Title: Not extending the standard cosmological model

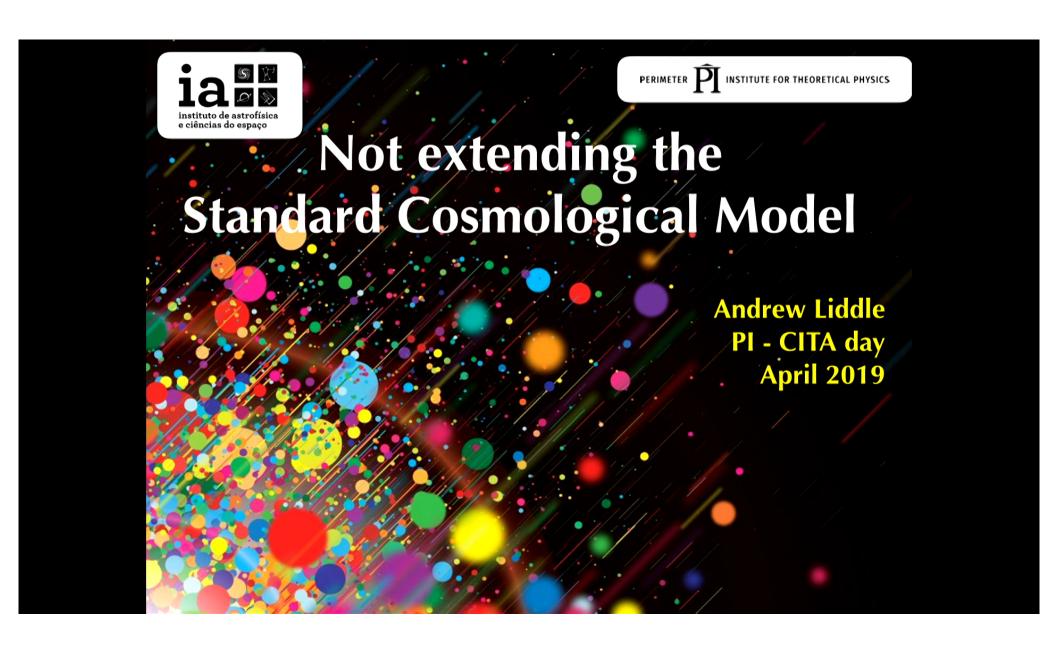
Speakers: Andrew Liddle

Collection: PI-CITA Day 2019

Date: April 02, 2019 - 2:00 PM

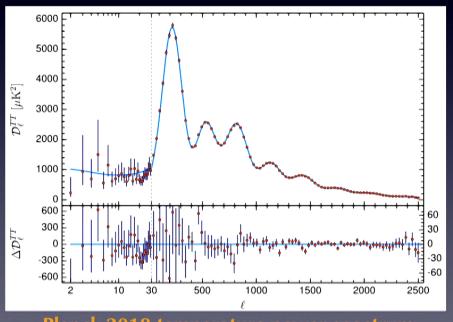
URL: http://pirsa.org/19020044

Pirsa: 19020044 Page 1/35



Pirsa: 19020044 Page 2/35

Starting point: the standard sixparameter cosmological model is extraordinarily successful and gives a precision description of our Universe.



Parameter	TT,TE,EE+lowE+lensing 68% limits
$\Omega_{\mathrm{b}}h^{2}$	0.02237 ± 0.00015
$\Omega_{\rm c}h^2$	0.1200 ± 0.0012
$100\theta_{\mathrm{MC}}$	1.04092 ± 0.00031
τ	0.0544 ± 0.0073
$\ln(10^{10}A_{\mathrm{s}})$	3.044 ± 0.014
n _s	0.9649 ± 0.0042

Planck 2018 temperature power spectrum and parameter constraints.

Pirsa: 19020044 Page 3/35

Once you've decided what observations you are going to try and explain, you have to figure out how to make predictions. You'll need physical laws and environmental descriptors. These will depend on a bunch of unknown parameters that you hope to measure.

Pirsa: 19020044 Page 4/35

Once you've decided what observations you are going to try and explain, you have to figure out how to make predictions. You'll need physical laws and environmental descriptors. These will depend on a bunch of unknown parameters that you hope to measure.

Starting with the cosmic microwave background, this needs a minimum of six parameters, which can be taken to be

- Densities of baryons and of cold dark matter.
- Expansion rate (Hubble constant).
- Density perturbation amplitude and scale-dependence (spectral index).
- Optical depth for rescattering of CMB photons.

These six turn out to be sufficient as well as necessary.

Pirsa: 19020044 Page 5/35

There might also be any number of so-called `nuisance' parameters to model various effects (eg subtraction of uncertain foreground contaminants from CMB maps, or intrinsic alignment of galaxies from weak lensing surveys).

These are eventually marginalized over and ignored.

Pirsa: 19020044 Page 6/35

There might also be any number of so-called `nuisance' parameters to model various effects (eg subtraction of uncertain foreground contaminants from CMB maps, or intrinsic alignment of galaxies from weak lensing surveys).

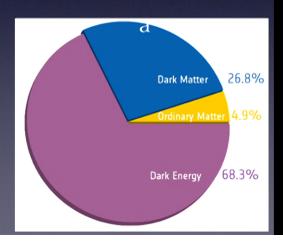
These are eventually marginalized over and ignored.

Of course, one person's nuisance may be another person's signal.

Pirsa: 19020044 Page 7/35

Status of observational cosmology

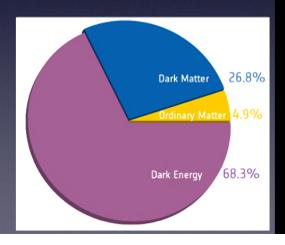
- We can make precision measurements, especially of the cosmic microwave background.
- A simple model, with just six parameters, is sufficient to explain all the major datasets.
- The composition of the Universe is accurately measured. It contains substantial amounts of dark matter and dark energy.
- There are strong indications that period of inflation took place in the very young Universe.



Pirsa: 19020044 Page 8/35

Status of theoretical cosmology

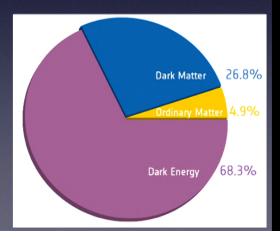
We don't know why the Universe is all matter and no anti-matter.



Pirsa: 19020044 Page 9/35

Status of theoretical cosmology

- We don't know why the Universe is all matter and no anti-matter.
- We don't know what the dark matter is.
- We don't know what the dark energy is.
- We don't know how inflation took place.



Pirsa: 19020044 Page 10/35

Cosmological parameters 2019			
PARAMETER	MEASURED ACCURACY	THEORETICAL ACCURACY	
Baryon density Hubble parameter Cold dark matter density Perturbation amplitude Optical depth Spectral index (dev. from	3% 13%	Many orders of magnitude Not fundamental None, despite plausible candidates Undetermined by theory Factor of 2, but not fundamental Plausibly explained by inflation	
Dark energy density (der	ived) 1%	60 or 120 orders of magnitude?!	
Principles Atomic and nuclear phys Neutrino physics General Relativity Quantum field theory	ics Very good Consistent but weal Approximately, mos Mildly tested		

Pirsa: 19020044 Page 11/35

Changing the six

The ideal outcome of any observational programme is to **break** the standard cosmological model.

New parameter
=
New physics to learn about!

Instead, we've been constraining the same six parameters better and better for over a decade.

Pirsa: 19020044 Page 12/35

Changing the six

The ideal outcome of any observational programme is to **break** the standard cosmological model.

New parameter
=
New physics to learn about!

Instead, we've been constraining the same six parameters better and better for over a decade.

But they were already measured far better than current theory demands.

Pirsa: 19020044 Page 13/35

Hidden parameter:

The radiation temperature is considered so well measured as not to be varied. However, it defines the present moment.

6

Pirsa: 19020044 Page 14/35

Hidden parameter:

The radiation temperature is considered so well measured as not to be varied. However, it defines the present moment.

6

On the way out:

The optical depth isn't really an independent parameter. It should be predictable from the others.

5

Pirsa: 19020044 Page 15/35

Hidden parameter:

The radiation temperature is considered so well measured as not to be varied. However, it defines the present moment.

6

On the way out:

The optical depth isn't really an independent parameter. It should be predictable from the others.

5

Coming and then going:

The helium fraction is needed for high-precision CMB calculations. But, it should be predictable from the others.

5

Pirsa: 19020044 Page 16/35

Hidden parameter:

The radiation temperature is considered so well measured as not to be varied. However, it defines the present moment.

6

On the way out:

The optical depth isn't really an independent parameter. It should be predictable from the others.

5

Coming and then going:

The helium fraction is needed for high-precision CMB calculations. But, it should be predictable from the others.

5

Coming soon:

Neutrino mass(es) is on the verge of inevitable detection.



Pirsa: 19020044 Page 17/35

Much more interesting is what *might* happen in the future.

Pirsa: 19020044 Page 18/35

Table 2. Candidate parameters: those which might be relevant for cosmological observations, but for which there is presently no convincing evidence requiring them. They are listed so as to take the value zero in the base cosmological model. Those above the line are parameters of the background homogeneous cosmology, and those below describe the perturbations.

$ \Omega_k \\ N_{\nu} - 3.04 \\ m_{\nu_i} $	spatial curvature effective number of neutrino species (CMBFAST definition) neutrino mass for species 'i'	
	[or more complex neutrino properties]	
$m_{ m dm}$	(warm) dark matter mass	
w+1	dark energy equation of state	
dw/dz	redshift dependence of w	
	[or more complex parametrization of dark energy evolution]	
$c_{\rm S}^2 - 1$	effects of dark energy sound speed	
$1/r_{ m top}$	topological identification scale	
	[or more complex parametrization of non-trivial topology]	
dlpha/dz	redshift dependence of the fine structure constant	
dG/dz	redshift dependence of the gravitational constant	
n-1	scalar spectral index	
$dn/d \ln k$	running of the scalar spectral index	
r	tensor-to-scalar ratio	
$r + 8n_{\mathrm{T}}$	violation of the inflationary consistency equation	
$dn_{\rm T}/d\ln k$	running of the tensor spectral index	
k_{cut}	large-scale cut-off in the spectrum	
$A_{ m feature}$	amplitude of spectral feature (peak, dip or step)	
$k_{ m feature}$	and its scale	
	[or adiabatic power spectrum amplitude parametrized in N bins]	
$f_{ m NL}$	quadratic contribution to primordial non-gaussianity	
	[or more complex parametrization of non-gaussianity]	
\mathcal{P}_S	CDM isocurvature perturbation	
n_S	and its spectral index	
$\mathcal{P}_{S\mathcal{R}}$	and its correlation with adiabatic perturbations	
$n_{SR} - n_S$	and the spectral index of that correlation	
	[or more complicated multi-component isocurvature perturbation]	
$G\mu$	cosmic string component of perturbations	From Liddle 2004

Pirsa: 19020044 Page 19/35

Table 2. Candidate parameters: those which might be relevant for cosmological observations, but for which there is presently no convincing evidence requiring them. They are listed so as to take the value zero in the base cosmological model. Those above the line are parameters of the background homogeneous cosmology, and those below describe the perturbations.

Ω_k	spatial curvature		
$N_{ u} - 3.04$	effective number of neutrino species (CMBFAST definition)		
$m_{ u_i}$	neutrino mass for species 'i'		
,	[or more complex neutrino properties]		
$m_{ m dm}$	(warm) dark matter mass		
w+1	dark energy equation of state		
dw/dz	redshift dependence of w		
,	[or more complex parametrization of dark energy evolution]		
$c_{\rm S}^2 - 1$	effects of dark energy sound speed		
$1/r_{\rm top}$	topological identification scale		
,	or more complex parametrization of non-trivial topology		
$d\alpha/dz$	redshift dependence of the fine structure constant		
$dG^{'}/dz$	redshift dependence of the gravitational constant		
n-1	scalar spectral index This is the only one convincingly		
$dn/d \ln k$	running of the scalar spectral index detected in the intervening 15 years.		
r	tensor-to-scalar ratio		
$r + 8n_{\mathrm{T}}$	violation of the inflationary consistency equation		
$dn_{\mathrm{T}}/d\ln k$	running of the tensor spectral index		
k_{cut}	large-scale cut-off in the spectrum		
$A_{ m feature}$	amplitude of spectral feature (peak, dip or step)		
$k_{ m feature}$	and its scale		
	or adiabatic power spectrum amplitude parametrized in N bins		
$f_{ m NL}$	quadratic contribution to primordial non-gaussianity		
	[or more complex parametrization of non-gaussianity]		
\mathcal{P}_S	CDM isocurvature perturbation		
n_S	and its spectral index		
$\mathcal{P}_{S\mathcal{R}}$	and its correlation with adiabatic perturbations		
$n_{SR} - n_S$	and the spectral index of that correlation		
	[or more complicated multi-component isocurvature perturbation]		
$G\mu$	cosmic string component of perturbations From Liddle 2004		

Pirsa: 19020044 Page 20/35

Table 2. Candidate parameters: those which might be relevant for cosmological observations, but for which there is presently no convincing evidence requiring them. They are listed so as to take the value zero in the base cosmological model. Those above the line are parameters of the background homogeneous cosmology, and those below describe the perturbations.

$\Omega_k \\ N_{\nu} - 3.04 \\ m_{\nu_i}$	spatial curvature effective number of neutrino species (CMBFAST definition) neutrino mass for species 'i' This is the only one whose		
$m_{ m dm}$	[or more complex neutrino properties] (warm) dark matter mass future detection seems inevitable.		
w+1	dark energy equation of state		
dw/dz	redshift dependence of w		
2	[or more complex parametrization of dark energy evolution]		
$c_{\rm S}^2 - 1$	effects of dark energy sound speed		
$1/r_{ m top}$	topological identification scale		
10/10	[or more complex parametrization of non-trivial topology]		
$rac{dlpha/dz}{dG/dz}$	redshift dependence of the fine structure constant		
	redshift dependence of the gravitational constant		
n-1	scalar spectral index This is the only one convincingly		
$dn/d \ln k$	running of the scalar spectral index detected in the intervening 15 years.		
r	tensor-to-scalar ratio		
$r + 8n_{\mathrm{T}}$	violation of the inflationary consistency equation		
$dn_{\mathrm{T}}/d\ln k$	running of the tensor spectral index		
$k_{ m cut}$	large-scale cut-off in the spectrum		
$A_{ m feature}$	amplitude of spectral feature (peak, dip or step)		
$k_{ m feature}$	and its scale		
C	[or adiabatic power spectrum amplitude parametrized in N bins]		
$f_{ m NL}$	quadratic contribution to primordial non-gaussianity		
D-	[or more complex parametrization of non-gaussianity]		
\mathcal{P}_S	CDM isocurvature perturbation and its spectral index		
$n_S \ \mathcal{P}_{S\mathcal{R}}$	and its correlation with adiabatic perturbations		
$n_{SR} - n_S$	and the spectral index of that correlation		
. SK .VS	[or more complicated multi-component isocurvature perturbation]		
$G\mu$	cosmic string component of perturbations From Liddle 2004		

Pirsa: 19020044 Page 21/35

Table 2. Candidate parameters: those which might be relevant for cosmological observations, but for which there is presently no convincing evidence requiring them. They are listed so as to take the value zero in the base cosmological model. Those above the line are parameters of the background homogeneous cosmology, and those below describe the perturbations.

$\Omega_k \\ N_{\nu} - 3.04 \\ m_{\nu_i} \\ m_{\rm dm}$	spatial curvature effective number of neutrino species (CMBFAST definition) neutrino mass for species 'i' [or more complex neutrino properties] (warm) dark matter mass This is the only one volume of the complex neutrino properties of the complex neutrino prope	
$w + 1$ dw/dz $c_{\rm S}^2 - 1$ $1/r_{\rm top}$ $d\alpha/dz$ dG/dz	dark energy equation of state redshift dependence of w [or more complex parametrization of dark energy evolution] effects of dark energy sound speed topological identification scale [or more complex parametrization of non-trivial topology] redshift dependence of the fine structure constant redshift dependence of the gravitational constant	lified that
n-1	scalar spectral index This is the only one convince	cingly
$dn/d \ln k$	running of the scalar spectral index detected in the intervening 15	0 ,
r	tensor-to-scalar ratio	years.
$r + 8n_{\mathrm{T}}$	violation of the inflationary consistency equation	
$dn_{\mathrm{T}}/d\ln k$	running of the tensor spectral index	
k_{cut}	large-scale cut-off in the spectrum	
$A_{ m feature}$	amplitude of spectral feature (peak, dip or step)	
$k_{ m feature}$	and its scale	
	[or adiabatic power spectrum amplitude parametrized in N bins]	
$f_{ m NL}$	quadratic contribution to primordial non-gaussianity	
	[or more complex parametrization of non-gaussianity]	
\mathcal{P}_S	CDM isocurvature perturbation	
n_S	and its spectral index	
$\mathcal{P}_{S\mathcal{R}}$	and its correlation with adiabatic perturbations	
$n_{SR} - n_S$	and the spectral index of that correlation	
$G\mu$	[or more complicated multi-component isocurvature perturbation] cosmic string component of perturbations From Liddle	2004
$G\mu$	cosmic string component of perturbations From Liddle	2004

Pirsa: 19020044 Page 22/35

Methods for model selection

Invent a threshold.

Archetype: the particle physicists' 5-sigma criterion.

Pirsa: 19020044 Page 23/35

Methods for model selection

Invent a threshold.
Archetype: the particle physicists' 5-sigma criterion.

Information theory methods.

These view a model as an algorithmic compression of the data.

A successful model optimizes the compression.

Archetype: Akaike information criterion.

Number of bits

DATA ALONE		
MODEL	DATA RESIDUALS	
WORSE MODEL	DATA RESIDUALS	
BETTER MODEL	DATA RESIDUALS	

Pirsa: 19020044 Page 24/35

Methods for model selection

Invent a threshold.

Archetype: the particle physicists' 5-sigma criterion.

Information theory methods.

These view a model as an algorithmic compression of the data.

BETTER MODEL

A successful model optimizes the compression.

Archetype: Akaike information criterion.

Number of bits

DATA ALONE		
MODEL	DATA RESIDUALS	
WORSE MODEL	DATA RESIDUALS	

DATA RESIDUALS

Bayesian methods.

These assign probabilities to all quantities of interest and update when new data comes in. For model selection we assign a probability to each *set* of parameters as well as to the parameter values. Archetype: Bayesian evidence (aka Bayes factor).

Pirsa: 19020044 Page 25/35

Dataset comparison

An interesting recent use of the Bayesian evidence is to test compatibility of two datasets D_1 and D_2 . We compare the evidences of a single-model fit with one assuming that each dataset measures independent parameters within the model:

$$R = \frac{P(\vec{D_1}, \vec{D_2}|M)}{P(\vec{D_1}|M)P(\vec{D_2}|M)}$$

Pirsa: 19020044 Page 26/35

Dataset comparison

An interesting recent use of the Bayesian evidence is to test compatibility of two datasets D_1 and D_2 . We compare the evidences of a single-model fit with one assuming that each dataset measures independent parameters within the model:

$$R = \frac{P(\vec{D_1}, \vec{D_2}|M)}{P(\vec{D_1}|M)P(\vec{D_2}|M)}$$

This is better understood via a rewriting using Bayes' Theorem:

$$R = \frac{P(\vec{D}_2 | \vec{D}_1, M)}{P(\vec{D}_2 | M)}$$

Pirsa: 19020044 Page 27/35

Dataset comparison

An interesting recent use of the Bayesian evidence is to test compatibility of two datasets D_1 and D_2 . We compare the evidences of a single-model fit with one assuming that each dataset measures independent parameters within the model:

$$R = \frac{P(\vec{D_1}, \vec{D_2}|M)}{P(\vec{D_1}|M)P(\vec{D_2}|M)}$$

This is better understood via a rewriting using Bayes' Theorem:

$$R = \frac{P(\vec{D}_{2}|\vec{D}_{1}, M)}{P(\vec{D}_{2}|M)}$$

Note that dataset consistency depends on the chosen model.

Pirsa: 19020044 Page 28/35

DES collaboration, arXiv:1708.01530, PRD

Dark Energy Survey on w

PHYSICAL REVIEW D 98, 043526 (2018)

Editors' Suggestion Featured in Physics

Dark Energy Survey year 1 results: Cosmological constraints from galaxy clustering and weak lensing

T. M. C. Abbott, ¹ F. B. Abdalla, ^{2,3} A. Alarcon, ⁴ J. Aleksić, ⁵ S. Allam, ⁶ S. Allen, ⁷ A. Amara, ⁸ J. Annis, ⁶ J. Asorey, ^{9,10} S. Avila, ^{11,12} D. Bacon, ¹¹ E. Balbinot, ¹³ M. Banerii, ^{14,15} N. Banik, ⁶ W. Barkhouse, ¹⁶ M. Baumer, ^{7,17,18} E. Bakter, ¹⁹ K. Bechtol, ²⁰ M. R. Becker, ^{7,17} A. Benoit-Lévy, ^{3,21,22} B. A. Benson, ^{6,23} G. M. Bernstein, ¹⁹ E. Bertin, ^{22,21} J. Blazek, ^{24,25} S. L. Bridle, ²⁶ D. Brooks, ³ D. Brout, ¹⁹ E. Buckley-Geer, ⁶ D. L. Burke, ^{17,18} M. T. Busha, ¹⁷ A. Campos, ^{27,28} D. Capozzi, ¹¹ A. Carrero Rosell, ^{2,2} M. Carrasco Kind, ^{3,3,1} J. Carretero, ⁵ F. J. Castander, ⁴ R. Cawthon, ² C. Chang, ² N. Chen, ³ M. Childress, ² A. Choi, ² C. Conselice, ³ R. Crittenden, ¹¹ M. Crocce, ⁴ C. E. Cunha, ¹⁷ C. B. D'Andrea, ¹⁹ L. N. da Costa ^{3,2,3} R. Das, ^{3,4} T. M. Davis, ^{3,1} C. Davis, ^{1,7} J. De Vicente, ^{3,6} D. L. DePoy, ^{3,6} J. DeRose, ^{3,7} S. Desai, ^{3,7} H. T. Diehl, ⁶ J. P. Dietrich, ^{38,30} S. Dodelson, ^{6,2,3} P. Doel, ³ A. Orlica-Wagner, ⁶ T. F. Eifler, ^{40,41} A. E. Elliott, ^{4,7} F. Esner, ¹ J. Elvin-Poole, ^{2,6} J. Estrada, ⁶ A. E. Evrard, ^{43,34} Y. Fang, ¹⁰ E. Fernandez, ⁵ A. Ferté, ⁴⁴ D. A. Finley, ⁶ B. Flaugher, ⁶ P. Fosalba, ⁴ O. Friedrich, ^{45,46} J. Frieman, 23.6 J. García-Bellido, 12 M. Garcia-Fernandez, 35 M. Gatti, E. Gaztanaga, D. W. Gerdes, 34.43 T. Giannantonio, ^{45,15,14} M. S. S. Gill, ¹⁸ K. Glazebrook, ⁴⁷ D. A. Goldstein, ^{48,49} D. Gruen, ^{30,17,18} R. A. Gruendl, ^{11,20} J. Gschwend, ^{28,29} G. Gutierrez, ⁶ S. Hamilton, ¹⁴ W. G. Hartley, ^{3,8} S. R. Hinton, ⁹ K. Honscheid, ^{25,42} B. Hoyle, ⁴⁵ D. Hutterer, ^{21,54} G. Grand, ^{25,29} G. Gutierrez, ⁶ S. Hamilton, ¹⁴ W. G. Hartley, ^{3,8} S. R. Hinton, ⁹ K. Honscheid, ^{25,29} B. Hoyle, ⁴⁵ D. Hutterer, ^{21,54} G. Grand, ^{25,20} G. Gutierrez, ⁶ S. Hamilton, ¹⁵ G. Grand, ^{25,20} G. B. Jain, D. J. James, J. M. Jarvis, T. Jeltema, S. M. D. Johnson, M. W. G. Johnson, T. Kacprzak, S. Kent, S. A. G. Kim, A. A. G. Kim, A. A. King, D. Kirk, N. Kokron, A. Kovacs, E. Krause, C. Krawiec, A. Kremin, K. Kuehn, K. Kuehn, A. G. Kim, D. Kirk, A. Kokron, A. Kovacs, E. Krause, C. Krawiec, A. Kremin, K. Kuehn, K. Kuehn, A. K. Kuehn, M. S. Kuhlmann. S. N. Kuropatkin, F. Lucasa, Z. O. Lahav, T. S. Li, A. R. Liddle, H. C. Lidman, 10.54 M. Lima, N. MacCrann, 12.25 M. A. G. Maia, 29.28 M. Makler, M. Manera, M. March, J. L. Marshall, P. Martini, 77.25 R. G. McMahon, 14, 15 P. Melchior, S. F. Menanteau, 30.31 R. Miquel, 5.59 V. Miranda, 19 D. Mudd, 7 J. Muir, 34 A. Möller, 66, 10 E. Neilsen, R. C. Nichol, B. Nord, P. Nugent, R. L. C. Ogando, A. Palmese, J. Peacock, H. V. Peiris, J. Peoples, W. J. Percival, D. Petravick, A. A. Plazas, A. Porredon, J. Prat, A. Pujol, M. M. Rau, A. Refregier, P. M. Ricker, 11,30 N. Roe, 40 R. P. Rollins, 26 A. K. Romer, 61 A. Roodman, 17,18 R. Rosenfeld, 27,28 A. J. Ross, 25 E. Rozo, E. S. Rykoff, ^{17,18} M. Sako, ¹⁰ A. I. Salvador, ³ S. Samuroff, ²⁶ C. Sánchez, ⁵ E. Sanchez, ⁵ B. Santiago, ^{3,28} V. Scarpine, ⁸ R. Schindler, ¹⁸ D. Scolnic, ²³ L. F. Secco, ¹⁹ S. Serrano, ⁴ I. Sevilla-Noarbe, ³⁵ E. Sheldon, ⁵⁴ R. C. Smith, ¹ M. Smith, ³² J. Smith, 65 M. Soares-Santos, 6 F. Sobreira, 28,66 E. Suchyta, 67 G. Tarle, 34 D. Thomas, 11 M. A. Troxel, 42,25 D. L. Tucker, 6 B. E. Tucker, ^(1,00) S. A. Uddin, ^(1,08) T. N. Varga, ^{46,45} P. Vielzeuf, ⁵ V. Vikram, ⁵⁵ A. K. Vivas, A. R. Walker, M. Wang, R. H. Wechsler, ^{18,17,7} J. Weller, ^{18,45,46} W. Wester, ⁶ R. C. Wolf, ¹⁹ B. Yanny, ⁶ F. Yuan, ^{10,60} A. Zenteno, ¹ B. Zhang, 60,10 Y. Zhang, and J. Zuntz44

(Dark Energy Survey Collaboration)

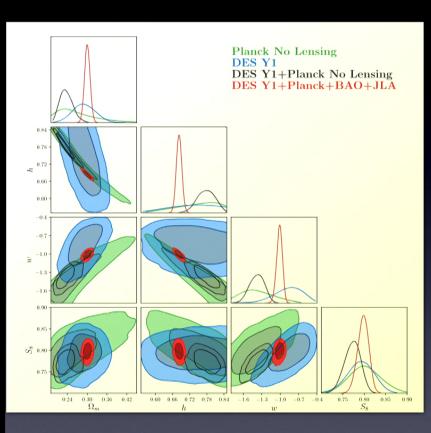
We used the first-year dataset from the Dark Energy Survey to constrain the equation of state w, through combination with Planck and other data. We found

$$W = -1.00^{+0.05}$$
-0.04

Pirsa: 19020044 Page 29/35

DES collaboration, arXiv:1708.01530, PRD

Dark Energy Survey on w



We used the first-year dataset from the Dark Energy Survey to constrain the equation of state w, through combination with Planck and other data. We found

$$W = -1.00^{+0.05}_{-0.04}$$

Clearly this does not indicate against the cosmological constant case w = -1.

This is borne out by the Bayesian evidence ratio $R_w = 0.1$.

Pirsa: 19020044 Page 30/35

Here are the combined constraints on the modified potentials.

Taking CFHTlens as our fiducial lensing dataset we have

$$\Sigma_0 = +0.05^{+0.05}_{-0.07}$$

$$\mu_0 = -0.10^{+0.20}_{-0.16}$$

Clearly there is no evidence here against GR / \(\Lambda\)CDM.

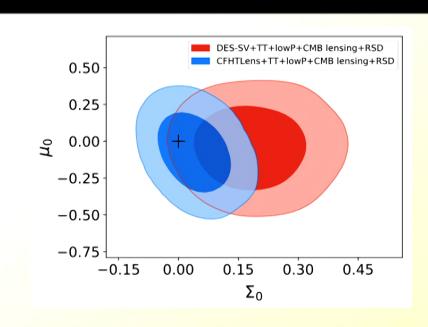


FIG. 8. 68% and 95% confidence contours on Σ_0 and μ_0 combining RSD data (BOSS DR12 + 6dFGS), CMB data (TT+lowP+CMB lensing from *Planck*) and cosmic shear data (CFHTLenS in blue and DES-SV in red).

Pirsa: 19020044 Page 31/35

DES collaboration, arXiv:1810.02499

Dark Energy Survey extensions paper

In this paper we used the Y1 DES data alongside the others to test for various extensions to the standard cosmological model.

Curvature	DES Y1 + External
Ω_k	$0.0020^{+0.0037}_{-0.0032}$
Number Rel. Species	DES Y1 + External
$N_{ m eff}$	< 3.28
Dynamical dark energy	DES Y1 + External
w_0	$-0.95^{+0.09}_{-0.08}$
w_a	$-0.28^{+0.37}_{-0.48}$
w_p	$-1.01^{+0.04}_{-0.04}$
Modified Gravity	DES Y1 + External
Σ_0	$0.06^{+0.08}_{-0.07}$
μ_0	$-0.11^{+0.42}_{-0.46}$

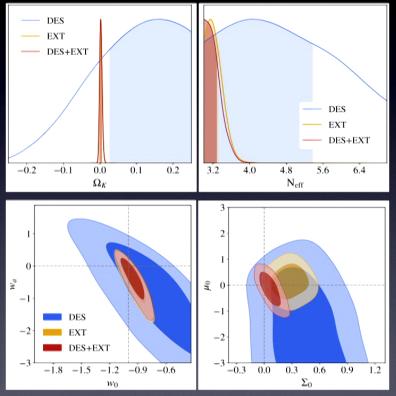
Pirsa: 19020044 Page 32/35

DES collaboration, arXiv:1810.02499

Dark Energy Survey extensions paper

In this paper we used the Y1 DES data alongside the others to test for various extensions to the standard cosmological model.

$egin{aligned} \mathbf{Curvature} \\ & \Omega_k \\ & \mathbf{Number\ Rel.\ Species} \\ & N_{\mathrm{eff}} \end{aligned}$	DES Y1 + External 0.0020 ^{+0.0037} DES Y1 + External < 3.28	— DES — EXT — DES+EXT
Dynamical dark energy w_0 w_a	DES Y1 + External $-0.95^{+0.09}_{-0.08}$ $-0.28^{+0.37}_{-0.48}$	-0.2 -0.1 0.0 0.1 Ω_K
w_p Modified Gravity Σ_0	$-1.01^{+0.04}_{-0.04}$ DES Y1 + External $0.06^{+0.08}_{-0.07}$ $0.11^{+0.42}_{-0.04}$	0 ≥ −1 DES EXT DES+EXT
μ_0	$-0.11^{+0.42}_{-0.46}$	$-3\frac{1}{-1.8}$ -1.5 -1.2 -0.9 w_0



Pirsa: 19020044 Page 33/35

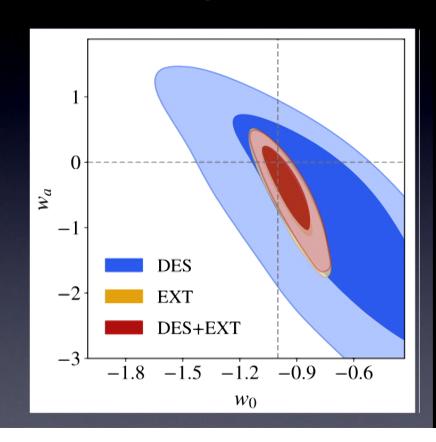
DES collaboration, arXiv:1810.02499

Dark Energy Survey extensions paper

In this paper we used the Y1 DES data alongside the others to test for various extensions to the standard cosmological model.

Curvature	DES Y1 + External
Ω_k	$0.0020^{+0.0037}_{-0.0032}$
Number Rel. Species	DES Y1 + External
$N_{ m eff}$	< 3.28
Dynamical dark energy	DES Y1 + External
w_0	$-0.95^{+0.09}_{-0.08}$
w_a	$-0.28^{+0.37}_{-0.48}$
w_p	$-1.01^{+0.04}_{-0.04}$
Modified Gravity	DES Y1 + External
Σ_0	$0.06^{+0.08}_{-0.07}$
μ_0	$-0.11^{+0.42}_{-0.46}$

 R_{waw_0} = 0.006 ! Less than 1% chance of dynamical dark energy being right!



Pirsa: 19020044 Page 34/35

Conclusion

The Universe is doing very well at resisting our attempts to find any interesting new physics beyond ΛCDM.

Pirsa: 19020044 Page 35/35