

Title: Higher operations in supersymmetric $\mathcal{N}=4$ theory

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Abstract: I will review the construction of "higher operations" on local and extended operators in topological $\mathcal{N}=4$ theory, and some applications of this construction in supersymmetric $\mathcal{N}=4$ theory. In particular, the higher operation on supersymmetric local operators in a 3d $\mathcal{N}=4$ theory turns out to be induced by the holomorphic Poisson structure on the moduli space of the theory. This leads to a new way of establishing the non-renormalization properties of this Poisson structure, and also to a simple topological reason for the appearance of its deformation quantization when the theory is placed in Omega-background. This is an account of joint work with Christopher Beem, David Ben-Zvi, Mathew Bullimore, and Tudor Dimofte.

Higher products in SUSY QFT

joint w/ C. Beem, D. Ben-Zvi,
M. Bullimore, T. Dimofte

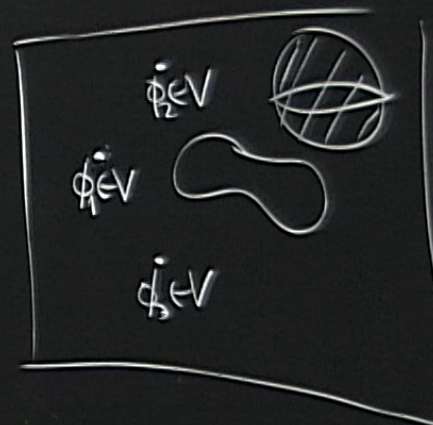
Aim: an overlooked structure in SUSY
QFT —

practitioners of (modern) TFT.

Companion to [Elliott-Sabner] on
 E_n -algebra structures in SUSY QFT

I'll focus on homology, not chain level.

In QFT, spacetime \mathbb{R}^n w/ various kinds of operators inserted, carrying labels, e.g. vector space V of local operators



translation invariance:



labels, e.g. vector space V of
local operators

Translation invariance: $V \ni \mathbb{R}^n$
by momentum operators P_μ

$\phi_1 \in V$

$\phi_2 \in V$

$\mu = 1, \dots, n$

$$\frac{\partial}{\partial x^\mu} \langle \phi(x) \dots \rangle = \langle (P_\mu \phi)(x) \dots \rangle$$

and P_μ is a symmetry:

$$\langle P_\mu(\dots) \rangle = 0$$

eg.

$$\langle P_\mu \phi(x_1) \phi(x_2) \rangle + \langle \phi(x_1) P_\mu \phi(x_2) \rangle = 0$$

Now suppose the theory has $(\mathbb{Z}_2\text{-odd})$ \swarrow symmetry operators

Q, Q_μ with $Q^2 = 0$
 $[Q, Q_\mu] = P_\mu$

(ex: Q is a "topological supercharge" in SUSY field theory)

If all $Q\phi_i = 0$

then $\langle Q\phi(x) \phi_1(x_1) \phi_2(x_2) \dots \phi_n(x_n) \rangle$

$$= \langle Q(\phi(x) \phi_1(x_1) \dots \phi_n(x_n)) \rangle = 0$$

$\Rightarrow \langle \prod \phi_i(x_i) \rangle$ depends only on $[\phi_i] \in H(Q)$

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!!
A

$$\begin{aligned}
 Q\phi = 0 &\Rightarrow P_\mu \phi = [Q, P_\mu] \phi \\
 &= Q(P_\mu \phi) \quad Q\text{-exact} \\
 &\Rightarrow \langle \prod \phi_i(x_i) \rangle \text{ is } \underline{\text{locally constant}} \text{ f}^n \text{ of } \{x_i\}
 \end{aligned}$$

Then given $\phi_1, \phi_2 \in A$

define $[(\phi_1 \cdot \phi_2)(x)] = \phi_1(x) \cdot \phi_2(x)$

$\lim_{x_2 \rightarrow x_1} [\phi_1(x_1) \phi_2(x_2)]$

$\phi_1 \cdot \phi_2 \in A$

If $n \geq 2$ this product is
commutative

If $n=1$ define $\phi_1 \cdot \phi_2(x) = \phi_1(x_1) \phi_2(x_2)$

$$\phi_2 \cdot \phi_1(x) = \phi_2(x_2) \phi_1(x_1)$$

This product is always associative —

in dim $n \geq 2$ that's because $\mathcal{C}_3(\text{ball})$ is connected
"topological ring"

Descent

Given ϕ with $Q\phi = 0$,

$$\phi^*(x) = \sum \phi^{(k)}(x)$$

$$\phi^{(0)}(x) = \phi(x)$$

$$\phi^{(1)}(x) = Q_{\mu} \phi(x) dx^{\mu}$$

$$\phi^{(2)}(x) = Q_{\mu} Q_{\nu} \phi dx^{\mu} dx^{\nu}$$

⋮

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$$\begin{aligned} d\phi^{(k)}(x) &= P_{\mu} \phi(x) dx^{\mu} \\ &= Q_{\mu} Q_{\nu} \phi(x) dx^{\mu} \end{aligned}$$

For any k -chain $\Gamma \subset \mathbb{R}^n$

define $\phi(\Gamma) = \int_{\Gamma} \phi^{(k)}$

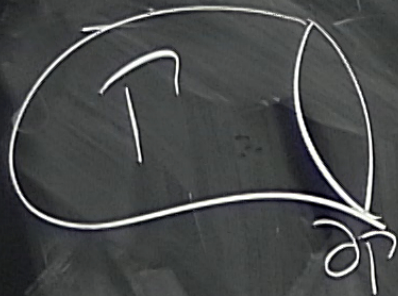
extended operator

$$d\phi^{(0)} = Q\phi^{(1)}$$

$$d\phi^{(1)} = Q\phi^{(2)}$$

$$d\phi^* = Q\phi^*$$

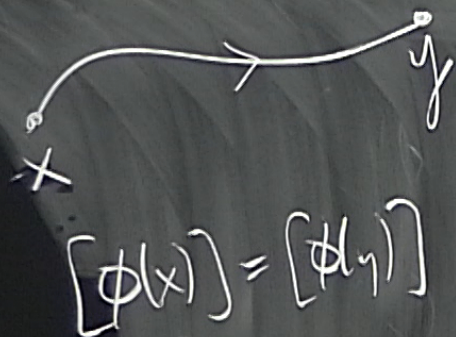
$$Q\phi(\Gamma) = \int_{\Gamma} Q\phi^{(k)} = \int_{\Gamma} d\phi^{(k-1)} = \int_{\partial\Gamma} \phi^{(k-1)} = \phi(\partial\Gamma)$$



$$Q(\phi(\Gamma)) = \phi(\partial\Gamma)$$

If $\partial\Gamma = 0$ then $Q(\phi(\Gamma)) = 0$

and $[\phi(\Gamma)]$ depends only on $[\Gamma]$



What to do with $\phi(\Gamma)$?

One application: [Witten]

build twisted version of SUSY
theory on n -manifold X

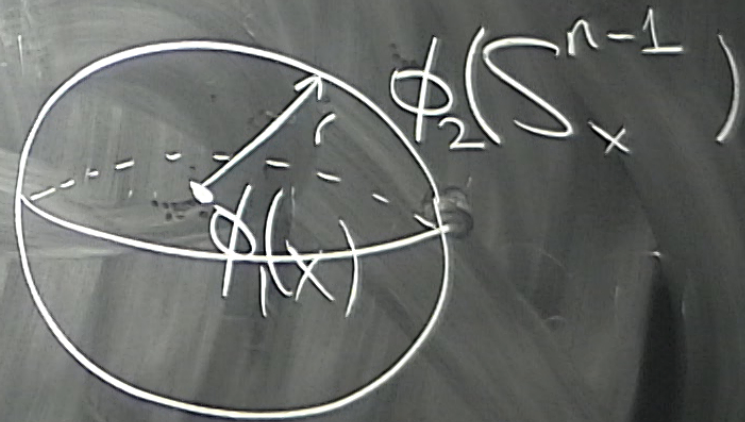
then insert $\phi(\Gamma)$ or number
cycles X .

Secondary product

$$\text{Say } Q\phi_1 = Q\phi_2 = 0$$

$$\{\phi_1, \phi_2\}(x) = \lim_{r \rightarrow 0} \phi_1(x) \phi_2(S_x^{n-1})$$

$$A \otimes A \rightarrow A$$



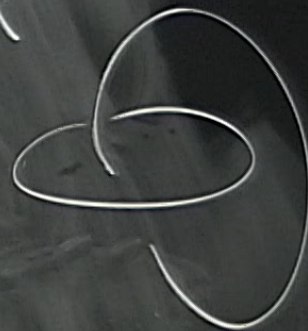
Then $\{ \cdot, \cdot \}$ gives a (degree $n-1$) Poisson bracket on A .

To prove $\{ \cdot, \cdot \}$ is a Poisson bracket:

(Skew) Symmetry: extend $\phi_1(x_1)\phi_2(x_2)$ to form-valued op
 $(\phi_1\phi_2)^*$ on $\mathcal{C}_2(\text{ball})$

\Rightarrow every class $\alpha \in H_*(\mathbb{C}_2(\text{ball}))$
gives a map $\star_\alpha: A \otimes A \rightarrow A$

and $H_{n-1}(\mathbb{C}_2(\text{ball})) \cong \mathbb{Z}$



\int^{n-1}
Jacobi comes from relation $H_{2n-2}(\mathbb{C}_3(\text{ball}))$

Rk All operations we get in this way
are built from \cdot and $\{, \}$

(Thm the homology of the little n -discs
operad is the degree $n-1$ Poisson operad)

(Cohen)

(Dez Sinhs)

Example

$n=3$, $\mathcal{N}=4$ SUSY σ -model into
hyperkähler space X .

choose $\omega \in \mathcal{Q} \leftrightarrow$ choose $\omega \in \mathcal{C}$ str on X

X is hol. symplectic

$$A = H_{\frac{\partial}{\partial}}^{0, \infty}(X)$$

gets Poisson bracket

$$\{f_1, f_2\} = \Omega^{-1}(df_1, df_2)$$

Applications

a) Non-renormalization

ξ, η can be computed exactly
either in the IR or the UV

✓ [Bullimore-Dimofte-Gaiotto, Gaiotto-Moore-N]

b) Ω -background

exactly

$$A_\varepsilon = H(Q_\varepsilon)$$

proposal: A_ε is def. quant of A
[Yagi, GMN, Itô Okubo-Tate]



actly

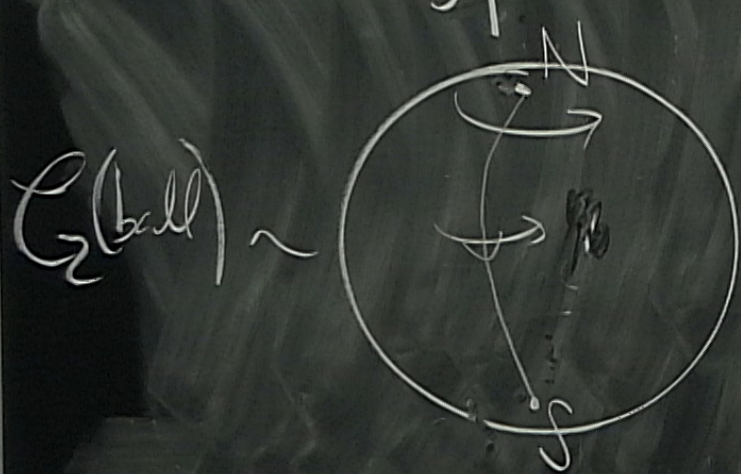
$$A_{\xi} = H(Q_{\xi})$$

proposal: A_{ξ} is def. quant of A

[Yagi, GMM, It's Okude-Take]

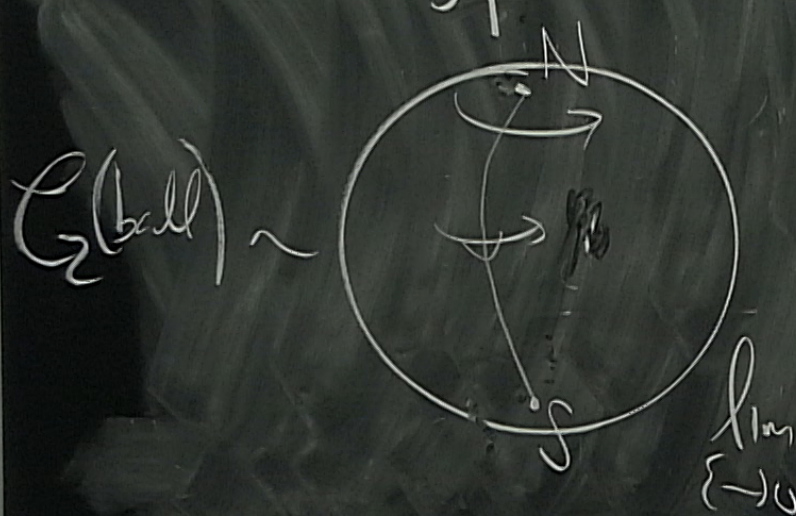
Why?

Operations on A_G come from $U(1)$ -equiv
 homology of $C_n(\text{ball})$



$$[N] - [S] = \varepsilon [S^2]$$

Operations on A_Σ come from $U(1)$ -equiv
 homology of $C_n(\text{ball})$



$$\begin{bmatrix} N \\ 0 \end{bmatrix} - \begin{bmatrix} S \\ 0 \end{bmatrix} = \varepsilon \begin{bmatrix} S^2 \end{bmatrix}$$

$$\lim_{\varepsilon \rightarrow 0} \frac{\phi_1 \phi_2 - \phi_2 \phi_1}{\varepsilon} = \left\{ \begin{matrix} \phi_1 \\ \phi_2 \end{matrix} \right\}^2$$