalized Penrose processes

$$0 \leq \Delta E_{\text{BU}} - \int_{\text{H}} \Delta T_{\text{BU}}$$

Superradiance

Scalar field:

Massless Klein-Gordon Eqn: $\square \psi = \nabla_a \nabla^a \psi = 0$

Stress-energy tensor:

$$T_{ab} = \nabla_a \psi \nabla_b \psi - \frac{1}{2} g_{ab} \nabla_c \psi \nabla^c \psi$$

$$\psi = \bar{\Psi}(r, \theta) \operatorname{Re} \left[e^{\underset{\substack{\uparrow \\ \text{frequency}}}{-i(\omega t - m\phi)}} \right]$$

azimuthal number

$$\nabla^a \psi = 0$$

$$\Delta E_{\text{BH}} = \int dt \int_{r=r_+} \chi^a T_{as} \hat{t}^b dS_{\text{BH}}$$

$$\chi^a = \hat{t}^a + \Omega_H \hat{\phi}^a$$

$$\chi^a \hat{t}^a = \chi^a (\chi^a - \Omega_H \hat{\phi}^a)$$

$$= \int dt \int_{r=r_+} \left[\chi^a (\nabla_a \psi) (\nabla_b \psi) \hat{t}^b - \frac{1}{2} \chi_a \hat{t}^a \nabla_c \psi \nabla^c \psi \right] dS_{\text{BH}}$$

$$= \int dt \int_{r=r_+} dS_{\text{BH}} (\partial_t \psi + \Omega_H \partial_\phi \psi) \partial_t \psi$$

$$= \int dt \int_{r=r_+} dS_{\text{BH}} (\omega - m \Omega_H) \omega \underbrace{|\psi|^2 \sin^2(\omega t - m\phi)}_{\geq 0}$$

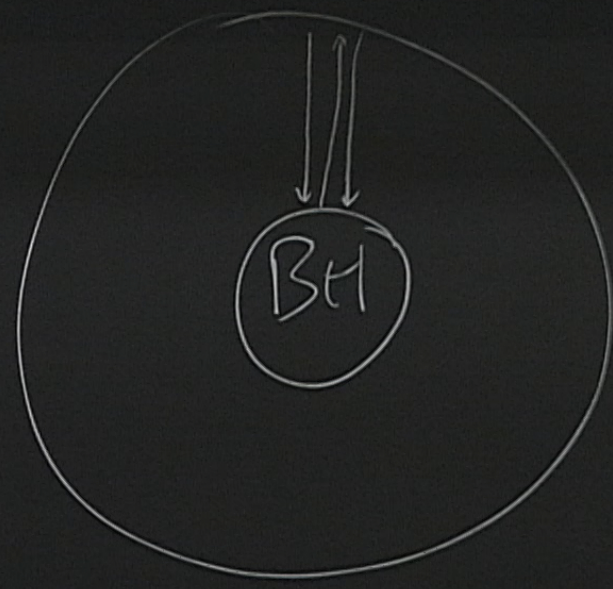
$$0 < \omega < m\Omega_H \Rightarrow \Delta E_{BH} < 0$$

$$\frac{\delta E}{\delta J} = \frac{\omega}{M}$$

Maximum Amplification factor	$S=0$	scalar	0.3%
	$S=1$	E-M	4.4%
	$S=2$	GWs	138%

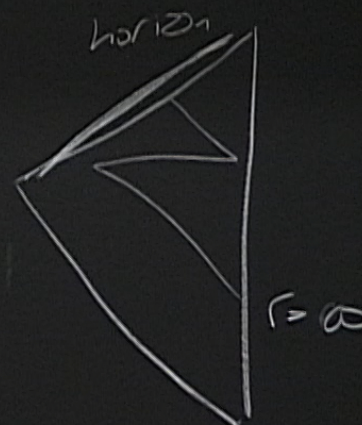
3%
4%
3%

mirror



Press & Teukolsky
BH bomb

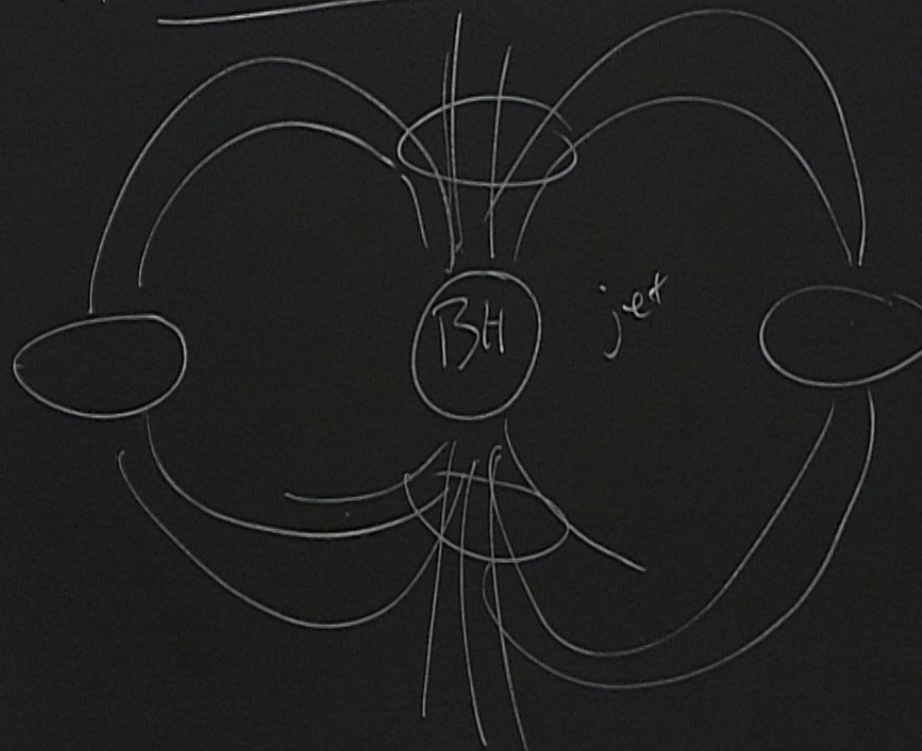
Press & Teukolsky
BH bomb



Generalized Ergoregion for charged BHs

$$E_k = -\hat{t}^a P_a = -m \hat{t}^a U_a - q \hat{t}^a A_a$$

Blandford Znajek



Massive

K-G Egn.

$$\mu' \gg \Lambda$$

$$\square \phi = \nabla$$

Massive scalar field

K-G Eqn. $\square\phi = \mu^2\phi$

$\mu \gg M$, $r \gg M$, non-relativistic limit

$$\square\phi = \nabla_a (g^{ab} \partial_b \phi) = \frac{1}{\sqrt{-g}} \partial_a (\sqrt{-g} g^{ab} \partial_b \phi)$$

For Schw $\sqrt{-g} = r \sin\theta$

$$g^{tt} \partial_t^2 \phi + \partial_r \partial_r \phi - \mu^2 \phi = 0$$

$$g^{tt} \approx \left(1 + \frac{2M}{r}\right)$$

Use ansatz: $\phi = \frac{1}{\sqrt{2\mu}} \left[\psi(x) e^{-i\omega t} + \psi(x) e^{i\omega t} \right]$

$$\left[-\left(1 + \frac{2M}{r}\right) (-\omega)^2 + \partial_r \partial_r - \mu^2 \right] \psi = 0$$

$$(\omega^2 - \mu^2) \psi = \left(-\partial_r \partial_r - \frac{2M\omega^2}{r} \right) \psi$$

$$(H - \mu) \psi = \left(-\frac{1}{2\mu} \nabla^2 - \frac{\mu M}{r} \right) \psi$$

$$\alpha = \mu M$$

$$\text{Energy} = H - \mu$$

$$\text{Bohr radius} = \mu^{-1} \alpha$$

$$(n, l, m)$$

$$l = n = 1$$

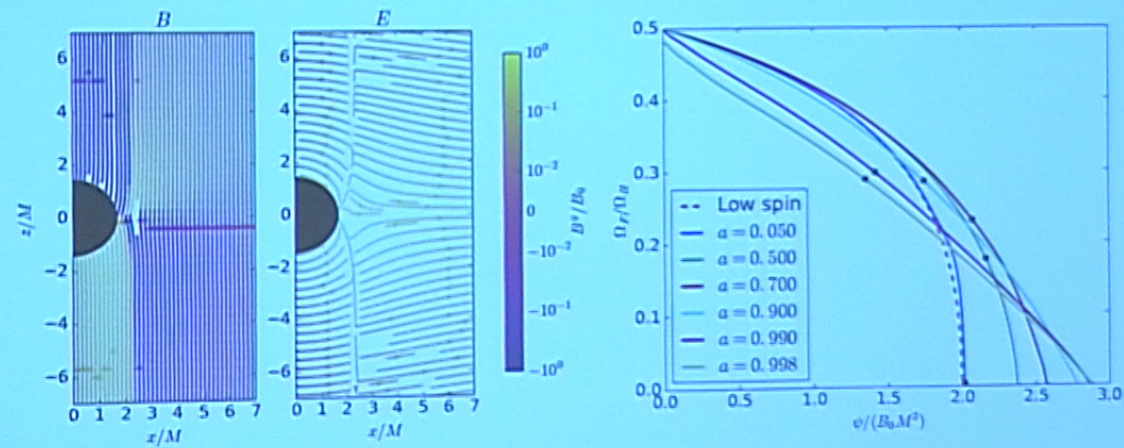
$$\phi \sim e^{-r/2r_0} Y_l^m(\vartheta, \phi) e^{-i(\omega t)} \quad \nwarrow_{\text{complex}}$$

$$\text{Re}(\omega) = \mu \left(1 - \frac{\alpha^2}{2n^2} \right)$$

$$|\text{Im}(\omega)| \ll |\text{Re}(\omega)|$$

$$\alpha = 0.1 \left(\frac{\mu}{10^{-12} \text{ eV}} \right) \left(\frac{M_{\text{BH}}}{10 M_{\odot}} \right)$$

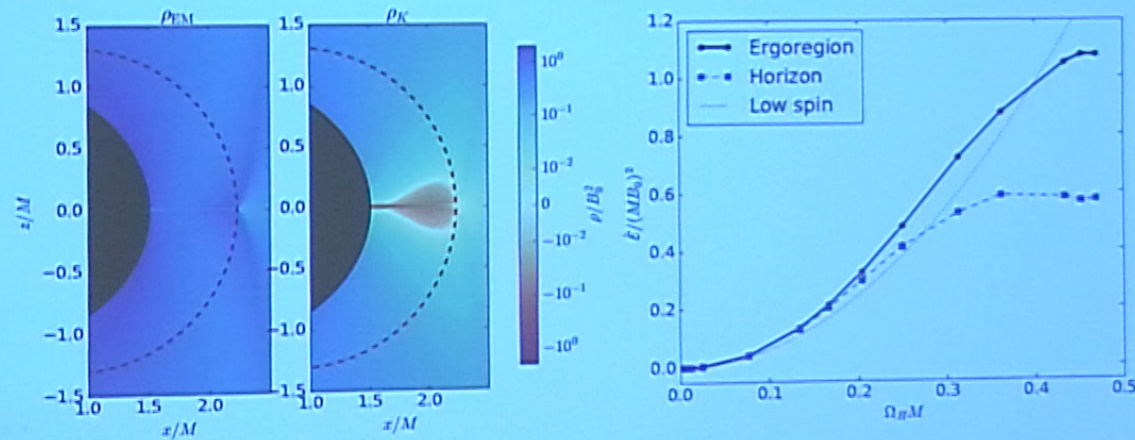
Black hole rotation powering jets



WE & Yang (2018)

Uniform field force-free jet solution (see also Komissarov 2004, Palenzuela et al. 2010, others)

Dissipating negative energy



WE & Yang (2018)

In ergosphere Killing energy density can be negative.

$$l=m=1$$

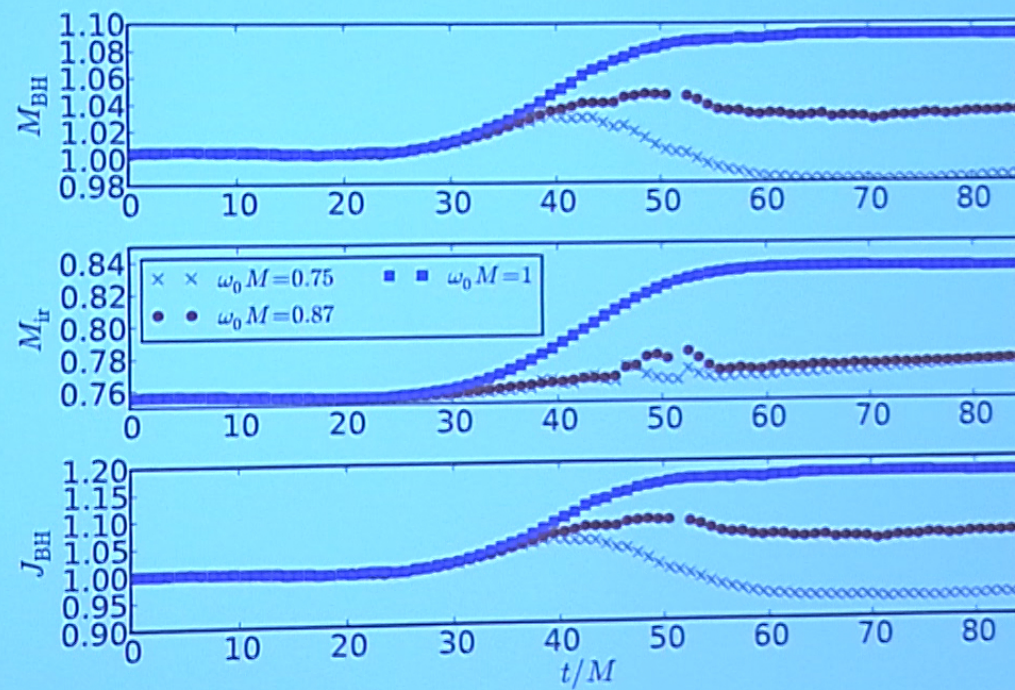
$$\phi \sim e^{-r/2r_0} Y_l^m(\theta, \phi) e^{-i(\omega t)} \quad \text{complex}$$

$$\left[\begin{array}{c} n_a n_b T^{ab} \\ n_a T^a_b \hat{T}_s \end{array} \right]$$

$$\text{Re}(w) = \mu \left(1 - \frac{\alpha^2}{2n^2} \right)$$

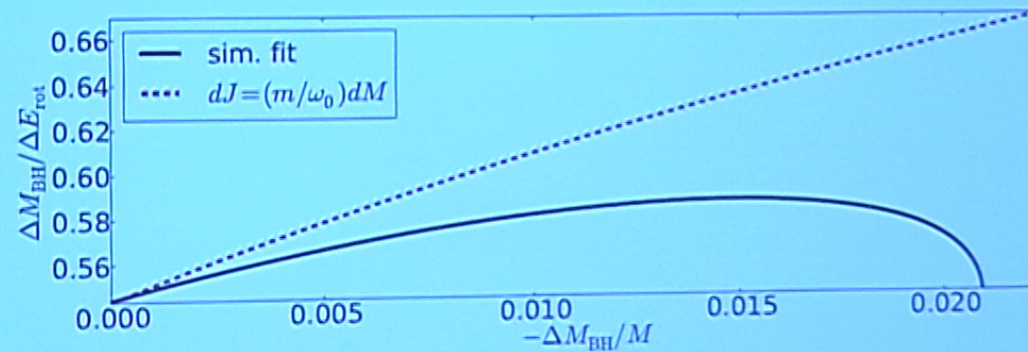
$$|T_m(w)| < | \text{Re}(w) |$$

Nonlinear gravitational wave superradiance



WE, Ramazanoglu & Pretorius (2013)

Liberating energy/angular momentum from a black hole

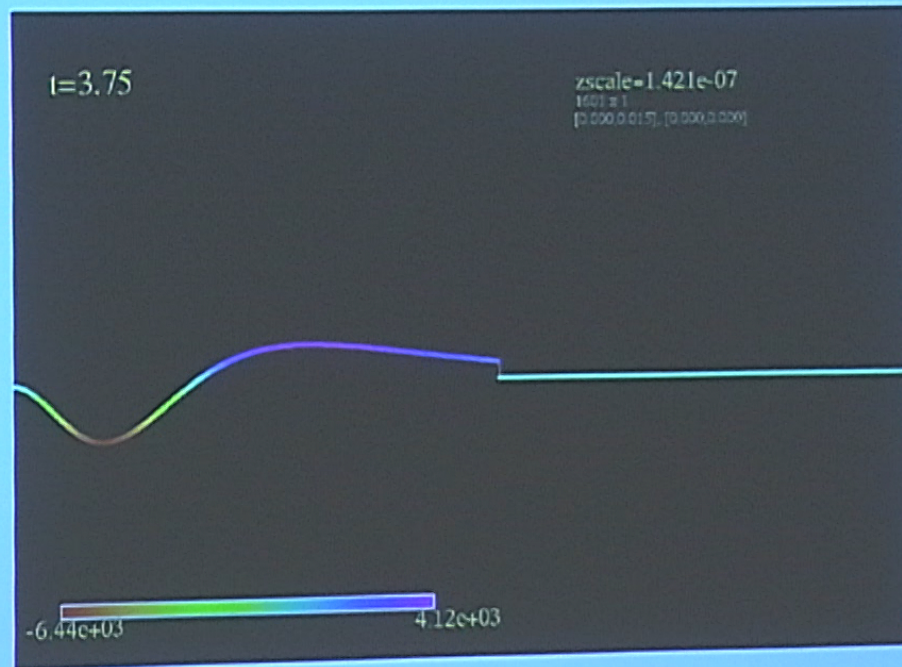


WE, Ramazanoglu & Pretorius (2013)

Efficiency of energy extraction reduced in nonlinear regime

Superradiant instability in Anti-de Sitter Space

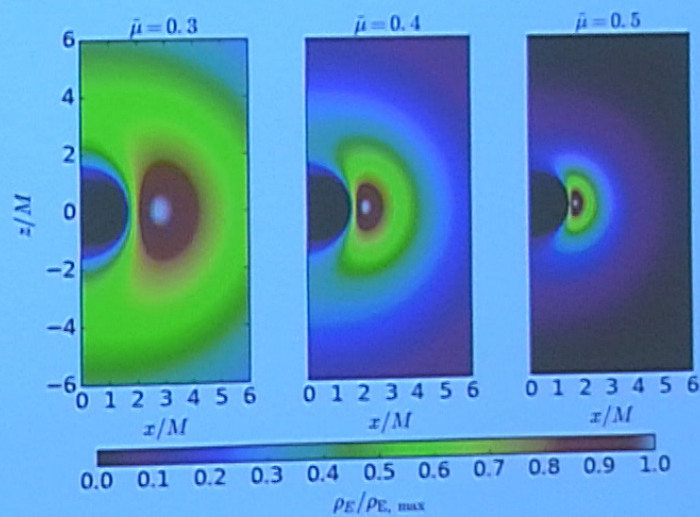
Charged BH/scalar field in AdS



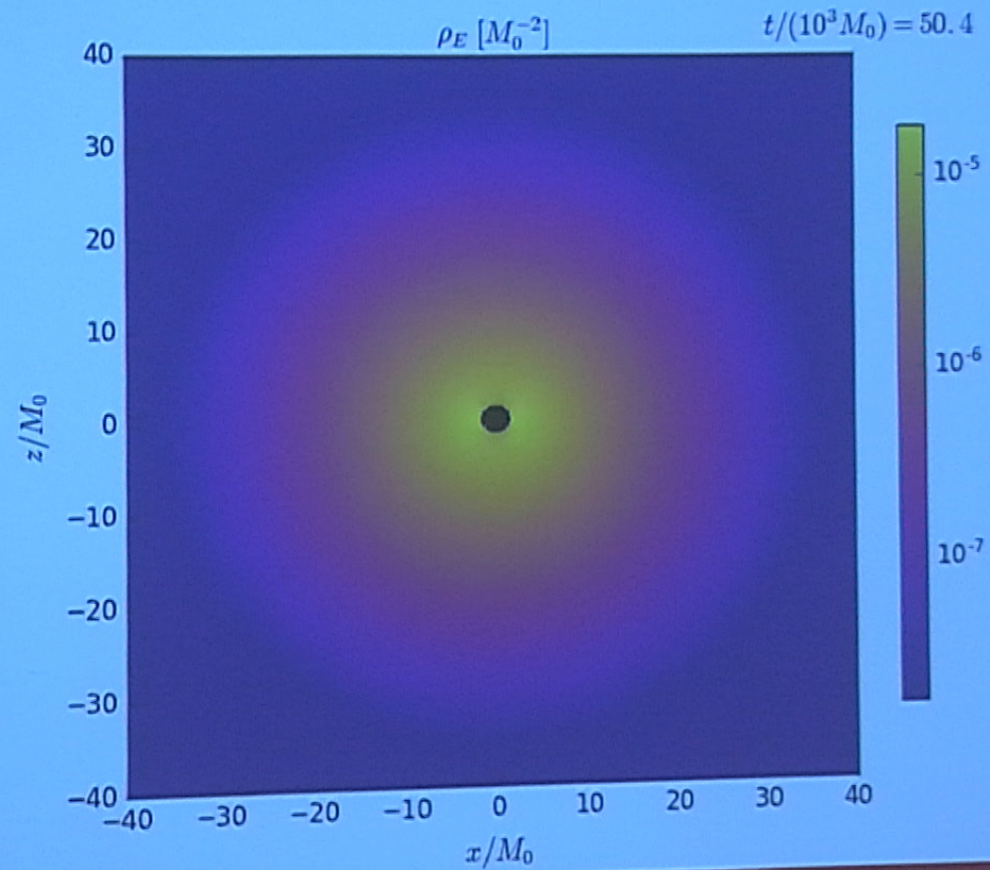
Bosch, Green, & Lehner (2016)

Superradiant instability of massive bosons

- **Massive** bosons with frequency $\omega < m\Omega_H$ grow exponentially in time.
- Significant when Compton wavelength comparable to black hole radius.

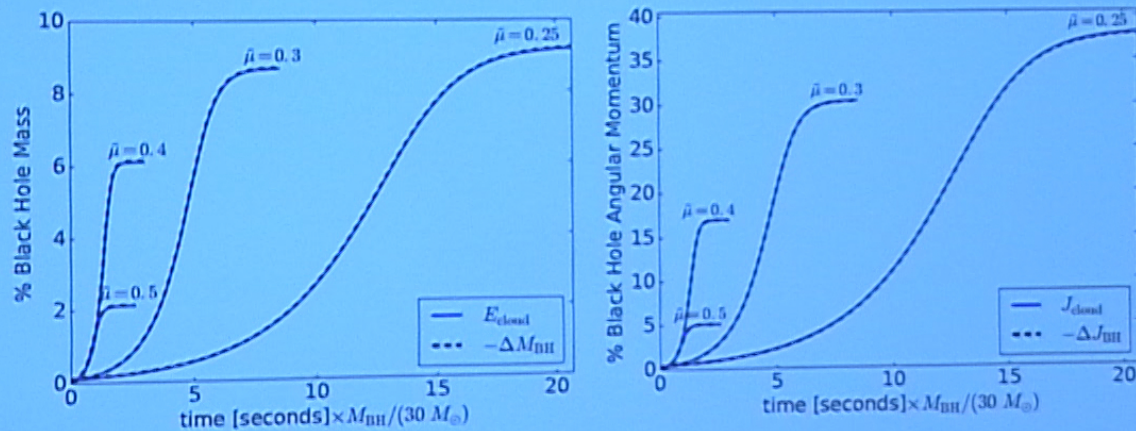


Massive vector field energy density



Superradiant instability: spinning down a black hole

Black hole with initial spin $a = 0.99$.

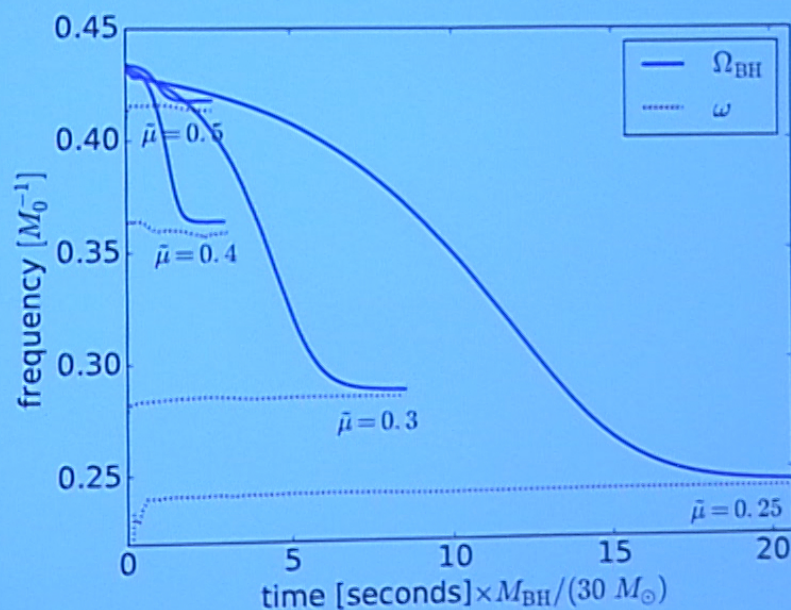


WE & Pretorius (2017)

Vector boson mass: $\approx 10^{-12} \text{ eV } (\tilde{\mu}/0.25)(30M_{\odot}/M_{\text{BH}})$

Superradiant instability: spinning down a black hole

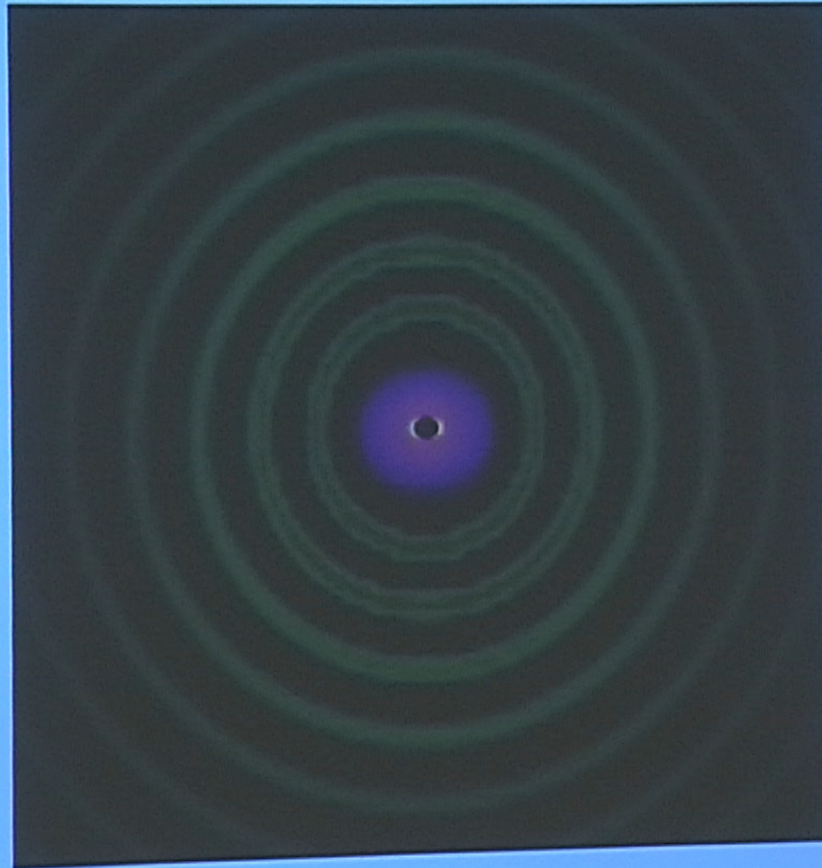
Black hole with initial spin $a = 0.99$.



WE& Pretorius (2017)

Vector boson mass: $\approx 10^{-12} \text{ eV } (\tilde{\mu}/0.25)(30M_{\odot}/M_{\text{BH}})$

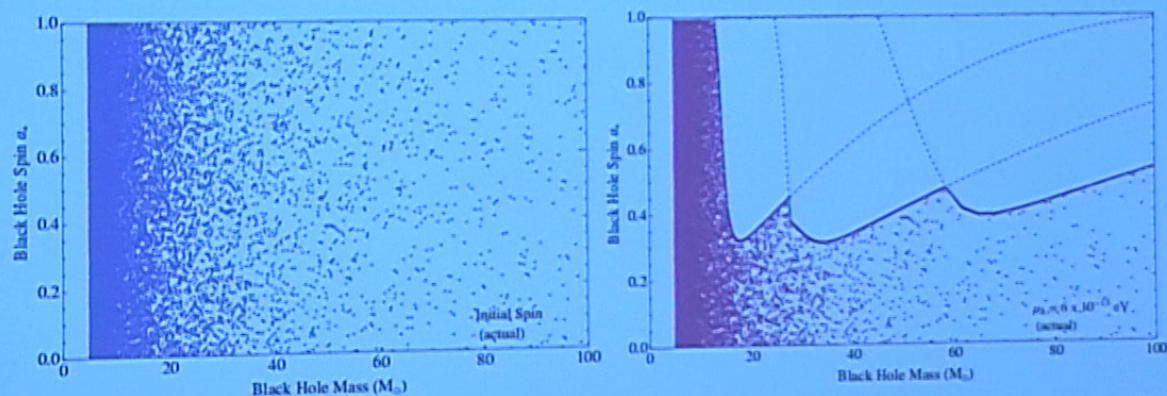
Massive vector field energy density and GWs



WE (2018)

Observing the superradiant instability

Indirect probe: Use black hole spin measurements from GWs or accretion disks (Arvanitaki et al. 2015, 2017)



Need reliable spin measurements