

Title: PSI 2018/2019 - Condensed Matter Review - Lecture 8

Date: Jan 16, 2019 11:30 AM

URL: <http://pirsa.org/19010039>

Abstract:

$$H = -J \sum_{\langle ij \rangle} s_i s_j - h \sum_i s_i$$

$$\{s_i\} \rightarrow \{-s_i\}$$

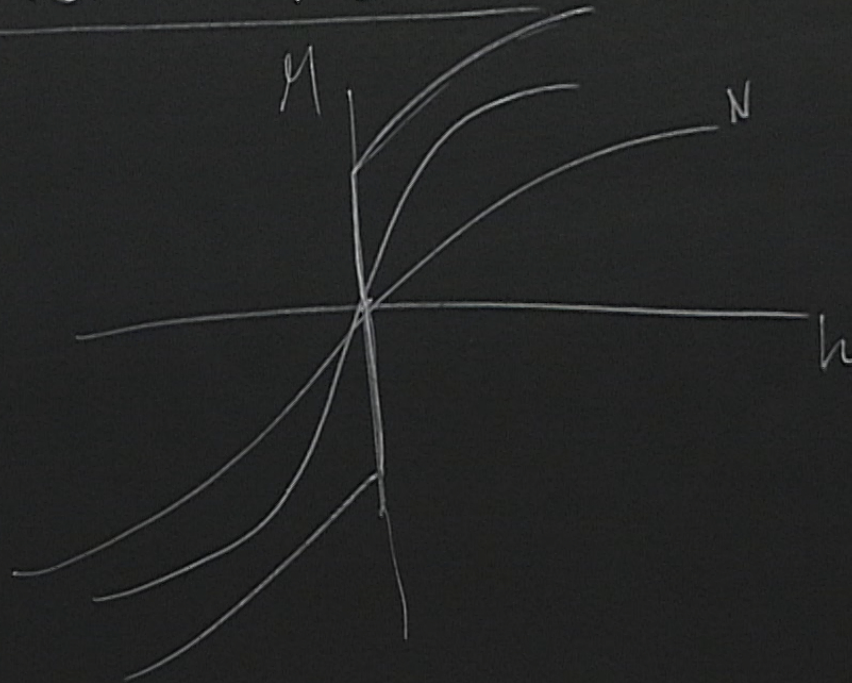
$$P(\{s_i\}) = \frac{1}{Z} e^{-\beta H(\{s_i\})}$$

$$M = \langle s_i \rangle = \sum_{\{s_i\}} P(\{s_i\}) s_i = 0!$$

In TDL the  
P is NOT Boltzmann

# Method of Sources

the  
of Boltzmann



$$= \frac{1}{\lambda} e^{-\beta H(\{s_i\})}$$

$= 0!$

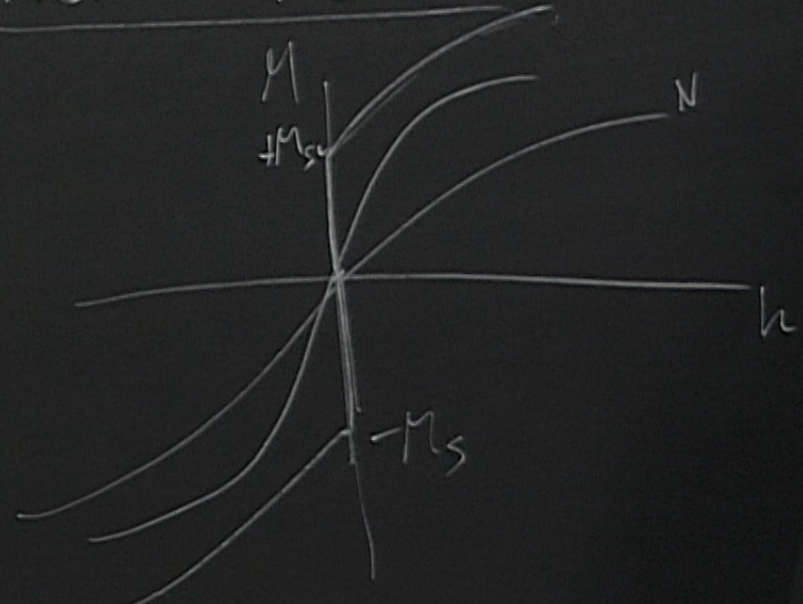
In TDL the

$P$  is NOT Boltzmann

$$M = \frac{1}{N} \frac{\partial F}{\partial h}$$

$$\lim_{h \rightarrow 0} \lim_{N \rightarrow \infty} M$$

### Method of Sauer



$$H = -J \sum_{\langle ij \rangle} s_i s_j - h \sum_i s_i$$

$$d\{s_i\} \rightarrow \{-s_i\}$$

$$P(\{s_i\}) = \frac{1}{Z} e^{-\beta_0 H(\{s_i\})}$$

$$M = \langle s_i \rangle = \sum_{\{s_i\}} P(\{s_i\}) s_i = 0!$$

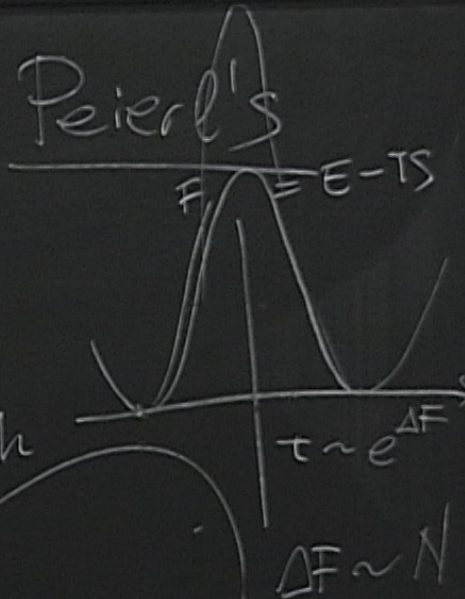
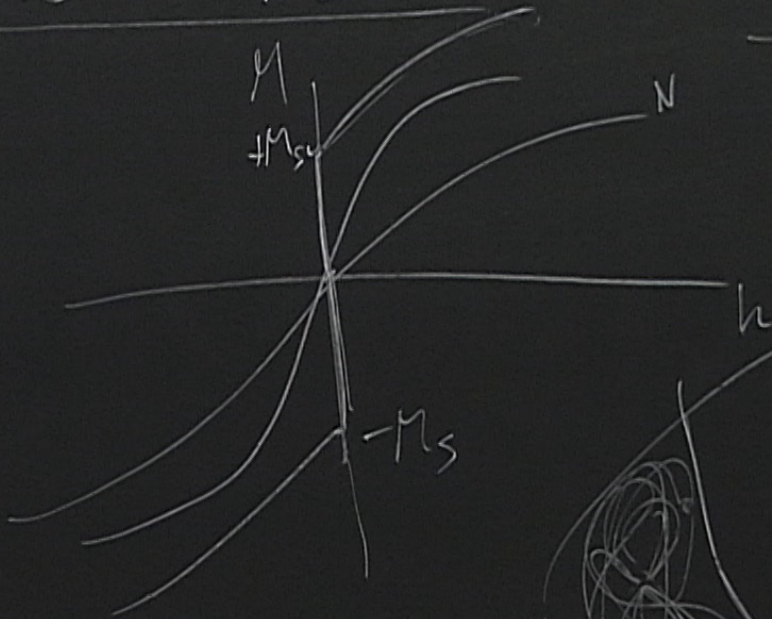
Ergodicity Broken

$$M = \frac{1}{N}$$

In TD

P is

# Method of Sources



Witzmann

$\lim_{N \rightarrow \infty} M$

QCD  $\rightarrow$  Lattice Gauge Theory

$Z_2$  symmetry

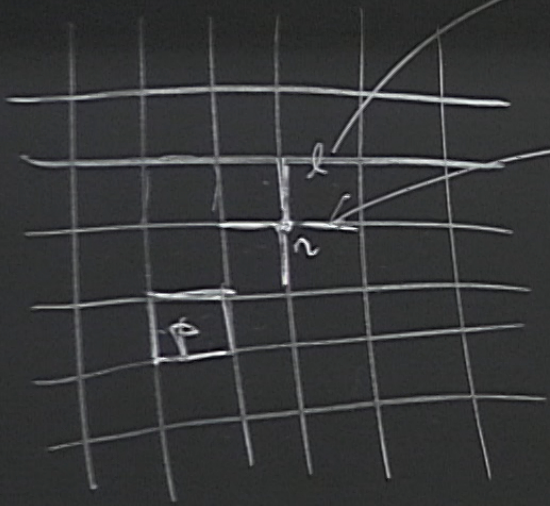
Globel

$Z_n$  Gauge Symmetry

local

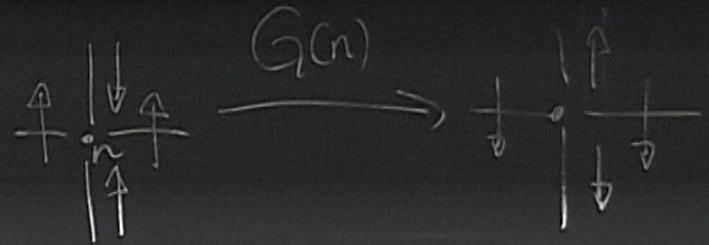
Wegner

xy



$$s_e = \pm 1$$

$$l_n \in \mathbb{Z}$$



$$H = -J \sum_p B_p$$

$$B_p = \underbrace{ssss}_{\text{around } p}$$

$\mathbb{Z}_2$  lattice gauge theory



$$H = \frac{-J \sum_p s_i s_{i+p} - h' \sum_e s_e}{\sum_{\{s_e\}} e^{-\beta H(\{s_e\})}} \quad h \equiv \beta h'$$

$$\begin{aligned} \langle s_e \rangle_h &= \frac{1}{Z} \sum_{\{s_e\}} s_e e^{-\beta H(\{s_e\})} \\ &= \frac{1}{Z} \sum_{\{s_e\}} s_e \exp[\beta J \sum_p s_i s_{i+p} + h \sum_e s_e] \end{aligned}$$

A gauge symm. CANNOT be spontaneously broken  
(local)

## ELITZUR'S THEOREM

$$\{S_e\} \xrightarrow{G(x)} \{S'_e\}$$

$$S(\varrho_n) \rightarrow S'(\varrho_n) = -S(\varrho_n)$$

A gauge symm. CANNOT be spontaneously broken  
(local)

$$H = -J \sum_p S$$

### ELITZUR'S THEOREM

$$\{s_e\} \xrightarrow{G(n)} \{s'_e\}$$

$$s(\ell_n) \rightarrow s'(\ell_n) = -s(\ell_n)$$

$$h \sum_e s_e = h \sum_e s'_e + h \sum_e \delta s_e \rightarrow -2S(\ell_n)$$

$$\langle s_e \rangle_n = \frac{1}{Z} = \frac{1}{Z} \sum_{\{s_e\}}$$

around 1

$$H = \frac{-J \sum_p s_i s_{i+p} - h \sum_e s_e}{k \equiv \beta h'}$$

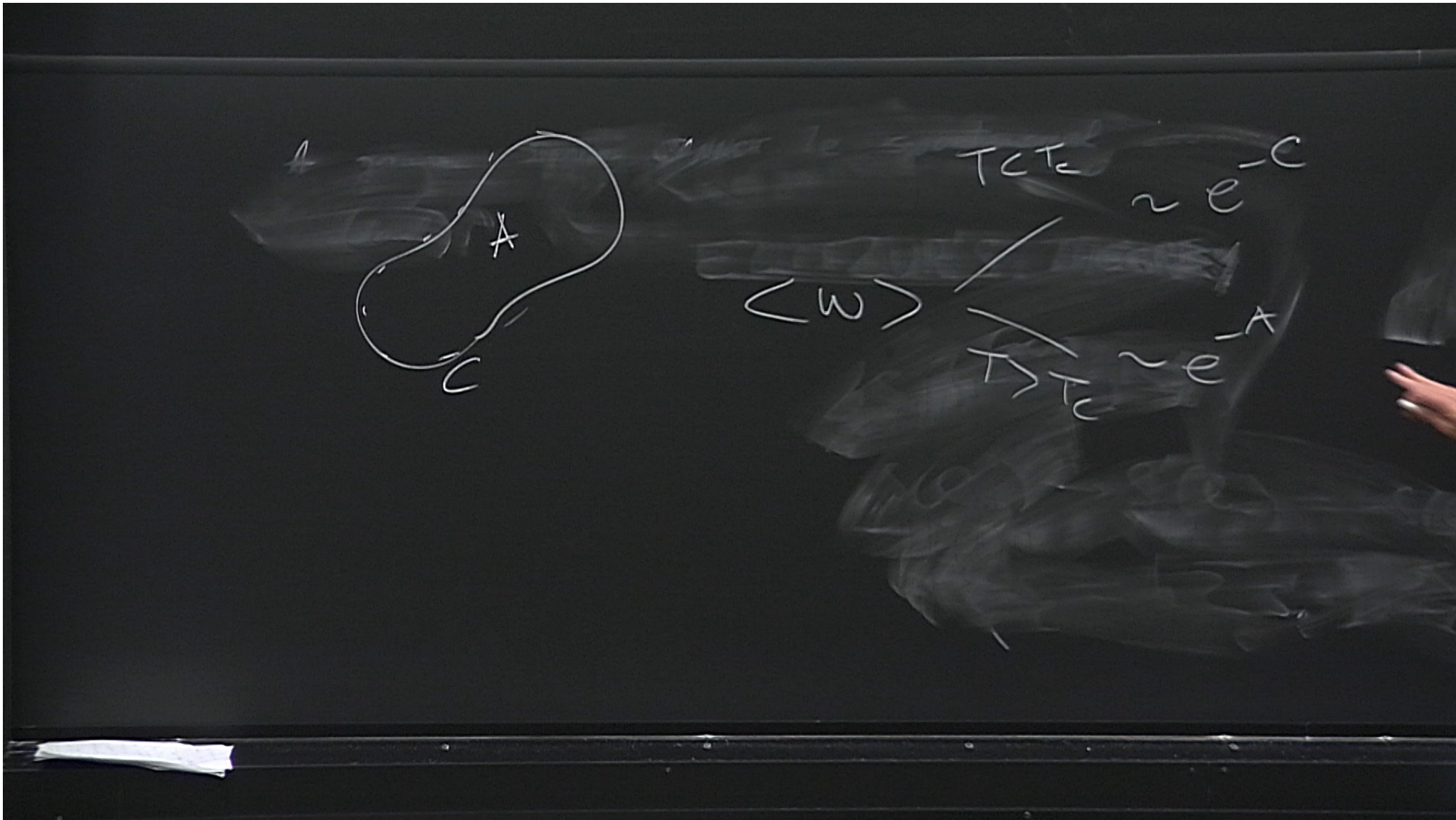
$$\langle s_e \rangle_h = \frac{1}{Z} \sum_{\{s_e\}} s_e e^{-\beta H(\{s_e\})}$$

$$= \frac{1}{Z} \sum_{\{s_e\}} s_e \exp \left[ \beta J \sum_p s_i s_{i+p} + h \sum_e s_e - h \sum_e s_i s_{i+1} \right]$$

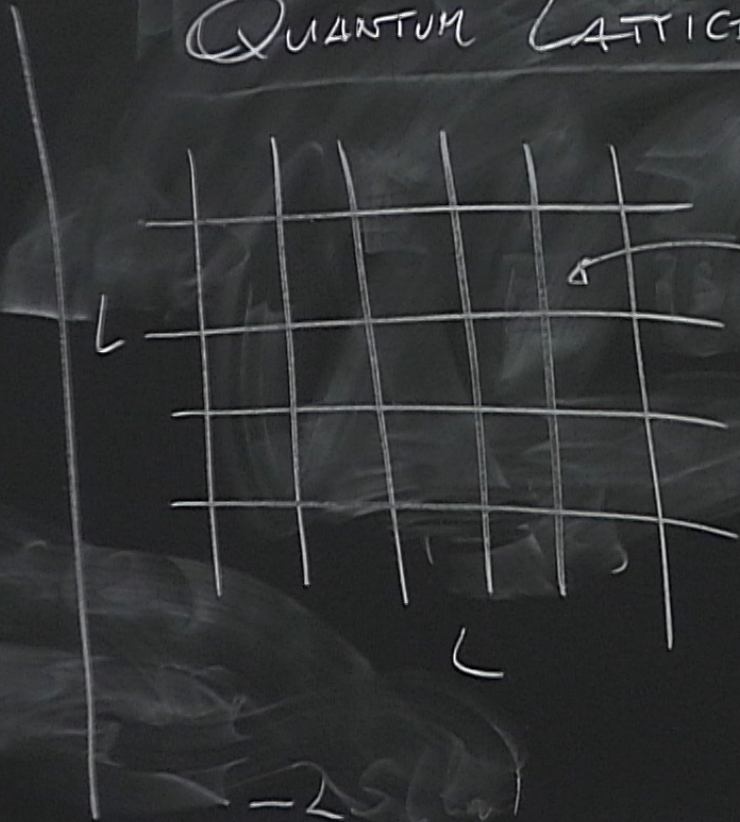
$$= -S(h)$$

$$+ h \sum_e s_i s_{i+1} - 2S(h)$$

$$\begin{aligned}
 & \left| \langle S_e \rangle_h - \langle -S_e \rangle_h \right| = \left| \langle -S_e \left( e^{\frac{-h \sum_{i=1}^S S_i}{2en^2} - 1} \right) \rangle_h \right| \\
 & \stackrel{||}{=} 0 \quad h \rightarrow 0 \leq (e^{4dh} - 1) \left| \langle -S_e \rangle_h \right| \leq e^{4dh} \\
 & \quad \quad \quad h \rightarrow 0
 \end{aligned}$$

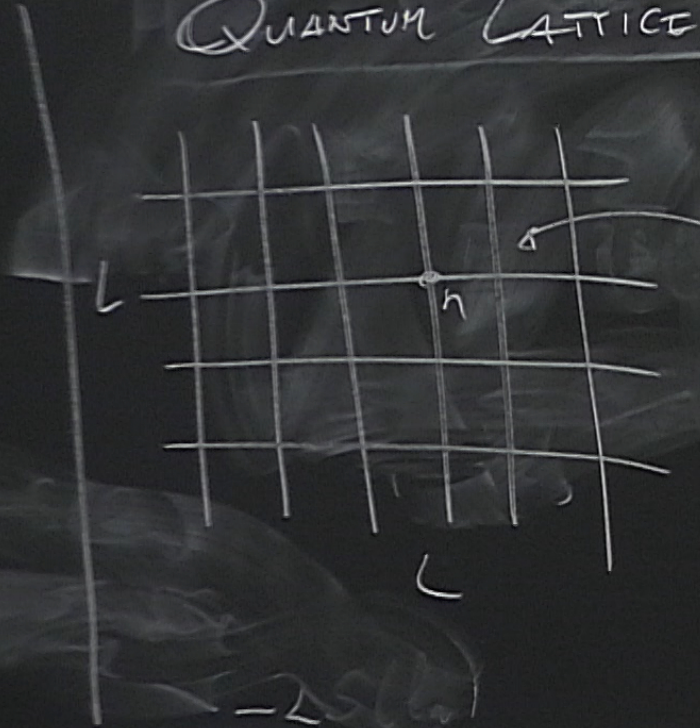


# QUANTUM LATTICE GAUGE THEORY



# links  $2L^2$   
 $\mathbb{C}^2 \cong \text{span} \{ |0\rangle, |1\rangle \}$

# QUANTUM LATTICE GAUGE THEORY



# links  $2L^2$

$$\mathbb{C}^2 \cong \text{span} \{ |0\rangle, |1\rangle \}$$

$$\mathcal{H} = \mathbb{C}^2 \otimes 2L^2$$

$$\dim \mathcal{H} = 2^{2L^2}$$

$$A(n) = \prod_{l \in n} \sigma_l^x$$

$$\mathcal{H}_{\text{gauge}} = \{ |\psi\rangle \in \mathcal{H} \mid A_n |\psi\rangle = |\psi\rangle \forall n \}$$



$$A_n P_n |X\rangle = A_n \frac{1 + A_n}{2} |X\rangle = P_n |X\rangle$$

$$P_n^2 \stackrel{?}{=} P_n$$

$$P_n^2 = \left( \frac{1 + A_n}{2} \right)^2 = \frac{1}{4} (1 + 2A_n + A_n^2)$$

$$P_k \dots P_n P_n |X\rangle$$

$$= \frac{1 + A_n}{2} = P_k$$

# GAUGE THEORY

$$\mathbb{C}^2 \cong \text{span} \{ |0\rangle, |1\rangle \}$$

# links  $2L^2$

$$\mathcal{H} = \mathbb{C}^2 \otimes 2L^2$$
$$\dim \mathcal{H} = 2^{2L^2}$$

$$A(n) = \prod_{l \in n} \sigma_l^x$$

$$\mathcal{H}_{\text{gauge}} = \left\{ |14\rangle \in \mathcal{H} \mid \underbrace{A_n |14\rangle = |14\rangle}_{L^2+1} \forall n \right\}$$
$$\dim \mathcal{H}_{\text{gauge}} = 2$$