

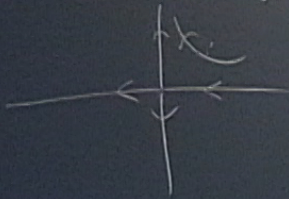
Title: Renormalization and Effective Field Theory - Lecture 5

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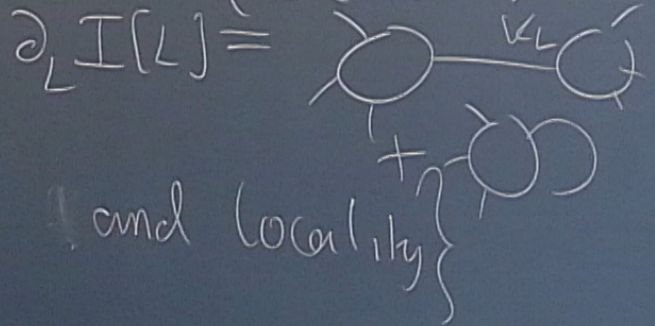
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Abstract:

Last time
RG trajectories



Theory = $\{ \mathbb{I}[L](\varphi),$
fncts of a field
 $\varphi \in C^\infty(\mathbb{R}^n)$



{ Theorem }

Theories

$\mathbb{R}_{>0}$ acting

NON CANONICAL SCHEME DEPENDENT

Local Lagrangians

$$I = \sum h^k I_k$$

$$I_k = \int \varphi D \varphi \dots D \varphi$$

$\mathbb{R}_{>0}$ action

Lagrangian for free theory

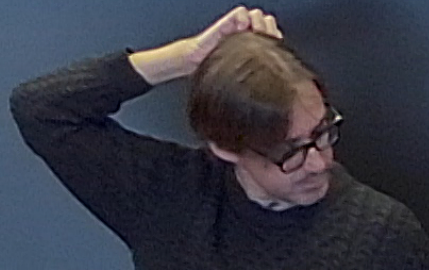
$$\int \psi \Delta \psi d^n x$$

$$\text{Let } R_\lambda(\psi)(x) = \lambda^{n/2-1} \psi(\lambda x)$$

R_λ acts on fns of ψ

$$(R_\lambda^* I)(\psi) = I(R_\lambda \psi)$$

$$R_\lambda^* \int \psi \Delta \psi = \int \psi \Delta \psi$$



$$R_\lambda \int \varphi \Delta \varphi - \int \varphi \Delta \varphi$$

$$K_t(x, x') = t^{-n/2} e^{-\|x-x'\|^2/t}$$

$$R_\lambda K_t(x, x') = t^{-n/2} e^{-\|x-x'\|^2/(t/\lambda^2)} \lambda^{n-2}$$

$$= \lambda^2 K_{t/\lambda^2}(x, x')$$

$$R_\lambda \int_\varepsilon^L K_t(x, x') dt = \int_\varepsilon^L K_{t/\lambda^2}(x, x') \lambda^{-2} dt = \int_{\varepsilon/\lambda^2}^{L/\lambda^2} K_t(x, x') dt$$

Claim
 If $\{I[L]\}$ defines a theory then so does

$$I_\lambda[L](\varphi) = R_\lambda^+ I[\lambda^2 L](\varphi)$$

Proof

$$\frac{d}{dL} I_\lambda[L/\lambda^2] = R_\lambda^+ \frac{d}{dL} I[L](\varphi) = R_\lambda^+ \left(\begin{array}{c} \text{Diagram 1} \\ + \\ \text{Diagram 2} \end{array} \right)$$

The diagrams are Feynman diagrams for the derivative of the action. Diagram 1 shows a circle labeled $I(L)$ with a vertical line passing through its center. Diagram 2 shows a circle labeled $I(L)$ with a horizontal line passing through its center. A plus sign is between the two diagrams, and they are enclosed in large parentheses.

$$R_\lambda^\nu \int \varphi \Delta \varphi = \int \varphi \Delta \varphi$$

Define the $\mathbb{R}_{>0}$ action on the space of theories

$$\text{by } \{\mathcal{I}[L]\} \longrightarrow \{\mathcal{I}_\lambda[L] = R_\lambda^\nu \mathcal{I}[\lambda^2 L]\}$$

... scheme, we can write

$$I[L] = \lim_{\varepsilon \rightarrow 0} W(P(\varepsilon, L), I - I^{CT}(\varepsilon))$$

I th dependent Lagrangian

$I^{CT}(\varepsilon)$ Counter-terms - uniquely determined if we ask

they are purely singular as fns of ε

$$\begin{aligned} R_\lambda^* I[\lambda^2 L] &= \lim_{\varepsilon \rightarrow 0} R_\lambda^* W(P(\varepsilon, \lambda^2 L), I - I^{CT}(\varepsilon)) \\ &= \lim_{\varepsilon \rightarrow 0} W(R_\lambda P(\varepsilon, \lambda^2 L), R_\lambda^* I - R_\lambda^* I^{CT}(\varepsilon)) \\ &= \lim_{\varepsilon \rightarrow 0} W(P(\lambda^{-2} \varepsilon, L), R_\lambda^* I - R_\lambda^* I^{CT}(\varepsilon)) \end{aligned}$$

$\varepsilon \leftarrow$ purely sing. $(\lambda^2 \varepsilon) \leftarrow$ still is
 $\log \varepsilon$ " " $\log \lambda^2 \varepsilon$ is not

$$A = \{ \text{real analytic fns of } \varepsilon \in \mathbb{R}_{>0} \}$$

$$A_+ = \{ f(\varepsilon), \lim_{\varepsilon \rightarrow 0} f(\varepsilon) \text{ exists} \}$$

Scheme = a choice of $A_- \subseteq A$, complementary to A_+

was in old scheme

ε^{-k} purely sing
 $\log \varepsilon$ " "

$(\lambda^2 \varepsilon)^{-k}$ still is
 $\log \lambda^2 \varepsilon$ is not

$$A_{\cup} = \{ \text{real analytic fns of } \varepsilon \in \mathbb{R}_{>0} \}$$

$$A_{+} = \{ f(\varepsilon), \lim_{\varepsilon \rightarrow 0} f(\varepsilon) \text{ exists} \}$$

Scheme = a choice of $A_{-} \subseteq A$, complementary to A_{+}

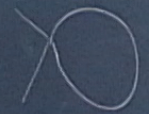
$\mathbb{R}_{>0}$ acts on A , $f \rightarrow f(\lambda^2 \varepsilon)$, preserves A_{+}

$\mathbb{R}_{>0}$ action on $\{\text{Lagrangians}\}$ sends
 $I \longrightarrow S_{R,\lambda}^{-1} S_{\psi} R_{\lambda}^* I$

φ^4 theory on \mathbb{R}^4
 $R_{\lambda}(\varphi) = \lambda \varphi(\lambda x)$, so $\int \varphi^4$ is fixed by R_{λ}
Classically, $\int \varphi \Delta \varphi + g \varphi^4$ is fixed

in old scheme

Scheme γ where ϵ^{-n} is singular, $\log \epsilon'$ is singular
One-loop counter-terms:



$$\epsilon^{-2} \int \varphi^2$$

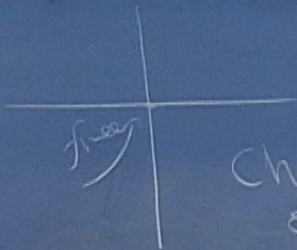


$$c^2 \log \epsilon \int \varphi^4$$

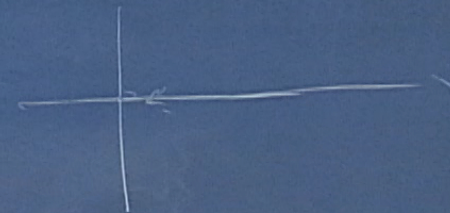
If we send $\varepsilon \rightarrow \lambda^2 \varepsilon$ and apply R_{λ}
then

$$\varepsilon^2 \int \varphi^2 \longrightarrow \varepsilon^2 \int \varphi^2 \quad \text{unchanged}$$

$$\log \varepsilon c^2 \int \varphi^4 \longrightarrow (\log \varepsilon c^2) \int \varphi^4 + 2c^2 \log \lambda \int \varphi^4$$



Choose a basis
of tangent space
corresponding to a
basis of Lagrangians



$$R_{s_0} = \text{v. field } \sum d_n c_n \frac{\partial}{\partial c_n} + \dots + \hbar c_0^2 \frac{\partial}{\partial c_0} + \hbar^2 c_0^3 \dots$$

$c_0 = \hbar^0$ Lagrangian

Say a theory is perturbatively renormalizable
if $\Gamma_\lambda[L]$ has log. growth as $\lambda \rightarrow 0$.

For φ^4 theory

One can show that the set of perturbatively renorm. theories
is $\left\{ \sum t^k (c_k \varphi^4 + d_k \varphi^6) \right\}$ (compatible all symmetries)