

Title: Renormalization and Effective Field Theory - Lecture 11

Date: Dec 05, 2018 10:30 AM

URL: <http://pirsa.org/18120005>

Abstract:

Chern-Simons

$$\mathcal{E}(M^3) = \Omega^2(M^3) \otimes \mathfrak{g}[1]$$

$$S_{CS} = \frac{1}{2} \int \langle \alpha, d\alpha \rangle + \frac{1}{6} \int \langle \alpha, [\alpha, \alpha] \rangle \quad \text{for } \alpha \in \mathcal{E}(M^3)$$

Gauge fix

Choose Lagrangian $\mathcal{L} \subseteq \Omega^2(M^3) \otimes \mathfrak{g}[1]$

$$\int_{\mathcal{L}} e^{S_S/\hbar}$$

Require $\langle \alpha, d\alpha \rangle$ is a non-degen. quadratic form on \mathcal{L}

Hodge Theory

$$d^{\sharp} : \Omega^k(M) \rightarrow \Omega^{k-1}(M)$$

[in local coords

$$d^{\flat} \sim g^{ij} \frac{\partial}{\partial x_i} (dx_j)$$

$(dx_i$ removes dx_j]

$$d^{\sharp} = \star d \star$$

$$\Delta = [d, d^{\sharp}]$$

$$= d d^{\flat} + d^{\sharp} d$$

any

$$\Omega^k(M) \rightarrow \Omega^{k-1}(M)$$

coords

$\frac{\partial}{\partial x_i} (dx_j)$
removes dx_j

$$d^* = \star d \star$$

$$\Delta = [d, d^*] = dd^* + d^*d$$

M compact

$$\Omega^k(M) = \text{Im } d^* \oplus \text{Im } d$$

$$\mathcal{H} \cong H^k(M) \oplus \mathcal{H} = \text{Ker } \Delta$$

finite dimensional

We will integrate over

$$\text{Im } d^{\nu} \subseteq \text{Im } d \oplus \text{Im } d^{\nu} = \mathcal{H}^{\perp}$$

Left with: an "effective measure" on space
of solⁿs to the EOM.
Tangent space of this is $\mathcal{H}(\otimes \mathfrak{g}[1])$

If $\alpha \in \text{Im } d^0 \oplus \mathfrak{g}(1)$

$\int \alpha d\alpha$ is a non-degenerate quadratic form

α a 1-form in $\text{Im } d^0$

$$\alpha = \sum \alpha_i dx_i$$

$$\text{Then } d^0 \alpha = \sum g_{ij} \frac{\partial}{\partial x_j} \alpha_i = 0$$

Lorenz gauge

$(x \sim \mathbb{R}^n)$ finite dimension

Heat Kernel

Simons

Define $K_t(x, y) \in \Omega^3(m \times m)$

So that:

1) $\partial_t K_t(x, y) + \Delta_x K_t(x, y) = 0$

2) $K_0(x, y) = \delta_{x=y}$

$$K_t(x, y) = \sum_{|I|+|J|=3} f_{I, J}(x, y, t) dx^I dy^J$$

With flat metric

$$K_t = e^{-\|x-y\|^2/t} t^{-3/2} (d(x_1-y_1) d(x_2-y_2) d(x_3-y_3))$$

With flat metric

$$K_t = e^{-\|x-y\|^2/t} t^{-3/2} (dx_1 - dy_1) dx_2 - dy_2 dx_3 - dy_3$$

$$- \Delta_x K_t = \Delta_y K_t$$

$$- (d_x + d_y) K_t = 0$$

$$\text{As, } \partial_t [(d_x + d_y) K_t] = (d_x + d_y) \Delta_x K_t = \Delta_x (d_x + d_y) K_t$$

$$/ [d, d'], d] = 0$$

$$t^{-3/2} \left(d(x_1 - y_1) d(x_2 - y_2) d(x_3 - y_3) \right)$$

Heat eqn
Initial condition

$$(d_x + d_y) \Delta_x K_t = \Delta_x (d_x + d_y) K_t$$

$$[d, d'], d = 0$$

Propagator

$$P \in \text{Im } d^v \otimes \text{Im } d^i$$

$$\text{is } P(x, y) = \int_0^\infty d_x^v K_t(x, y)$$

Key properties

1) $d_x^v P(x, y) = 0$ and $d_y^v P(x, y) = 0$ ✓

2) If $\alpha \in \text{Im } d^i$, then $d_y^i P(x, y) \alpha(y) = \alpha(x)$

$$dx \int_y^\infty d_x^\dagger K_t(x, y) \wedge \alpha(y)$$

$$= \int_0^\infty \int_y^\infty ([d_x, d_x^\dagger] - d_x^\dagger dx) K_t(x, y) \wedge \alpha(y)$$

$$= \int_y^\infty d_x^\dagger dx K_t(x, y) \alpha(y) - \int_0^\infty \int_y^\infty d_t K_t(x, y) \alpha(y)$$

$$d_x \int_y^\infty d_x^* K_t(x, y) \wedge \alpha(y)$$

$$= \int_0^\infty \int_y^\infty ([d_x, d_x^*] - d_x^* d_x) K_t(x, y) \wedge \alpha(y)$$

$$= \underbrace{\int_y^\infty d_x^* d_x K_t(x, y) \wedge \alpha(y)}_{\rightarrow \text{problematic?}} - \int_0^\infty \int_y^\infty d_t K_t(x, y) \wedge \alpha(y)$$

$K_0 = \delta_{x=y}$
 and $K_\infty = \text{proj onto } \alpha$
 $= \int \delta_{x=y} \alpha(y) = \alpha(x) \checkmark$

$$\Delta \int_0^\infty K_t(x,y) \alpha(y) = \alpha(x)$$

$$\int K_t(x,y) \alpha = \Delta^{-1} \alpha$$

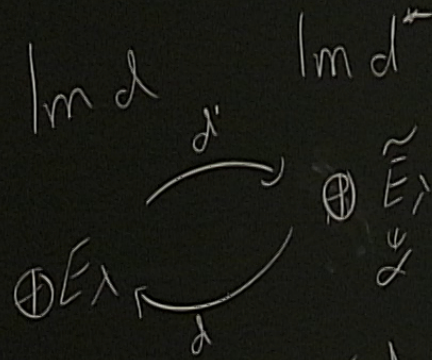
$$d(d^* \Delta^{-1} \alpha)$$

$$= (\Delta - d^* d) \Delta^{-1} \alpha$$

$$\int_x K_t(x,y) d_y \alpha = \alpha$$

$$= - \int d_x^* (d_y K_t(x,y)) \alpha$$

$$= \int d_x d_y K_t(x,y) \alpha = \alpha + \text{stuff in } \text{Im } d$$



Simons

$$(x, y) d_y \alpha$$

$$= \pm \int d_x d_x K_\epsilon(x, y) \alpha$$

$$= \pm \int \Delta_x K_\epsilon(x, y) \alpha + \text{things in Im } \alpha$$

$$\int d_x d_x' K_\epsilon(x, y) \alpha_y$$

$$= 0 \text{ using } (d_x + d_y') K_\epsilon = 0 \text{ and } d_y \alpha = 0 \text{ and IBP}$$

Propagation

$$P(\varepsilon, L) = \int_{\varepsilon}^L d^4x K_x dt$$

We can use the same Feynman rules as before

- Vertex $\int d^3x$
- edge : put

$$P(\varepsilon, L) = \int_{\varepsilon}^L d^4x K_x$$

$$W(P(\varepsilon, L), I^{CS}) (\alpha)$$

$$= \sum_{\text{conn graphs}} \frac{1}{h^{\# \text{ loops}}} \cdot \text{amplitude}$$

fn of $\alpha \in \Omega(M^3)_{\text{reg}}(1)$

Just as before

$\exists!$ set of counter terms, "purely singular" as a fn of ε

So $W(P(\varepsilon, L), I^{CS} - I^{\text{ct}}(\varepsilon))$ has an $\varepsilon \rightarrow 0$ limit

View

$P(\varepsilon, L)$ as in $\Omega^2(m \times m)$

Then

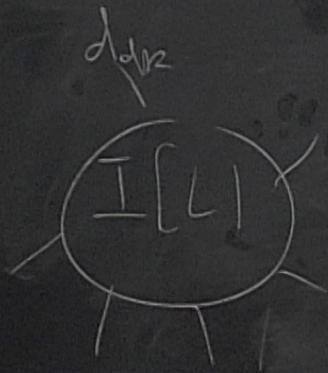
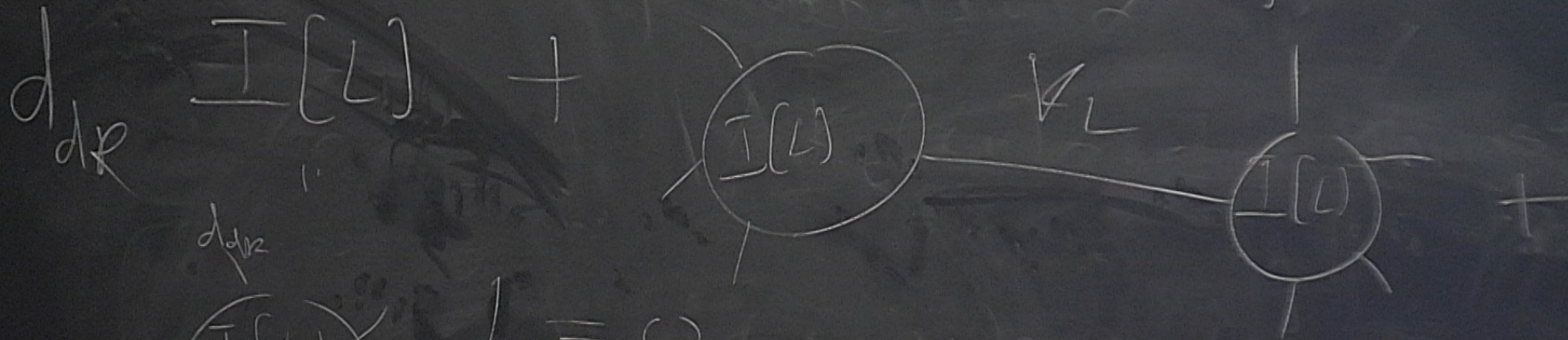
KEY FORMULA

$$\begin{aligned} & (d_x + d_y) P(\varepsilon, L) = -K_\varepsilon + K_L \\ \text{As, } & (d_x + d_y) \int_\varepsilon^L K_t(x, y) = \int_\varepsilon^L (\Delta_x - d_x^2 d_x - d_x^2 d_y) K_t(x, y) \end{aligned}$$

$$K^L - \left(\frac{d^2}{dx^2} - \frac{d^2}{dy^2} \right) K_t(x, y) = - \int_{\mathcal{E}}^L \partial_t K_t$$

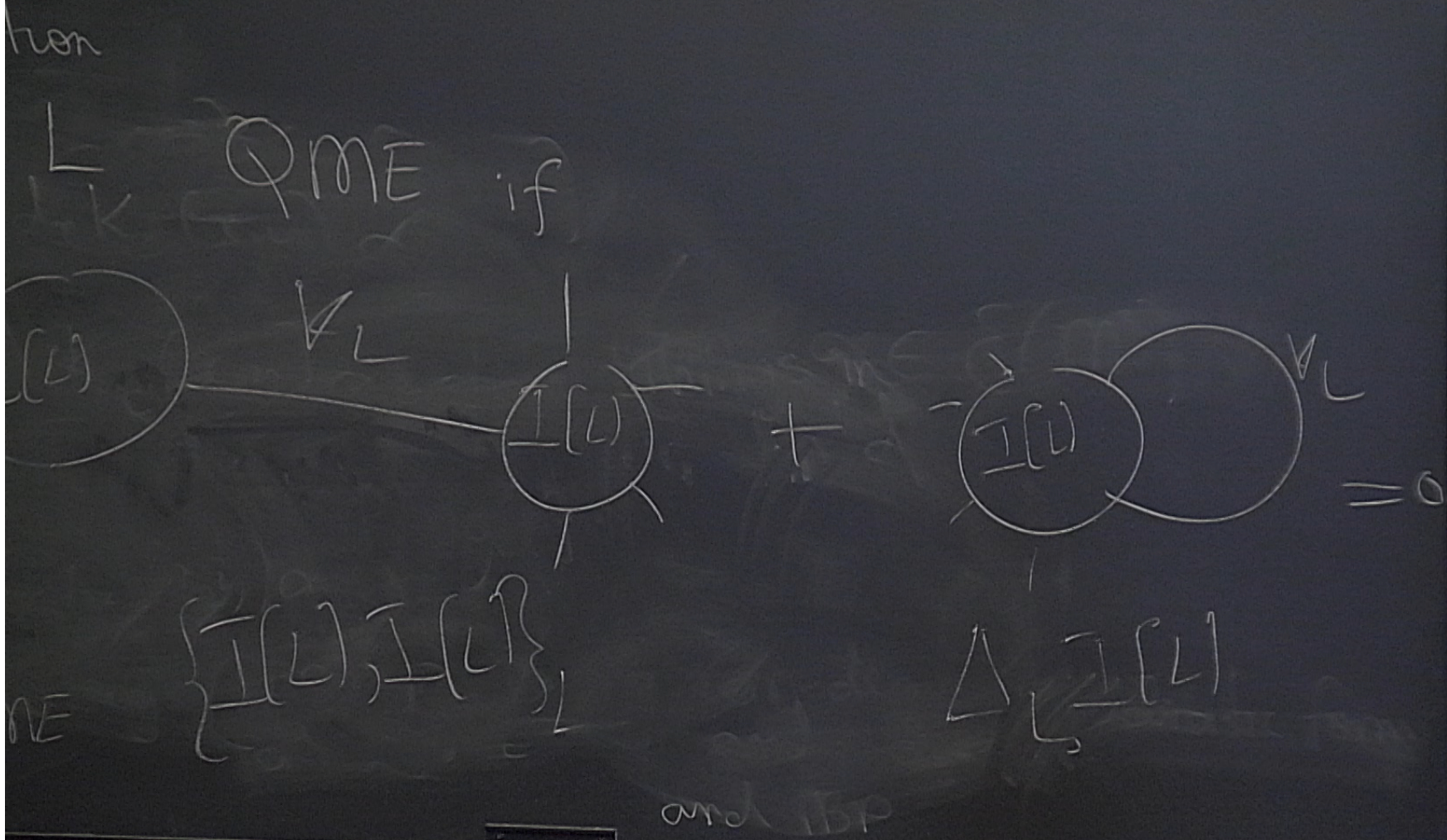
We say some interaction

$I[L]$ satisfies scale L QME if



$L = 0$
 This is ordinary QME

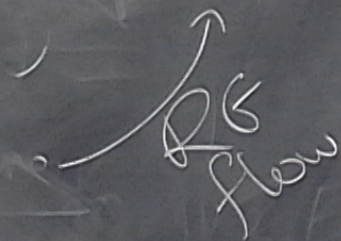
$\{I[L], I[L]\}$



Lemma

If $I[\varepsilon]$ satisfies scale ε QME

Then $W(P(\varepsilon, L), I[\varepsilon])$ satisfies scale L QME

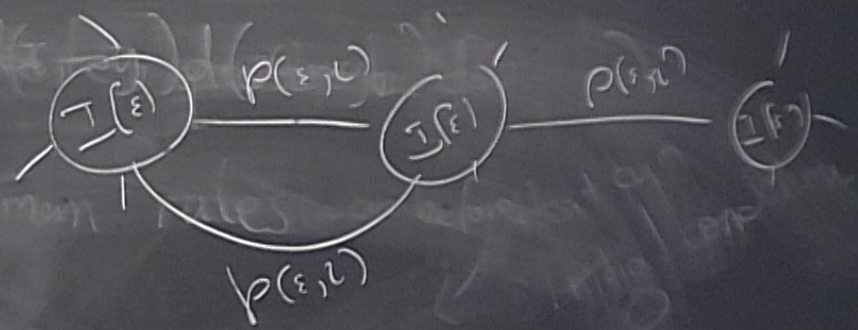


Proof:

Diagrams

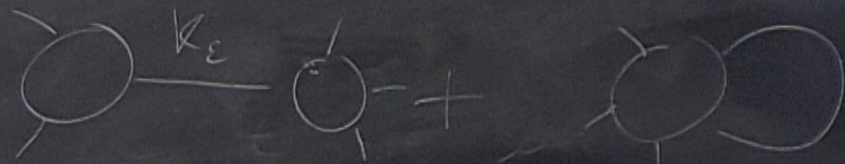
$$W(p(\varepsilon, L), I[\varepsilon]) = \sum_{\text{connected}} \dots$$

$$dW(p(\varepsilon, L), I[\varepsilon]) = \sum \dots$$

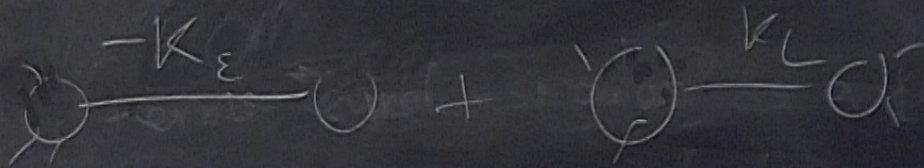


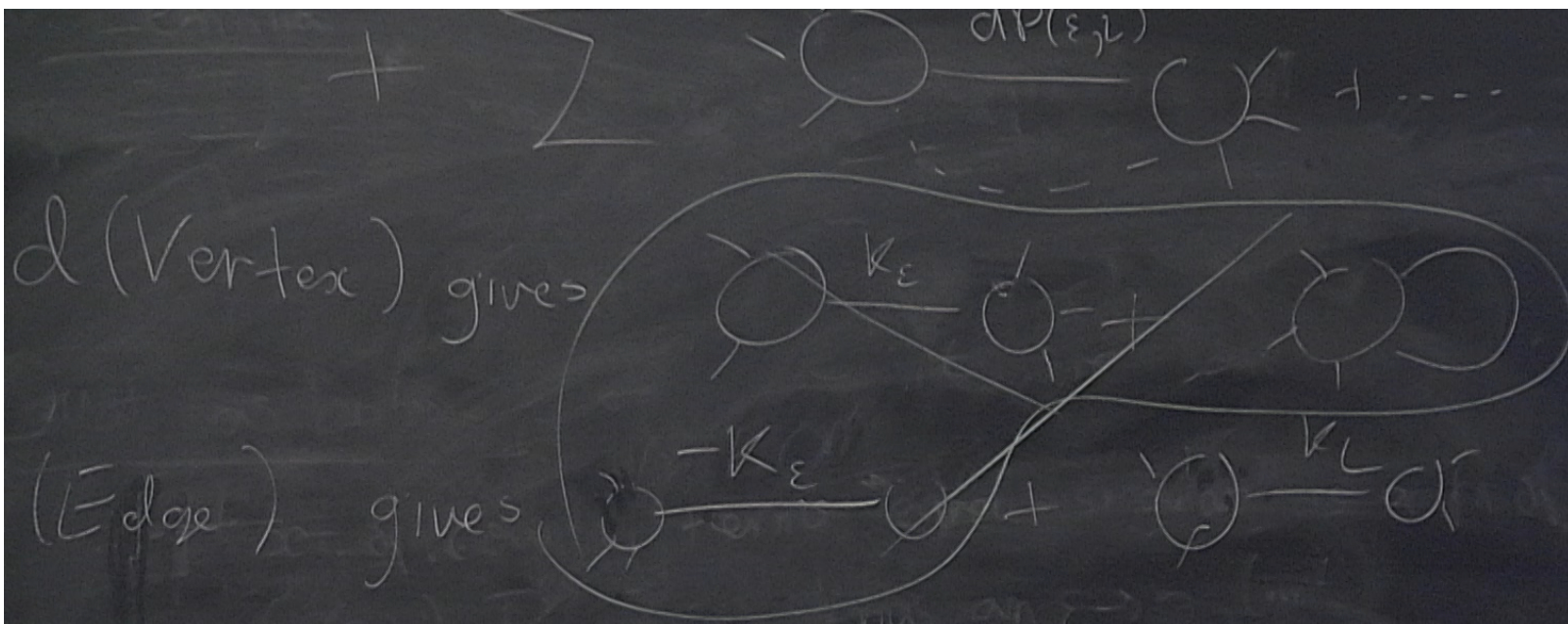
$$+ \sum \text{diagram} \xrightarrow{dP(\epsilon, L)} \text{diagram} + \dots$$

$d(\text{Vertex})$ gives

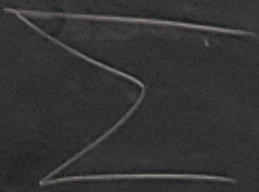


$d(\text{Edge})$ gives



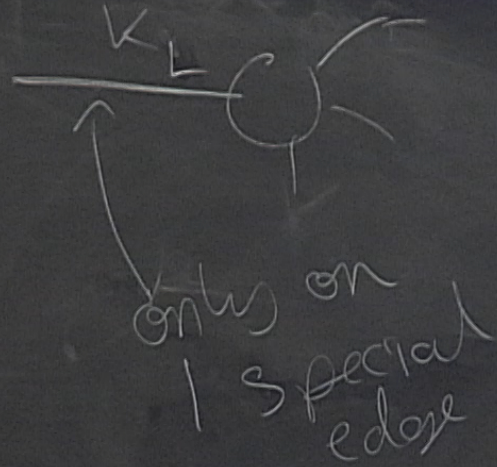
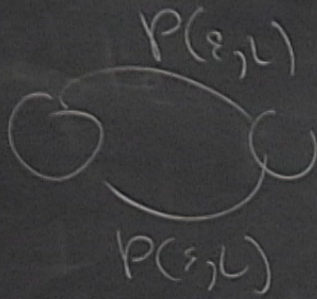


Left with



Conn. graphs

+ 1 special edge



Left with

