

Title: A new probe to axion-like particles from upcoming CMB experiments

Date: Jan 14, 2019 11:00 AM

URL: <http://pirsa.org/19010069>

Abstract: <p>Cosmic Microwave Background (CMB) is a powerful probe to the Universe which carries signatures of cosmic secrets over a vast range of redshifts. Along with spatial fluctuations, spectral distortions of CMB blackbody are also a rich source of cosmological information. In my talk, I will introduce a new kind of spectral distortion of CMB which can arise due to the conversion of CMB photons into Axion-Like Particles (ALPs) in the presence of an external magnetic field. This effect leads to both polarized and unpolarized spatially varying spectral distortion signal with a unique spectral shape when CMB photons undergo resonant and non-resonant conversion into ALPs in the presence of the magnetic field of Milky Way, galaxy clusters and voids. I will discuss the spatial structure of this distortion which can arise from Milky Way and galaxy clusters and will show its uniqueness from other known cosmological and astrophysical signals. I will present the first all-sky map of this distortion which is obtained using the data of Planck satellite and will discuss the capability of this new cosmological window to probe ALPs from the upcoming ground-based and space-based CMB experiments such as CMB-S4, LiteBIRD, Simons Observatory, etc.</p>

A new probe to axion-like particles from upcoming CMB experiments

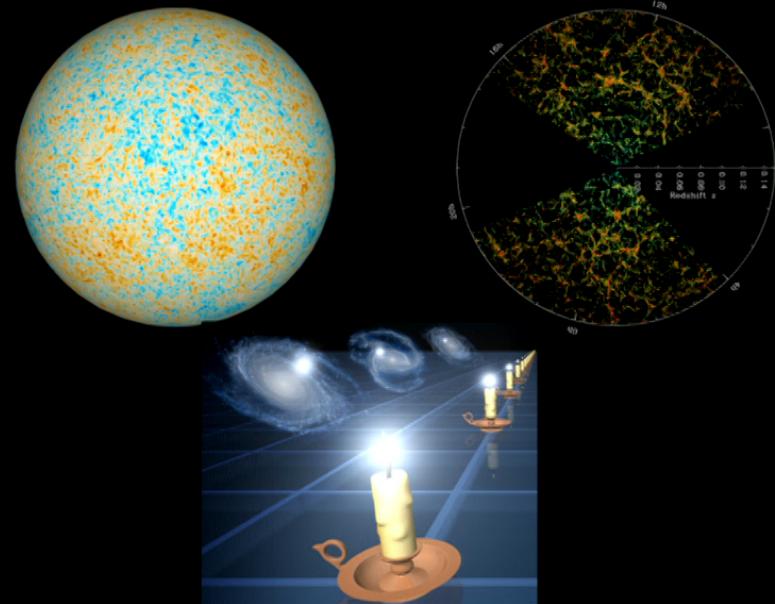
Suvodip Mukherjee

Lagrange Postdoctoral Fellow
IAP, Paris

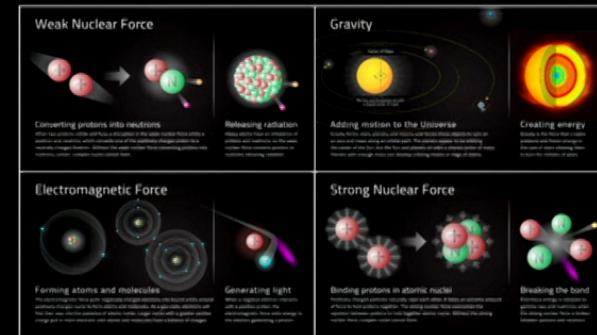
Perimeter Institute



Observational Probes

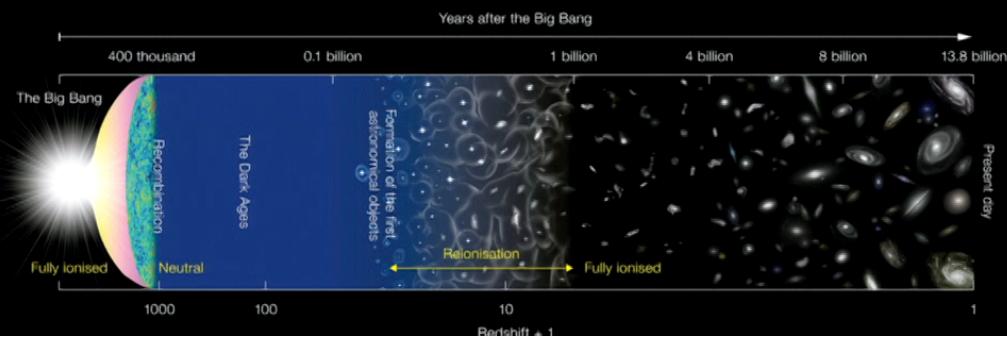
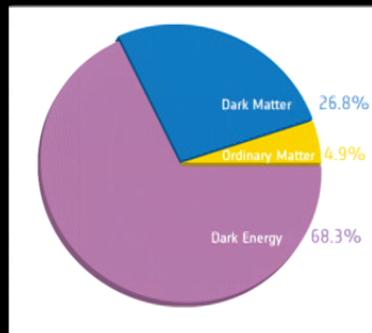


Fundamental forces



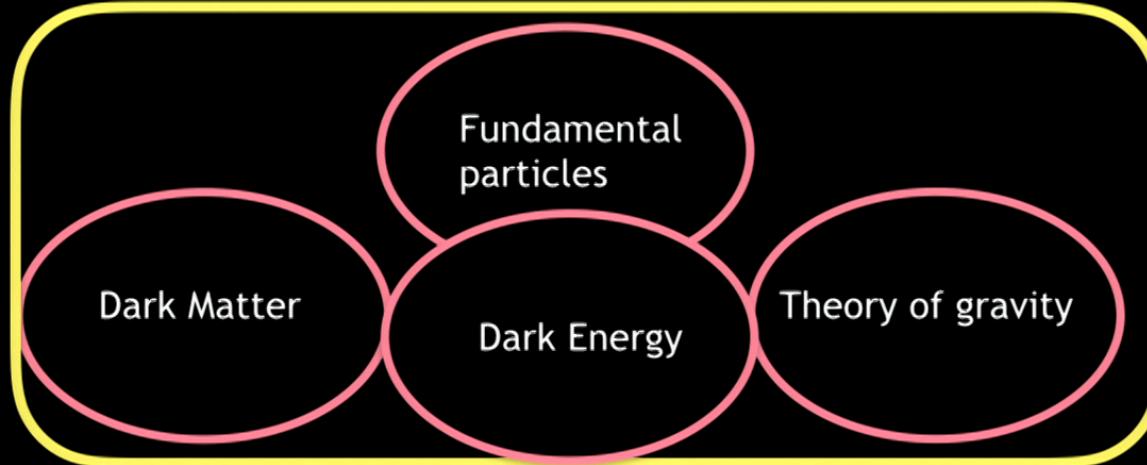
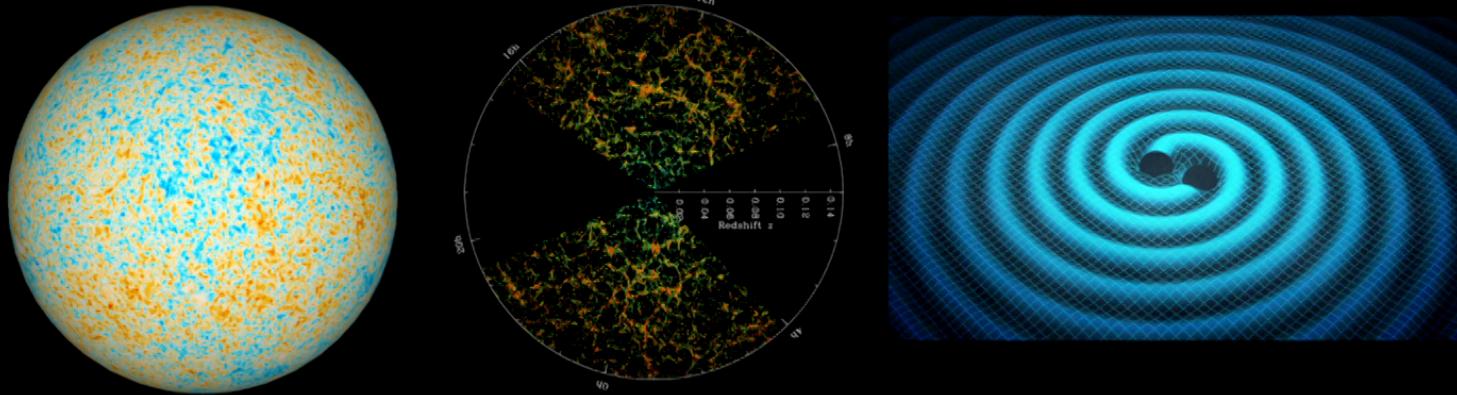
The Standard Model of Cosmology

Figure courtesy: ESO



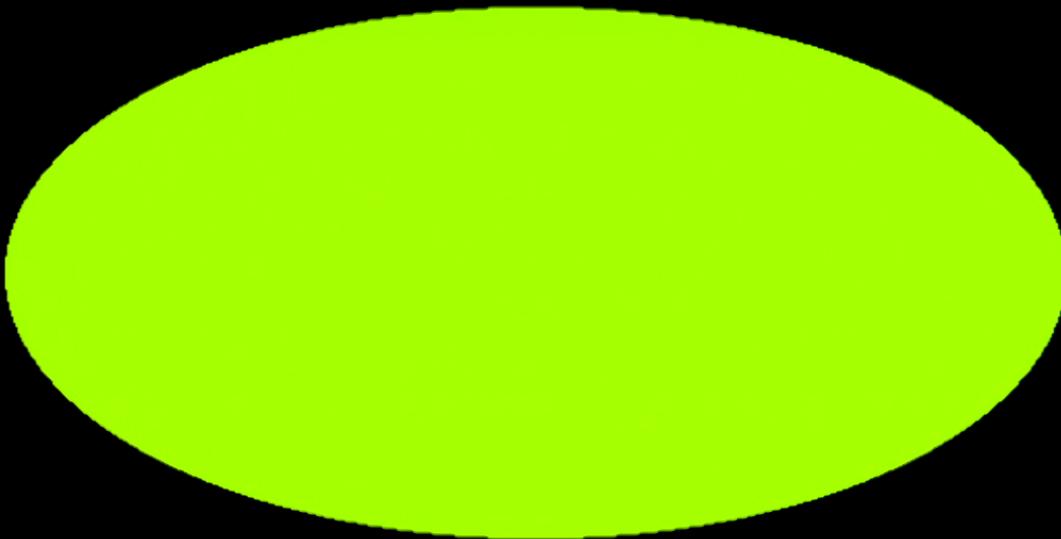
My research interest

Develop connections between cosmological probes and fundamental physics



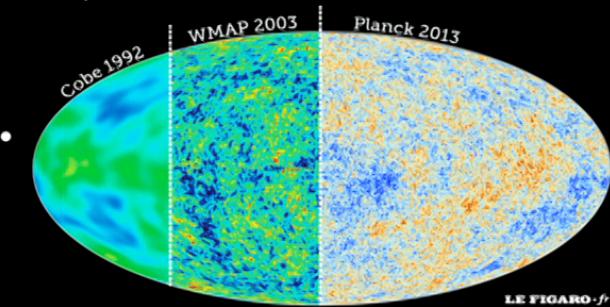
3

CMB Sky discovered in 1965

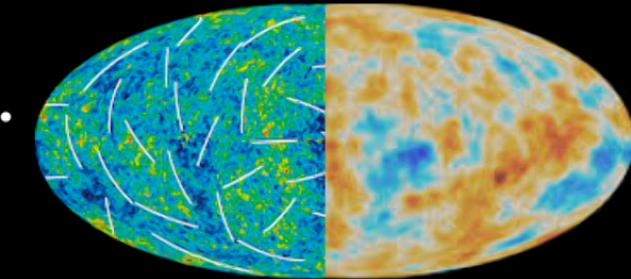


Blackbody with
spatial fluctuations

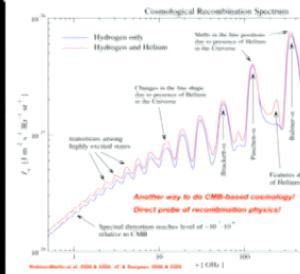
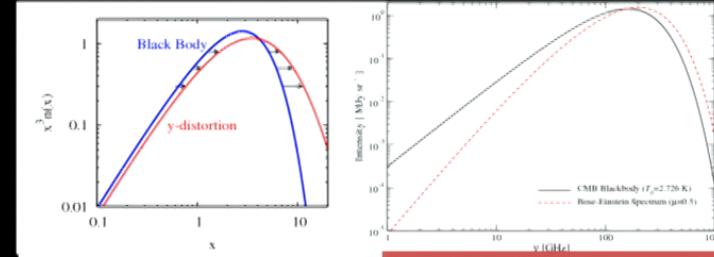
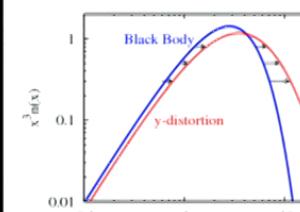
Temperature



Polarization



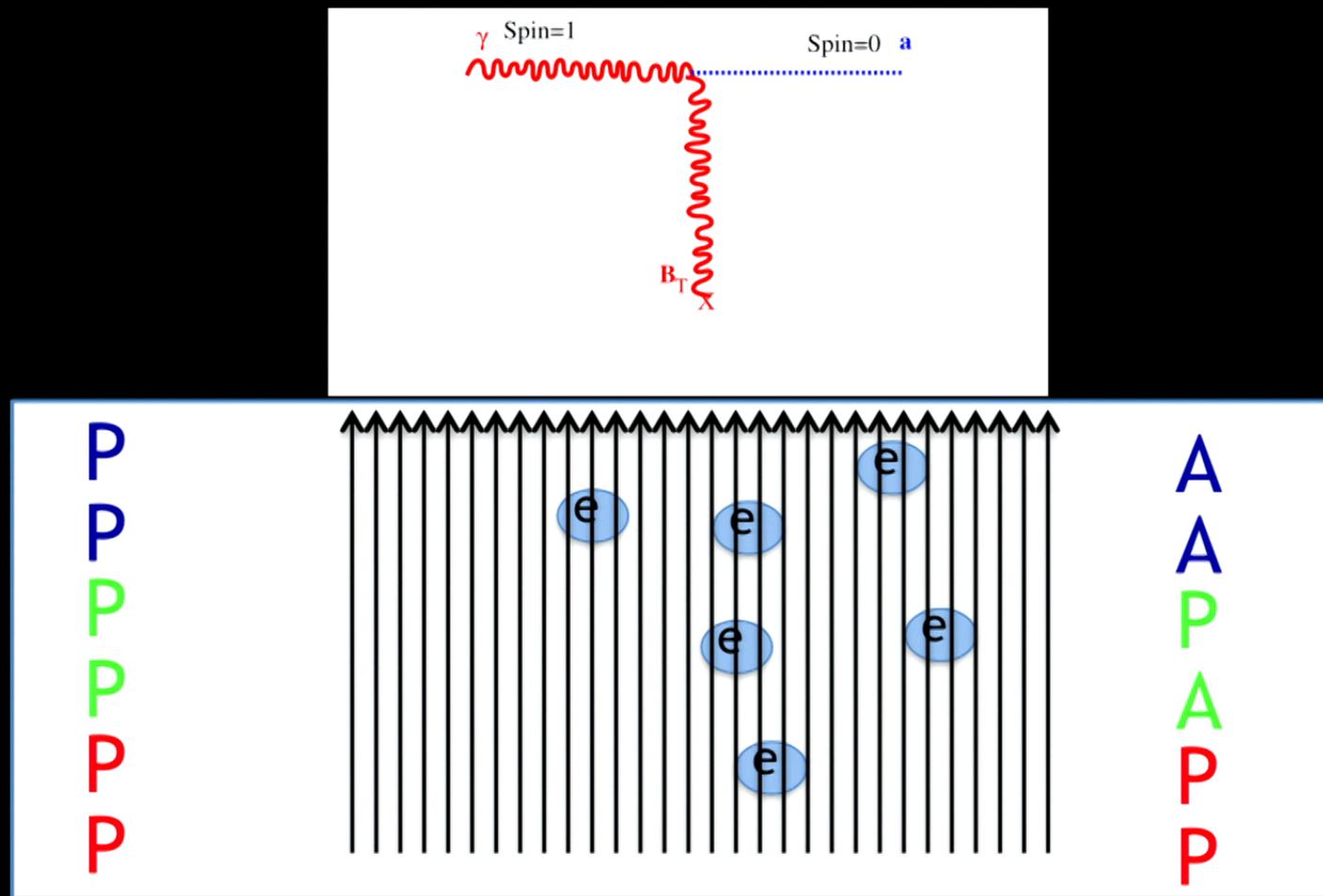
Non-Blackbody with
(or without)
spatial fluctuations



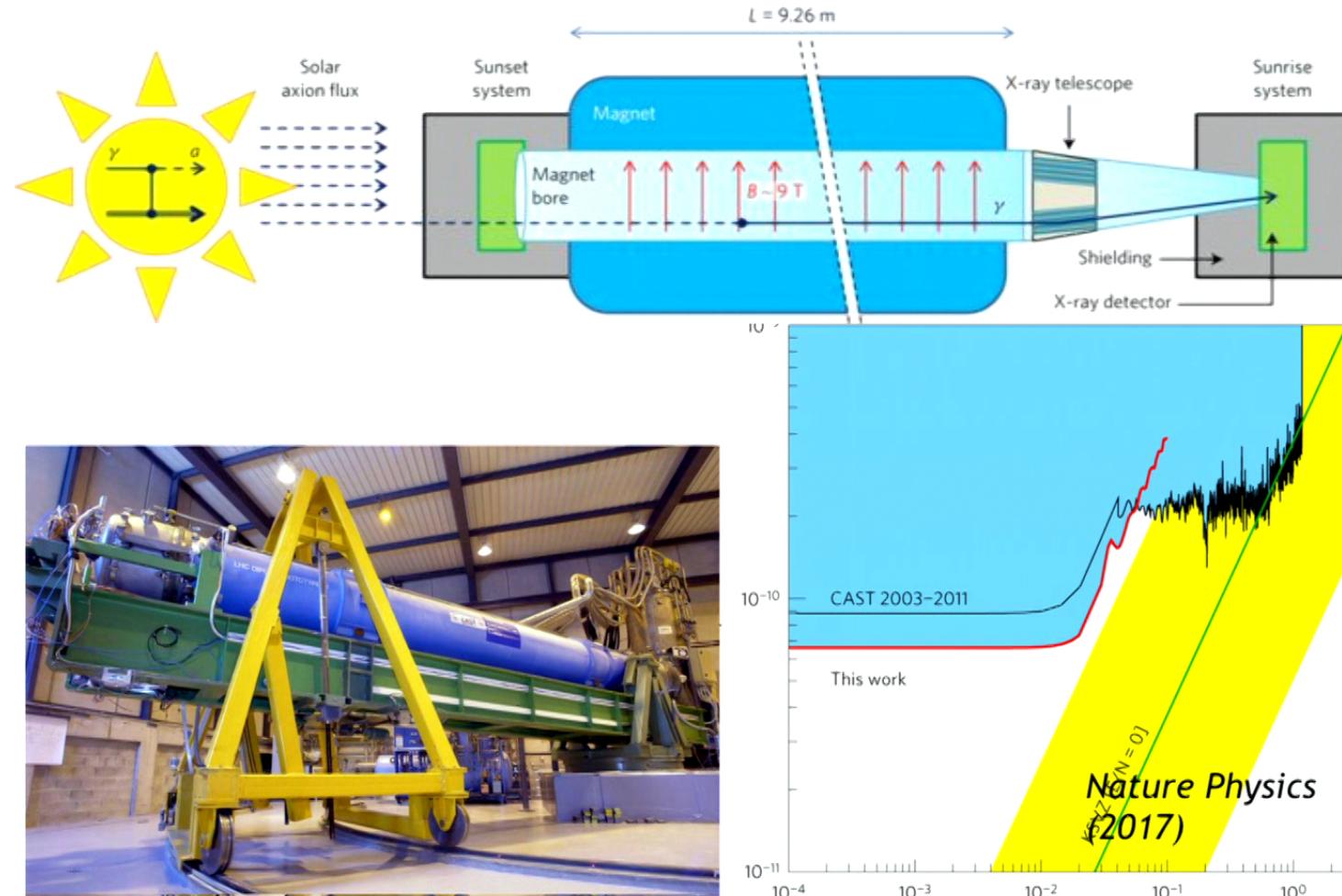
ALPs

Chluba 1405.6938,
Khatri & Sunyaev JCAP 09 (2012) 016
Hu & Silk Phys. Rev. Lett. 1993

Aspects of Spectral Distortions from ALPs



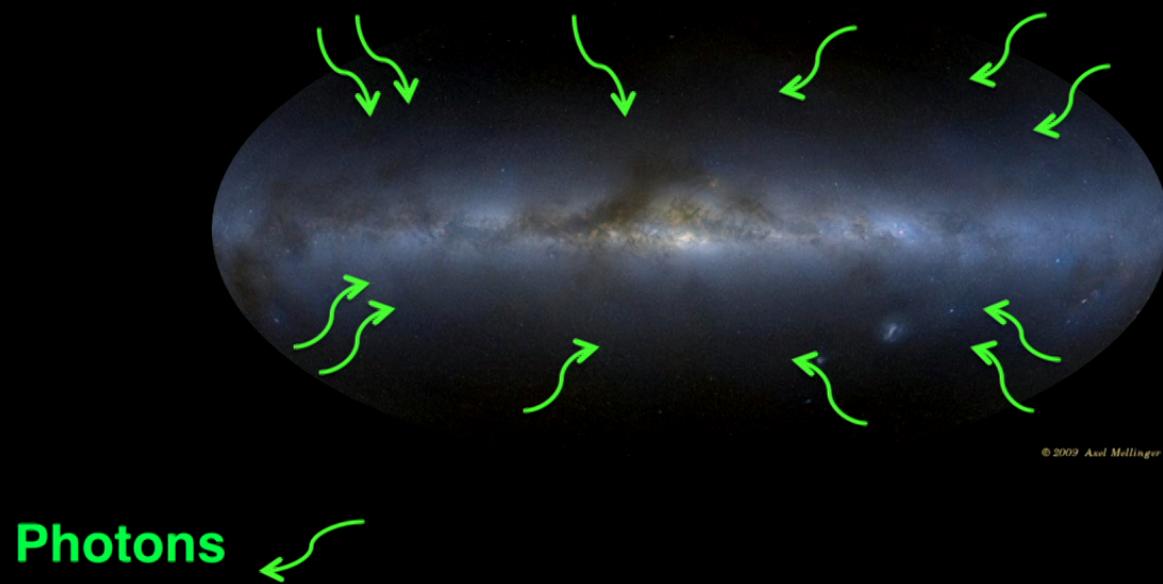
ALPs are also searched by particle physicists



Signature from the Milky Way

Mukherjee, Khatri, Wandelt
JCAP 04 (2018) 045

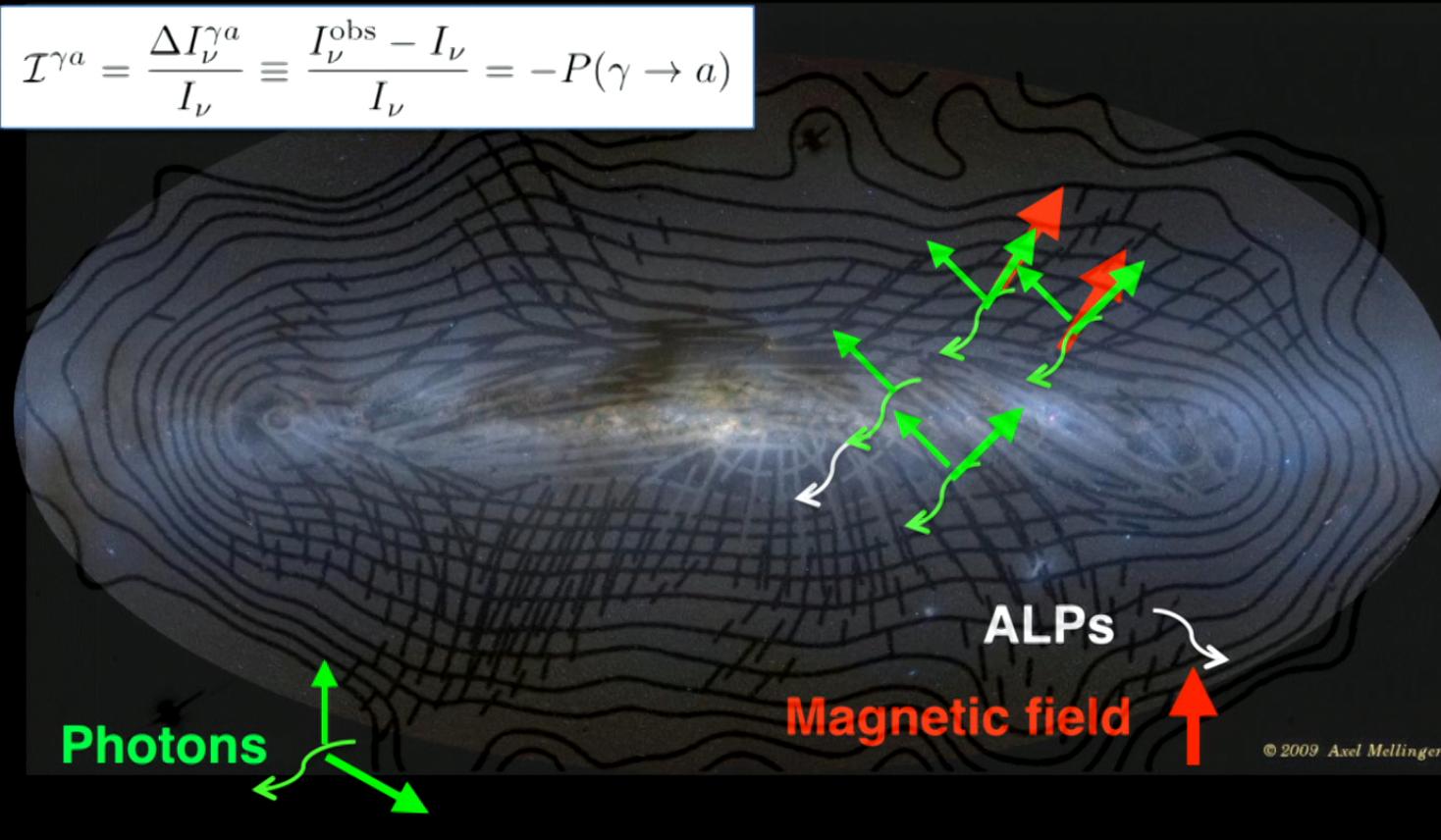
Milky Way scenario



© 2009 Axel Mellinger

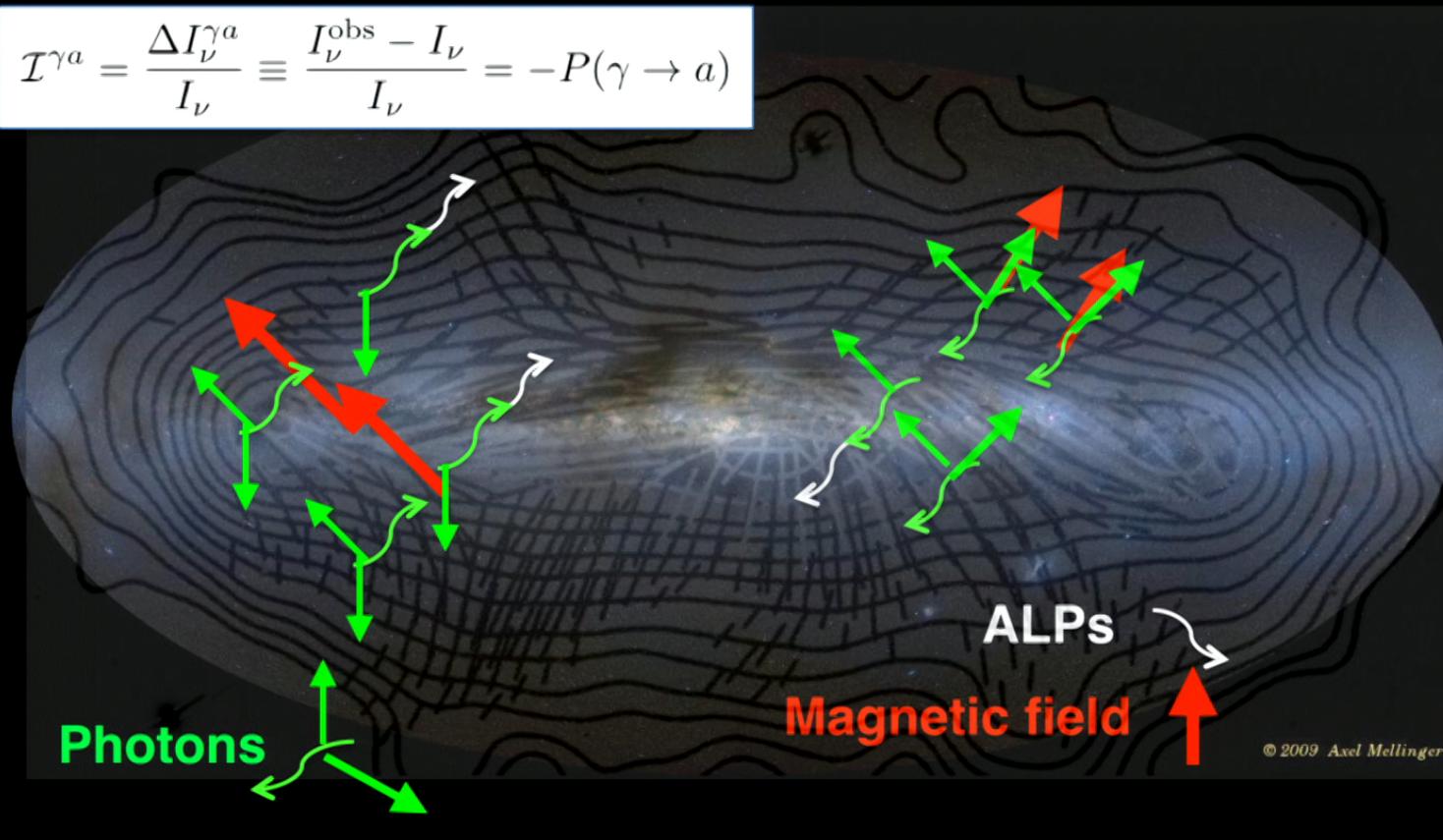
Photon to ALPs in Milky Way

$$\mathcal{I}^{\gamma a} = \frac{\Delta I_{\nu}^{\gamma a}}{I_{\nu}} \equiv \frac{I_{\nu}^{\text{obs}} - I_{\nu}}{I_{\nu}} = -P(\gamma \rightarrow a)$$



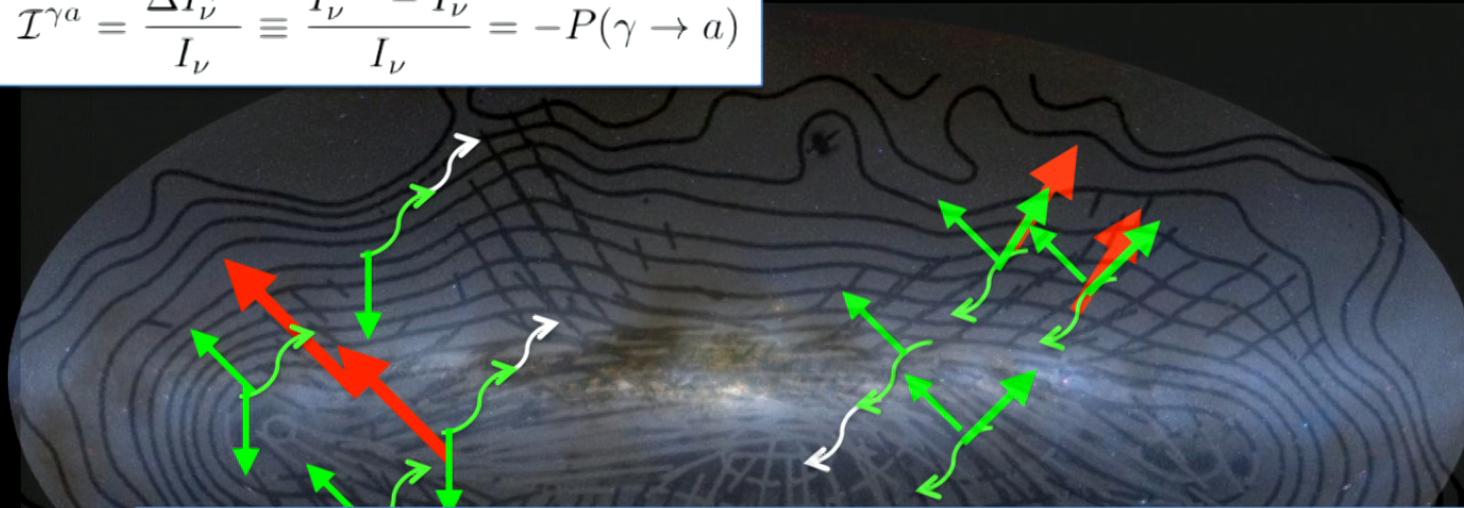
Photon to ALPs in Milky Way

$$\mathcal{I}^{\gamma a} = \frac{\Delta I_{\nu}^{\gamma a}}{I_{\nu}} \equiv \frac{I_{\nu}^{\text{obs}} - I_{\nu}}{I_{\nu}} = -P(\gamma \rightarrow a)$$



Photon to ALPs in Milky Way

$$\mathcal{I}^{\gamma a} = \frac{\Delta I_{\nu}^{\gamma a}}{I_{\nu}} \equiv \frac{I_{\nu}^{\text{obs}} - I_{\nu}}{I_{\nu}} = -P(\gamma \rightarrow a)$$



Photo

$$\left(\omega + \begin{pmatrix} \Delta_e & \Delta_f & \Delta_{\gamma a}^x \\ \Delta_f & \Delta_e & \Delta_{\gamma a}^y \\ \Delta_{\gamma a}^x & \Delta_{\gamma a}^y & \Delta_a \end{pmatrix} + i\partial_z \right) \begin{pmatrix} A_x \\ A_y \\ a \end{pmatrix} = 0,$$

$$\Delta_{\gamma a} \equiv \frac{g_{\gamma a}|B_T|}{2}$$

$$\Delta_a \equiv -\frac{m_a^2}{2\omega}$$

$$\Delta_e \equiv -\frac{m_\gamma^2}{2\omega}$$

Two kinds of conversion

Resonance conversion

Happen at places where ALPs
mass equals
photon mass in the plasma

Non-Resonance conversion

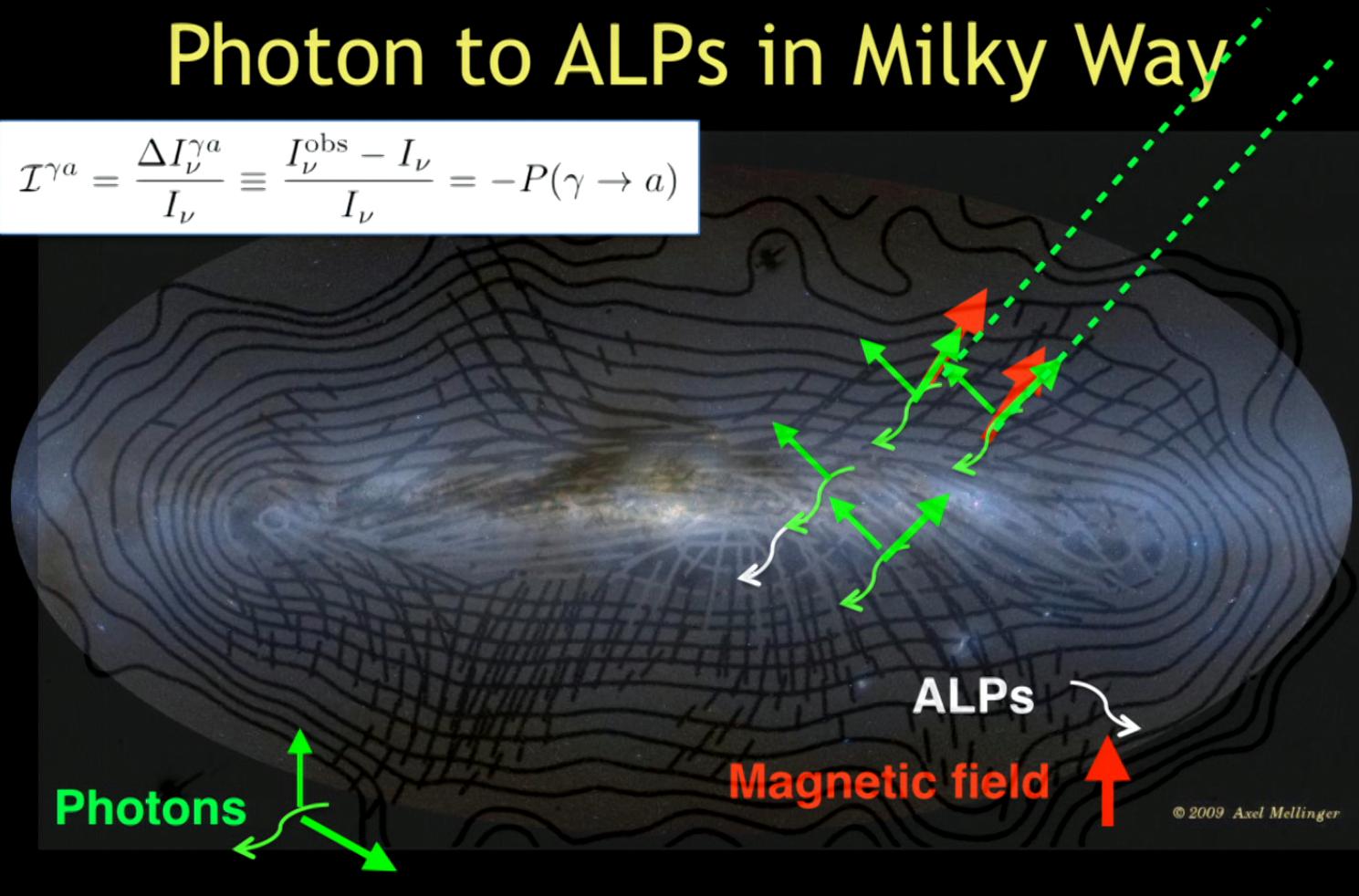
Happens throughout the
line of sight

Non-Resonant Conversion

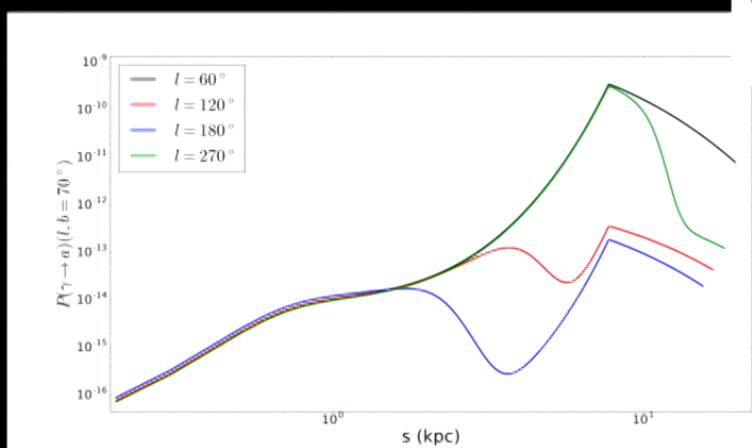
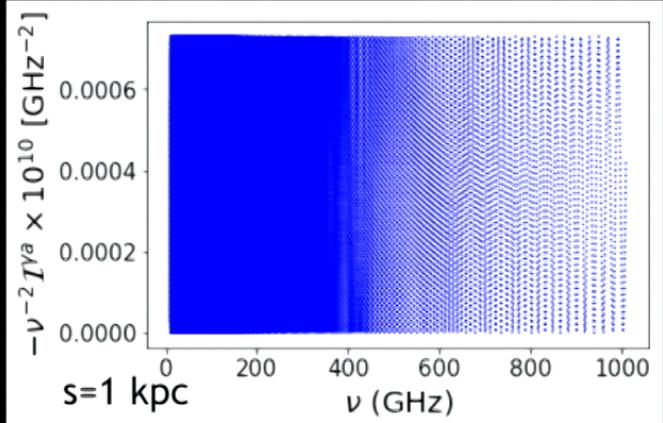
15

Photon to ALPs in Milky Way

$$\mathcal{I}^{\gamma a} = \frac{\Delta I_{\nu}^{\gamma a}}{I_{\nu}} \equiv \frac{I_{\nu}^{\text{obs}} - I_{\nu}}{I_{\nu}} = -P(\gamma \rightarrow a)$$



$$\mathcal{I}^{\gamma a} = \frac{\Delta I_{\nu}^{\gamma a}}{I_{\nu}} \equiv \frac{I_{\nu}^{\text{obs}} - I_{\nu}}{I_{\nu}} = -P(\gamma \rightarrow a)$$



$$\left(\omega + \begin{pmatrix} \Delta_e & \Delta_f & \Delta_{\gamma a}^x \\ \Delta_f & \Delta_e & \Delta_{\gamma a}^y \\ \Delta_{\gamma a}^x & \Delta_{\gamma a}^y & \Delta_a \end{pmatrix} + i\partial_z \right) \begin{pmatrix} A_x \\ A_y \\ a \end{pmatrix} = 0,$$

$$\begin{aligned} \left(\frac{\Delta_{\gamma a}^{x,y}}{\text{Mpc}^{-1}} \right) &\equiv \frac{g_{\gamma a}|B_{x,y}|}{2} = 15.2 \left(\frac{g_{\gamma a}}{10^{-11}\text{Gev}^{-1}} \right) \left(\frac{B_{x,y}}{\mu\text{G}} \right), \\ \left(\frac{\Delta_a}{\text{Mpc}^{-1}} \right) &\equiv -\frac{m_a^2}{2\nu} = -1.9 \times 10^4 \left(\frac{m_a}{10^{-14}\text{eV}} \right) \left(\frac{100 \text{ GHz}}{\nu} \right), \\ \left(\frac{\Delta_e}{\text{Mpc}^{-1}} \right) &\approx \frac{\omega_p^2}{2\nu} \left[-1 + 7.3 \times 10^{-3} \frac{n_H}{n_e} \left(\frac{\omega}{\text{eV}} \right)^2 \right] \\ &= -2.6 \times 10^6 \left(\frac{n_e}{10^{-5}\text{cm}^{-3}} \right) \left(\frac{100 \text{ GHz}}{\nu} \right) \times \left[1 - 7.3 \times 10^{-3} \frac{n_H}{n_e} \left(\frac{\omega}{\text{eV}} \right)^2 \right] \end{aligned}$$

Electron density model :
J.M.Cordes & T.J.W.Lazio
arXiv: astro-ph/0207156
Gaensler et al.
arXiv0808.2550[astro-ph]

Magnetic field model:
Jansson & Farrar
Astrophys.J. 757 (2012) 14

Non-Resonant conversion from the random magnetic field

$$\bar{P}(\gamma \rightarrow a)(r) = \frac{1}{3} \left(1 - e^{(-3P(\gamma \rightarrow a)r/2d_0)} \right) \quad r \gg d_0$$

A. Mirizzi, G. G. Raffelt, and P. D. Serpico
Phys. Rev. D, 72(2):023501
Y. Grossman, S. Roy, and J. Zupan.
Physics Letters B, 543:23-28

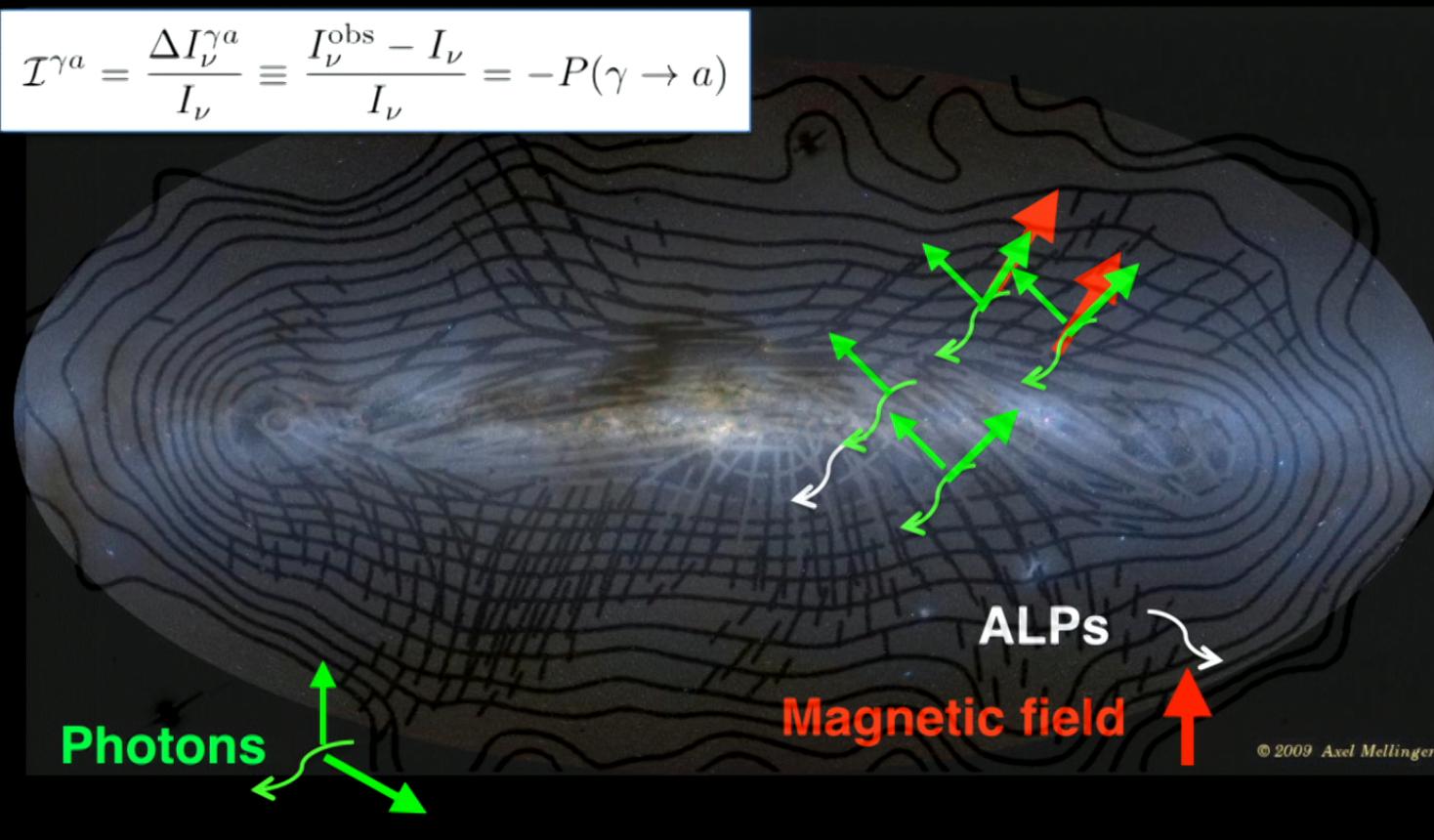
$$\begin{aligned} \bar{P}(\gamma \rightarrow a) &\approx \frac{P(\gamma \rightarrow a)R}{2s} \\ &\approx \frac{\Delta_{\gamma a}^2 R s}{2} = 10^{-9} \left(\frac{g_{\gamma a}}{10^{-10} \text{ GeV}^{-1}} \right)^2 \left(\frac{B_T}{1 \mu\text{G}} \right)^2 \left(\frac{R}{1000 \text{ pc}} \right) \left(\frac{s}{10^{-4} \text{ pc}} \right) \end{aligned}$$

Resonant Conversion

19

Photon to ALPs in Milky Way

$$\mathcal{I}^{\gamma a} = \frac{\Delta I_{\nu}^{\gamma a}}{I_{\nu}} \equiv \frac{I_{\nu}^{\text{obs}} - I_{\nu}}{I_{\nu}} = -P(\gamma \rightarrow a)$$

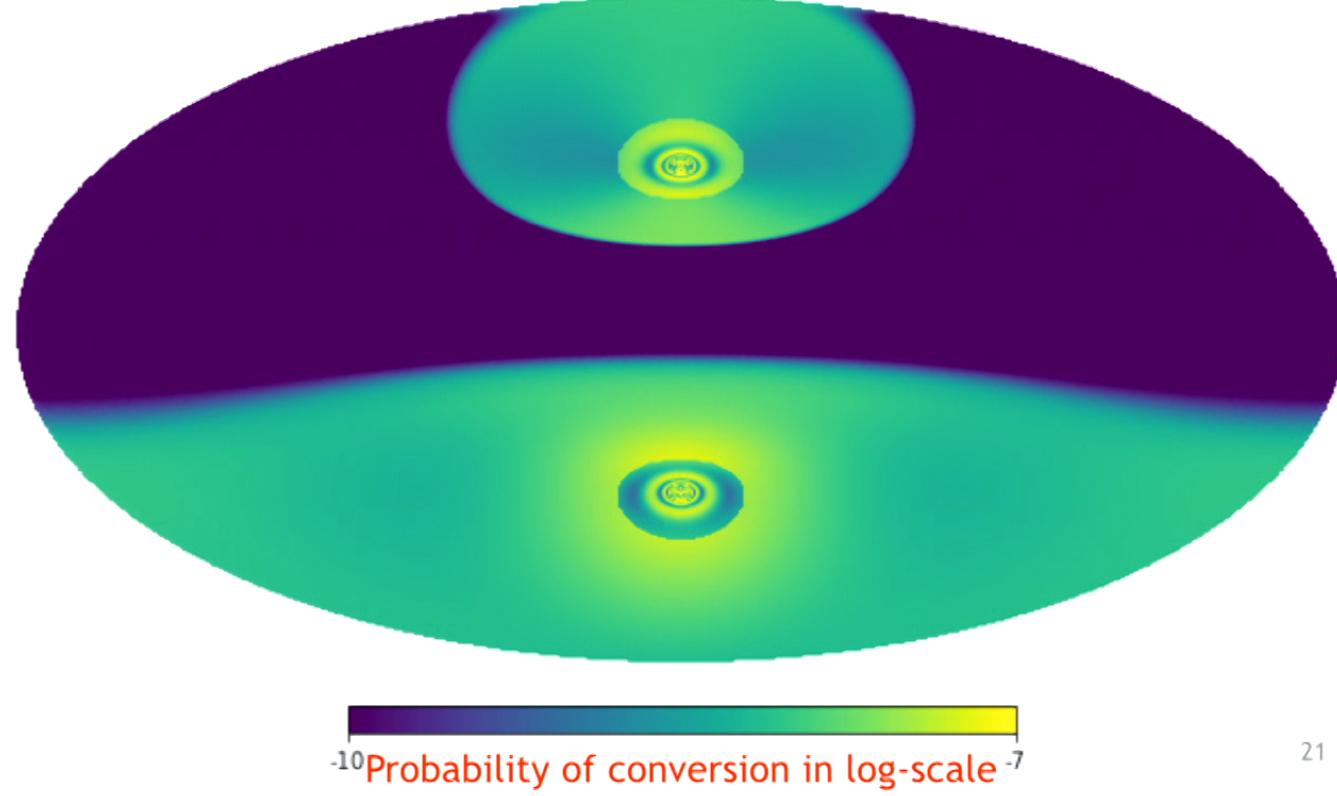


Resonant Conversion

ALPs mass: 5×10^{-13} eV

Solved for 3 Million sky pixels at 150 GHz

Mukherjee, Khatri, Wandelt
JCAP 04 (2018) 045



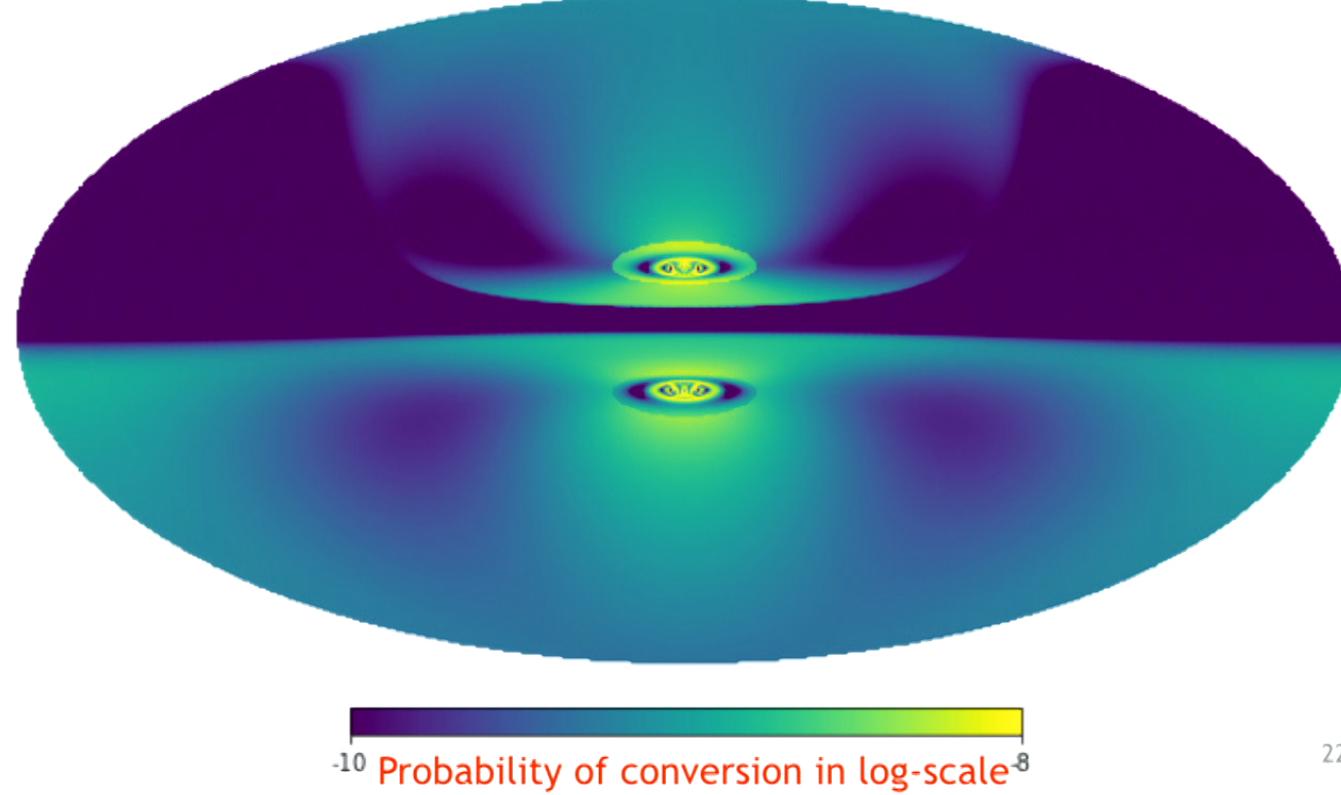
21

Resonant Conversion

ALPs mass: 5×10^{-12} eV

Mukherjee, Khatri, Wandelt
JCAP 04 (2018) 045

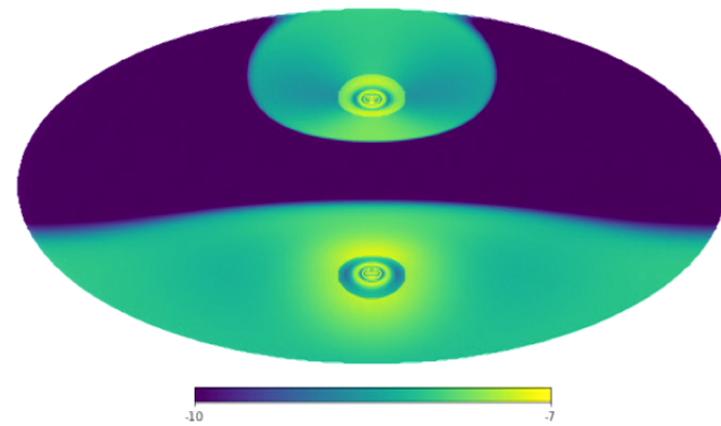
Solved for 3 Million sky pixels at 150 GHz



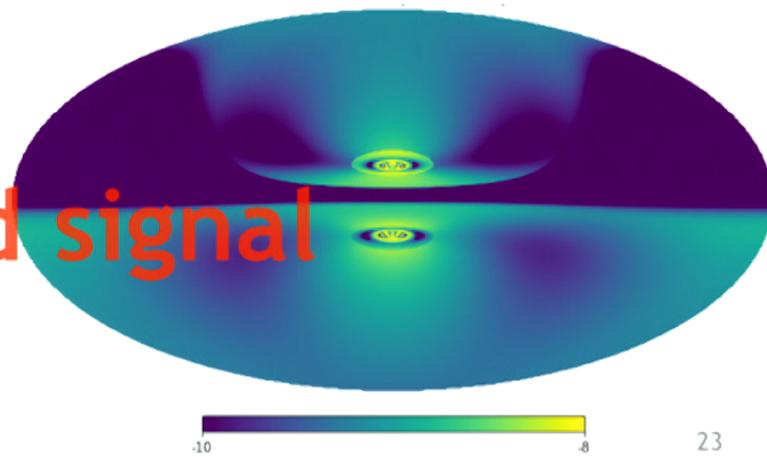
22

Resonant Conversion

ALPs mass: 5×10^{-13} eV



ALPs mass: 5×10^{-12} eV

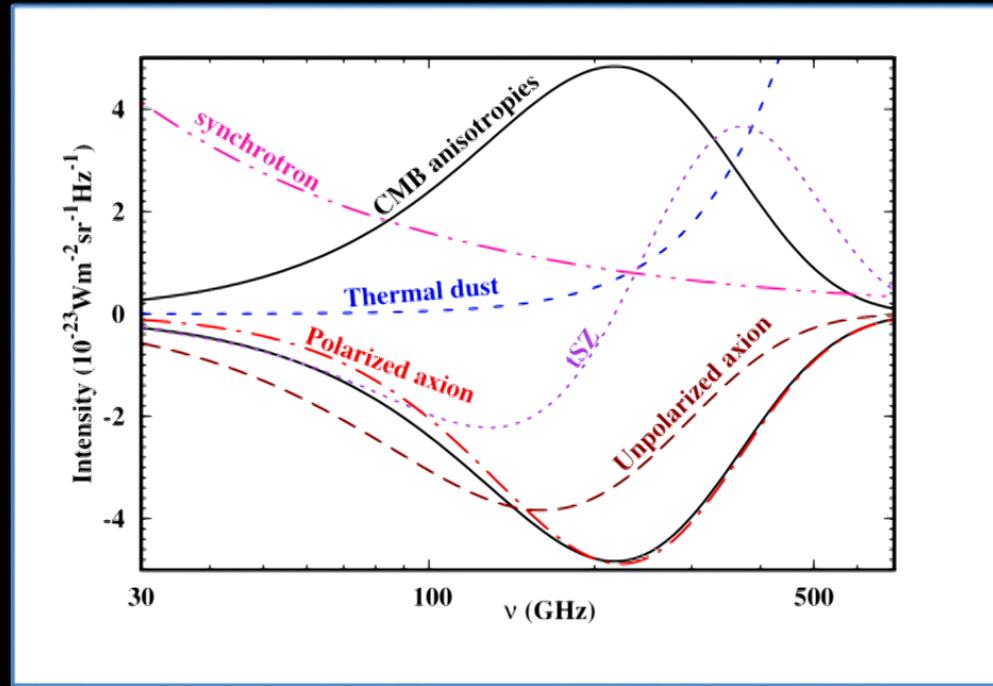


100 % Polarized signal

Spatially fluctuating
a unique structure

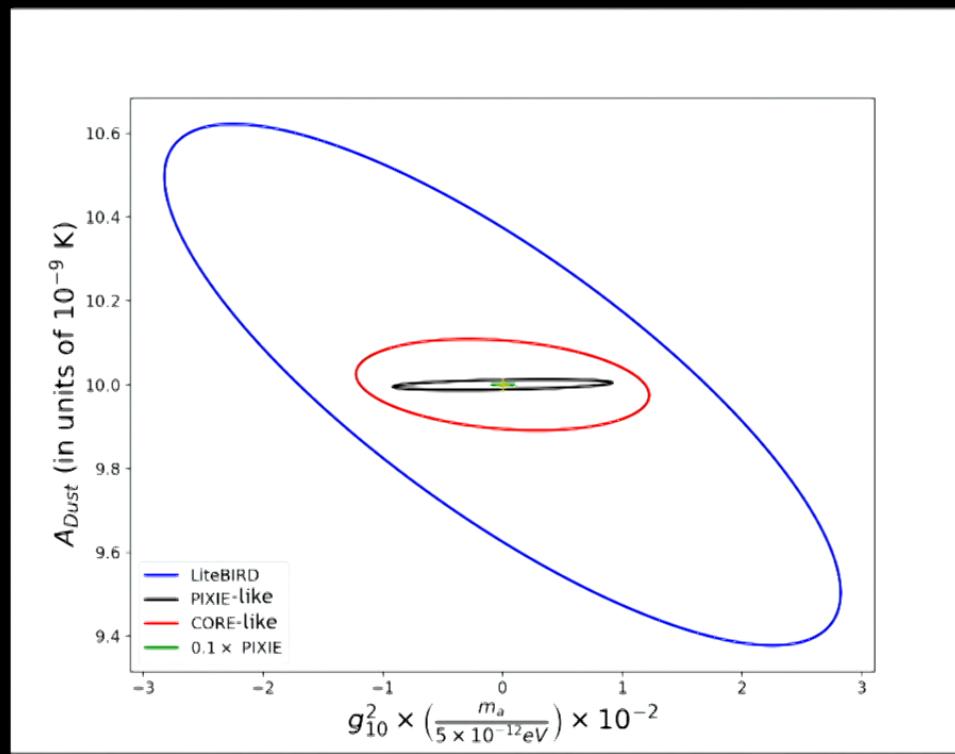
What about other signals (contaminations)

Two new signals of spectral distortions



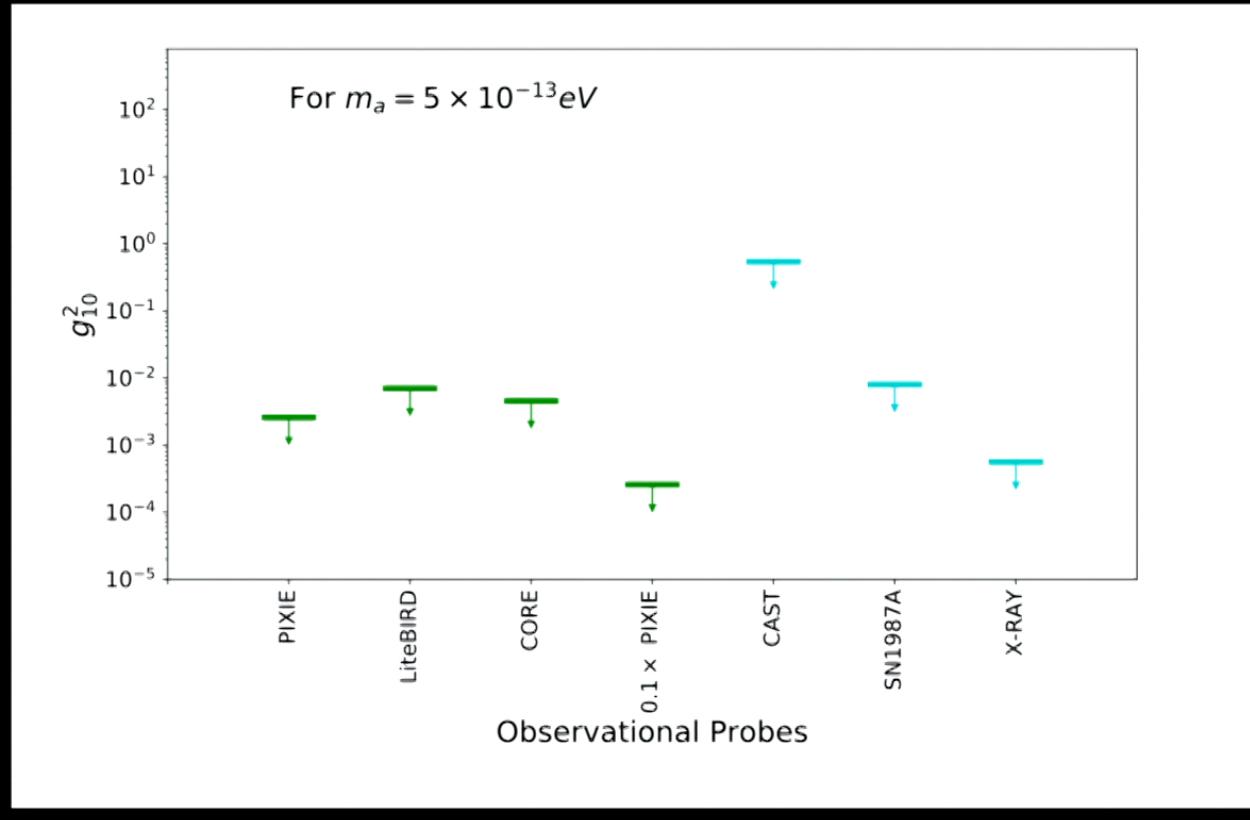
25

Resonant case



26

Resonant case

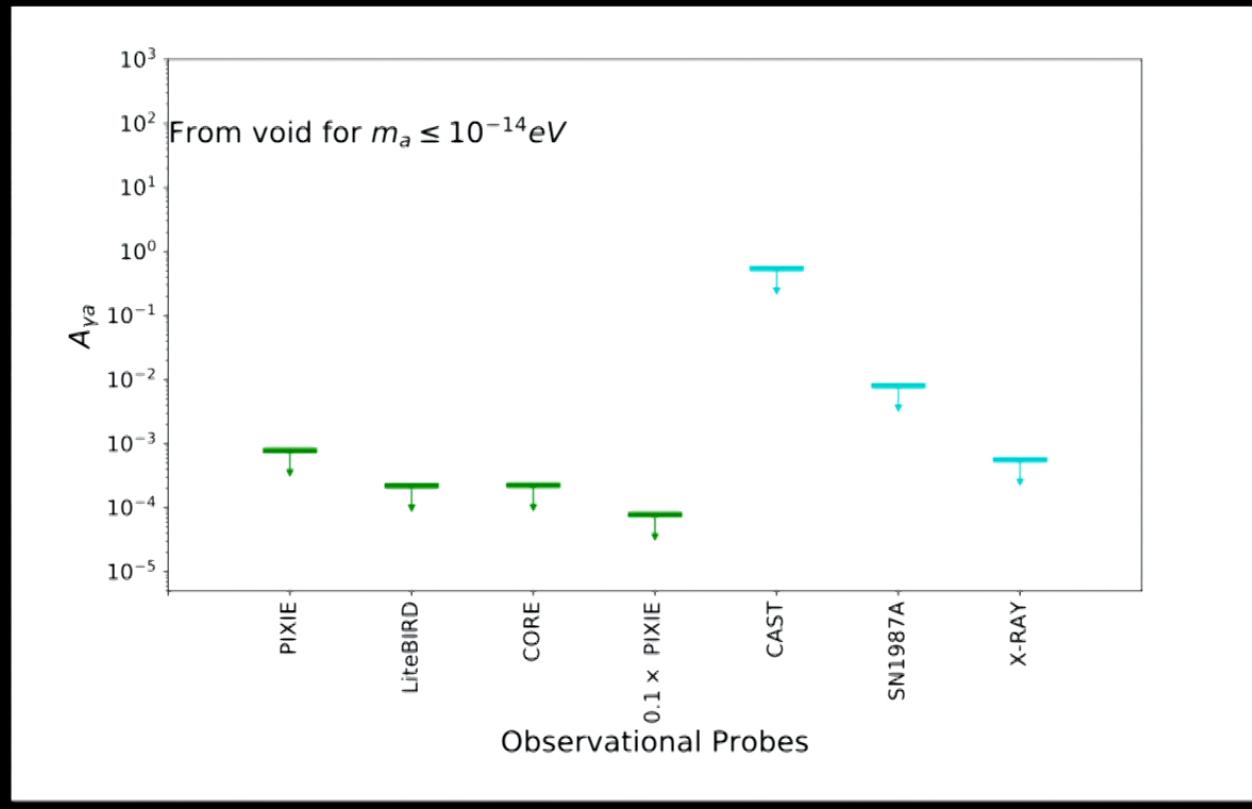


27

Extragalactic sources: Voids

Mukherjee, Khatri, Wandelt
[JCAP 04 \(2018\) 045](#)

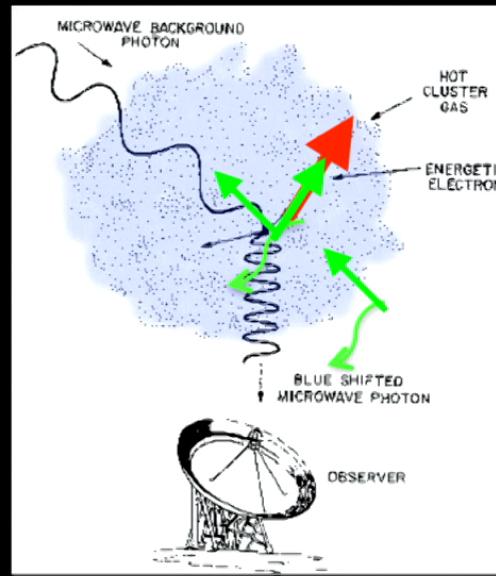
Forecast for future CMB mission after marginalizing other contaminations



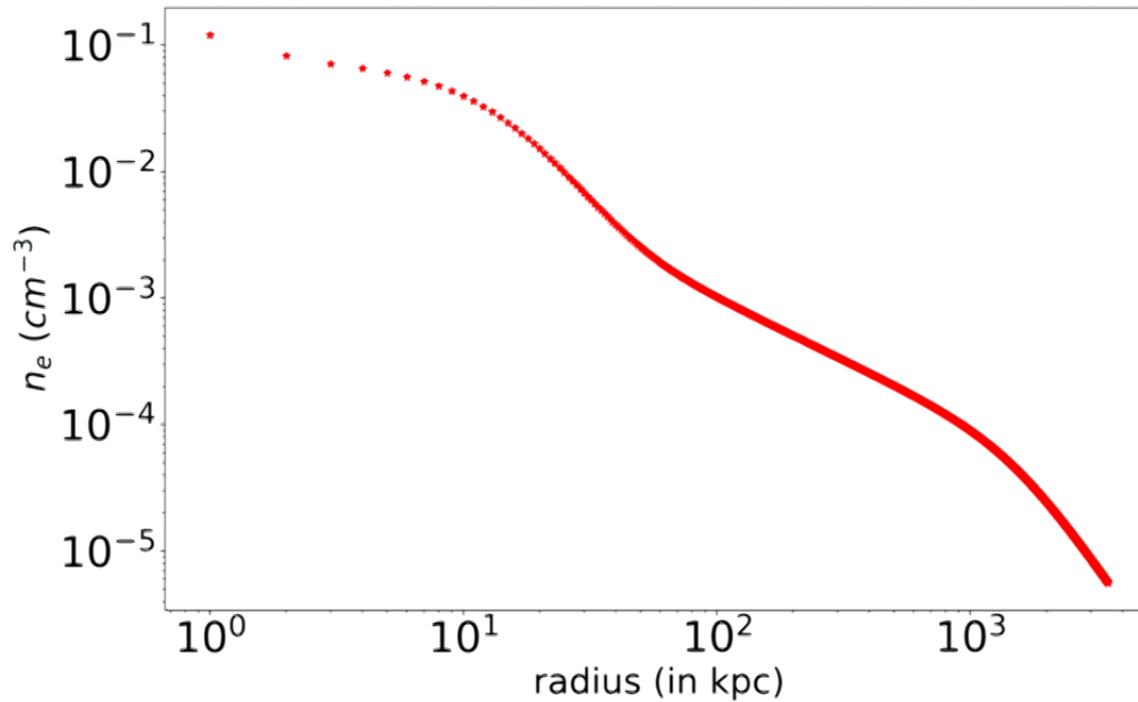
29

ALPs constraints from Simons Observatory and CMB-S4 using Galaxy Clusters

Mukherjee, Spergel, Khatri, Wandelt.
To be submitted



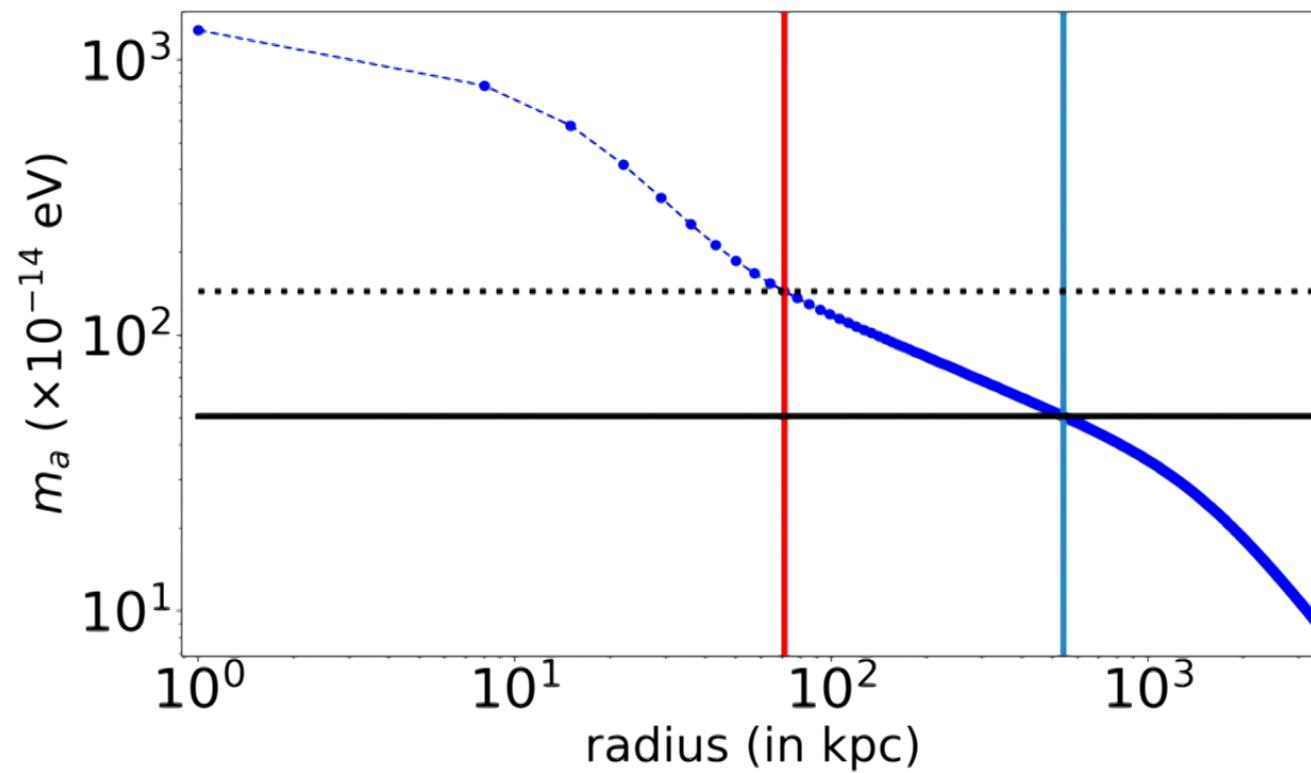
Electron density in galaxy clusters



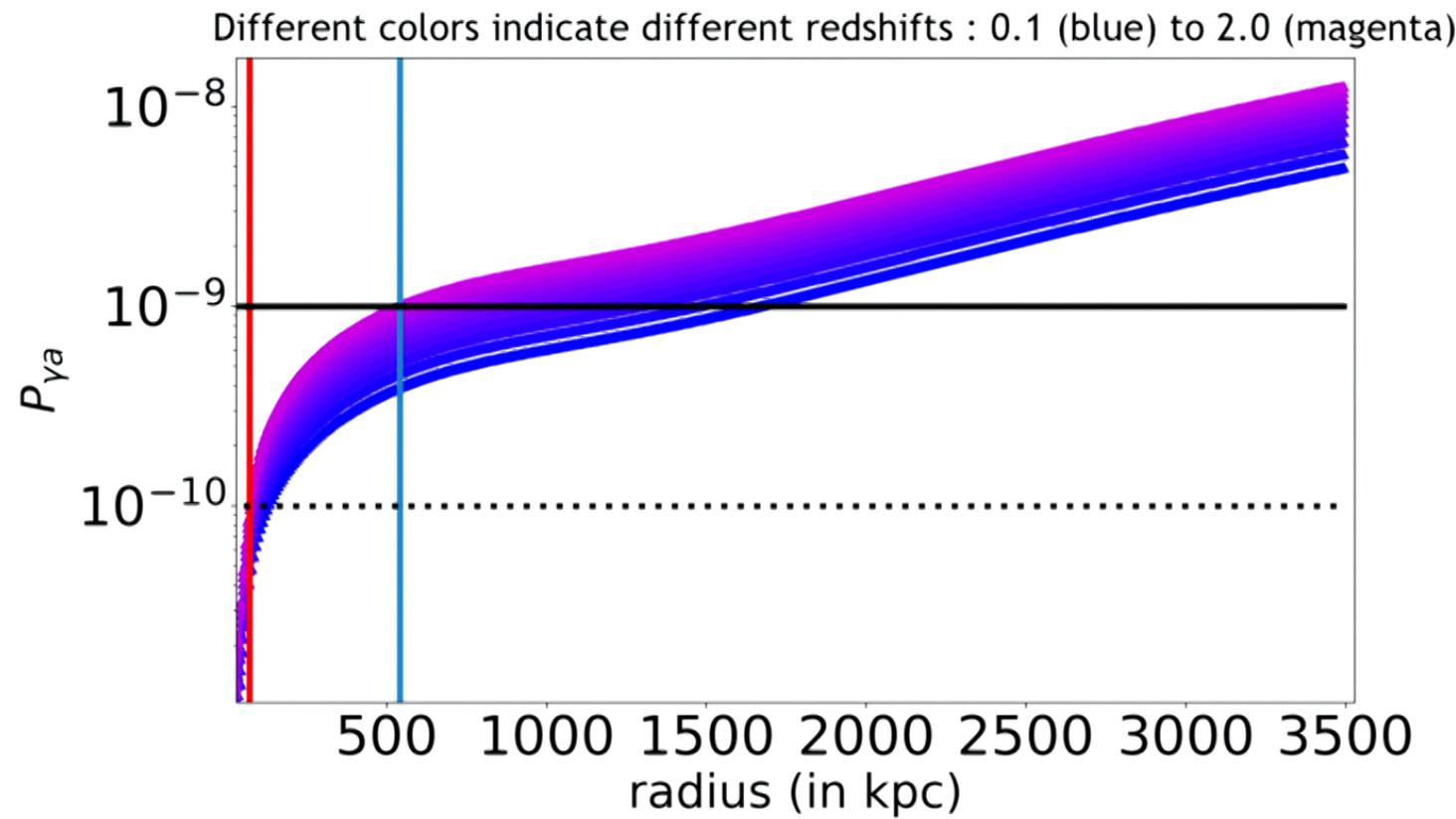
Vikhlinin et al. 2006,
Bartalucci et al. 2017

31

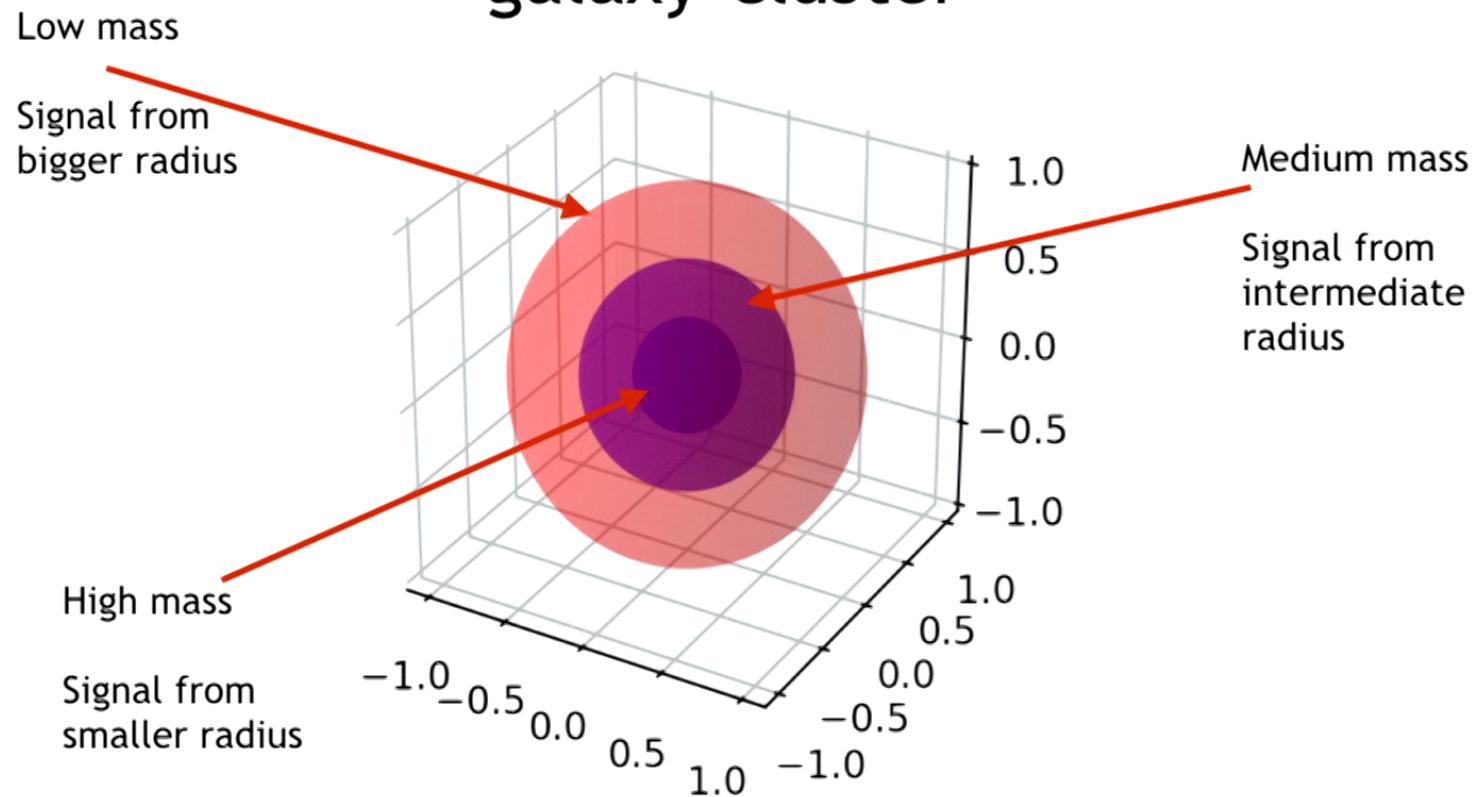
which ALP mass at which radius?



Signal profile with radius in galaxy cluster



Spatial shape of the distortion around a galaxy-cluster



34

A simplistic simulated
sky map of ALPs signal
around galaxy cluster

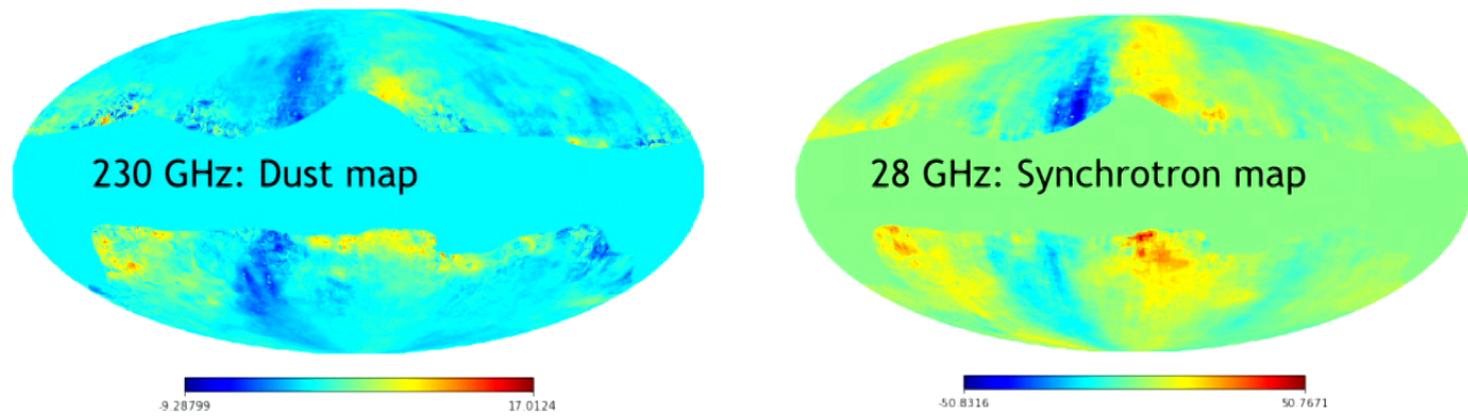
35

Model of the signal in the Sky

$$S_{\nu_i}(\hat{p}) = A_{\nu_i j} x_j(\hat{p}) + n_{\nu_i}(\hat{p}),$$

$$\mathbf{S}(\hat{p}) = \mathbf{A}\mathbf{x}(\hat{p}) + \mathbf{n}(\hat{p}),$$

We added simulated maps of synchrotron, dust, CMB over the ALPs signal at randomly selected cluster locations.

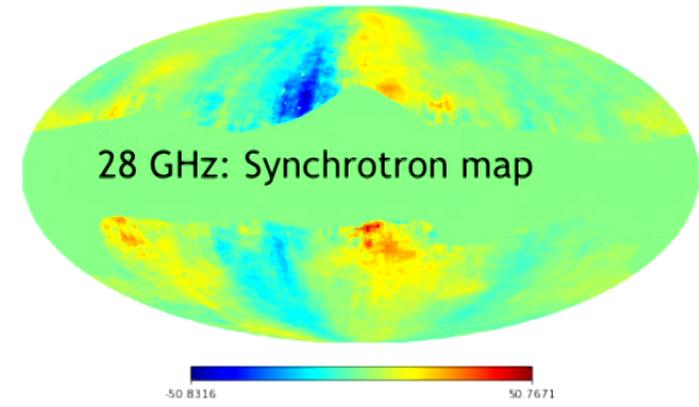
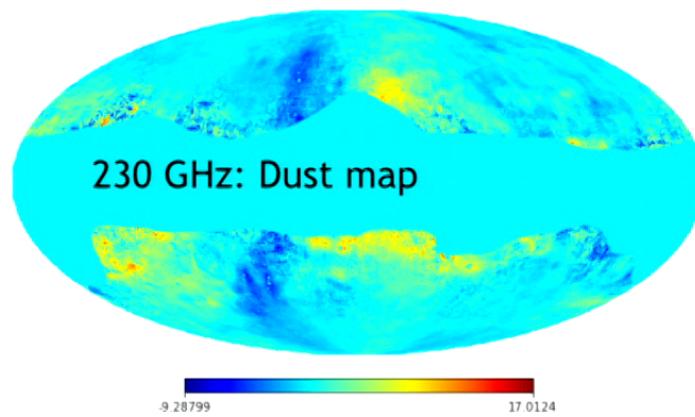


Python Sky Model (PYSM) code
Thorne et al. (2018)

Model of the signal in the Sky

$$S_{\nu_i}(\hat{p}) = A_{\nu_i j} x_j(\hat{p}) + n_{\nu_i}(\hat{p}),$$
$$\mathbf{S}(\hat{p}) = \mathbf{A}\mathbf{x}(\hat{p}) + \mathbf{n}(\hat{p}),$$

We added simulated maps of synchrotron, dust, CMB over the ALPs signal at randomly selected cluster locations.

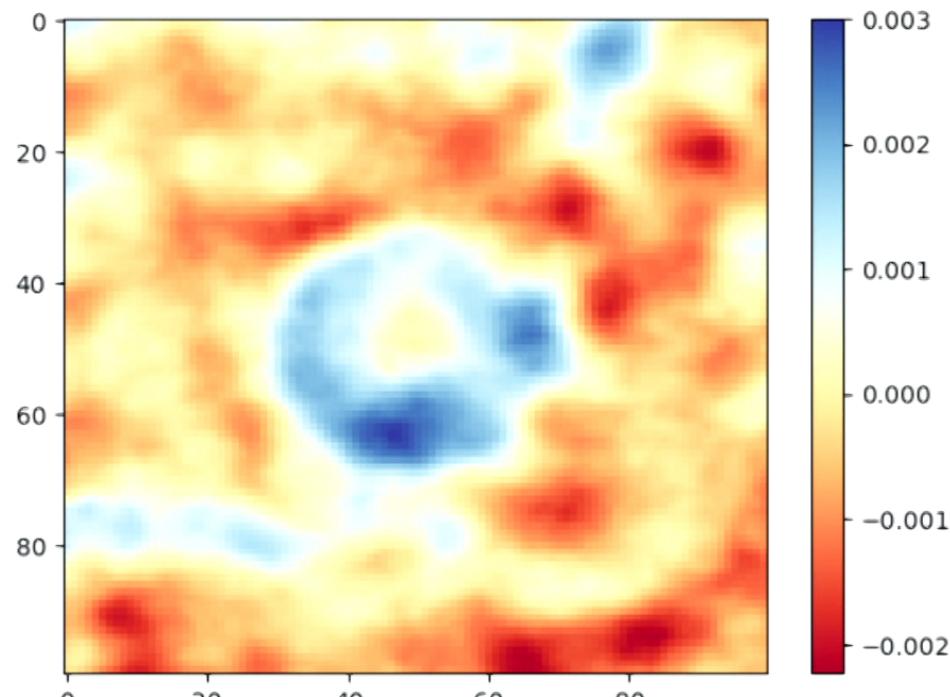


Then we clean the foregrounds using component separation method:
Internal Linear Combination (ILC)

$$\hat{x}_{\gamma a}(\hat{p}) = \hat{\mathcal{W}}_{\gamma a}^T \mathbf{S}(\hat{p}),$$

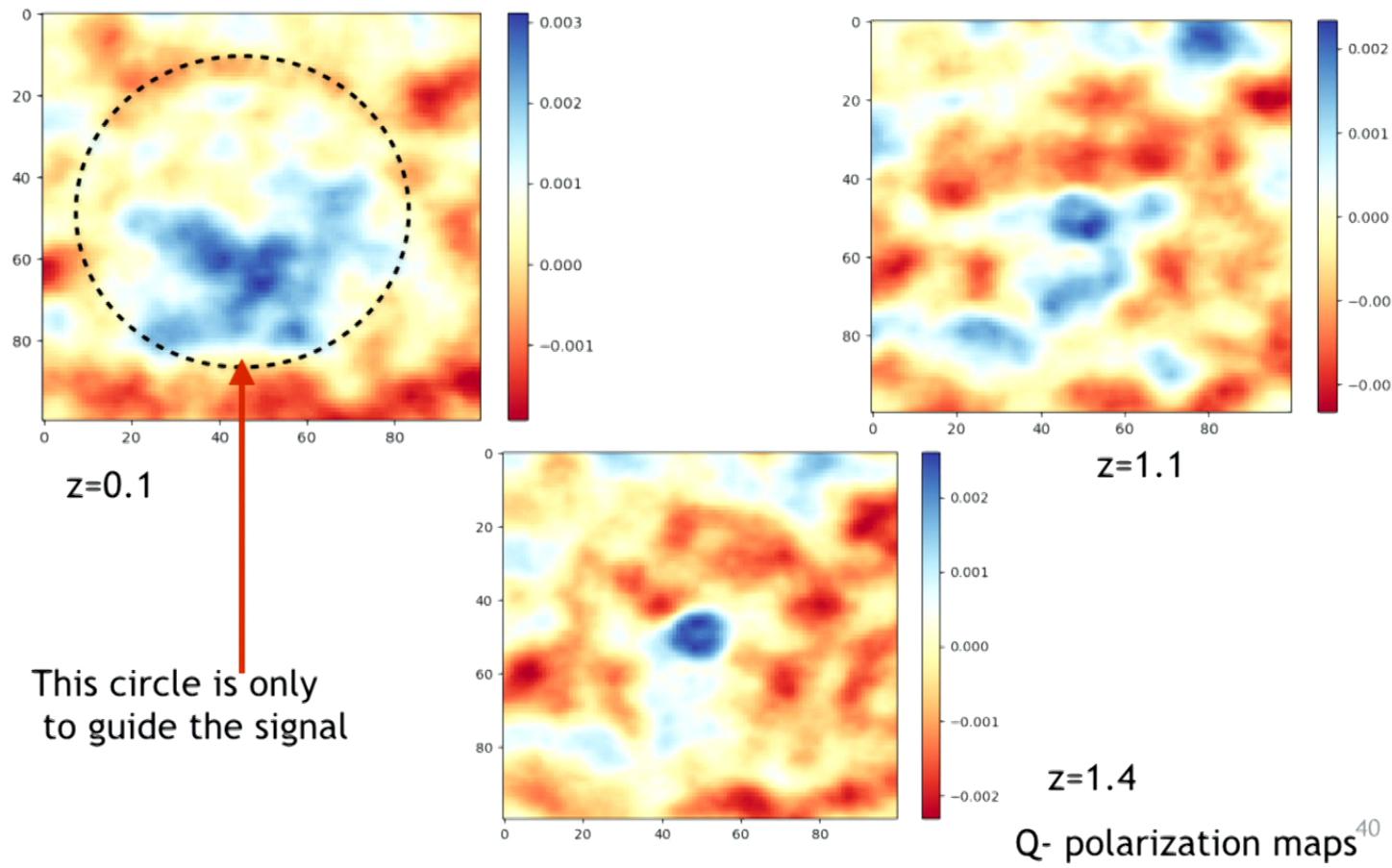
$$\hat{\mathbf{W}}_{\gamma a}(\nu) = \mathbf{C}_S^{-1} \mathbf{a}_{\gamma a} (\mathbf{a}_{\gamma a}^T \mathbf{C}_S^{-1} \mathbf{a}_{\gamma a})^{-1}$$

Adding 1000 clusters



Q-polarization map₃₈

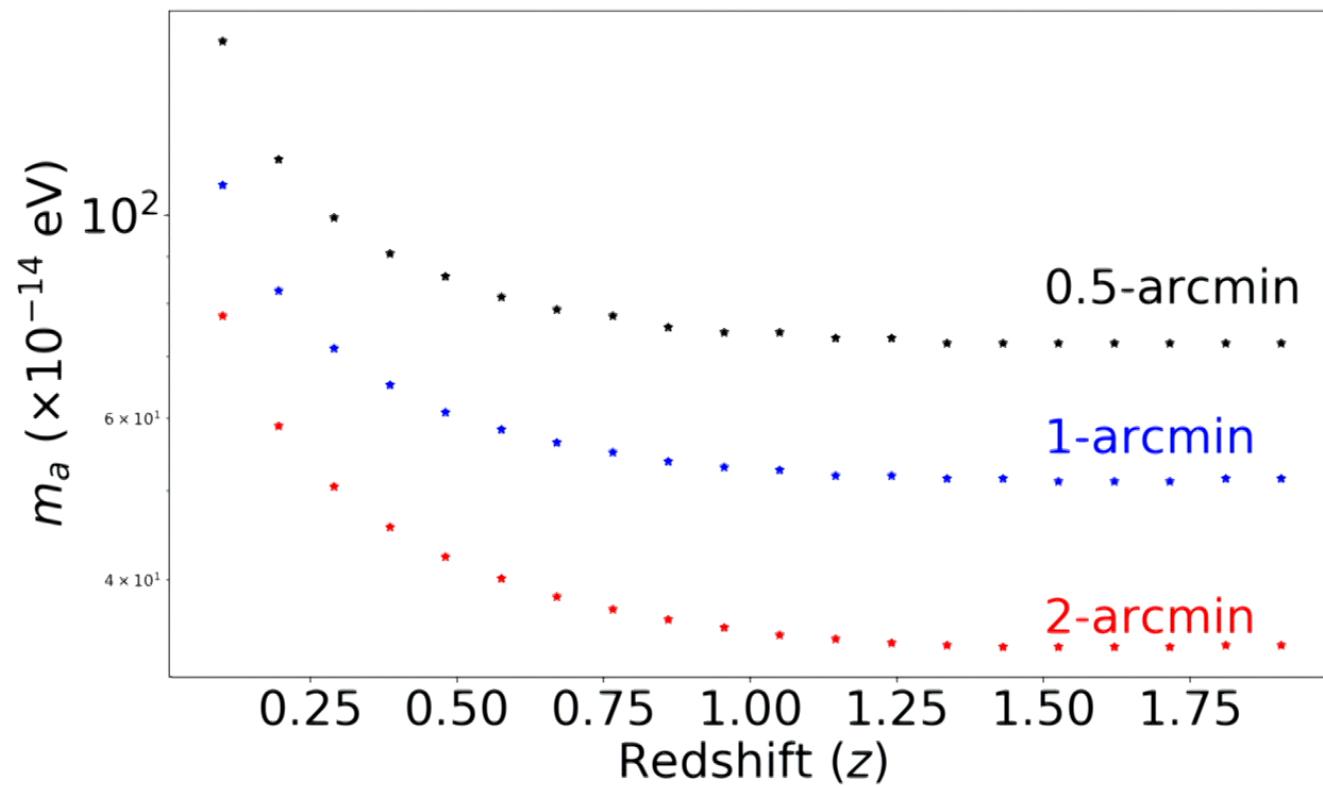
At different redshifts



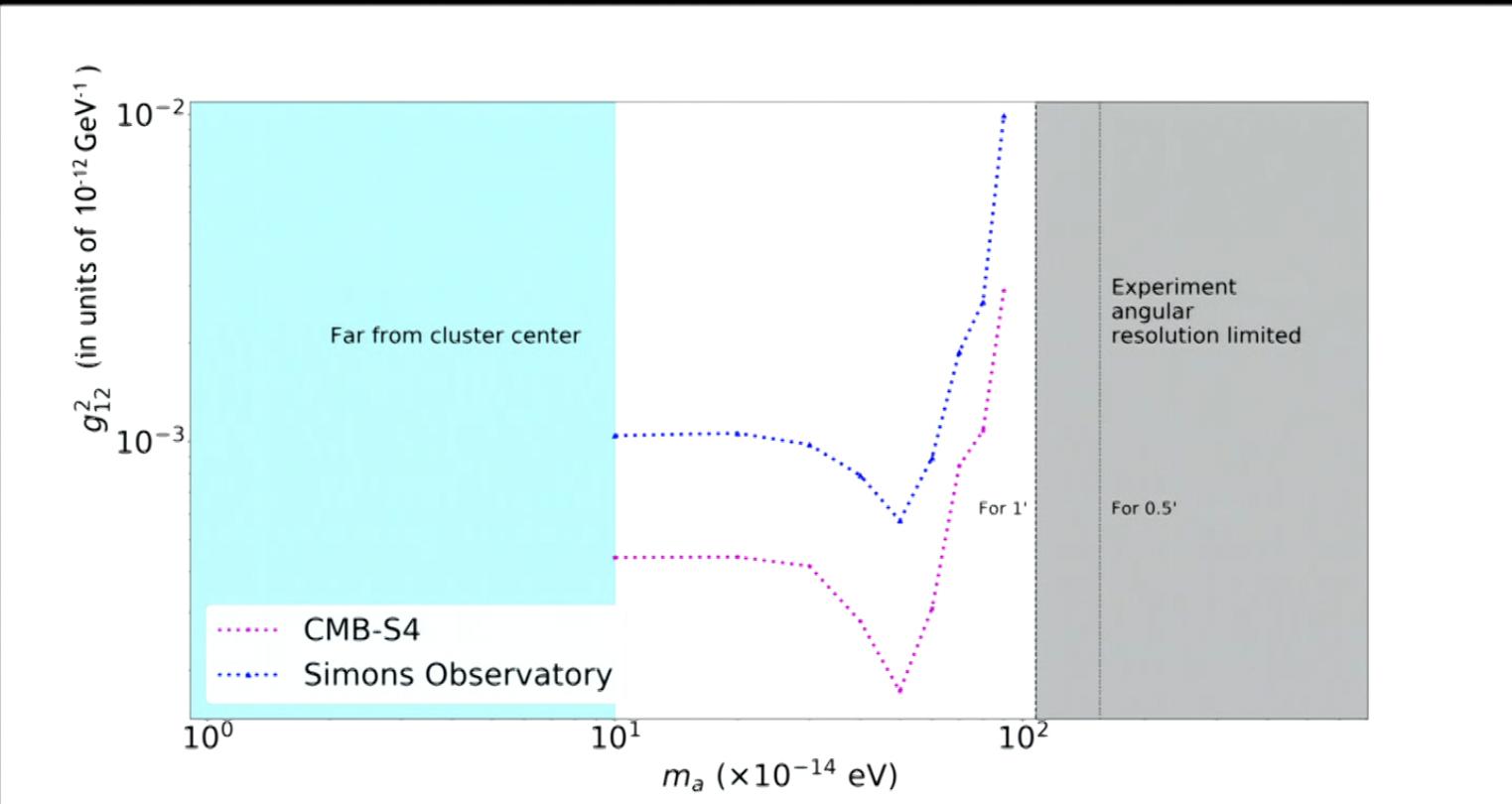
Detectability of the signal

41

ALP mass spatially resolved

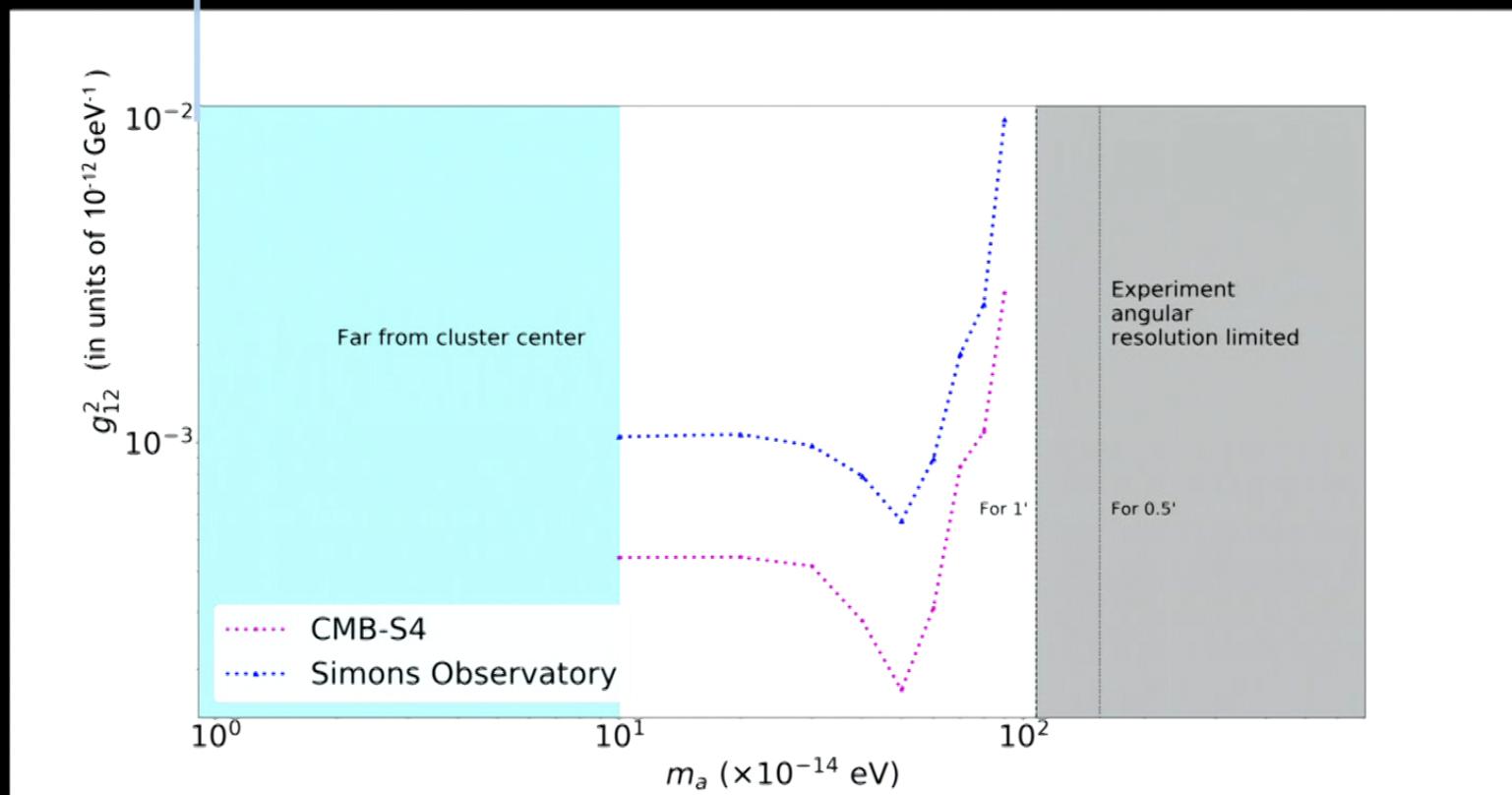


Bound achievable from CMB Ground based experiments (after including foregrounds, CMB, instrument noise)



4.4 X 10³ Current Bound from CERN Axion Solar Telescope (CAST)

Bound achievable from CMB Ground based experiments
(after including foregrounds, CMB, instrument noise)



ALPs distortion map

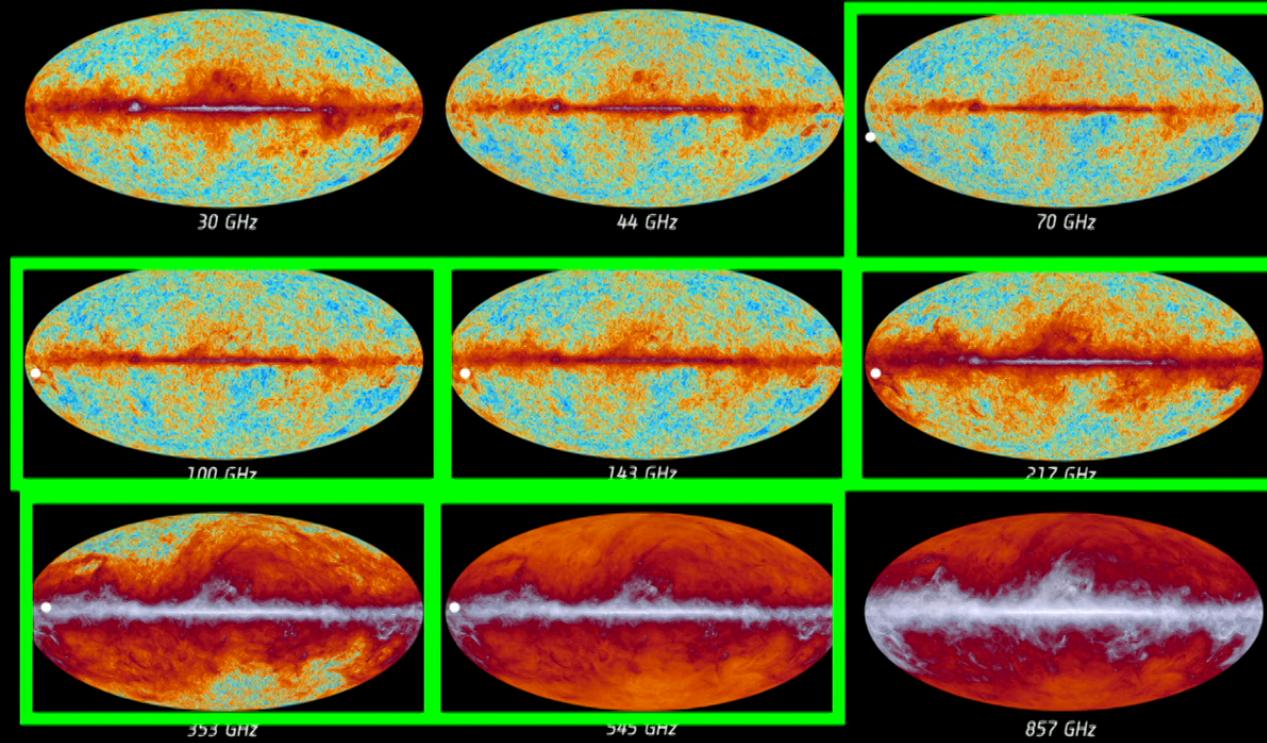
Mukherjee, Khatri, Wandelt

arXiv:1811.11177

Using the Planck temperature data
(For the non-resonance case)

From Polarization data: In preparation
(For the resonance case)

45

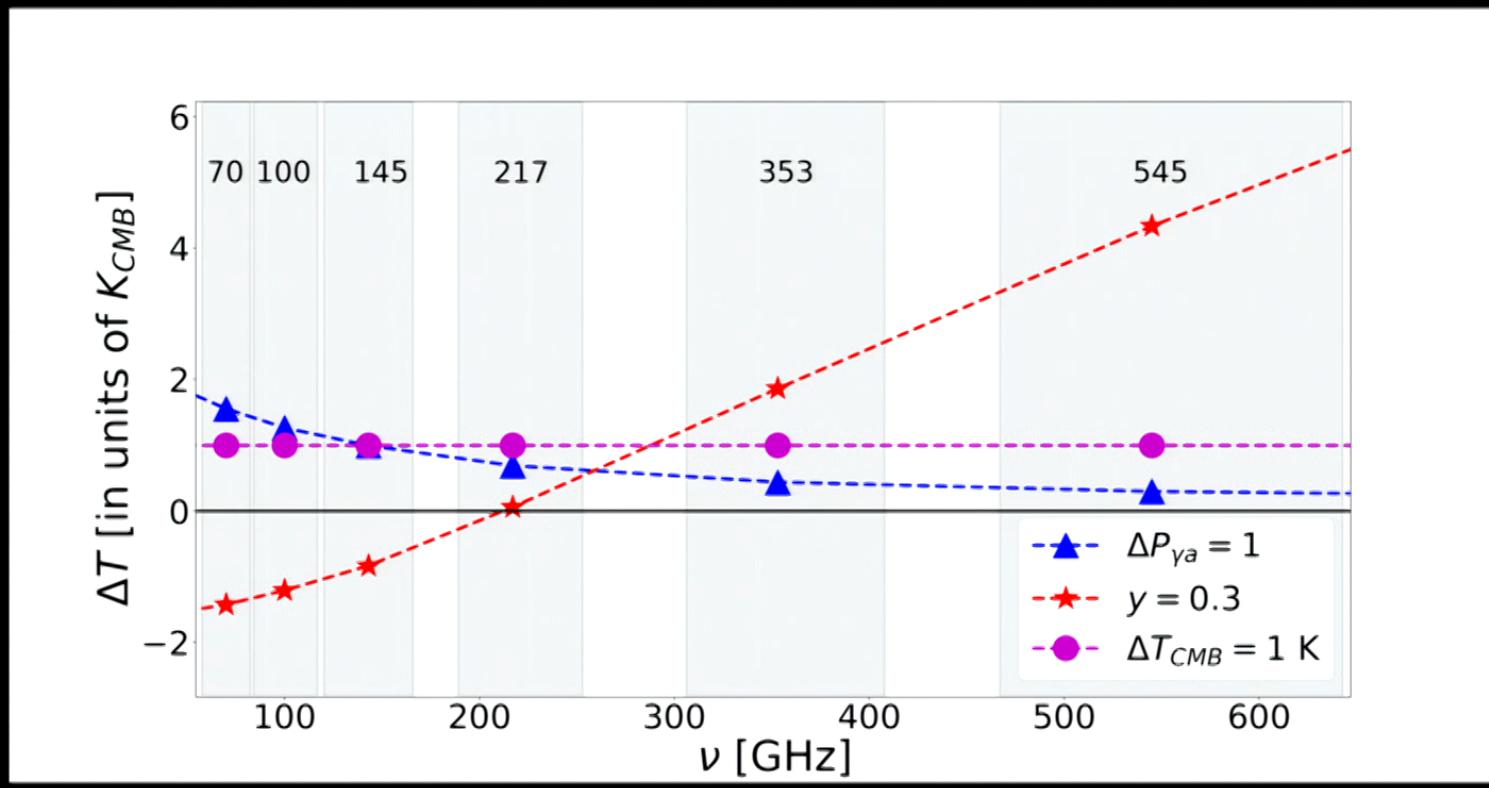


$$S_{\nu_i}(\hat{p}) = A_{\nu_i j} x_j(\hat{p}) + n_{\nu_i}(\hat{p}),$$

$$\mathbf{S}(\hat{p}) = \mathbf{Ax}(\hat{p}) + \mathbf{n}(\hat{p}),$$

46

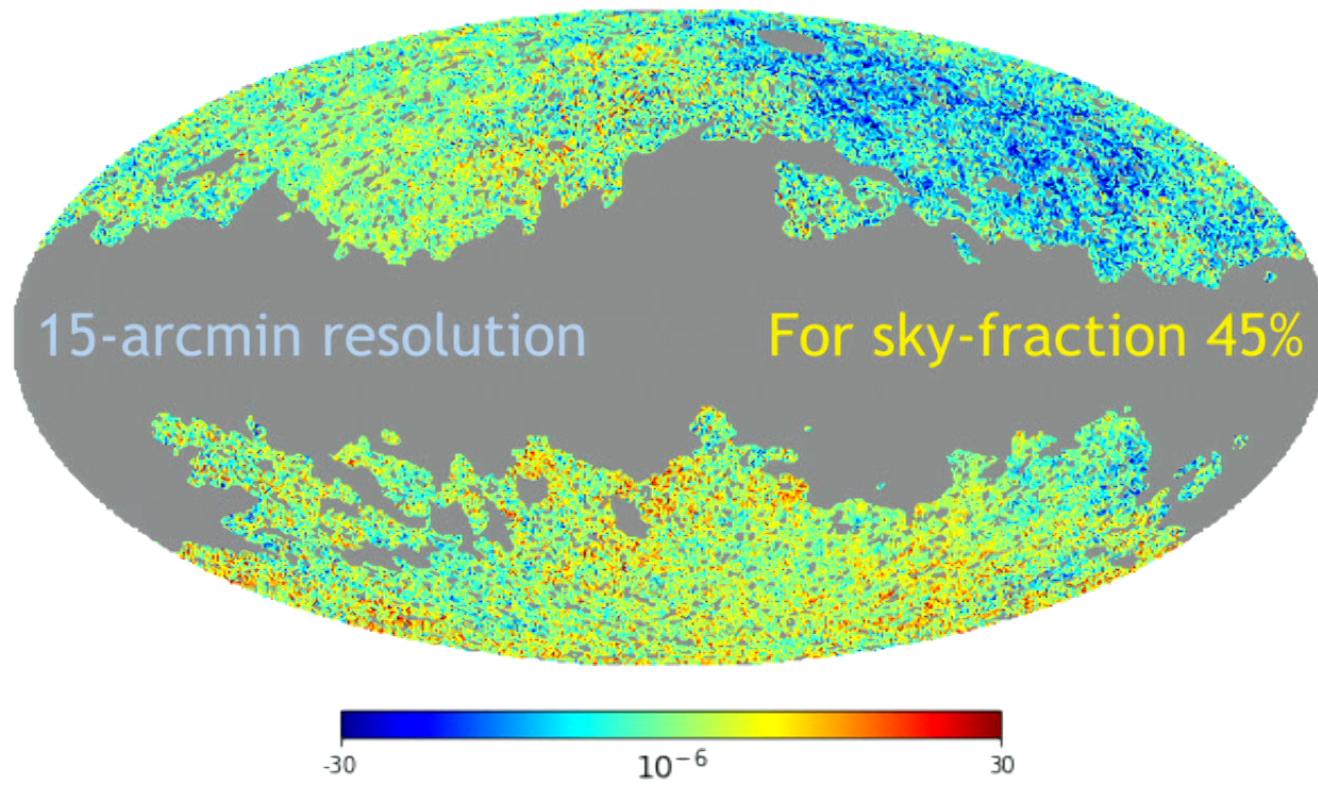
Spectrum of the distortion in the Planck frequency bands



47

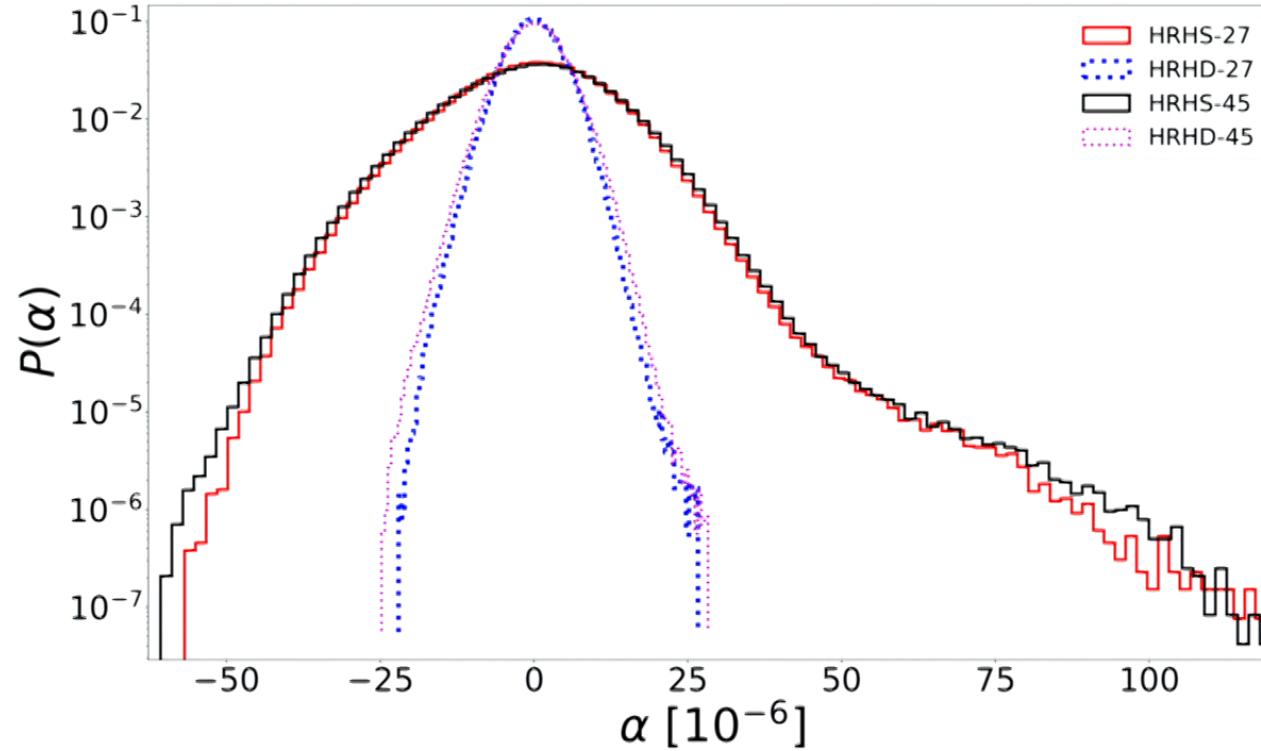
First all-sky constraints on Non-Resonant ALPs signal from Planck-data

$$\hat{x}_{\gamma a}(\hat{p}) = \hat{\mathcal{W}}_{\gamma a}^T \mathbf{S}(\hat{p}),$$

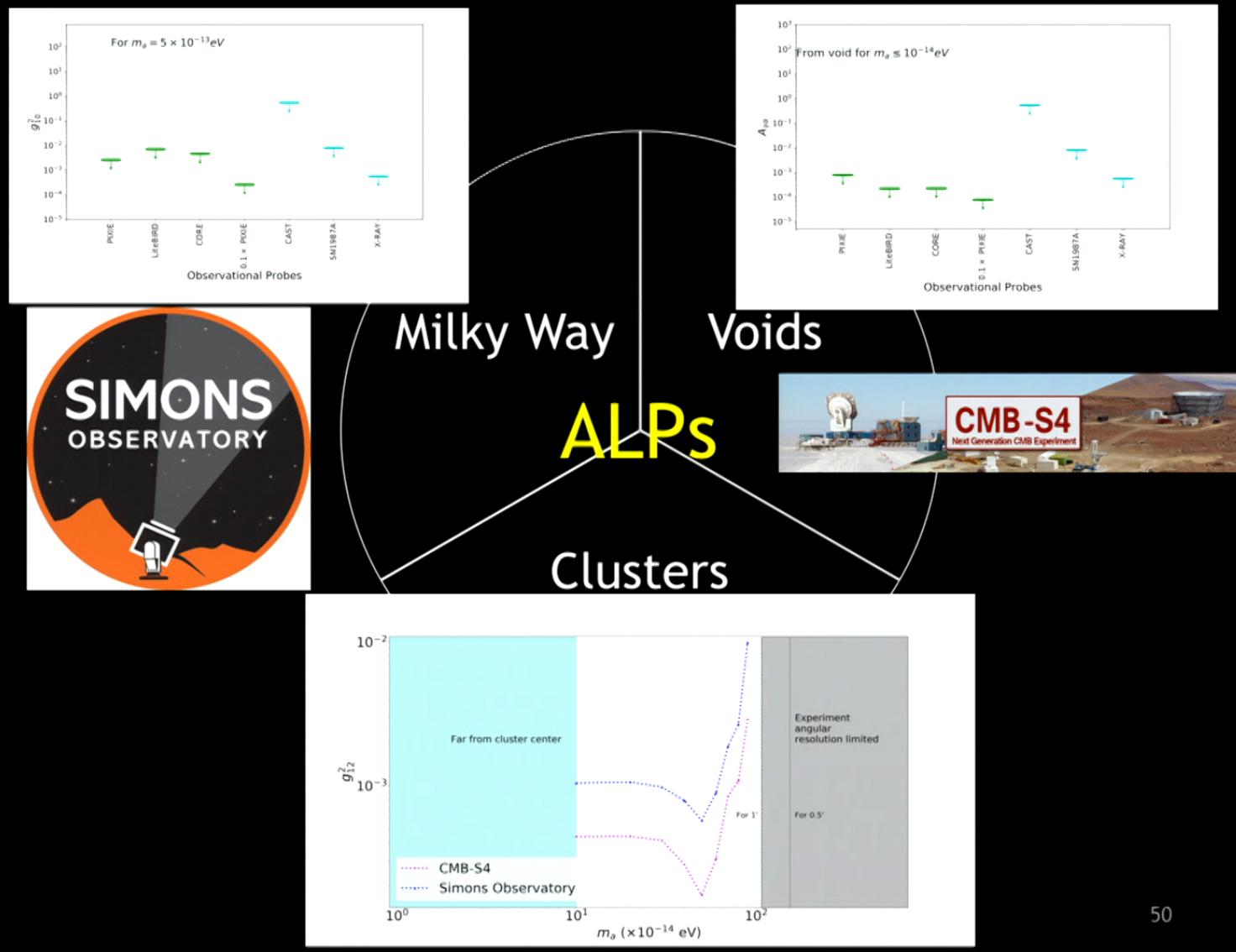


48

First all-sky constraints on Non-Resonant ALPs signal from Planck-data

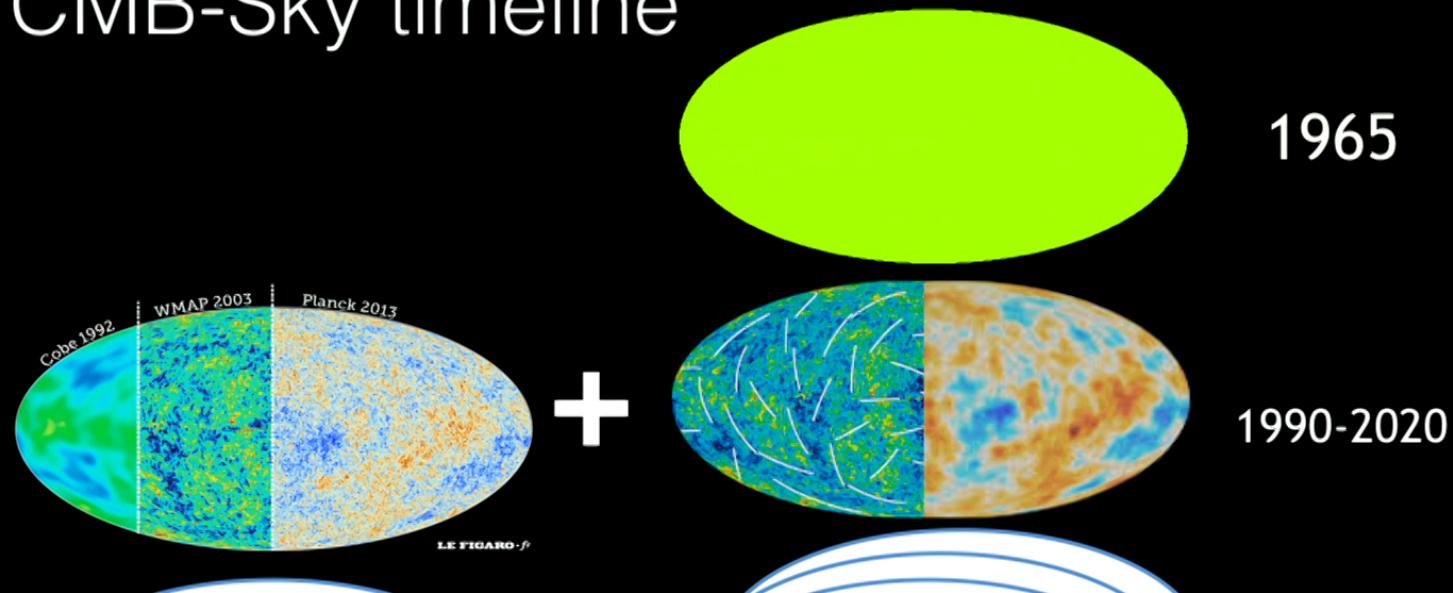


49



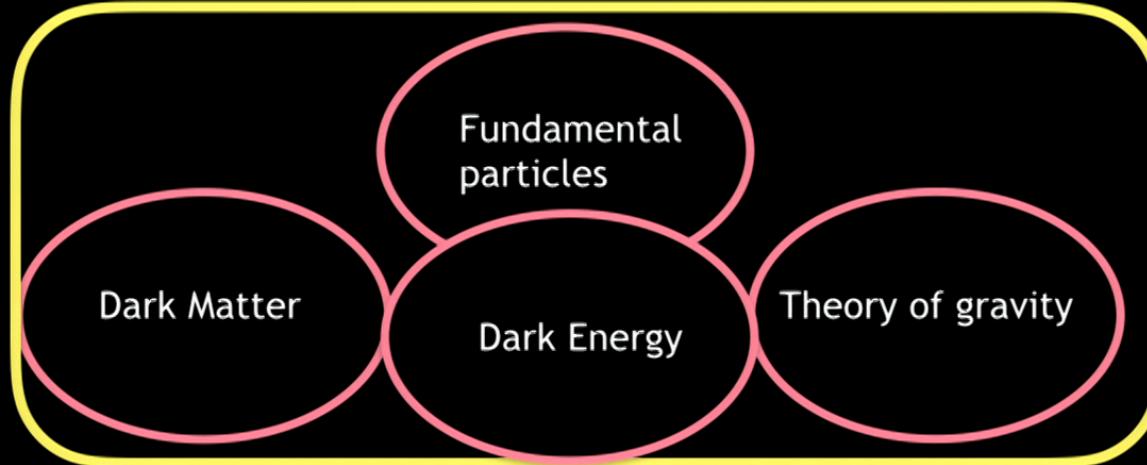
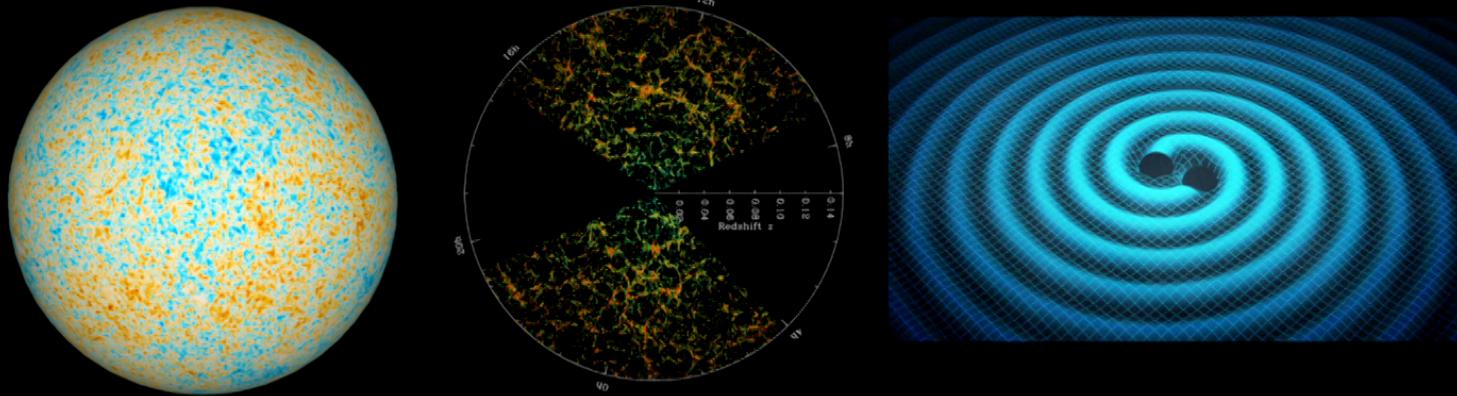
50

CMB-Sky timeline



My research interest

Develop connections between cosmological probes and fundamental physics



52

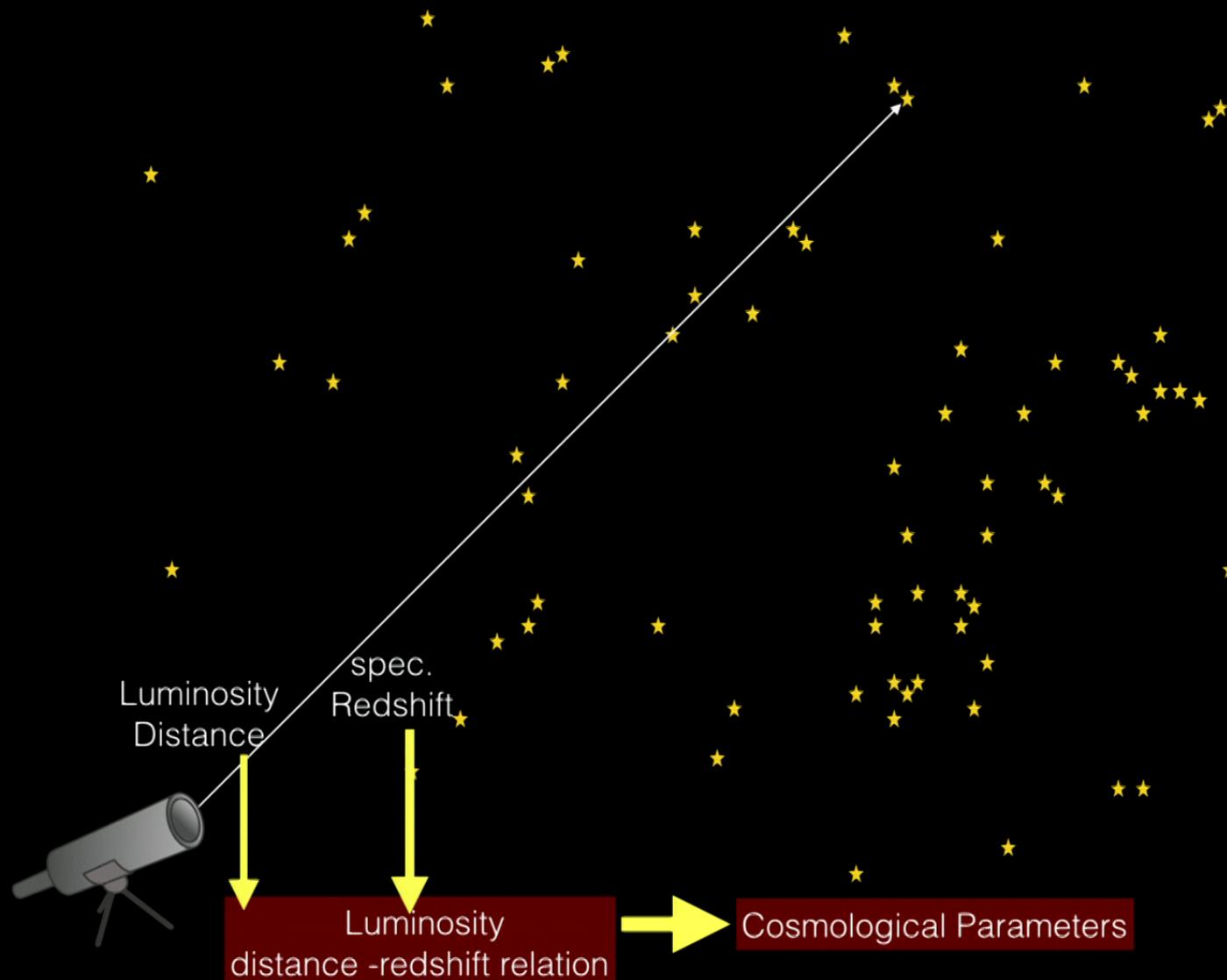
Cosmology with redshift unknown supernovae

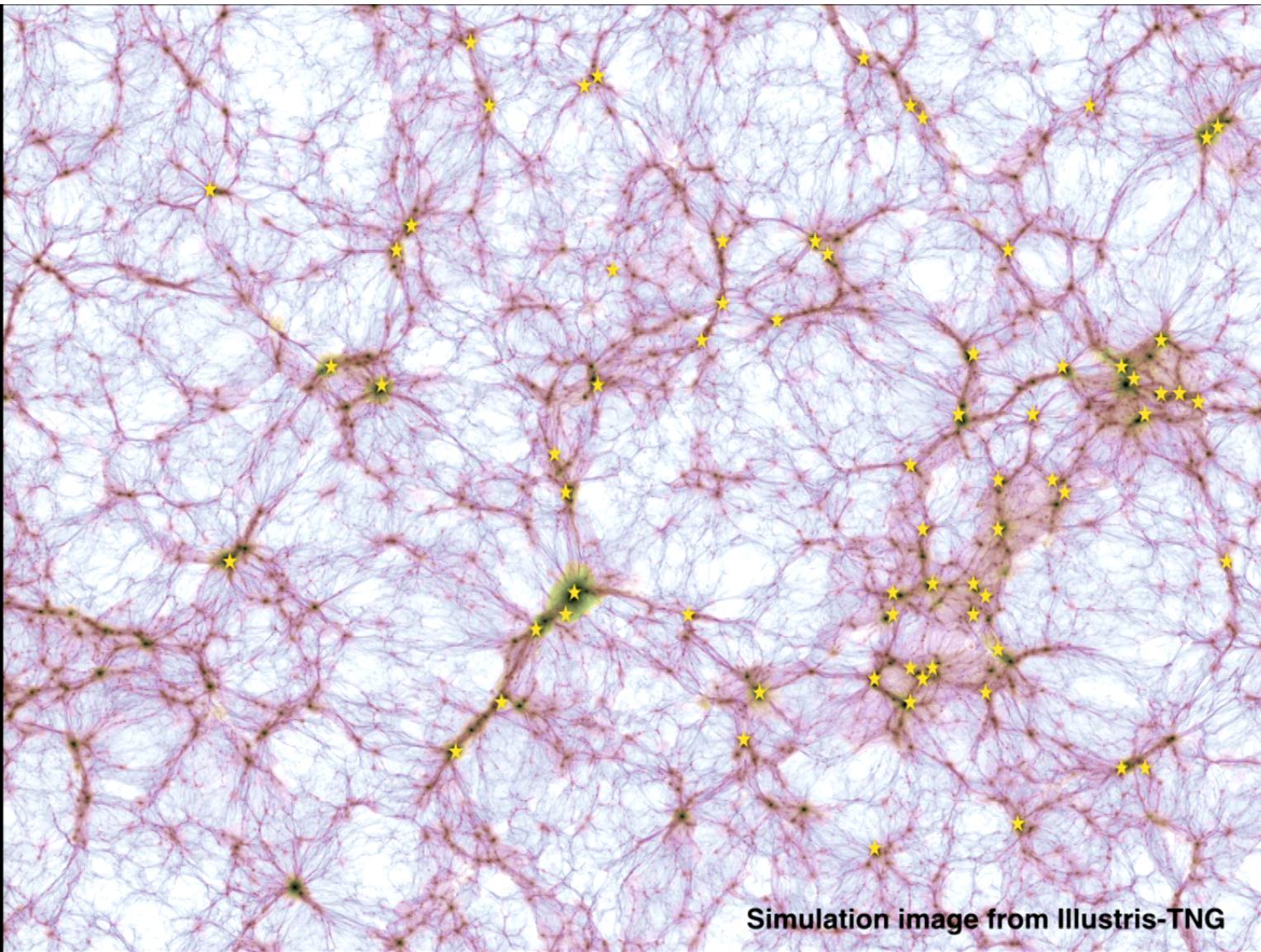
Mukherjee and Wandelt
arXiv:1808.06615

Upcoming surveys are going to measure

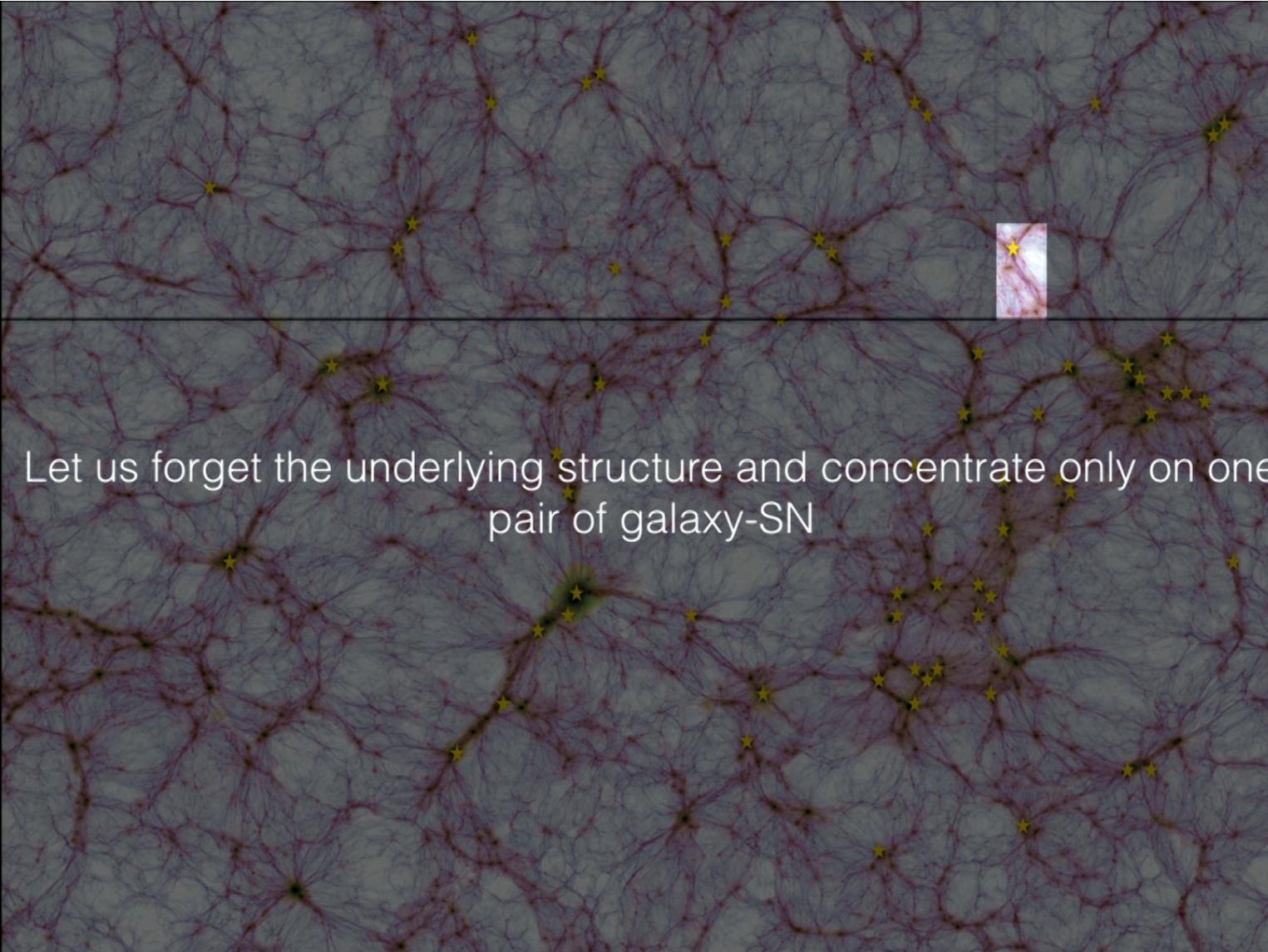
10^4 - 10^5 SN every year

These will primarily rely on photometric redshift measurements

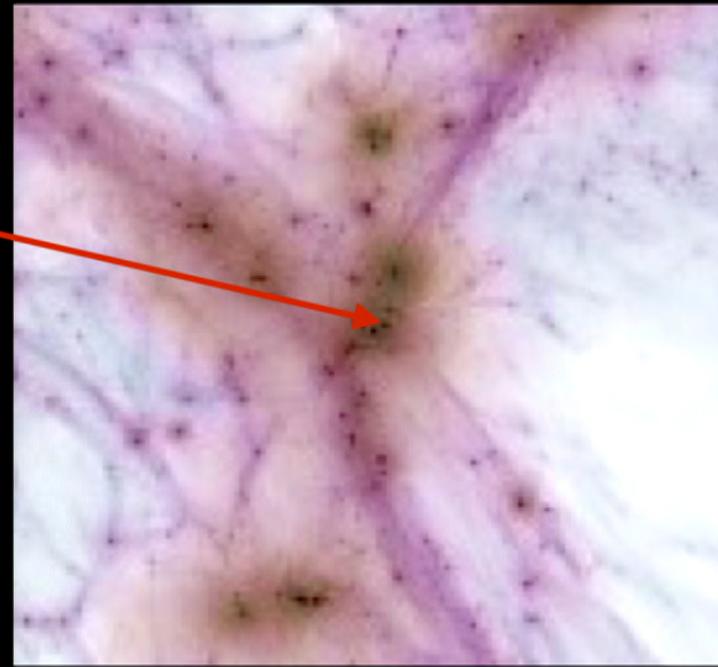




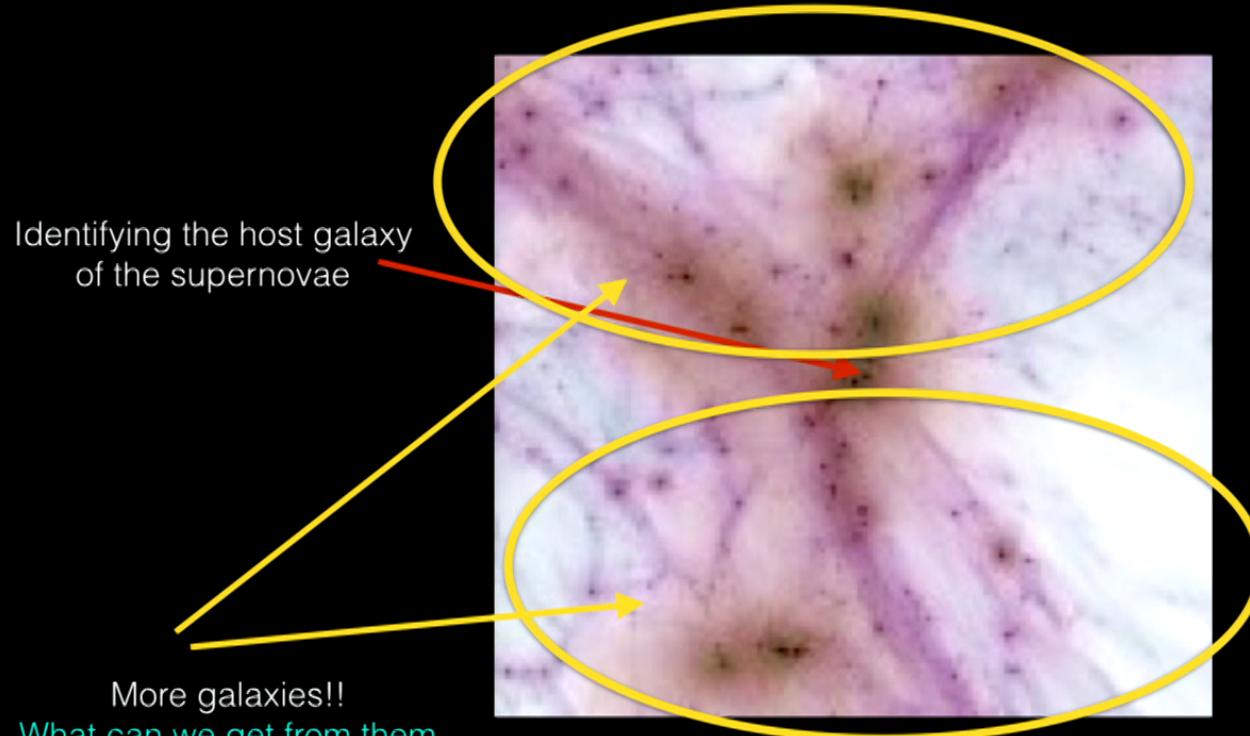
Simulation image from Illustris-TNG

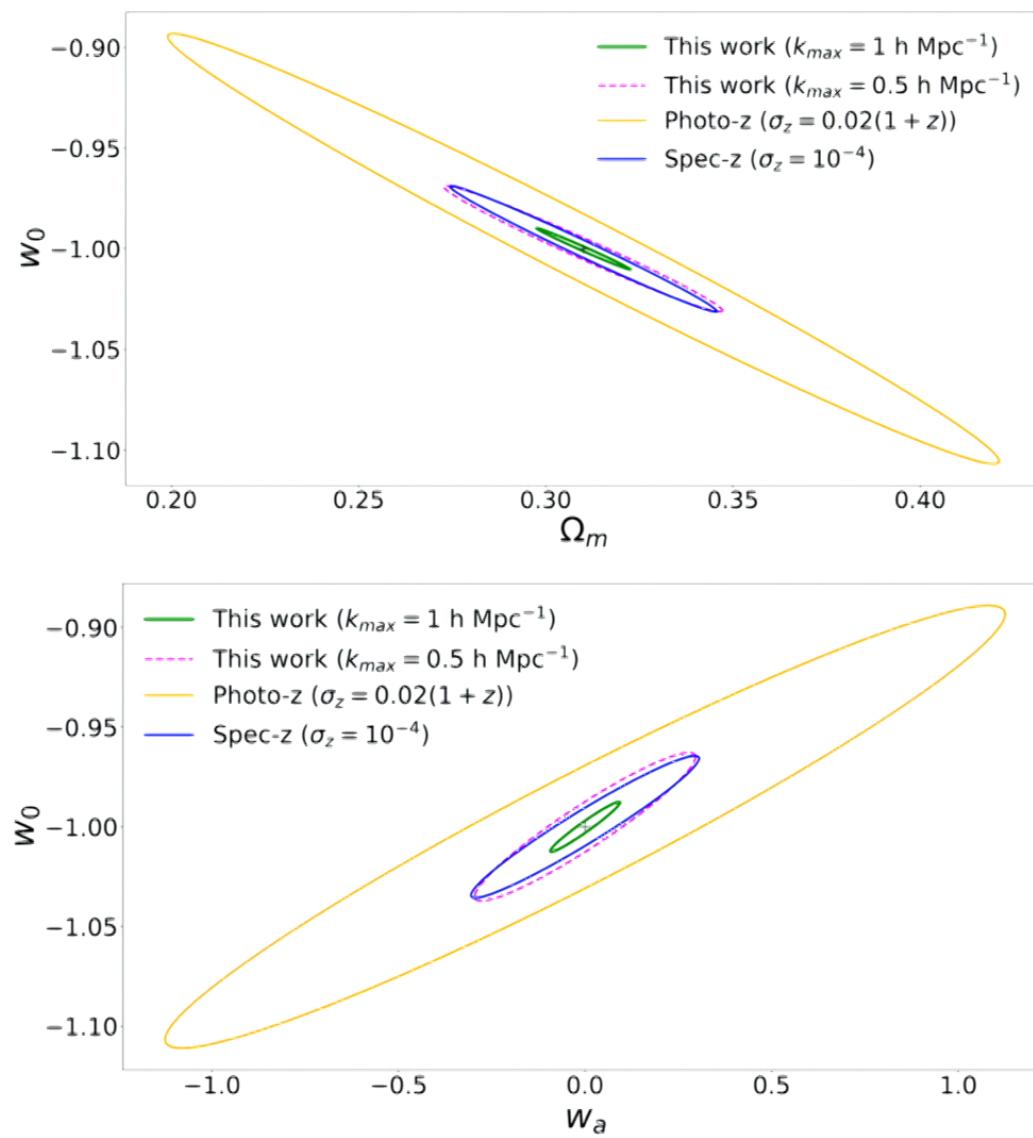


Identifying the host galaxy
of the supernovae



Cross-correlation of galaxies and supernovae





Cosmology with Gravitational Waves : Going beyond standard sirens

Mukherjee, Wandelt, Silk
under review in
Nature-Astronomy

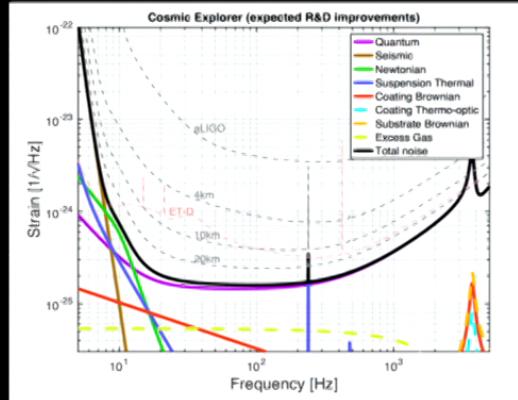


Image: LIGO collaboration
1607.08697

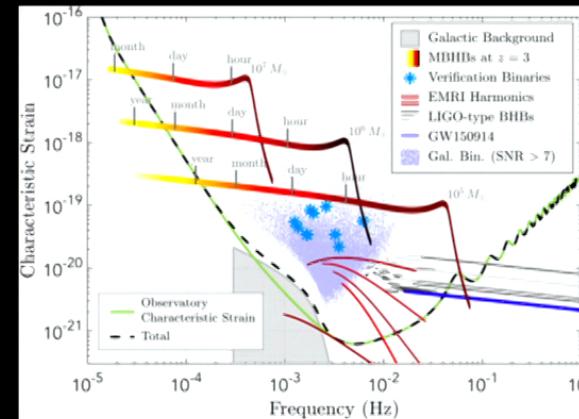
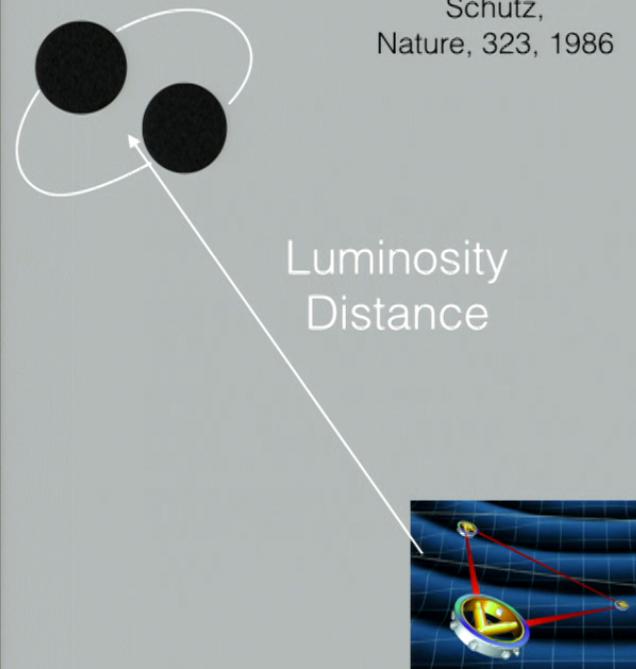


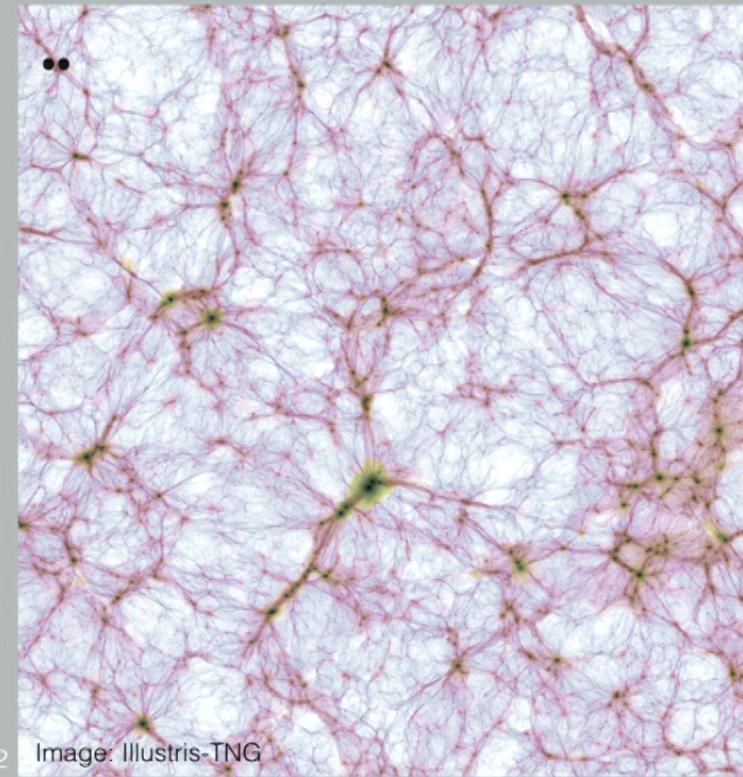
Image: LISA Science Proposal 61

Cosmology with GW

Standard Sirens: Distance Ladder

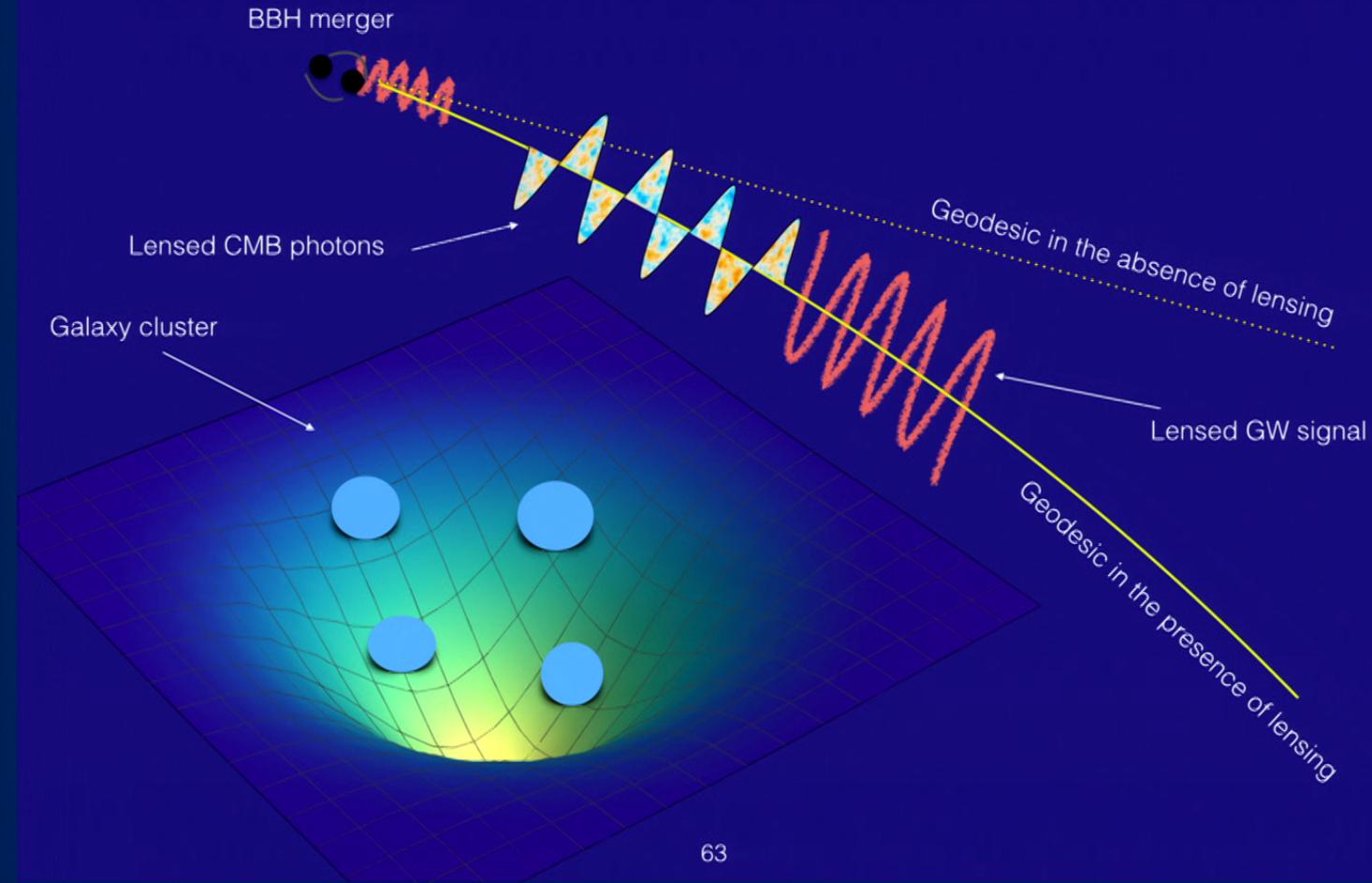


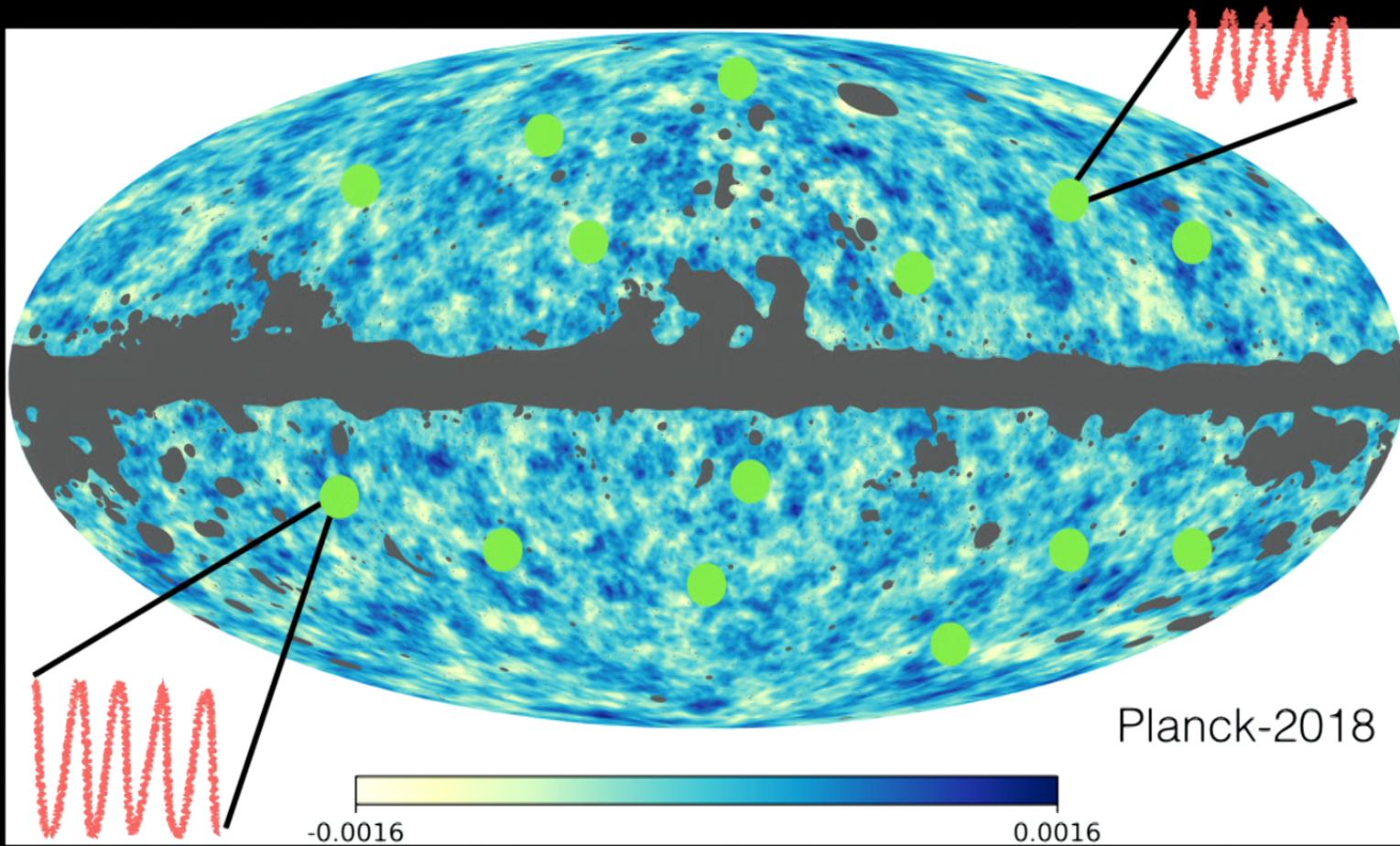
Cosmic Density field



Probing GW lensing using CMB-GW correlation

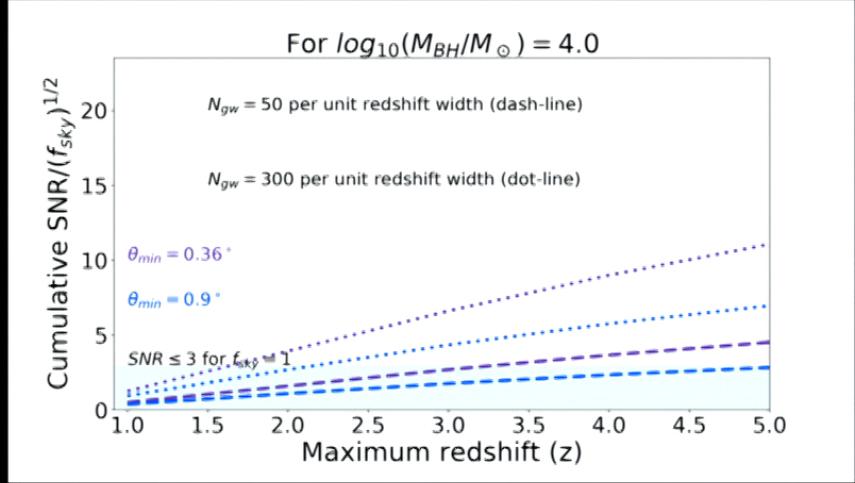
Mukherjee, Wandelt, Silk (2018)





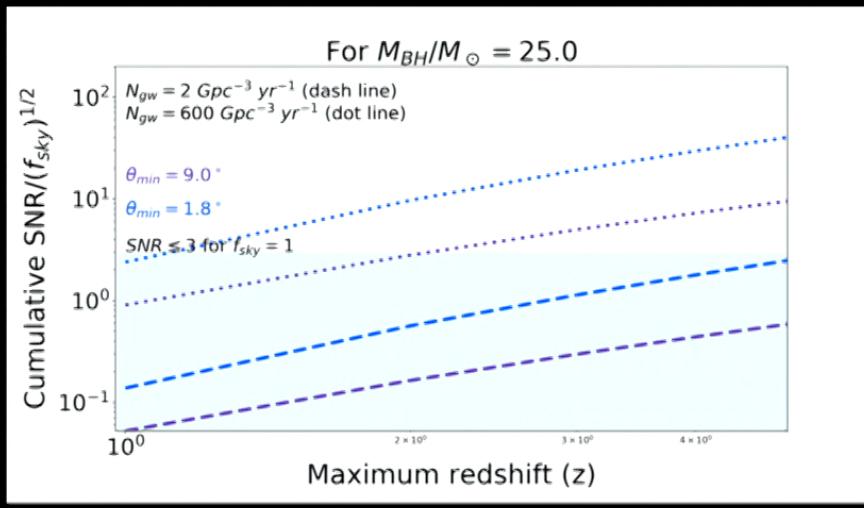
64

Discovery space



Mukherjee, Wandelt, Silk (2018)

LISA

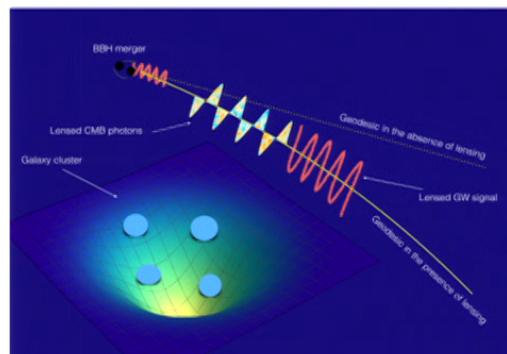
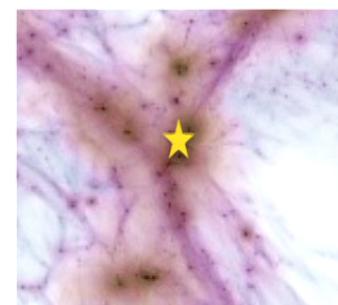
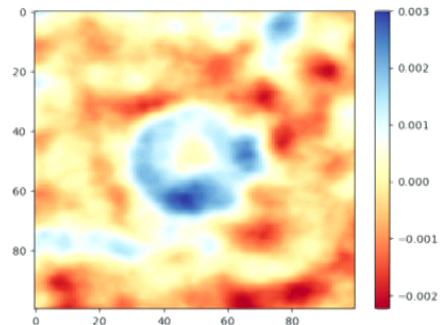


Cosmic Explorer

Diverse scientific scopes

- Concordant trajectory between electromagnetic waves and gravitational waves
- A probe to the theories of extra-dimensions of spacetime
(Deffayet & Menou 2007, LIGO-VIRGO Coll. (2018))
- A probe to the alternative theories of gravity
(Saltas et al. 2014,Nishizawa 2018)

A Few New Frontiers in Cosmology



Simons Observatory
CMB-S4

LSST
EUCLID
DESI

adv-LIGO
LISA
Cosmic Explorer

67