

Title: Fundamental physics and cosmology from gravitational-wave observations

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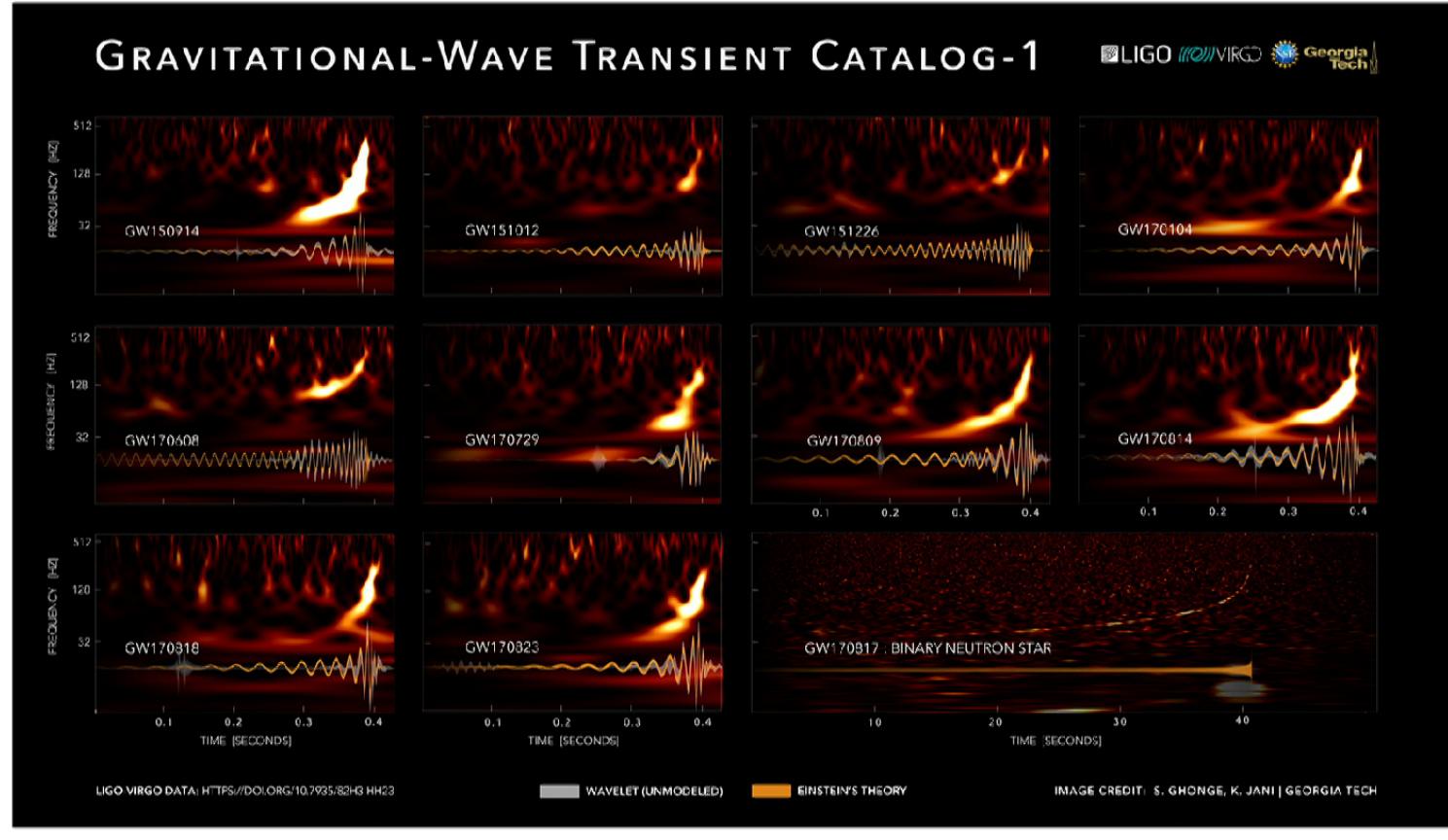
Abstract: <p>The detection of gravitational waves from mergers of compact binaries in the first two runs of the Advanced LIGO-Virgo have brought in valuable insights into fundamental physics and astrophysics. The coalescence process sweeping the components through a range of frequencies at highly relativistic velocities, have enabled some of the first tests of general relativity in its highly dynamical and extremely strong field regime. The recent detection of the binary neutron star merger has shed first light on the elusive neutron star equation of state. Furthermore, with its coincident electromagnetic counterpart, a first "standard-siren" measurement of the Hubble parameter has been made independent of a cosmic distance ladder. Subsequent detections are expected to give us a more direct understanding of the nature of the black holes and the composition of neutron stars, as well as bring in a plethora of results in astrophysics and possibly new insights into cosmology. I will go over some of these exciting new results and talk about future prospects of fundamental physics and cosmology from gravitational-wave observations.</p>

# Fundamental physics and cosmology from gravitational-wave observations

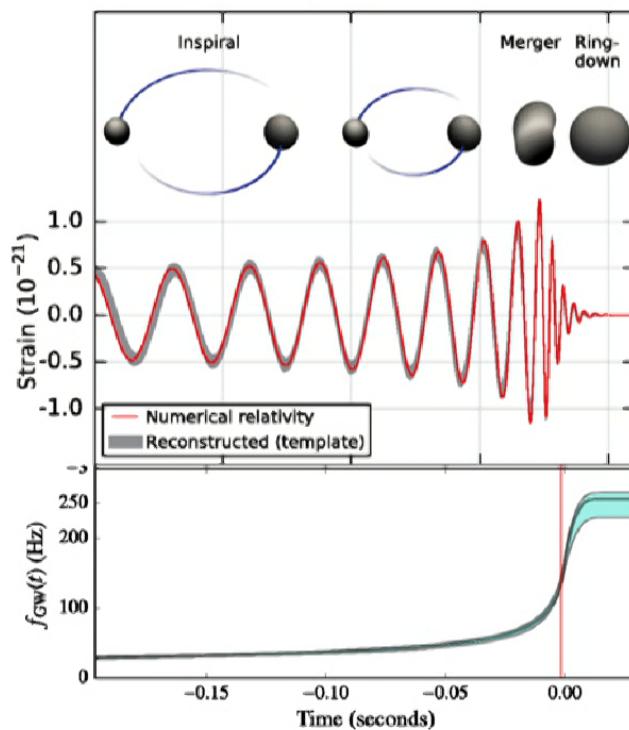
Archisman Ghosh  
Nikhef, Amsterdam

Perimeter Institute  
2019 January 10





Abbott et al., PRL 116, 061102 (2016)



Abbott et al., PRL 116, 221101 (2016)

PRL 116, 221101 (2016) Selected for a Viewpoint in Physics PHYSICAL REVIEW LETTERS week ending 3 JUNE 2016

### Tests of General Relativity with GW150914

B. P. Abbott et al.<sup>\*</sup>

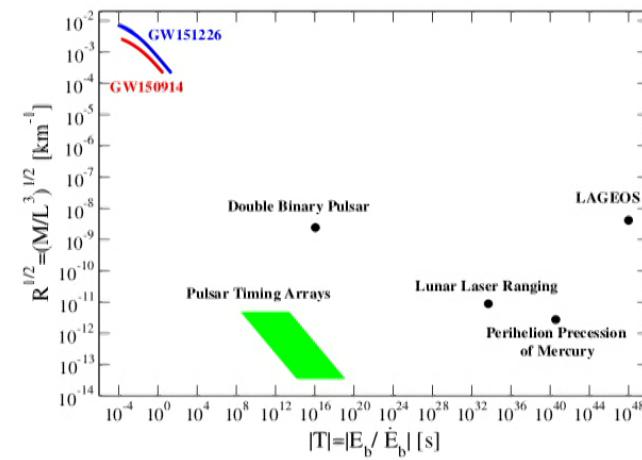
(LIGO Scientific and Virgo Collaborations)

(Received 26 March 2016; revised manuscript received 9 May 2016; published 31 May 2016)

## Probing strong-field gravity

First probes into the dynamical regime of strong field general relativity (GR).

Yunes et al. (2016)

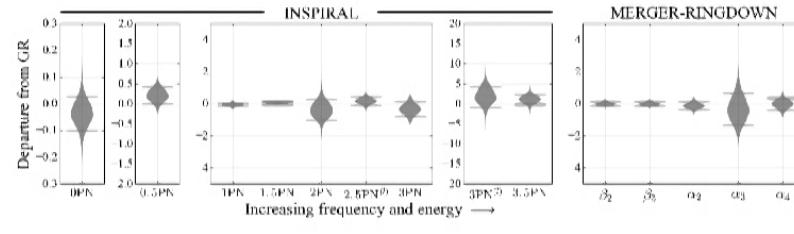


## Constraints on parameterized deformations from GR

Allowing coefficients in waveform models to deviate from their GR values, the deviation parameters do not show any departure from their GR values.

Li *et al.* (2011); Agathos *et al.* (2013); Meidam (PhD thesis, 2017); Meidam *et al.* (2017)

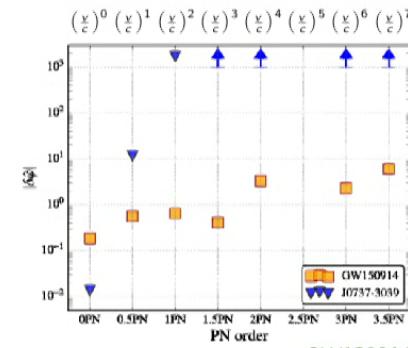
GW150914 + GW151226 + GW170104



Deviation in  $\left(\frac{v}{c}\right)^3$  coefficient constrained to  $\mathcal{O}(10\%)$

Dynamical self-interaction of spacetime

## Spin-orbit interaction



GWI56914

Abbott et al., PRL 118, 221101 (2017)

## First-ever measurement of orbital dynamics beyond leading order in $v/c$ .

## Constraints from modified dispersion

Will (1998); Mirshekari *et al.* (2012)

Modified dispersion relation:

(different frequencies travel with different speeds)

$$E^2 = p^2 c^2 + \mathbb{A} p^\alpha c^\alpha$$

$$\lambda_{\mathbb{A}} \equiv hc\mathbb{A}^{1/(\alpha-2)}$$

$\alpha \neq 0 \rightarrow$  local Lorentz invariance violation

$\alpha = 0 \rightarrow$  massive graviton (for  $\Lambda > 0$ )

**GW150914 + GW151226 + GW170104**

$$\lambda_g \equiv \frac{h}{m_g c} > 1.6 \times 10^{13} \text{ km}$$

$$m_g < 7.7 \times 10^{-23} \text{ eV}/c^2$$

Effect gets enhanced with propagation over a distance!

Agathos (PhD thesis, 2016); Samajdar (PhD thesis, 2017); Samajdar & Arun (2017)

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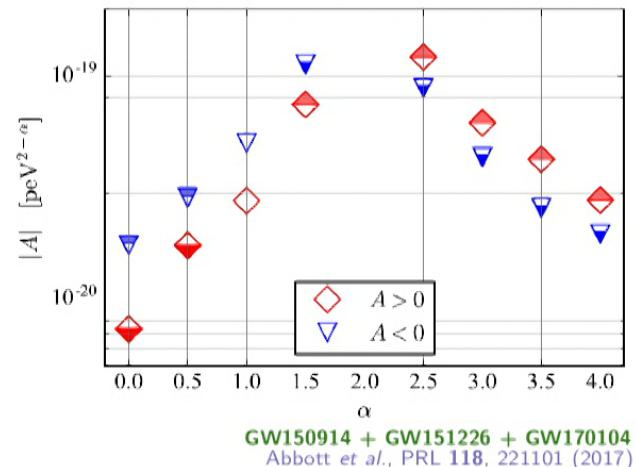
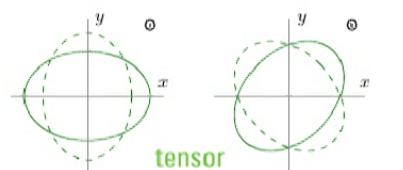


TABLE III. 90% credible level lower bounds on the length scale  $\lambda_A$  for Lorentz invariance violation test using GW170104 alone. **GW170104**

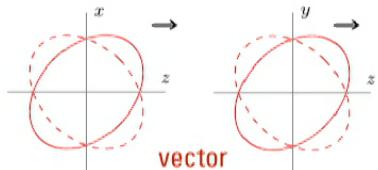
	$A > 0$	$A < 0$
$\alpha = 0.0$	$1.3 \times 10^{13}$ km	$6.6 \times 10^{12}$ km
$\alpha = 0.5$	$1.8 \times 10^{16}$ km	$6.8 \times 10^{15}$ km
$\alpha = 1.0$	$3.5 \times 10^{22}$ km	$1.2 \times 10^{22}$ km
$\alpha = 1.5$	$1.4 \times 10^{41}$ km	$2.4 \times 10^{40}$ km

Hubble scale  $\approx 1.3 \times 10^{23}$  km

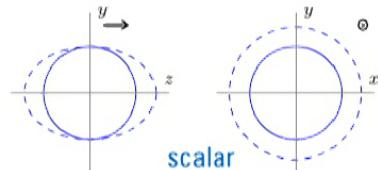
# Polarization from 3-detector observation of GW170814



tensor

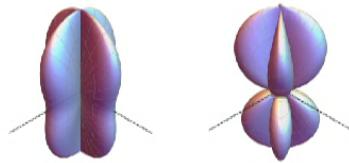


vector



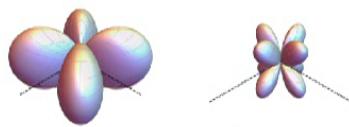
scala

six polarizations → distinct antenna patterns



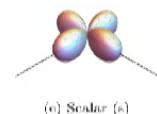
(a) Plus (+)

(b) Cross ( $\times$ )



(c) Vector-x ( $x$ )

(d) Vector- $y$  ( $y$ )



(e) Scalar ( $\phi$ )

$$|F_t^I(\alpha, \delta)| \equiv \sqrt{F_+^I(\alpha, \delta)^2 + F_\times^I(\alpha, \delta)^2}.$$

$$|F_v^I(\alpha, \delta)| \equiv \sqrt{F_x^I(\alpha, \delta)^2 + F_y^I(\alpha, \delta)^2},$$

$$|F_s^I(\alpha, \delta)| \equiv \sqrt{F_b^I(\alpha, \delta)^2 + F_1^I(\alpha, \delta)^2}$$

In GR: GW are **transverse**, **traceless**  
only **tensor** polarizations

pure tensor / pure scalar = 1000 / 1

pure tensor / pure vector = 200 / 1

Isi & Weinstein (2017)

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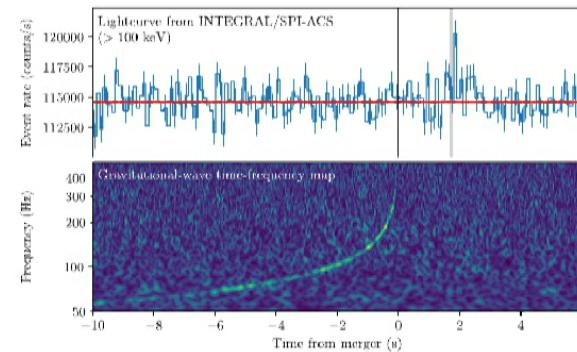
Need multiple detectors: **thanks to Virgo!**

Abbott *et al.*, PRL 119, 141101 (2017)

## Constraints from GW170817+GRB

Delay of only a few seconds after a propagation over one hundred million light years.

$$t_{\text{EM}} - t_{\text{GW}} = 1.74 \pm 0.05 \text{ s}$$



Constraints on speed of gravity

assuming GRB emitted within 10s of GW

$$-3 \times 10^{-15} \leq \frac{v_{\text{GW}} - v_{\text{EM}}}{v_{\text{EM}}} \leq +7 \times 10^{-16}$$

"Shapiro time delay" of GW and EM in the gravitational potential of our galaxy:

$$-2.6 \times 10^{-7} \leq \gamma_{\text{GW}} - \gamma_{\text{EM}} \leq 1.2 \times 10^{-6}$$

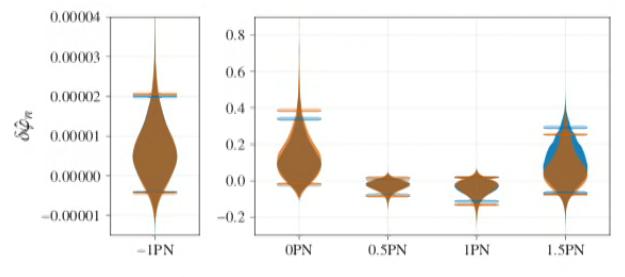
Test of the equivalence principle.

Abbott *et al.* *Astrophys. J.* 848 #2, L13 (2017)

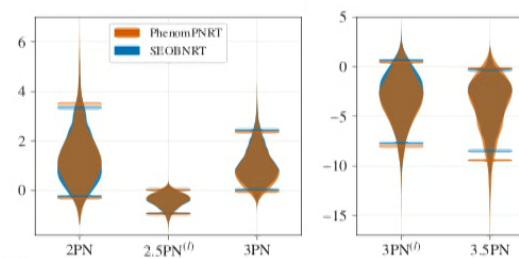
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## Tests of general relativity with GW170817

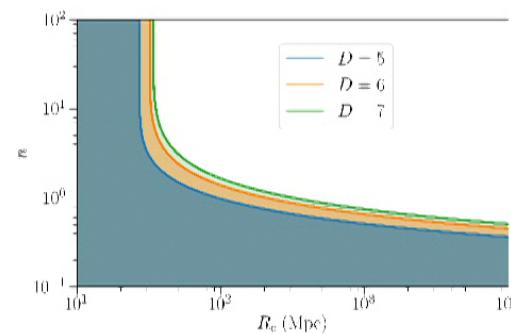
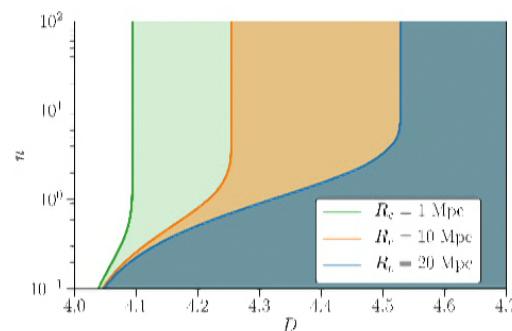
### Dipole radiation



Abbott et al. arXiv:1811.00364 [gr-qc]



- Parameterized deviations do not show any departures from GR values.
  - “Inverse square law” → constraints on extra dimensions.



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## Constraints from modified dispersion

Will (1998); Mirshekari *et al.* (2012)

Modified dispersion relation:

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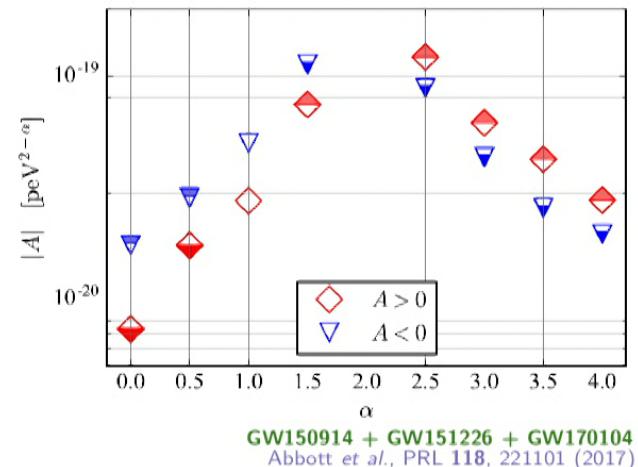


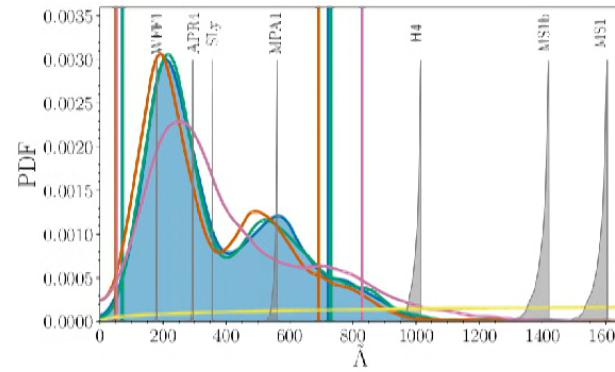
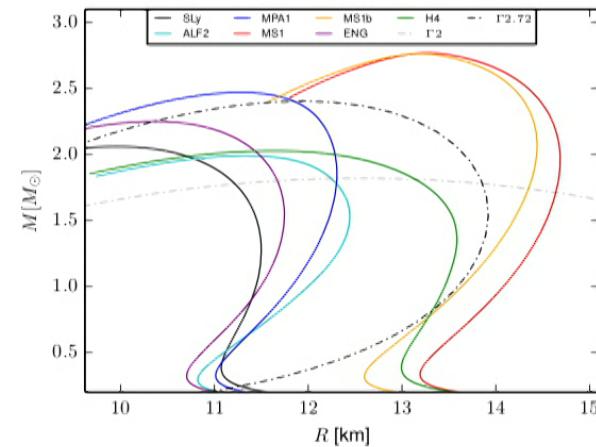
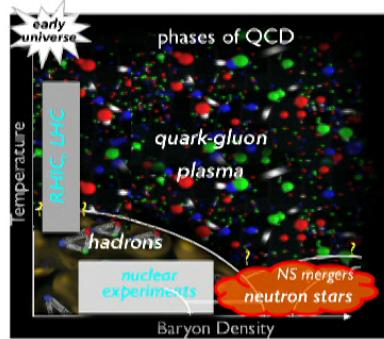
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Hubble scale  $\approx 1.3 \times 10^{23}$  km

## GW170817: measurement of properties of the neutron star

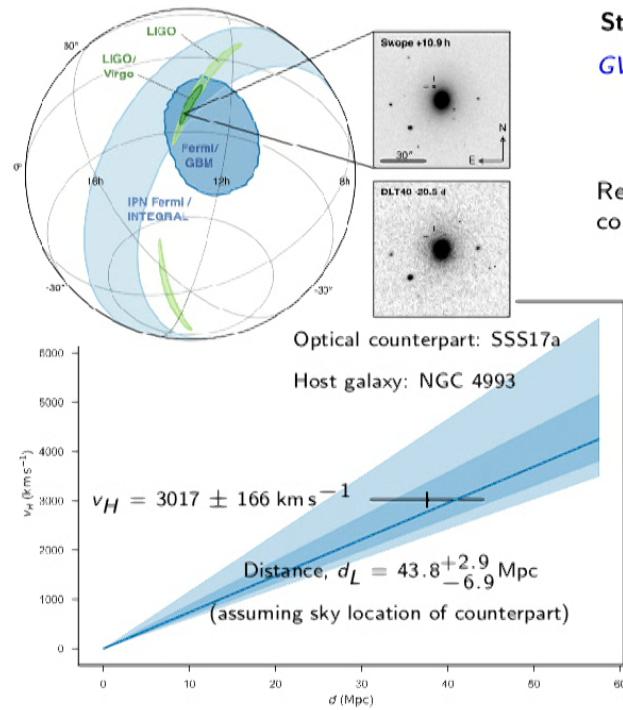
Figure from: Dietrich et al. (2015)



Abbott et al. arXiv:1805.11579 [gr-qc]

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## Cosmology: Hubble parameter with GW170817



Independent of any distance ladder!

Abbott et al. *Astrophys. J.* **848** #2, L12 (2017); LSC-EPO

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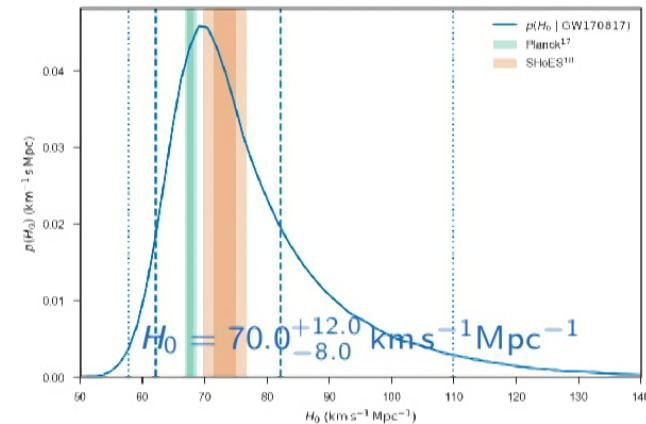
### Standard siren

Schutz (1986), Holz & Hughes (2005)

GWs provide a direct measurement of the luminosity distance!

$$v_H = H_0 d_L$$

Recession velocity (or redshift) can come from a transient EM counterpart or an identified host galaxy.



Abbott et al. *Nature* **551** #7678, 85-88 (2017)

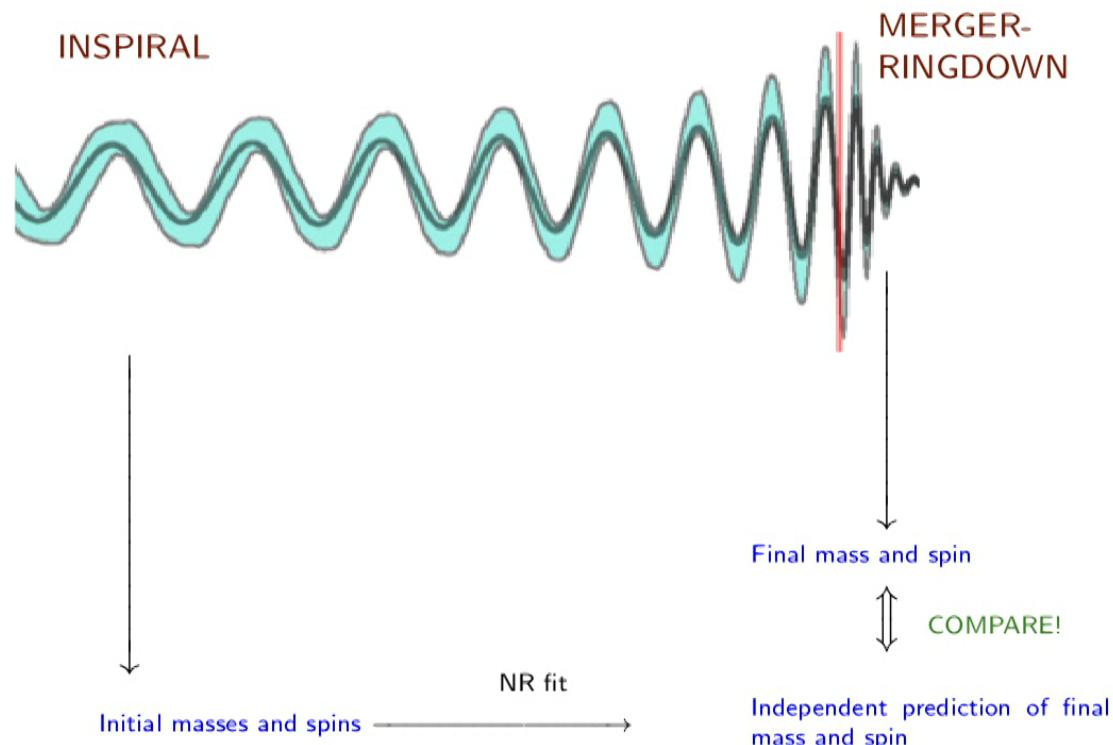
## Plan of the talk

- Overview of LIGO-Virgo detections and science-results ✓
- Probing strong-field gravity
  - The inspiral-merger-ringdown consistency test
  - Probing into the nature of compact objects
    - Ringdown
    - Echoes
- Cosmology
  - A “standard siren” measurement of  $H_0$
  - “ $H_0$ -statistical” and beyond
- Future prospects

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## Strong-field gravity

## Inspiral-merger-ringdown consistency test



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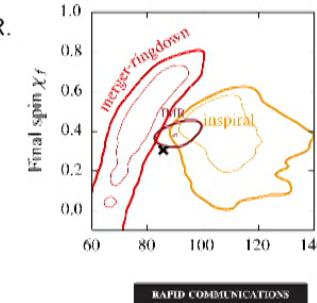
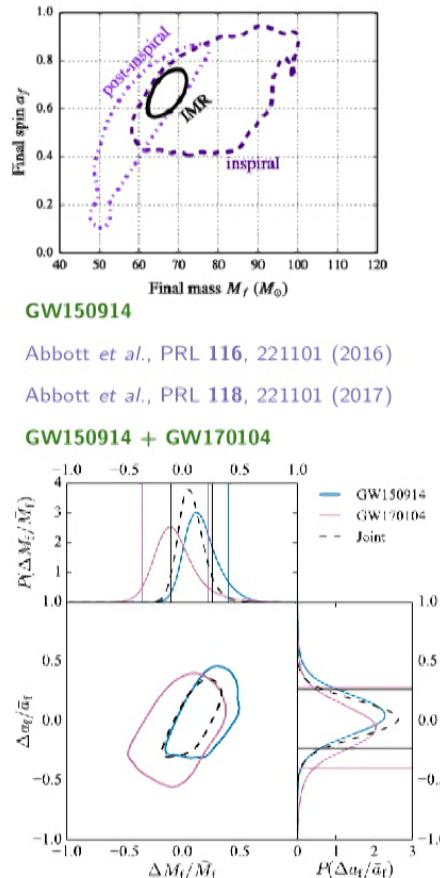
# Inspiral-merger-ringdown consistency test

One of the first tests of GR carried out with GW150914

Mass and spin of the remnant object estimated from the inspiral and merger-ringdown parts agree with each other given GR predictions.

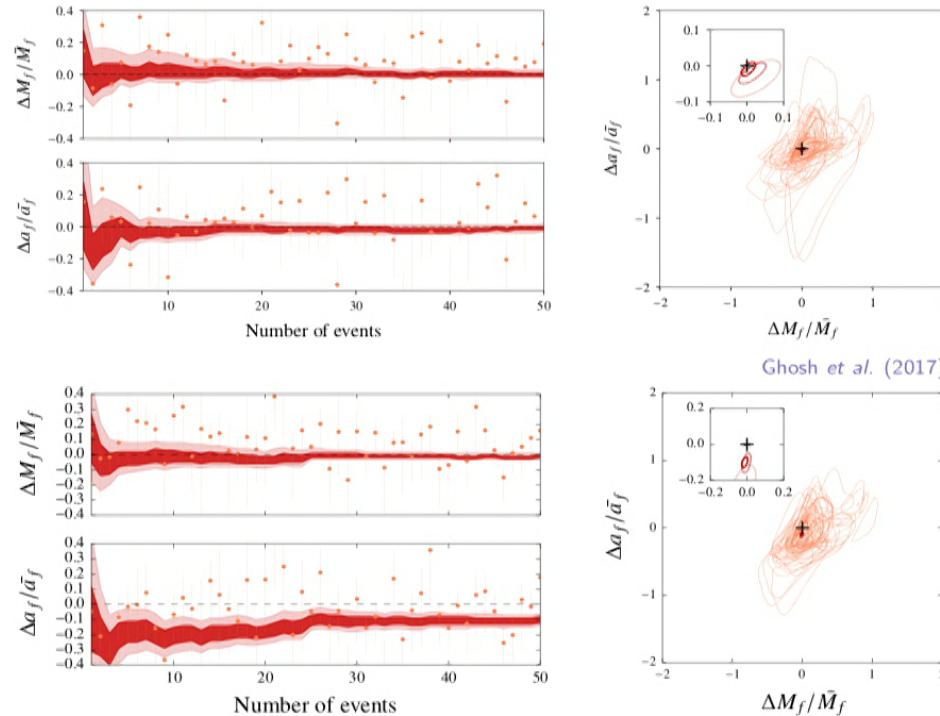
Ghosh *et al.* (2016); Ghosh *et al.* (2017)

Might not have been true in modified GR



Stronger constraints on systematic departures from GR combining information from multiple detections.

## Inspiral-merger-ringdown consistency test



Stronger constraints on systematic departures from  
GR combining information from multiple detections.

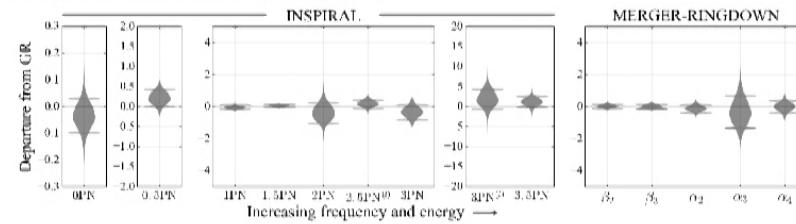
## Some prospects of doing better ...

- Searching for systematic **systematic** departures.
- Is this the optimal way of testing?
- Use theoretically motivated combinations.
- Use correlations between parameters.

Multipole moments: Kastha et al. (2018)

Insights from simulations in modified gravity?

GW150914 + GW151226 + GW170104



Abbott et al., PRL 118, 221101 (2017)

## Probing the nature of the progenitor and remnant compact objects

Are they really black holes, or exotic compact objects mimicking black holes?

Boson stars, dark matter stars, gravastars, shells, wormholes

Three “complementary” ways in three different regimes:

- Finite size effects during inspiral.
- No-hair conjecture with quasinormal modes.
- Search for post-merger oscillations or “echoes”.

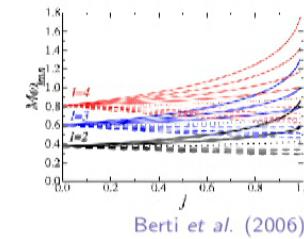
## Testing the no-hair conjecture with ringdown quasinormal modes

### No-hair conjecture:

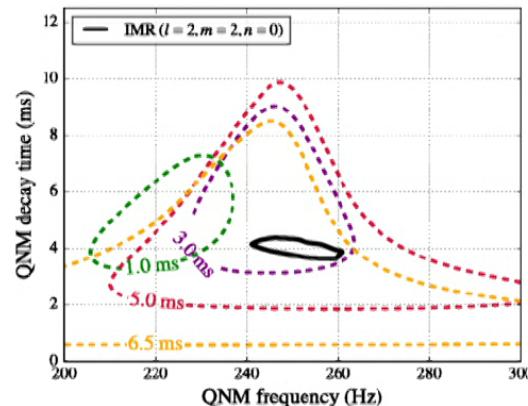
A stationary black hole in Einstein's general relativity is described only by its mass and spin.

During ringdown, the **quasinormal mode frequencies and damping times** will depend only on the **mass and spin of the remnant black hole**, which can be obtained from linearized Einstein equations on Kerr background.

⇒ Test for dependences  $\omega_{lmn}(M_f, J_f)$ ,  $\tau_{lmn}(M_f, J_f)$ .



Berti et al. (2006)



Difficult to measure leading QNM for GW150914.

Design sensitivity  $\sim 3$  times higher.

## Testing the no-hair conjecture with ringdown quasinormal modes

- Even where one is not able to isolate the individual modes, one can look for systematic departures in the QNM frequencies and damping times from their GR values.

Gossan *et al.* (2011) Meidam *et al.* (2014)

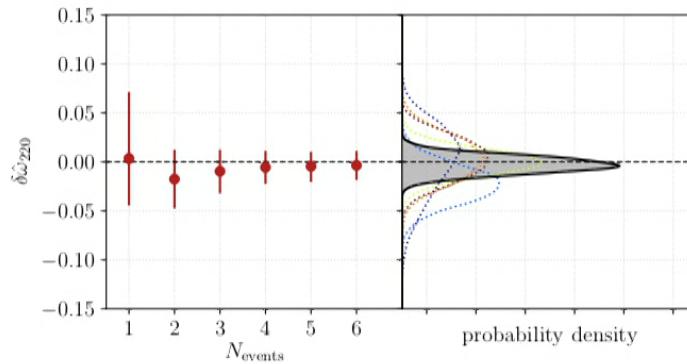
$$\omega_{lmn} = \omega_{lmn}^{GR}(1 + \delta\omega_{lmn}), \quad \tau_{lmn} = \tau_{lmn}^{GR}(1 + \delta\tau_{lmn})$$

- The general expectation was that such tests would become effective only for sources detected by third generation or space-based detectors.

**Empirical tests of the black hole no-hair conjecture using gravitational  
-wave observations**

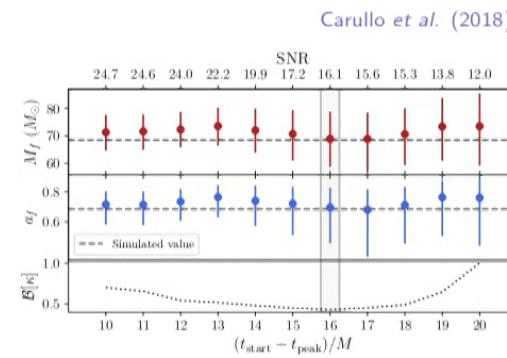
Gregorio Carullo,<sup>1,2,\*</sup> Laura van der Schaaf,<sup>2</sup> Lionel London,<sup>3</sup> Peter T. H. Pang,<sup>4</sup> Ka Wa Tsang,<sup>2</sup> Otto A. Hannuksela,<sup>4</sup>  
Jeroen Meidam,<sup>2</sup> Michalis Agathos,<sup>5</sup> Anuradha Samajdar,<sup>2</sup> Archisman Ghosh,<sup>2</sup> Tjonnie G. F. Li,<sup>4</sup>  
Walter Del Pozzo,<sup>1,6</sup> and Chris Van Den Broeck<sup>2,7</sup>

- With  $\mathcal{O}(5)$  BBH sources similar to GW150914, the systematic departures can be measured with an accuracy of  $\sim 1.5\%$  by the Adv LIGO-Virgo at design sensitivity.



- Effective criterion for “start of ringdown” from point of view of parameter estimation.

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## Search for “echoes” after the merger

In a large class of exotic compact objects,

Horizon-scale corrections  $\Rightarrow$  secondary bursts of radiation.

Modulated and distorted train of “echoes”.

$n=8$ : wormholes

$n=4$ : empty shell

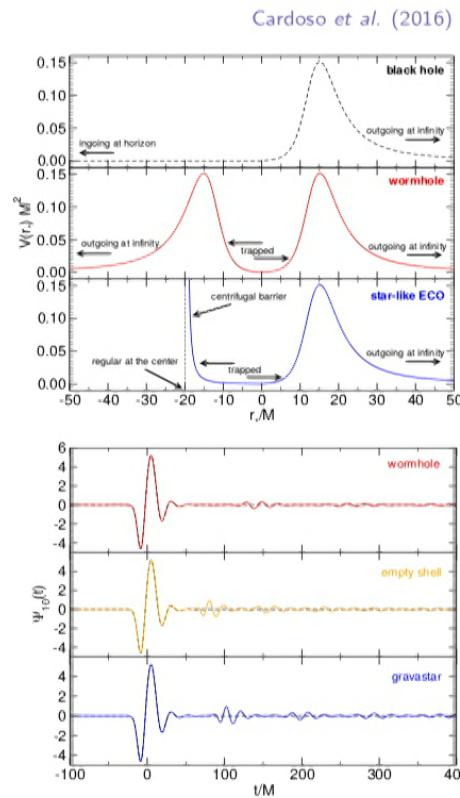
$n=6$ : thin-shell gravastars

$$\Delta t = nM \log(M/l)$$

Planck-scale corrections can appear relatively soon.

For an event like GW150914,  $\Delta t = \mathcal{O}(100\text{ ms})$ , at aLIGO design can hope to see first few echoes.

Can search for “echoes” immediately following the binary-merger detection.



Not sufficiently modelled;

Exotic objects not envisaged in literature.

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One feature expected to be reasonably robust: constancy of time difference between the subsequent echoes.

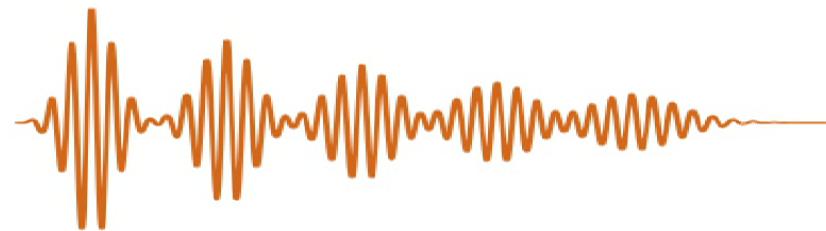
## A model-agnostic coherent search for echoes

- Use wavelets that are trains of sine-Gaussians to reconstruct the signal

$$\Psi(t; A_n, f_0, \tau, t_n, \phi_n) = \sum_{n=0}^{N_{\text{echoes}}} A e^{-(t-t_n)^2/\tau_n^2} \cos(2\pi f_0(t - t_n) + \phi_n)$$

With:

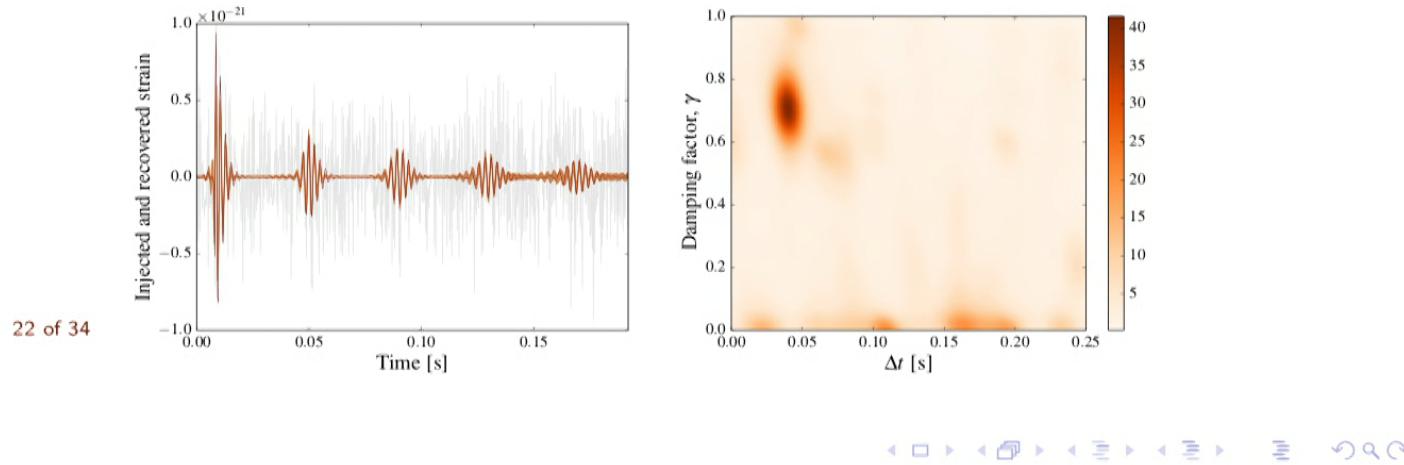
$A_n = \gamma^n A$	damping
$\tau_n = w^n \tau$	widening
$t_n = t_0 + n\Delta t$	time between subsequent echoes
$\phi_n = \phi_0 + 2\pi f_0 n \Delta t + n \Delta \phi$	phase shift subsequent echoes



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**A morphology-independent data analysis method for detecting and characterizing gravitational wave echoes**

Ka Wa Tsang,<sup>1</sup> Michiel Rollier,<sup>1</sup> Archisman Ghosh,<sup>1</sup> Anuradha Samajdar,<sup>1</sup> Michalis Agathos,<sup>2</sup> Katerina Chatzioannou,<sup>3</sup> Vitor Cardoso,<sup>4</sup> Gaurav Khanna,<sup>5</sup> and Chris Van Den Broeck<sup>1,6</sup>



## Runs with O1 data

We extend our analysis to O1 C02 data.  
The segments being used have good data quality (passed CBC CAT1,2,3)

Runs with O1 data:

- ① Simulated signal with O1 noise
- ② Background distribution
- ③ Echo searches at O1 BBH events

Abedi et al. results are within the prior ranges.

	GW150914	GW151012	GW151226
dt (s)	0.30068	0.09758	0.19043

Prior:

- $dt = [0., 0.7]s$
- $d\phi = [0., 2\pi]$
- $\gamma = [0., 1.]$
- $w = [1., 2.]$
- $T = 2s$
- $f_{low} = 20Hz$
- $srate = 2048Hz$

## Echoes search at O1 BBH detections

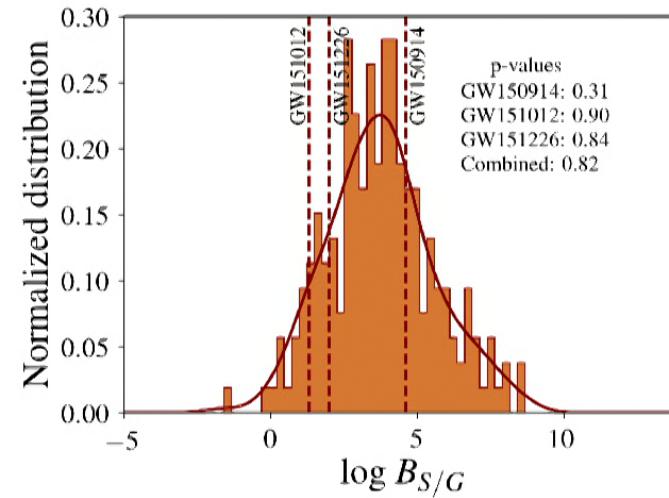
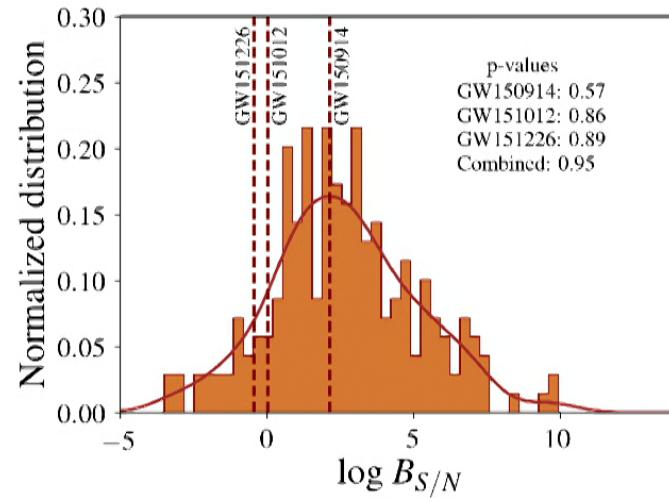
- $t_0$  prior is  $[t_{\text{trig}} + 4\tau_{220}, t_{\text{trig}} + 4\tau_{220} + 0.5]$
- QNM decay time  $\tau_{220}(M_f, a_f, z)$  can be obtained from fitting formula  
[Berti et al., Class. Quantum Grav. 26, 163001 \(2009\)](#)
- Take  $M_f, a_f, z$  to be upper bound of 90 % confidence interval to have conservatively large value for  $\tau_{220}$
- All events info are obtained from O2 catalog paper

Event	$m_1^{\text{src}}/M_\odot$	$m_2^{\text{src}}/M_\odot$	$\mathcal{M}^{\text{src}}/M_\odot$	$\chi_{\text{eff}}$	$M_f^{\text{src}}/M_\odot$	$a_f$	$E_{\text{rad}}/(M_\odot c^2)$	$\ell_{\text{peak}}/(\text{erg s}^{-1})$	$D_L/\text{Mpc}$	$z$	$\Delta\Omega/\text{deg}^2$
GW150914	$35.4^{+1.8}_{-3.0}$	$30.3^{+2.9}_{-4.3}$	$28.5^{+1.5}_{-1.4}$	$-0.03^{+0.10}_{-0.12}$	$62.7^{+3.2}_{-2.9}$	$0.68^{+0.05}_{-0.05}$	$3.1^{+0.4}_{-0.4}$	$3.6^{+0.4}_{-0.4} \times 10^{56}$	$410^{+150}_{-170}$	$0.09^{+0.03}_{-0.04}$	194.01
GW151012	$22.6^{+15.0}_{-4.9}$	$13.9^{+3.8}_{-5.1}$	$15.2^{+1.7}_{-1.1}$	$0.04^{+0.27}_{-0.19}$	$35.4^{+10.4}_{-3.5}$	$0.67^{+0.11}_{-0.10}$	$1.6^{+0.5}_{-0.5}$	$3.3^{+0.8}_{-1.8} \times 10^{56}$	$1030^{+510}_{-480}$	$0.20^{+0.09}_{-0.09}$	1490.79
GW151226	$14.1^{+9.1}_{-3.5}$	$7.5^{+2.2}_{-2.5}$	$8.9^{+0.3}_{-0.3}$	$0.19^{+0.21}_{-0.12}$	$20.7^{+6.8}_{-1.6}$	$0.71^{+0.03}_{-0.05}$	$1.0^{+0.1}_{-0.2}$	$3.3^{+0.7}_{-1.7} \times 10^{56}$	$440^{+180}_{-190}$	$0.09^{+0.04}_{-0.04}$	1074.50

## Echoes search at O1 BBH detections - Bayesian evidences

Background calculation:

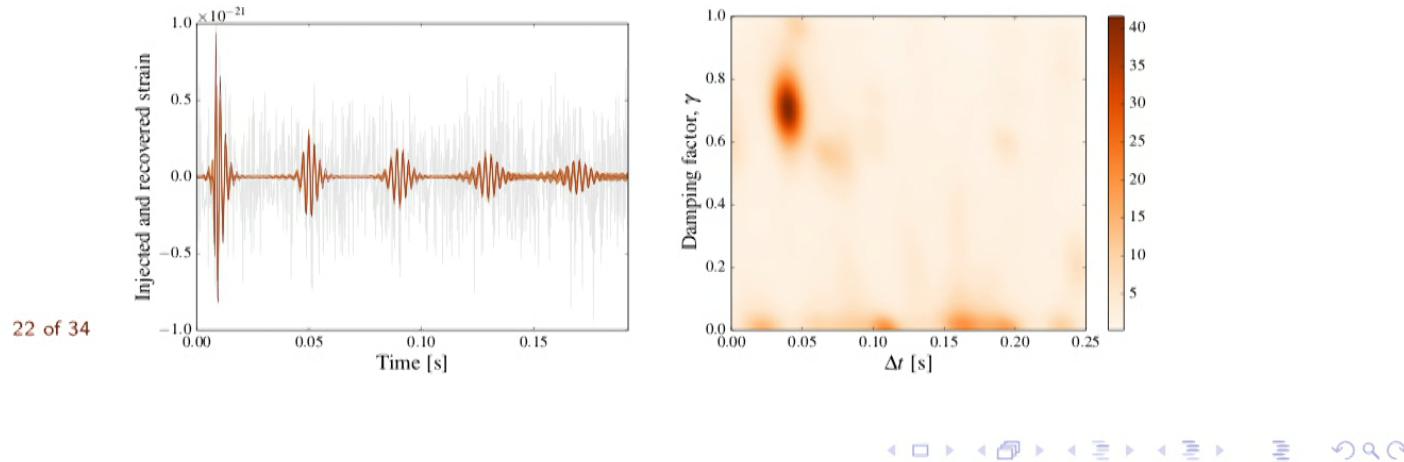
Analyze  $\sim 200$  8s-segments from GPSTime 1126073529 to 1126075217



All three events are well within the background.

**A morphology-independent data analysis method for detecting and characterizing gravitational wave echoes**

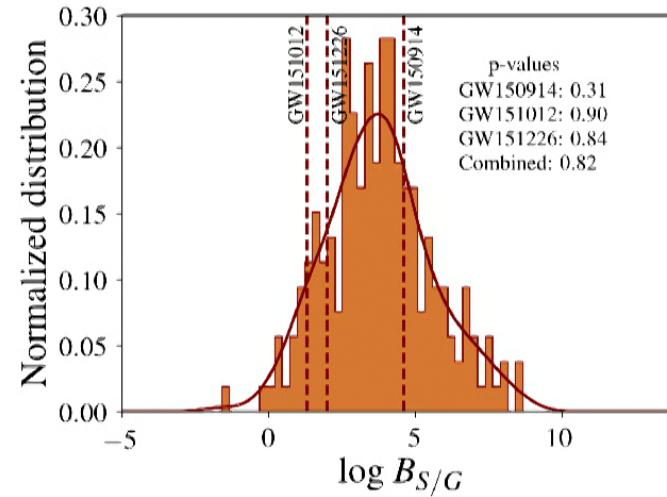
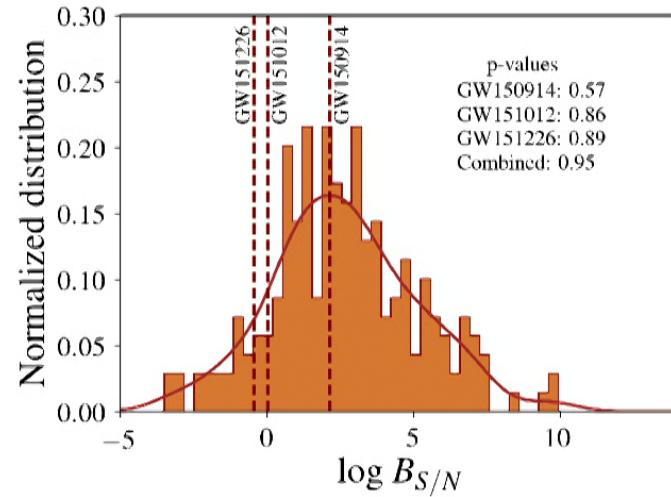
Ka Wa Tsang,<sup>1</sup> Michiel Rollier,<sup>1</sup> Archisman Ghosh,<sup>1</sup> Anuradha Samajdar,<sup>1</sup> Michalis Agathos,<sup>2</sup> Katerina Chatzioannou,<sup>3</sup> Vitor Cardoso,<sup>4</sup> Gaurav Khanna,<sup>5</sup> and Chris Van Den Broeck<sup>1,6</sup>



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# Cosmology



## Cosmology: Hubble's law

recession velocity of a galaxy in the local universe



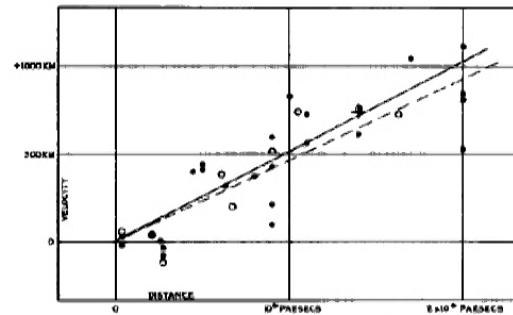
$$v_H = H_0 d$$



distance to the galaxy

Hubble parameter

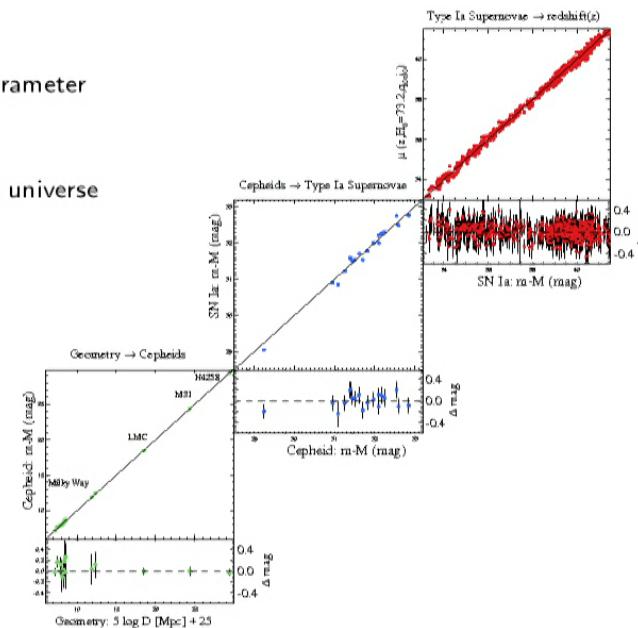
recession  $\rightarrow$  stretching of spacetime itself  $\rightarrow$  expansion of the universe  
usually measured as a cosmological redshift  $v_H = c z$



Edwin Hubble, *Proc. Nat. Acad. Sciences.* (1929)

Note: significant overestimate!

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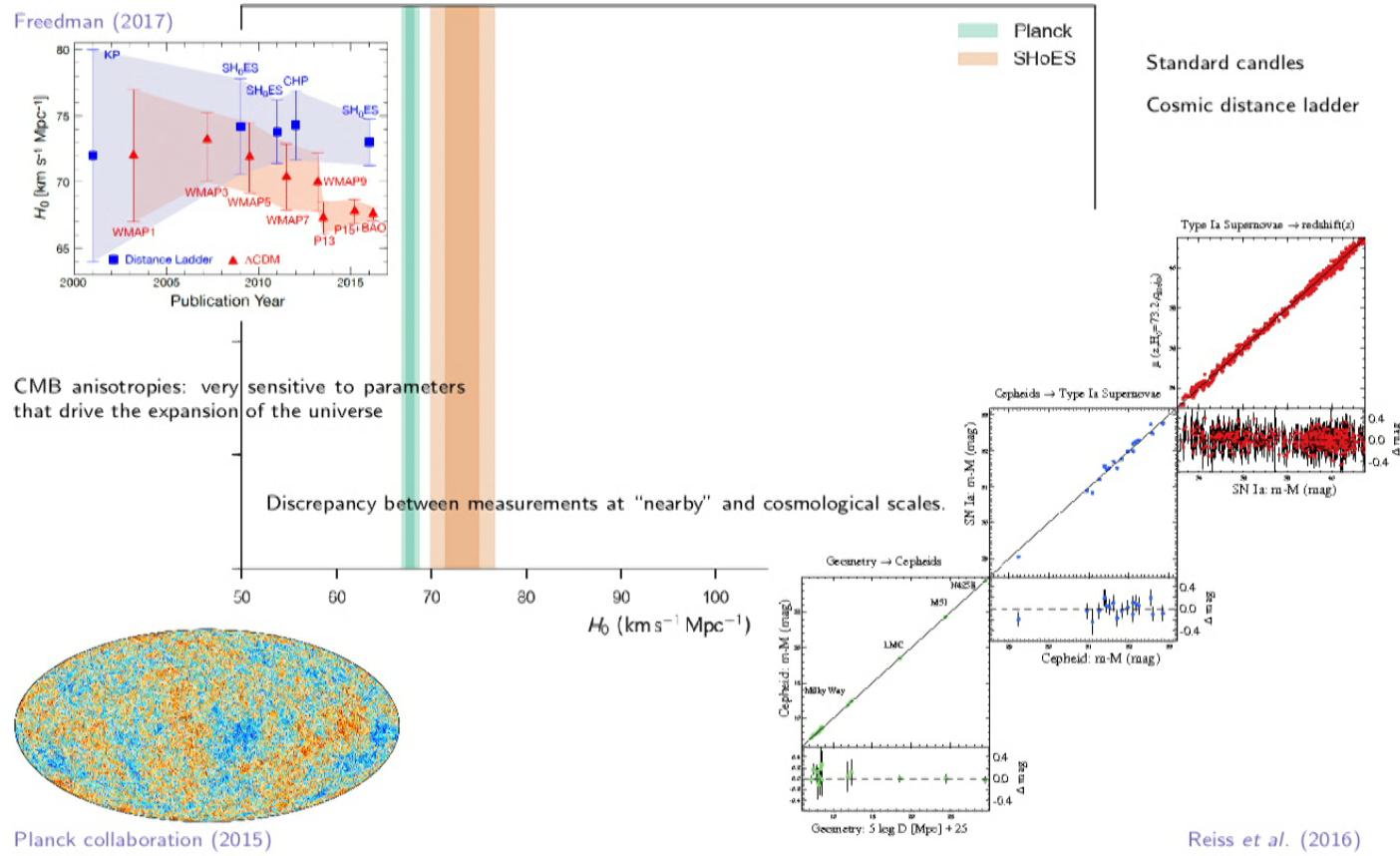


Cosmic distance ladder: Reiss et al. (2016)

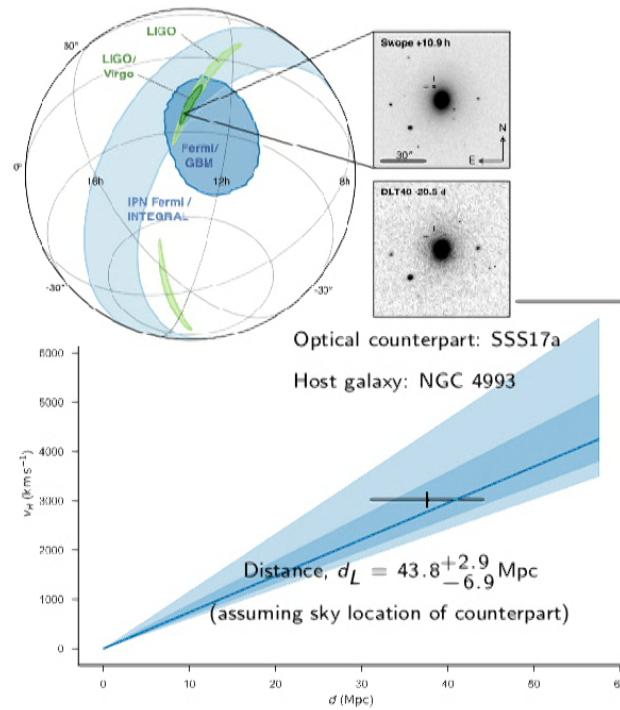


## State-of-the-art measurements of $H_0$

Two contrasting methods applied on nearby and very distant cosmological scales



## Hubble parameter with GW170817



Independent of any distance ladder!

Abbott et al. *Astrophys. J.* **848** #2, L12 (2017); LSC-EPO  
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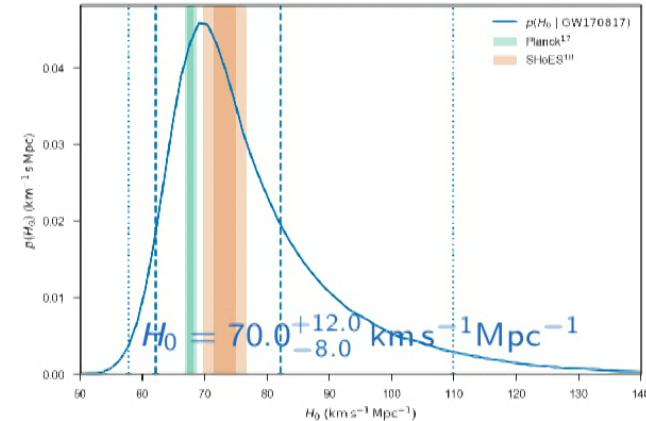
$$\text{observed } v_{\text{recession}} = v_H + v_{\text{peculiar}}$$

universe is not homogeneous at small scales:  
galaxies attracted towards local matter overdensities

$$\text{NGC 4993: } v_{\text{recession}} = 3327 \pm 72 \text{ km s}^{-1}$$

Correct for peculiar velocity of group of galaxies

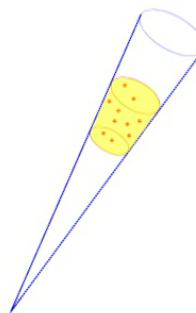
$$v_H = 3017 \pm 166 \text{ km s}^{-1}$$



Abbott et al. *Nature* **551** #7678, 85-88 (2017)

## Era of precision GW-cosmology ahead?

Multiple observations with transient counterparts.

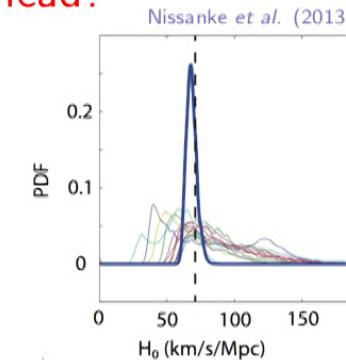


A fully statistical analysis using cross-correlation with a galaxy catalog in absence of a transient optical counterpart.

applicable also for binary black holes

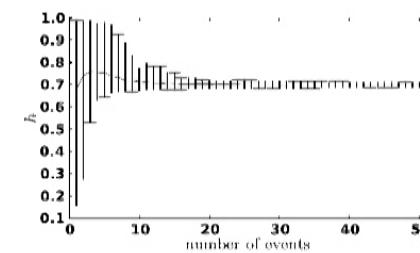
extension to other cosmological parameters?

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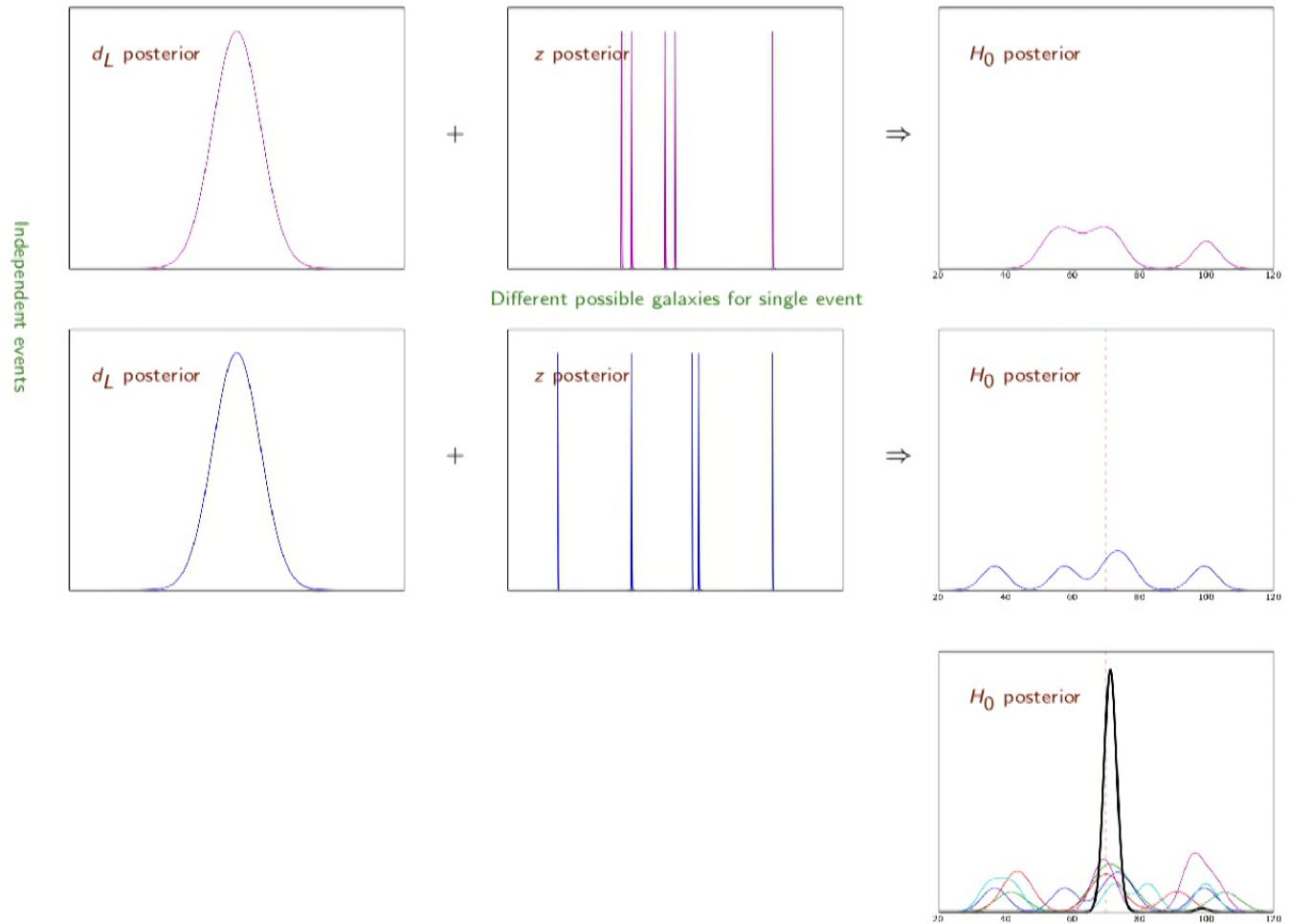


Narrow beam with potential host galaxies around optical counterpart if host galaxy not uniquely identified.

Schutz (1986)



Del Pozzo (2012)



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## $H_0$ -statistical: selection effects

$$d_L H_0 \approx zc$$

### GW selection effects

threshold SNR → interferometer horizon  
only nearby signals detected

### EM selection effects

depth of telescope  
incomplete galaxy catalogues

Correct for / take into account possible contribution of galaxies missing from catalogue.

Detection efficiency:

$$\mathcal{N}_{\text{eff}}(\Omega) = \int_{\mathcal{E}_{\text{det}}} d\mathcal{E} \int d\theta p(\mathcal{E}|\theta, \Omega, \mathcal{H}, \mathcal{I}) p(\theta|\Omega, \mathcal{H}, \mathcal{I})$$

Integrated method of taking into account both effects.

Abbott *et al.* Nature 551 #7678, 85-88 (2017)

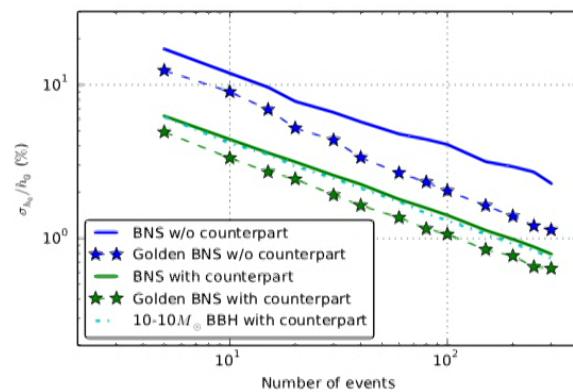
Mandel, Farr, Gair (2018); Chen *et al.* (2017)

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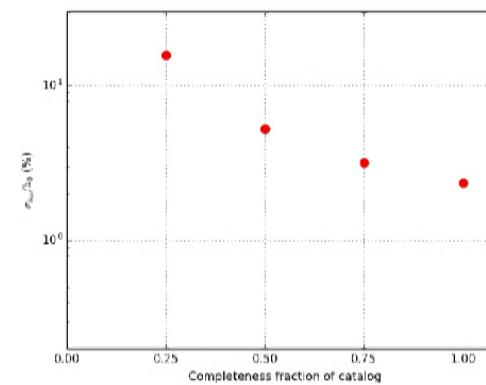
Messenger & Veitch (2013); Gray *et al.* (in prep.)

## $H_0$ -statistical: results on simulations

Chen et al. (2017)



Sur (2017, Masters thesis), Gray et al. (in prep.)

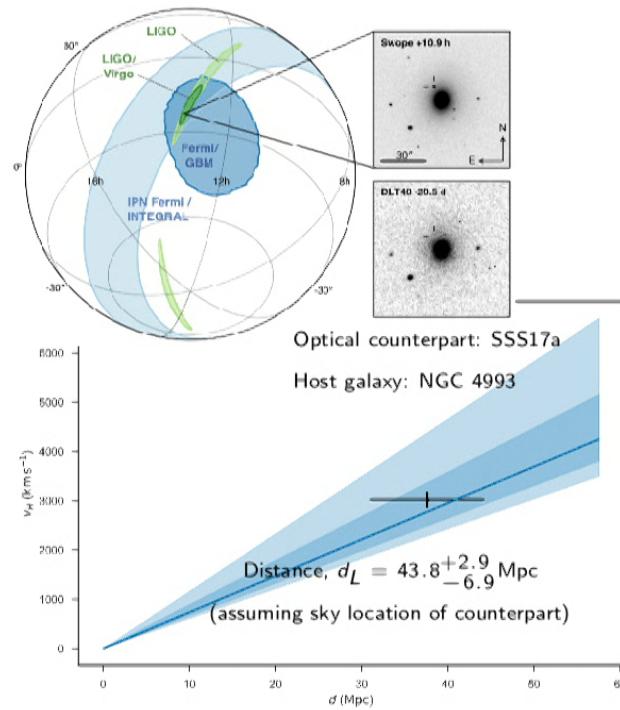


Incomplete galaxy catalogue

Ajith, Brady, Chen, Datrier, Del Pozzo, Fishbach, Gair, Ghosh, Gray, Hendry, Holz, Magaña-Hernandez, Messenger, Qi, Samajdar, Sur, Van Den Broeck, Veitch

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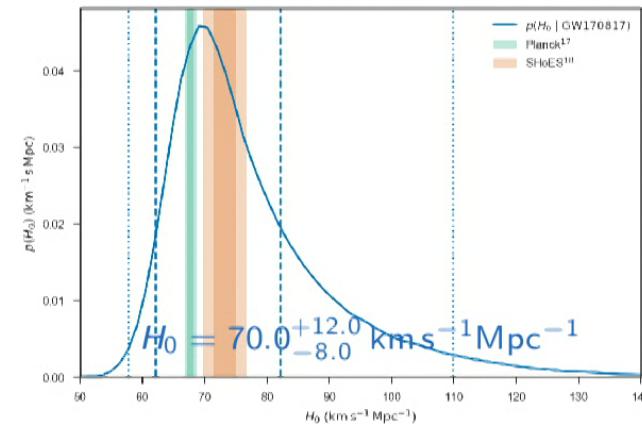
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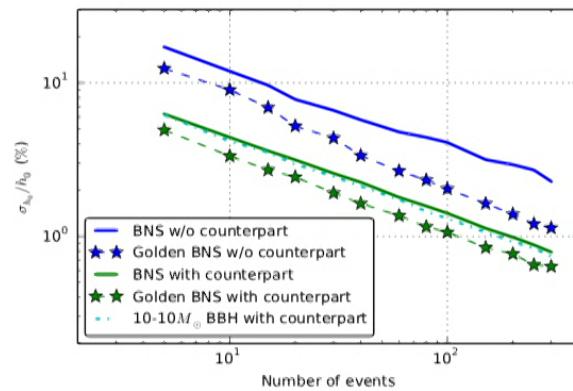
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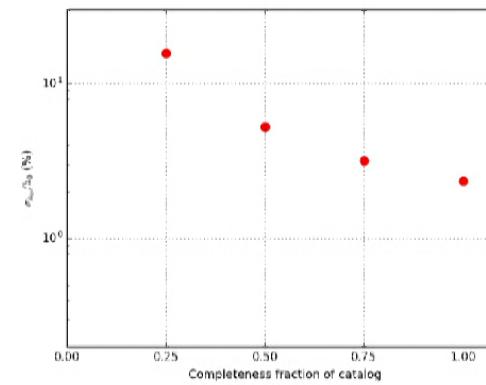
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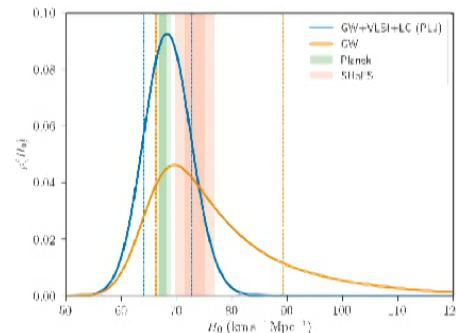
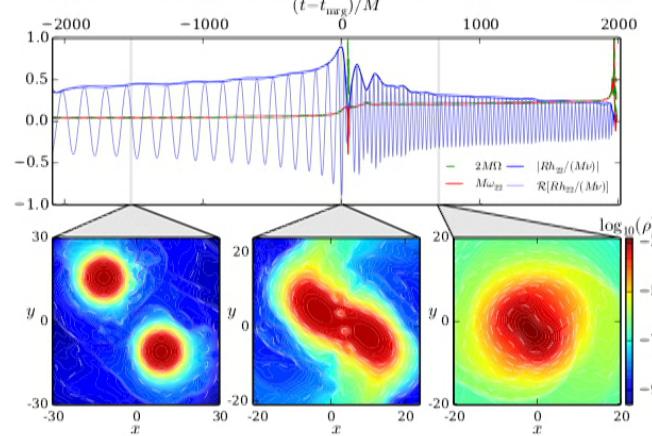
## Ongoing and future work

- Fold in probabilities of galaxies hosting the sources.
  - Luminosity weighting.
  - Astrophysically-motivated weighting of host galaxies?
- Going beyond  $H_0$ ?
  - Caveat: incompleteness of galaxy catalogues.
  - Sources correlated with visible matter distribution?
  - Cluster catalogues  $\Rightarrow$  probability density in redshift space?
- Cross-correlation with clustering more effective?

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## Synergetic multimessenger science: road to the future

Bernuzzi et al. (2015): NR simulations → postmerger templates?

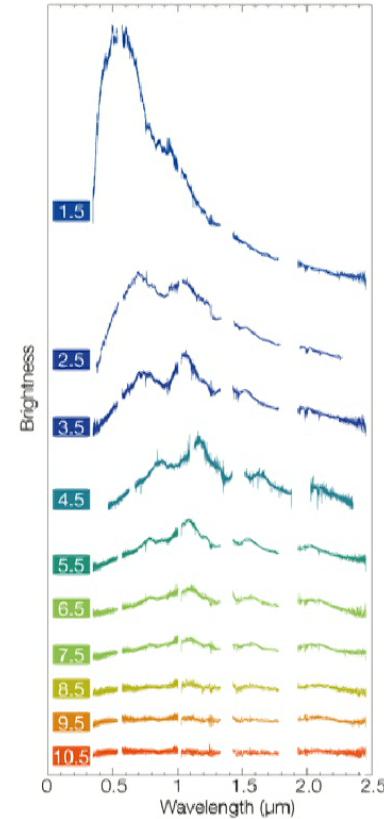


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Hotokezaka et al. (2018): jet → inclination!

**GW & EM!**

Kilonova → NS-EoS?



Pian et al. (2017)

THANK YOU!



