

Title: Cosmos, the beginnings...

Date: Jul 24, 2018 10:30 AM

URL: <http://pirsa.org/18070060>

Abstract: <div>Abstract: â€œHow did our universe begin?â€• is possibly one of the oldest questions that have</div>
<div>bewildered humans throughout history. As a theoretical cosmologist, our job is to find a</div>
<div>mathematically consistent picture for early universe that could explain observations, from</div>
<div>the largest to the smallest scales. The past thirty years have witnessed amazing progress,</div>
<div>both in developing technology for precision cosmological observations, and in perfecting</div>
<div>mathematical methodology to explain them. For example, ripples in cosmic geometry are</div>
<div>now measured with the precision of one part in a million. We also have sophisticated</div>
<div>mathematical frameworks such as general relativity and quantum theories that describe the</div>
<div>origin of these ripples in early universe. However, with all of these extraordinary</div>
<div>achievements, some old and new puzzles remain unsolved. For example we still have not</div>
<div>resolved the most crucial puzzle about the origin of cosmos, namely the Big Bang Singularity</div>
<div>problem. We will take a journey back in time to explore the fascinating realm of early</div>
<div>universe and some of its mysteries.</div>

Cosmic Observations Meeting Fundamental Theory



Hubble Law

$$a(t) = 1$$

$$a(t) \ll 1$$

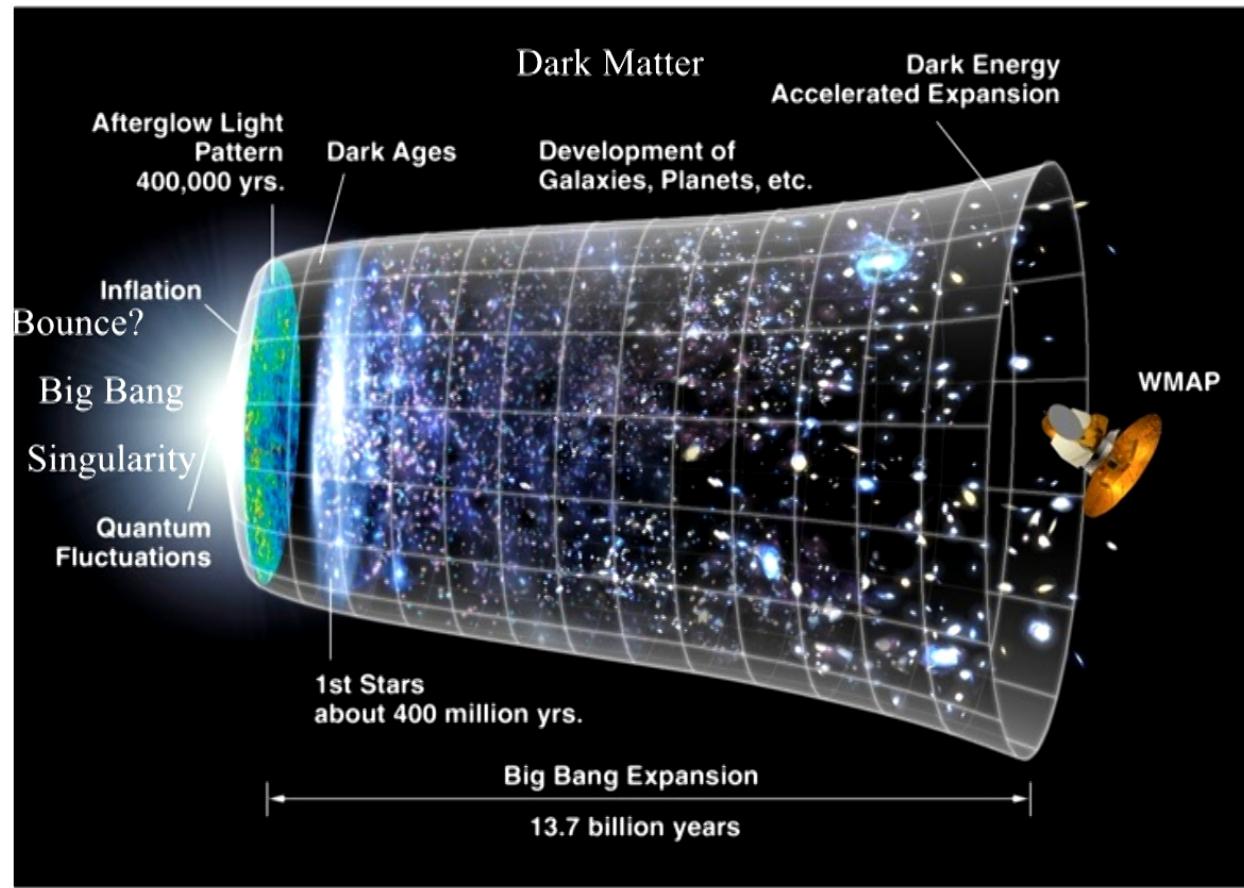


By [Ricardo García](#)

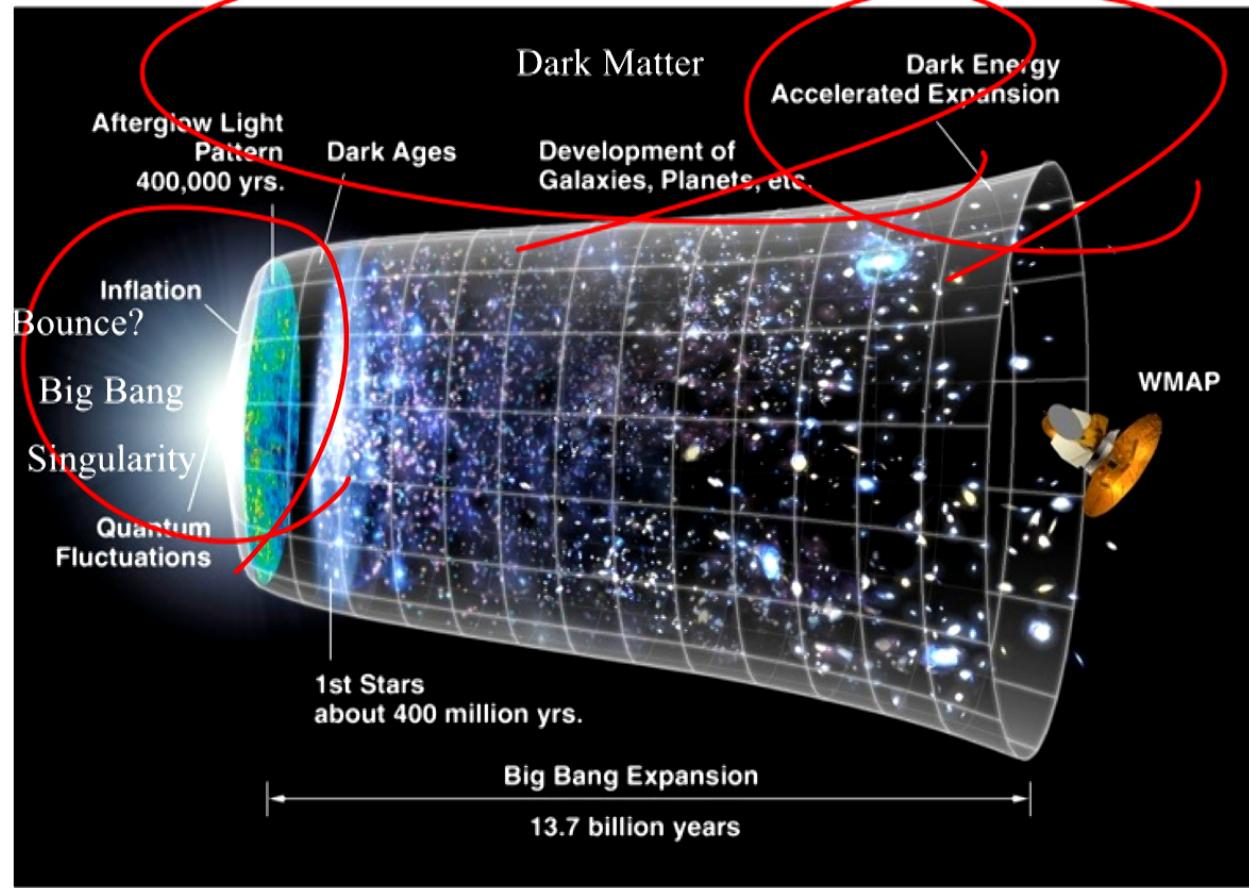
$$H = \frac{\dot{a}}{a}$$

$$V = Hr$$

The Jigsaw Puzzle



The Jigsaw Puzzle



- **Early Time:**

Big Bang Singularity, Quantum fluctuations, Generating seeds of Large scale structure, Big Bang nucleosynthesis, ...

Inflation, Bounce, Quantum gravity, New theories of particle Physics

Observation: CMB and Matter Power Spectrum, Non-Gaussianity, Gravitational waves? Neutrinos?

- **Intermediate time:**

Observations: Galaxy curves, lensing signal, CMB ...

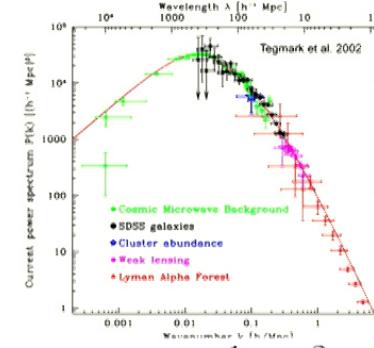
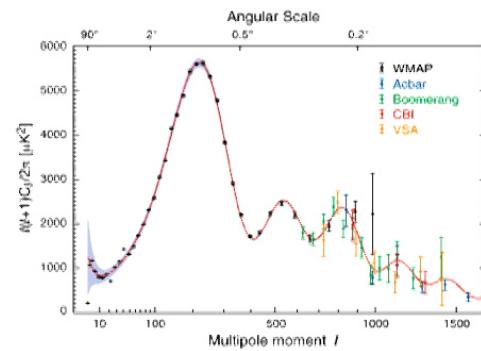
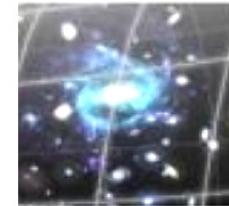
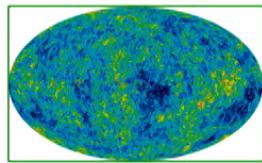
Missing Matter: Dark Matter, Mond/TEVES/MUG,
first stars, galaxy formation, re-ionization ...

- **Late Time Acceleration:**

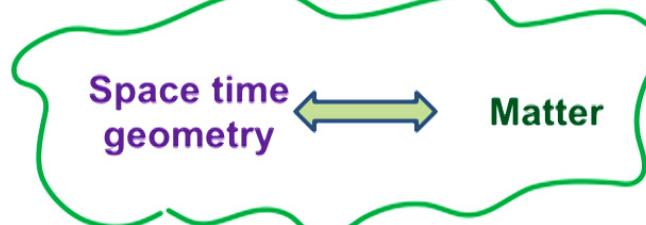
Cosmological constant problem, Lambda/Dark Energy/Modified Gravity

Observation: CMB, Supernovae, Gravitational waves

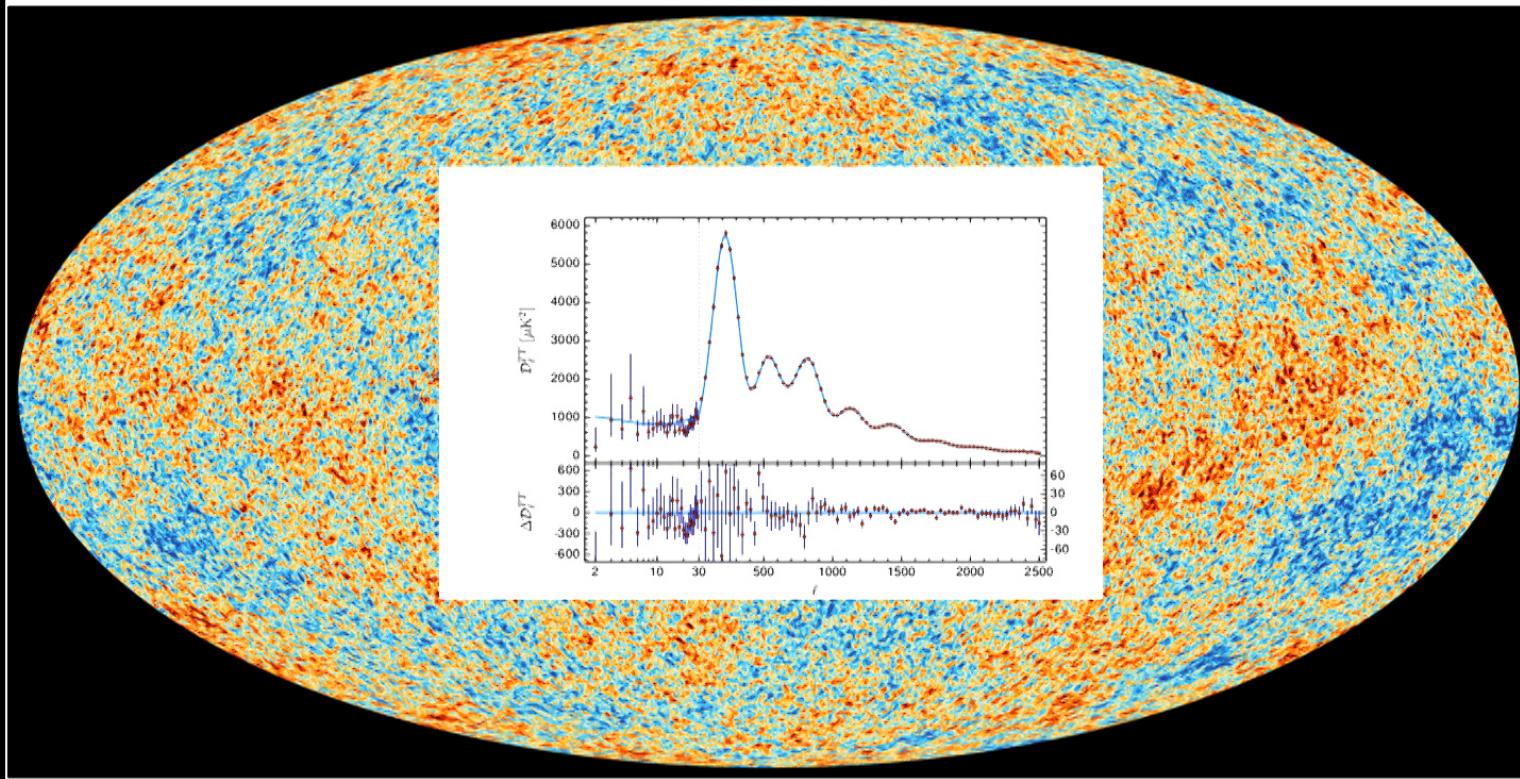
Anisotropies of order 10^{-5} ! Scale Invariant Power Spectrum



$$\langle \zeta(\mathbf{x})\zeta(\mathbf{x} + \mathbf{r}) \rangle \longleftrightarrow \mathcal{P}_k^\zeta \equiv |\zeta_k|^2 k^3 \longleftrightarrow k^{-1} |\delta\rho_k^2|$$

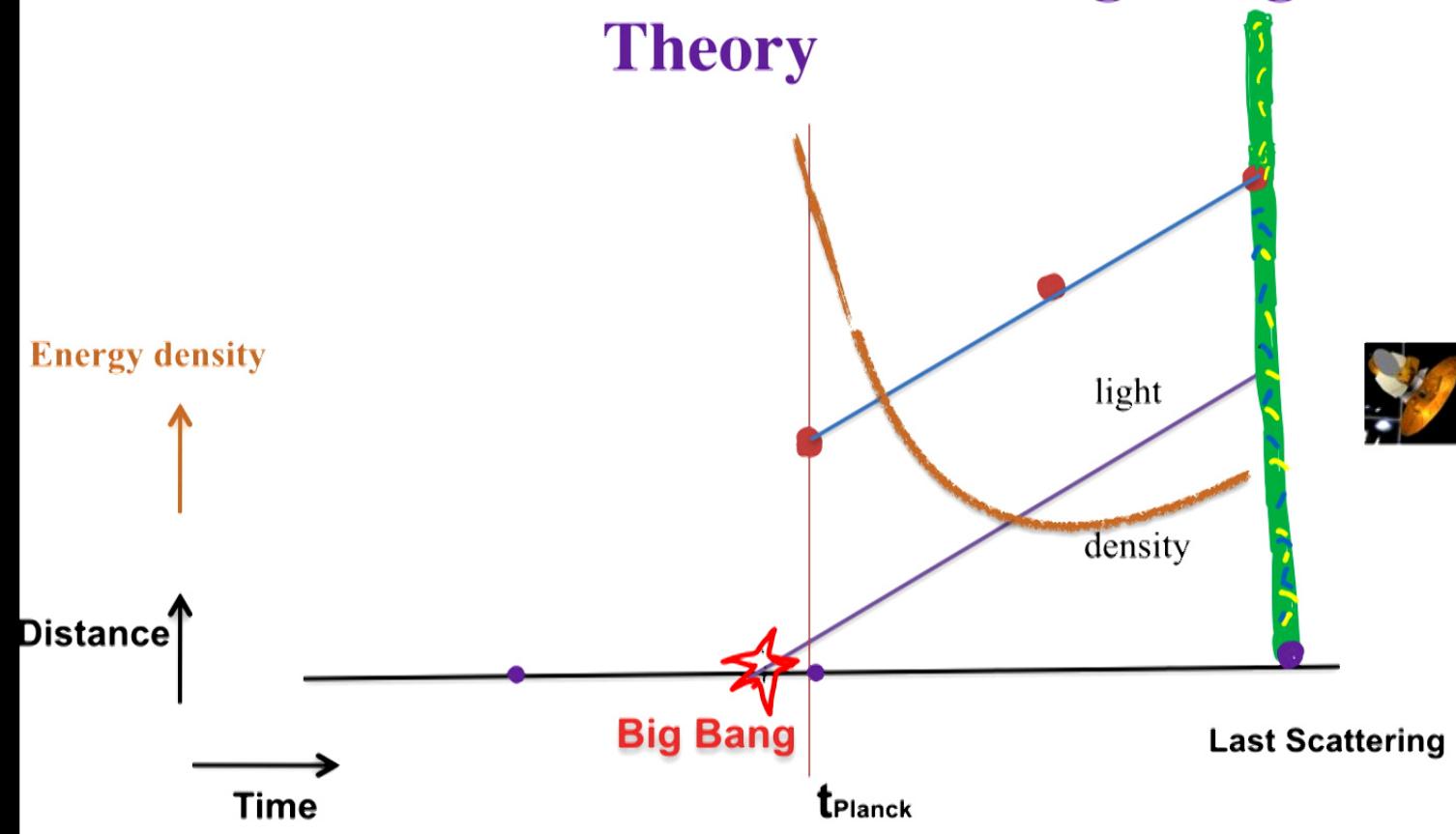


Most recent Picture of baby universe, released last week

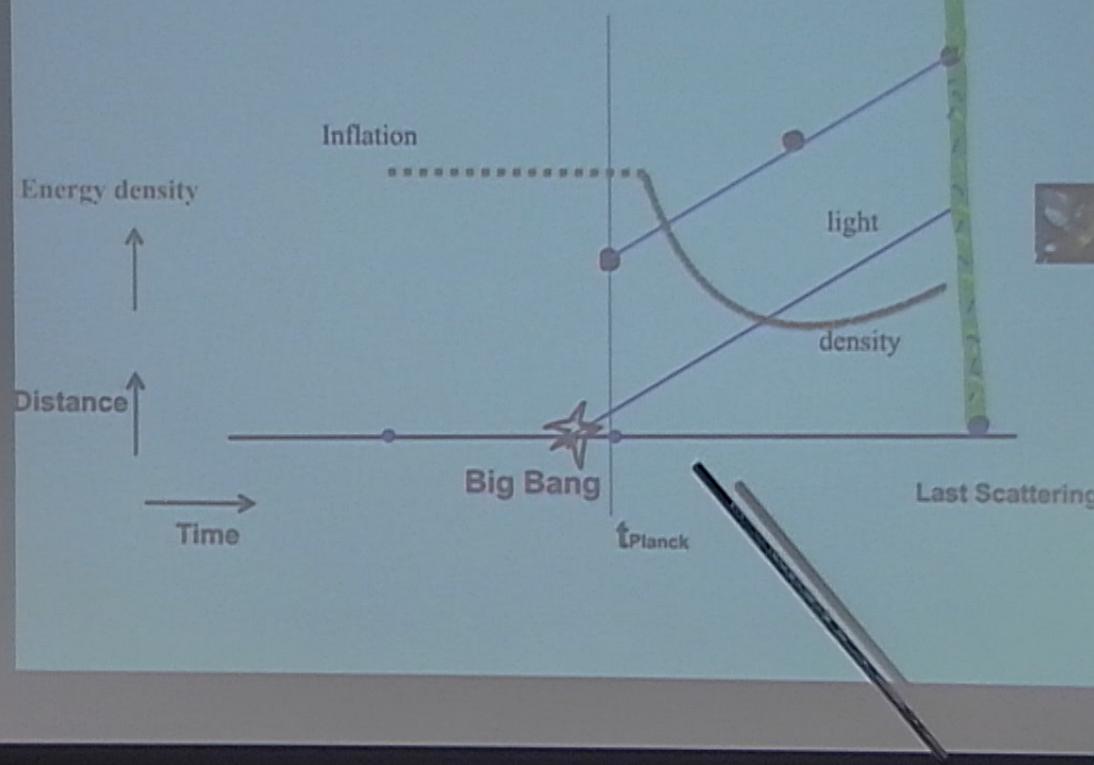


- **Title** Planck's view of the cosmic microwave background
- **Released** 17/07/2018 3:00 pm
- **Copyright** ESA/Planck Collaboration

Horizon Problem of Standard Big Bang Theory



Horizon Problem of Standard Big Bang Theory



Motivation from Particle Physics and String Theory

Standard Action of a Scalar Field:

$$S_\phi = \int dx^4 \sqrt{-g} \left[\frac{1}{2} (\partial_\mu \phi \partial^\mu \phi) - V(\phi) \right]$$

However at high energies Scalar Fields may have non-canonical kinetic terms and they could also inflate the universe!

$$S = \int dx^4 \sqrt{-g} \mathcal{L}(X, \phi) \quad X = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi$$

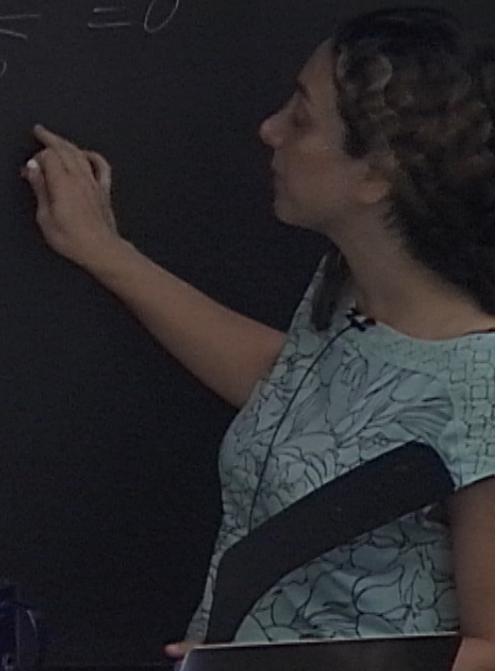
Warped D-Brane
Inflation

$$S_{DBI} = \int dx^4 \sqrt{-g} [-f(\phi)^{-1} \sqrt{1 - 2f(\phi)X} - (V(\phi) - f(\phi)^{-1})]$$

What are the general
criteria other fields
and particles

$$S = \int dx^4 L(\phi, \pi, \psi^i, \dots)$$

$$\frac{\delta S}{\delta \dot{\varphi}} \rightarrow \frac{d}{dt} \frac{\delta L}{\delta \dot{\varphi}} - \frac{\delta L}{\delta \varphi} = 0$$



Generalized Slow Roll Inflation

$$\mathcal{L}(X, \phi) \xleftarrow{\quad} X = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi$$

The scalar field is very similar
to a hydrodynamical fluid:

$$p(X, \phi) \quad \rho(X, \phi)$$



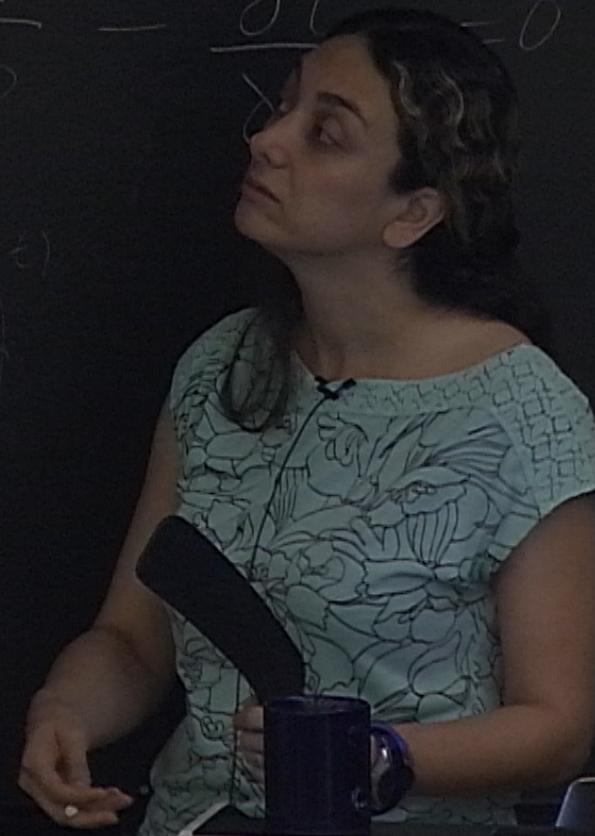
Gravity

1. Friedmann equations for homogenous background $a(t), \quad H = \frac{\dot{a}}{a} \longleftrightarrow \rho(\phi(t))$
2. Linear perturbation for quantum fluctuations



$$\zeta(x, t) \longleftrightarrow \delta\phi(x, t) \xrightarrow{\quad} \delta\rho$$

$$\frac{\delta S}{\delta \dot{\varphi}} \rightarrow \frac{d}{dt} \frac{\delta L}{\delta \dot{\varphi}} - \frac{\delta L}{\delta \varphi} = 0$$

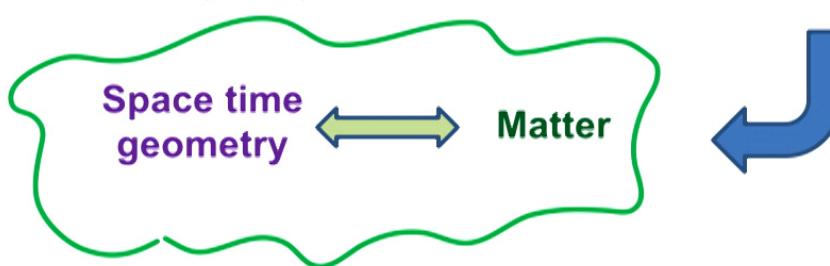


Generalized Slow Roll Inflation

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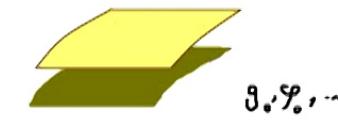
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Gravity

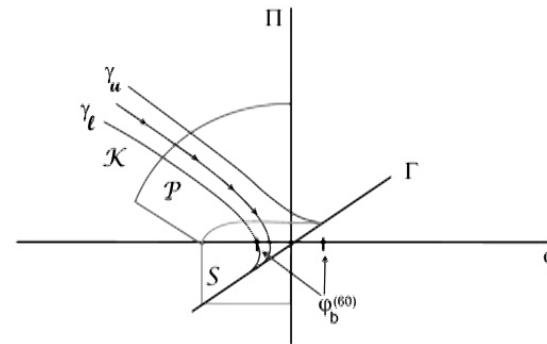
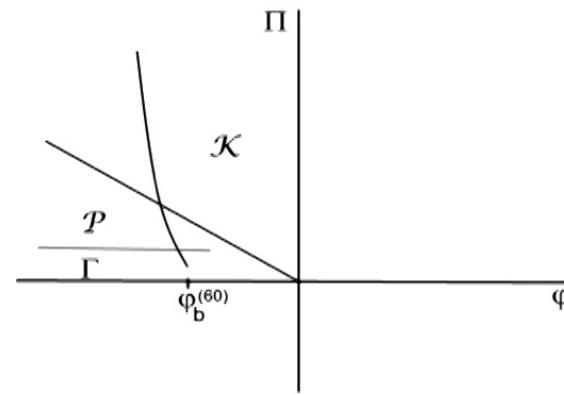
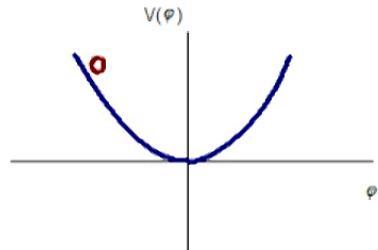
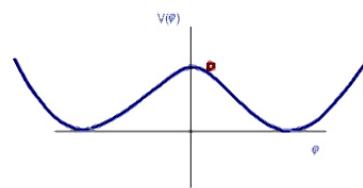
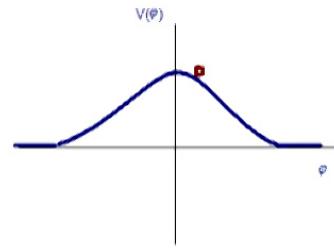
1. Friedmann equations for homogenous background $a(t), H = \frac{\dot{a}}{a} \longleftrightarrow \rho(\phi(t))$
Accelerated Expansion
2. Linear perturbation for quantum fluctuations



$$\zeta(x, t) \longleftrightarrow \delta\phi(x, t) \xrightarrow{\quad} \delta\rho$$

**CMB anisotropies and
Large scale structure**

Examples:



$$\frac{d^2 u_k}{d\tau^2} + (c_s^2 k^2 - 2(aH)^2 \{(1 + \frac{\eta}{2} + \kappa)(1 - \frac{\epsilon}{2} + \frac{\eta}{4} + \frac{\kappa}{2}) + \frac{\dot{\eta}}{2H} + \frac{\dot{\kappa}}{H}\}) u_k = 0$$

$$\eta \equiv \frac{\dot{\epsilon}}{H\epsilon}$$

$$\kappa \equiv \frac{\dot{c}_s}{H c_s}$$

slow roll: $\epsilon = -\frac{\dot{H}}{H^2}, \eta, \frac{\dot{\eta}}{H}, \kappa, \frac{\dot{\kappa}}{H} \ll 1$



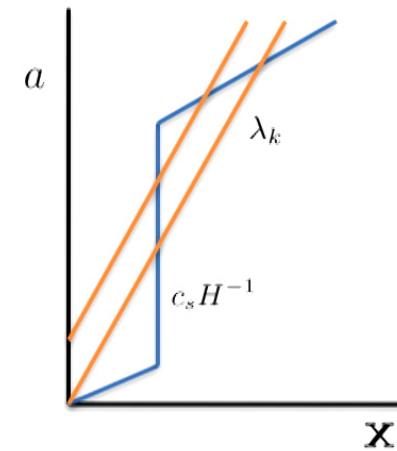
A Bessel equation

Scalar Power Spectrum $\mathcal{P}_k^\zeta \sim (\frac{1}{8\pi^2 M_p^2}) \frac{H^2}{c_s \epsilon} \Big|_{c_s k = aH}$

Scalar spectral index:

$$n_s - 1 \equiv \frac{d \ln \mathcal{P}_R}{d \ln k} \Big|_{k=k_s}, \\ \approx -(2\epsilon + \eta + \kappa) + O(\epsilon^2, \epsilon\eta, \kappa_N, \dots)$$

$$n_s = 0.9665 \pm 0.0038$$



$$\frac{d^2 u_k}{d\tau^2} + (c_s^2 k^2 - 2(aH)^2 \{(1 + \frac{\eta}{2} + \kappa)(1 - \frac{\epsilon}{2} + \frac{\eta}{4} + \frac{\kappa}{2}) + \frac{\dot{\eta}}{2H} + \frac{\dot{\kappa}}{H}\}) u_k = 0$$

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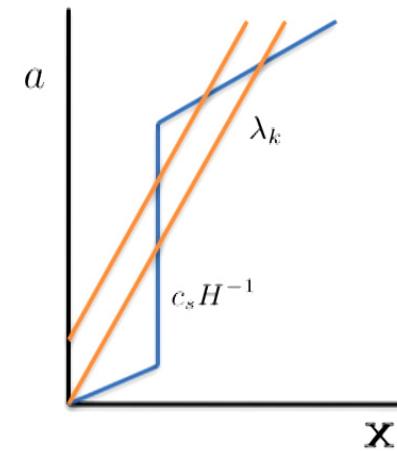
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$$n_s = 0.9665 \pm 0.0038$$



Alternative observations will result in getting complementary information:

Primordial Gravity Waves

Similar procedure for tensor power spectrum
Leave subdominant but distinct imprints in
the CMB Polarization

$$P_h = \frac{2H^2}{M_{pl}^2 \pi^2} \Big|_{k=aH}$$

Tensor spectral indexes:

$$n_t \equiv \left. \frac{d \ln P_h}{d \ln k} \right|_{k=k_t} \approx -2\epsilon + O(\epsilon^2, \dots)$$

Three point function, Non-Gaussianity :

Parameterizing non-gaussianities as

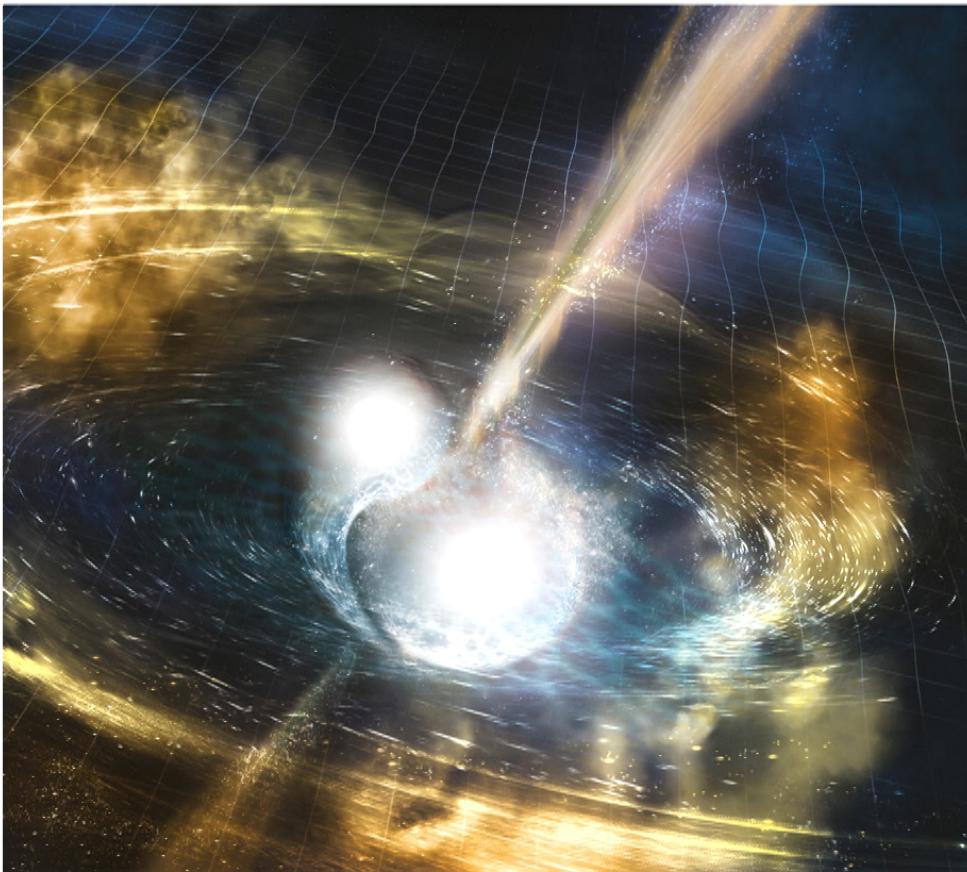
$$\zeta = \zeta_G - \frac{3}{5} f_{NL} \zeta_G^2$$

(Chen, Huang, Kachru, and Shiu)

$$f_{NL}^{\text{equil}} \approx (-0.26 + 0.12 c_s^2) \left(1 - \frac{1}{c_s^2}\right) - 0.08 \left(\frac{c_s^2}{\epsilon}\right) \frac{X^3 \mathcal{L}_{XXX}}{M_{pl}^2 H^2},$$



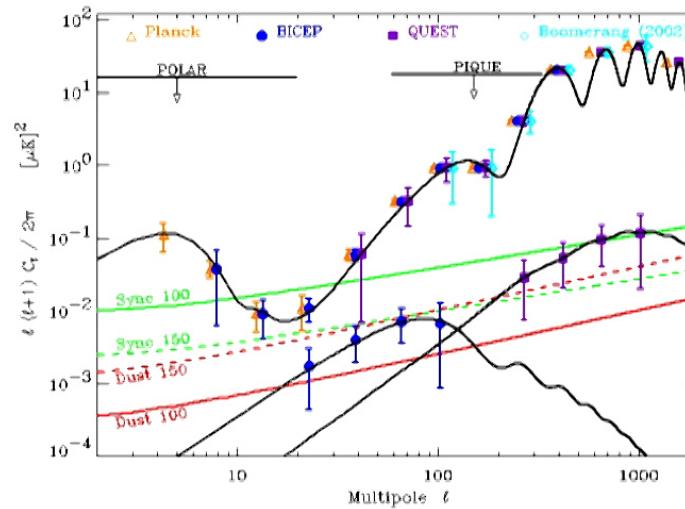
Gravitational waves



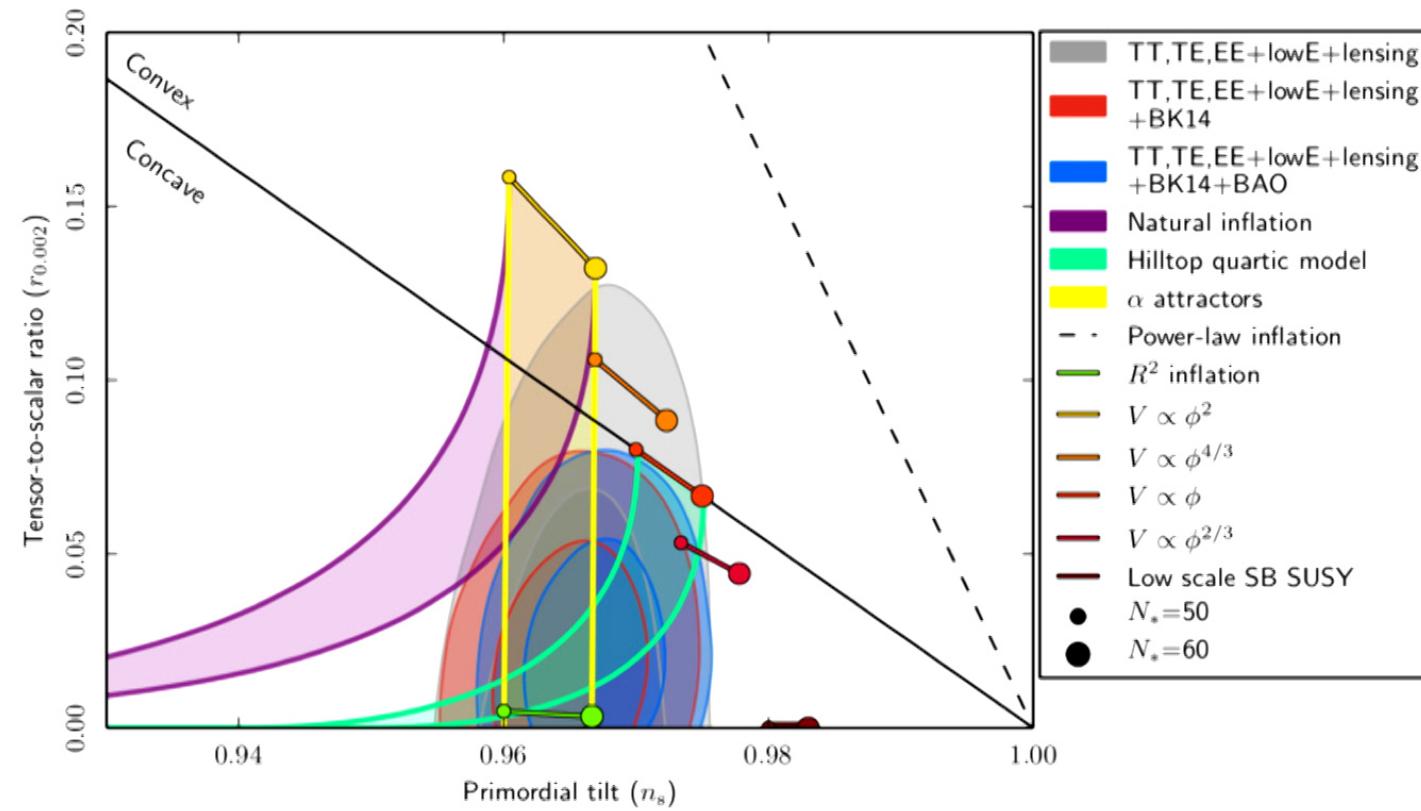
Credit

NSF/LIGO/Sonoma State University/A. Simonne

BICEP Experiment



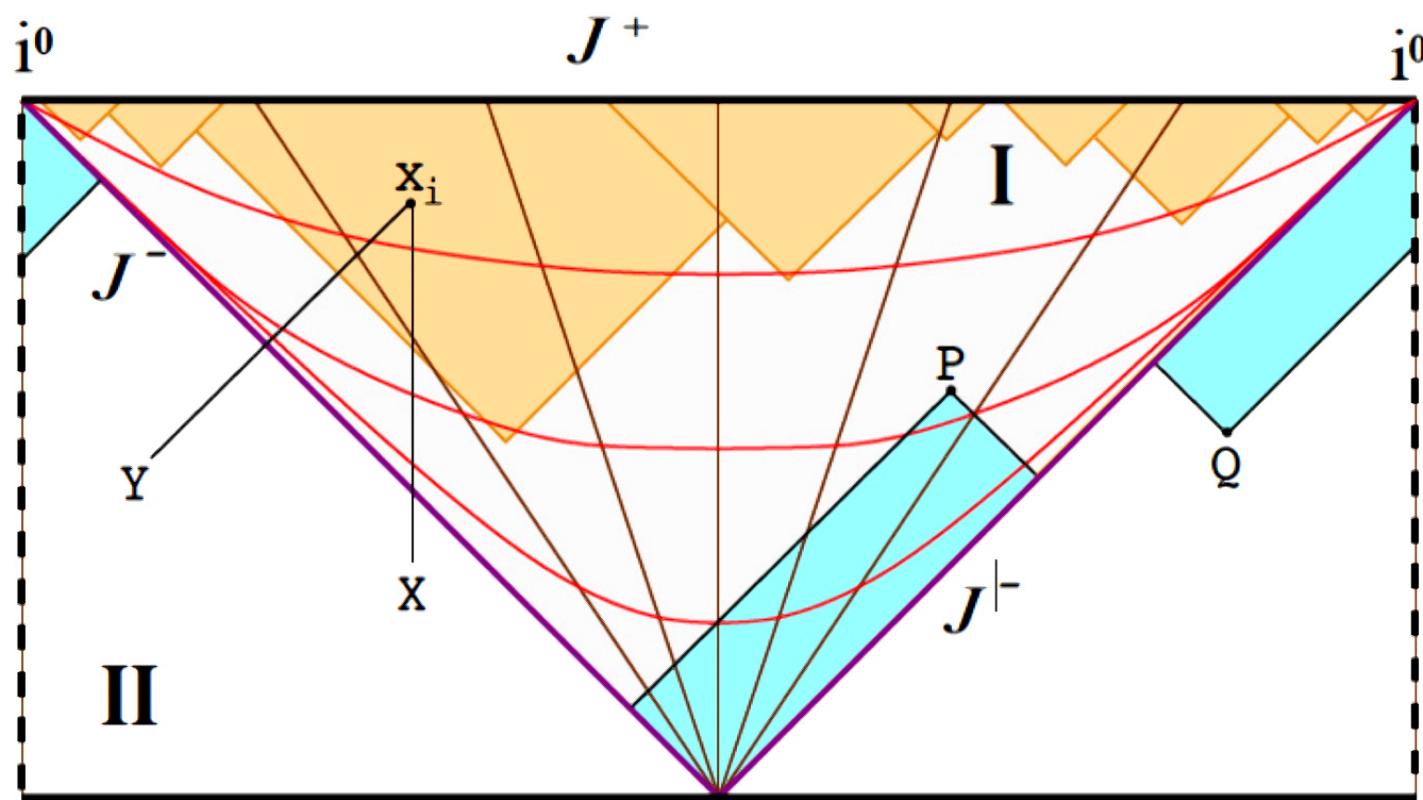
Planck Collaboration: Constraints on Inflation



What then?

- Falsifiable?
- What is Inflaton fi
- Eternal inflation/F

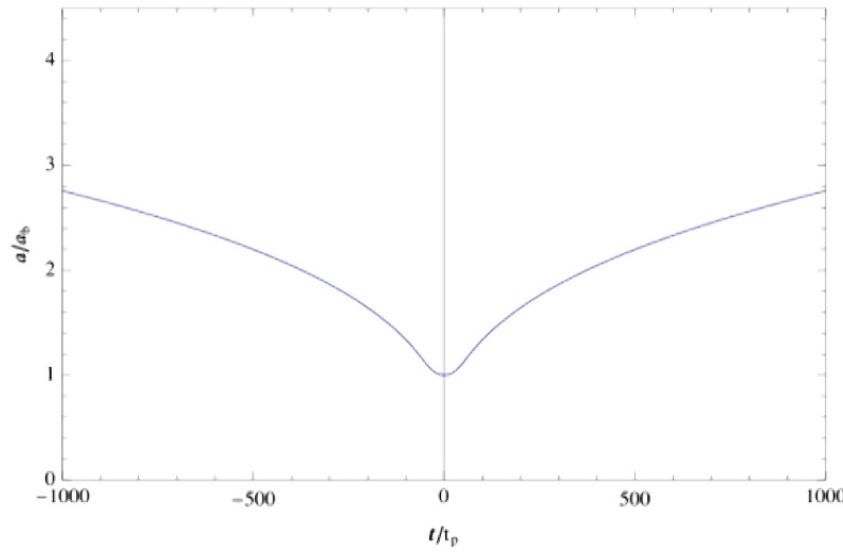




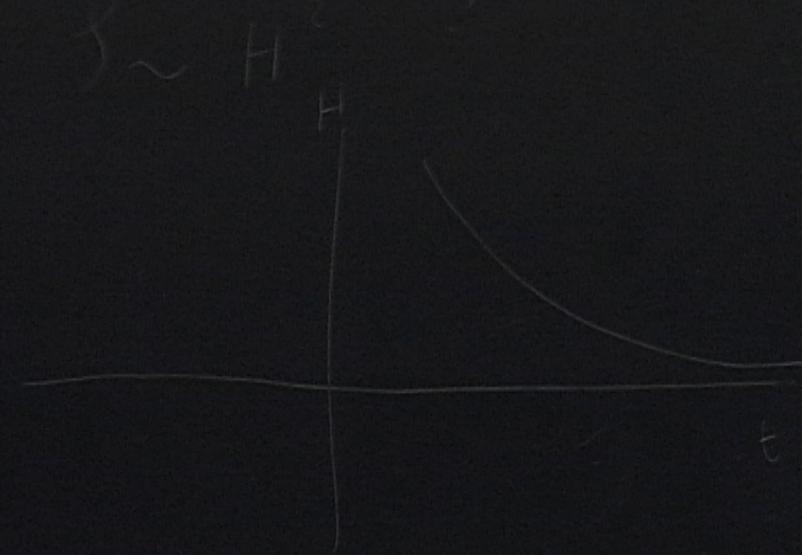
Aguirre & Gratton

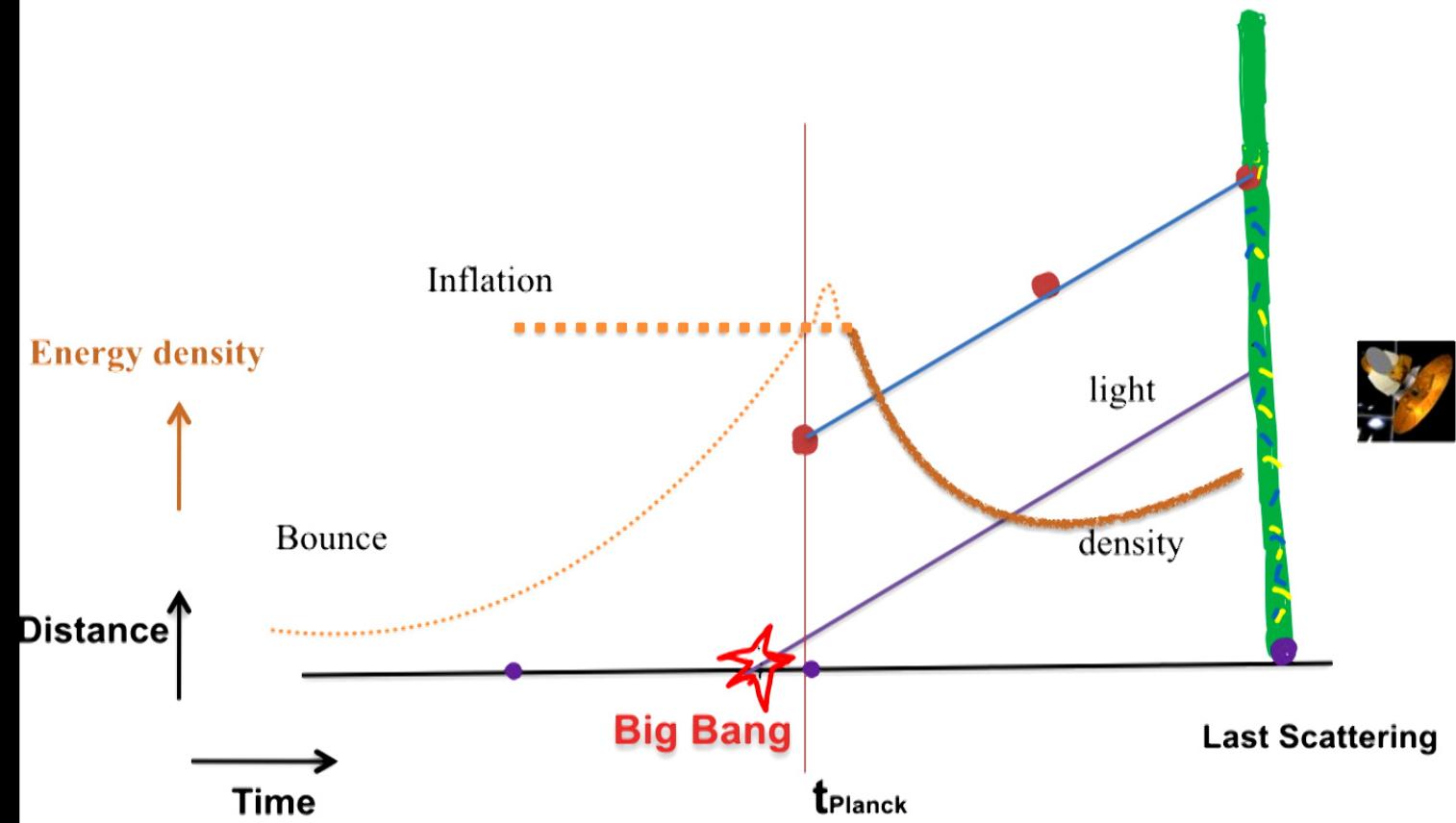
Alternatives: Bounce?

- Instabilities
- Getting the right spectrum
- What field?



$$\dot{\alpha}_B = 0 \rightarrow H = \frac{\dot{\psi}}{q} = 0$$





Getting a bounce cosmology through Cuscuton

Cuscuton (G.G., N. Afshordi, D. Chung, 07; G.G., N. Afshordi, D. Chung, M. Doran, 07)

Specific limit of K-essence Action: $c_s = \infty$.

$$\mathcal{L} = -\mu^2 \underbrace{\sqrt{|\partial^\mu \varphi \partial_\mu \varphi|}}_{2X} - V(\varphi) \quad c_s^2 = \frac{p_{,X}}{\rho_{,X}} = \frac{1}{1 + 2 \frac{X \mathcal{L}_{,XX}}{\mathcal{L}_{,X}}}$$

- Field equation becomes a **constraint equation** that uniquely determines Cuscuton as a function of metric
- No internal dynamics; only follows what it couples to
- *Cuscuton* (käs-kü-tän): derived from the Latin name for the parasitic plant of dodder, “Cuscuta”



Cusciton Bounce: No Instabilities!

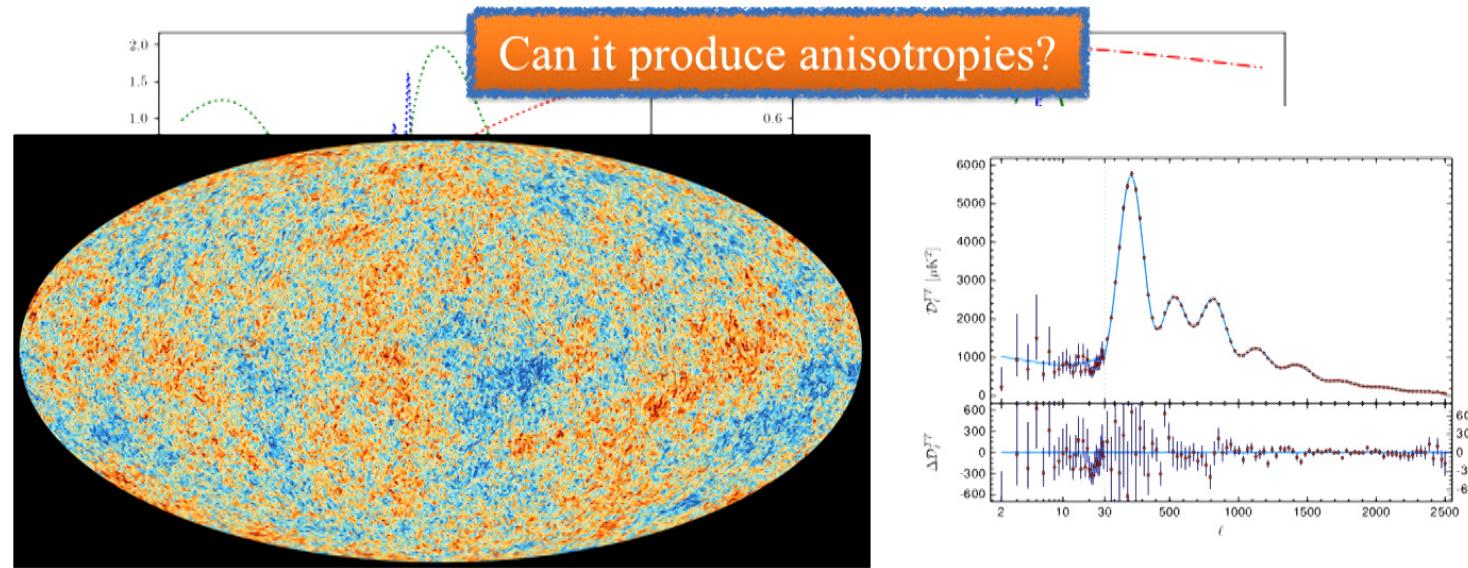


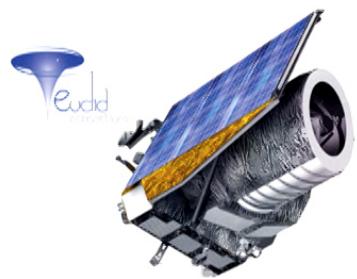


Illustration of NASA's Wide Field Infrared Survey Telescope (WFIRST).

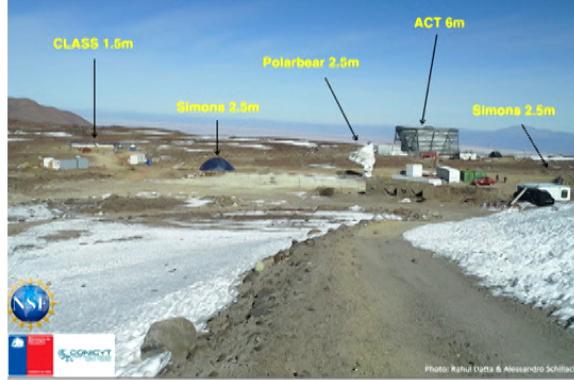
Credits: NASA/GSFC/Conceptual Image Lab



SKA
SQUARE KILOMETRE ARRAY
CANADA



Recent Atacama CMB experiments



Recent South Pole CMB experiments

