

Title: LIGO and Virgo continuous wave searches - Overview and all-sky searches

Date: May 10, 2018 02:30 PM

URL: <http://pirsa.org/18050033>

Abstract: The Continuous Waves (CW) Search Group of the LIGO Scientific Collaboration and Virgo Collaboration carries out a diverse suite of searches for a diverse set of possible CW sources. Assumptions underlying these searches will be discussed, along with strategies used so far to keep our eyes wide open while also giving due attention to the most promising sources. One important assumption to date has been that fast-spinning, non-axisymmetric neutron stars are the most promising class of CW sources. If the most promising source class is, instead, black hole super-radiance from axion clouds, then current searches may not be optimal. Potential modifications to current all-sky searches to address super-radiance will be discussed.

LIGO and Virgo continuous wave searches - Overview and all-sky searches



**Workshop on Searching for New
Particles with Black Hole
Superradiance**

**Perimeter Institute
May 10, 2018**

**Keith Riles
University of Michigan**

**LIGO Scientific Collaboration
and the Virgo Collaboration**



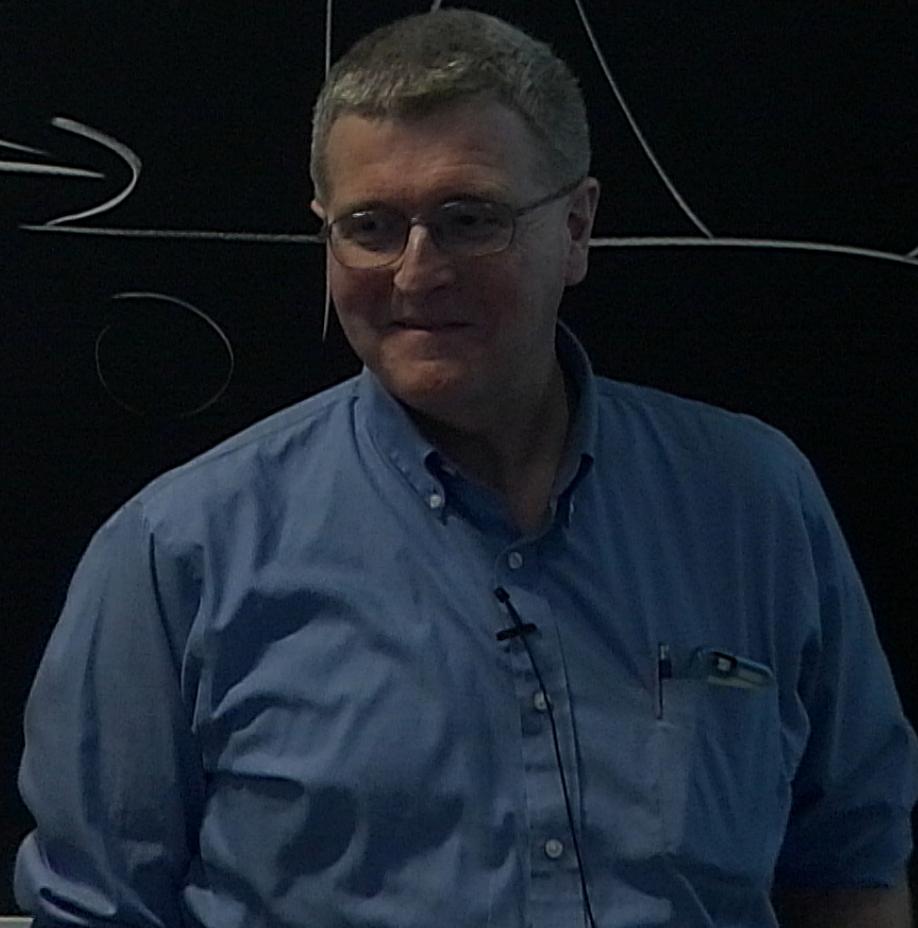
LIGO-G1800909

$$P \propto \gamma A^2$$

$$\propto \frac{F^2}{\gamma}$$

$$A \propto \frac{\bar{F}}{m \sqrt{\gamma \omega}}$$

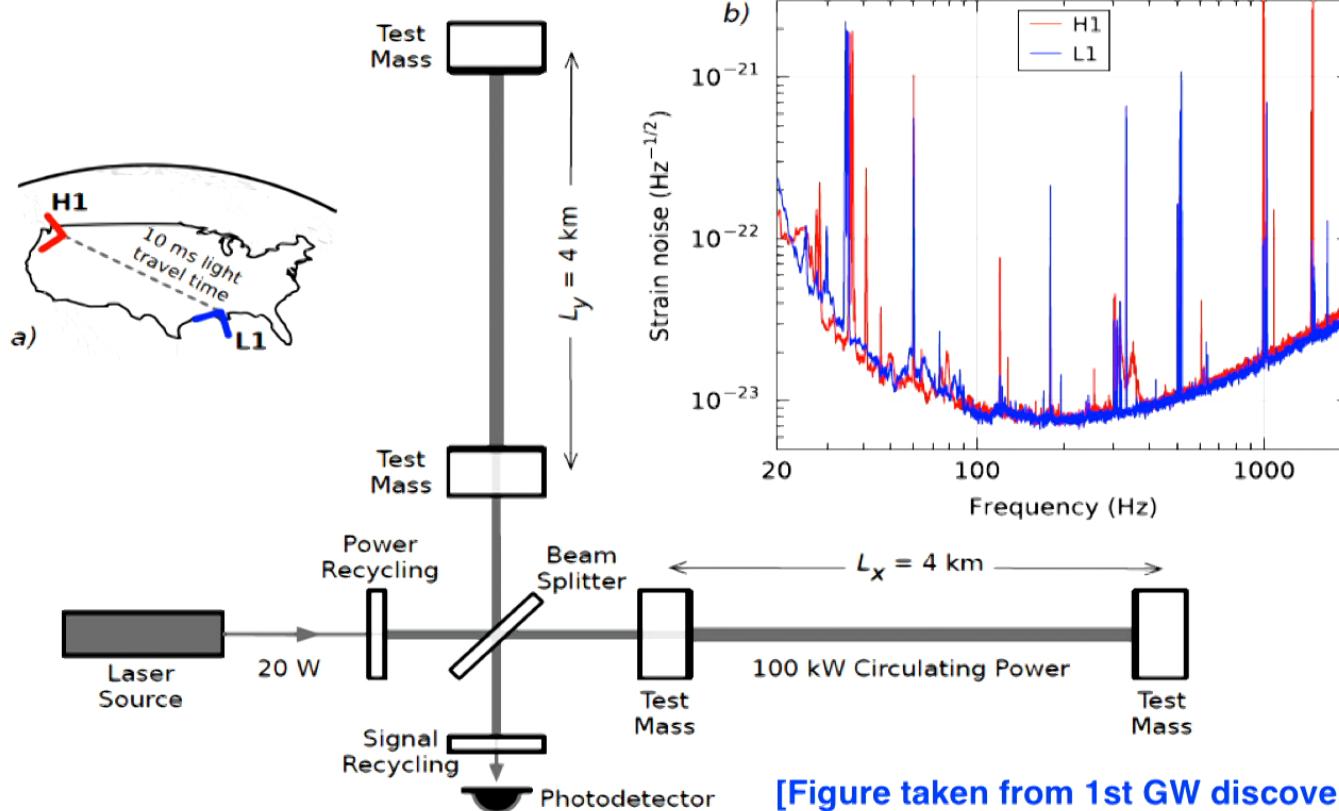
$\sim 1 - 2$



Outline

- **LIGO and Virgo – Where do we stand?**
 - ◆ Sensitivity progression to date
 - ◆ Prospects in near and far future
- **Continuous Wave (CW) search assumptions**
 - ◆ Emission mechanisms
 - ◆ Promising sources & sky directions
- **CW search methods & results to date**
 - ◆ Challenges – coping with unknown source parameters
 - ◆ Categories – focus on all-sky searches
 - ◆ Superradiance – what changes?

O1 Data Run (Sept 12, 2015 – January 19, 2016)



[Figure taken from 1st GW discovery paper:
B.P. Abbott *et al.*, PRL 116 (2016) 061102]

Translating 1-sided amplitude spectral noise density (ASD) to GW signal amplitude sensitivity: (back of the envelope level)

$$\text{Detectable signal amplitude} \approx \frac{\text{Amplitude spectral density}}{\sqrt{\text{[Coherent observation time]}}} \times \text{Statistical / geometry trials factor}$$

Translating 1-sided amplitude spectral noise density (ASD) to GW signal amplitude sensitivity: (back of the envelope level)

$$\text{Detectable signal amplitude} \approx \frac{\text{Amplitude spectral density}}{\sqrt{\text{[Coherent observation time]}}} \times \text{Statistical / geometry trials factor}$$

Examples:

BBH like GW150914:

$$10^{-23} \text{ Hz}^{-1/2} / [0.2 \text{ s}]^{1/2} \times \sim 10 \approx \text{several} \times 10^{-22}$$

Actual amplitude = 10^{-21} → Very loud signal

Translating 1-sided amplitude spectral noise density (ASD) to GW signal amplitude sensitivity: (back of the envelope level)

$$\text{Detectable signal amplitude} \approx \frac{\text{Amplitude spectral density}}{\sqrt{\text{Coherent observation time}}} \times \text{Statistical / geometry trials factor}$$

Examples:

BBH like GW150914:

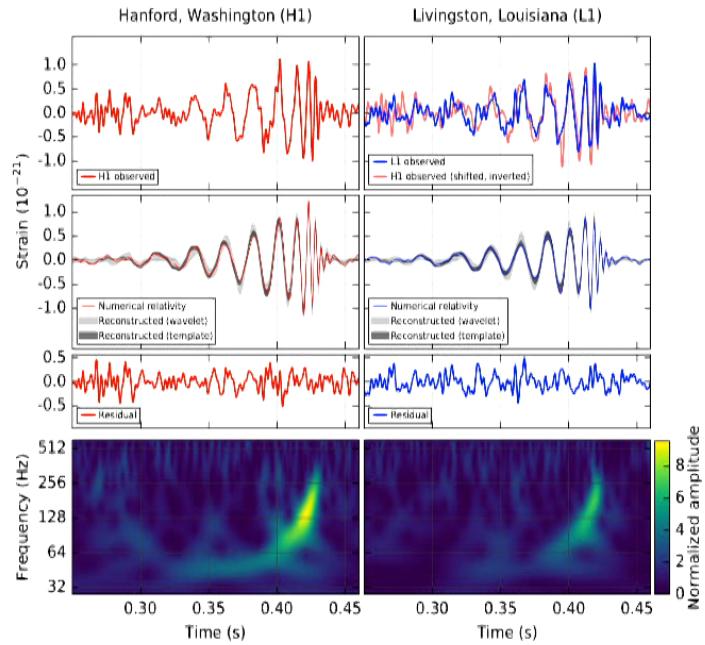
$$10^{-23} \text{ Hz}^{-1/2} / [0.2 \text{ s}]^{1/2} \times \sim 10 \approx \text{several} \times 10^{-22}$$

Actual amplitude = 10^{-21} → Very loud signal

CW signal from known pulsar measured coherently over 1 year:

$$10^{-23} \text{ Hz}^{-1/2} / [3 \times 10^7 \text{ s}]^{1/2} \times \sim 10 \approx 2 \times 10^{-26}$$

O1 “5-sigma” Detections

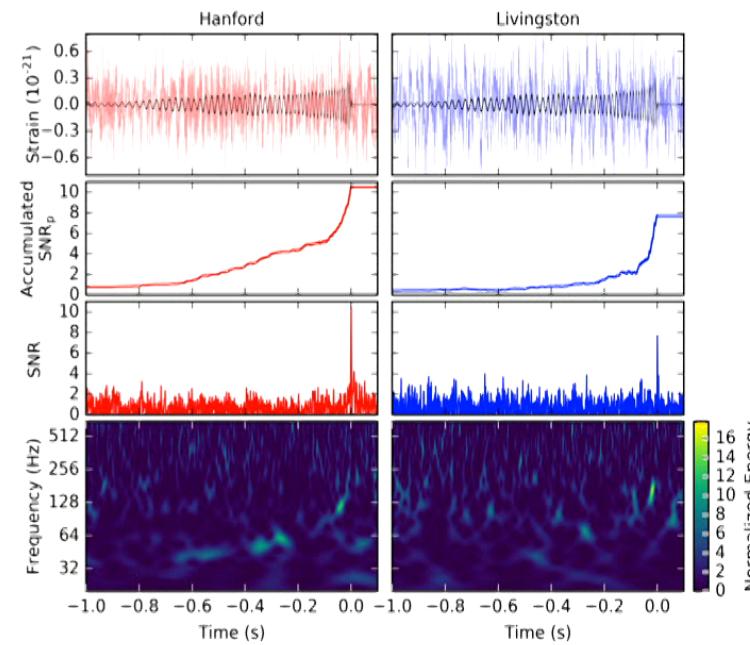


B.P. Abbott *et al.*, PRL 116 (2016) 061102

GW150914

Two binary black hole mergers

(louder in **Hanford (H1)** than **Livingston (L1)** interferometer)



B.P. Abbott *et al.*, PRL 116 (2016) 241103

GW151226

The O2 Run

After O1 completed in January 2016, both observatories began preparations for the the O2 run planned for the fall:

- ❑ Mitigate some non-fundamental noise sources seen in O1
- ❑ Raise laser powers to reduce fundamental noise and demonstrate mitigation of parametric instability (PI) associated with high power

The O2 Run

After O1 completed in January 2016, both observatories began preparations for the the O2 run planned for the fall:

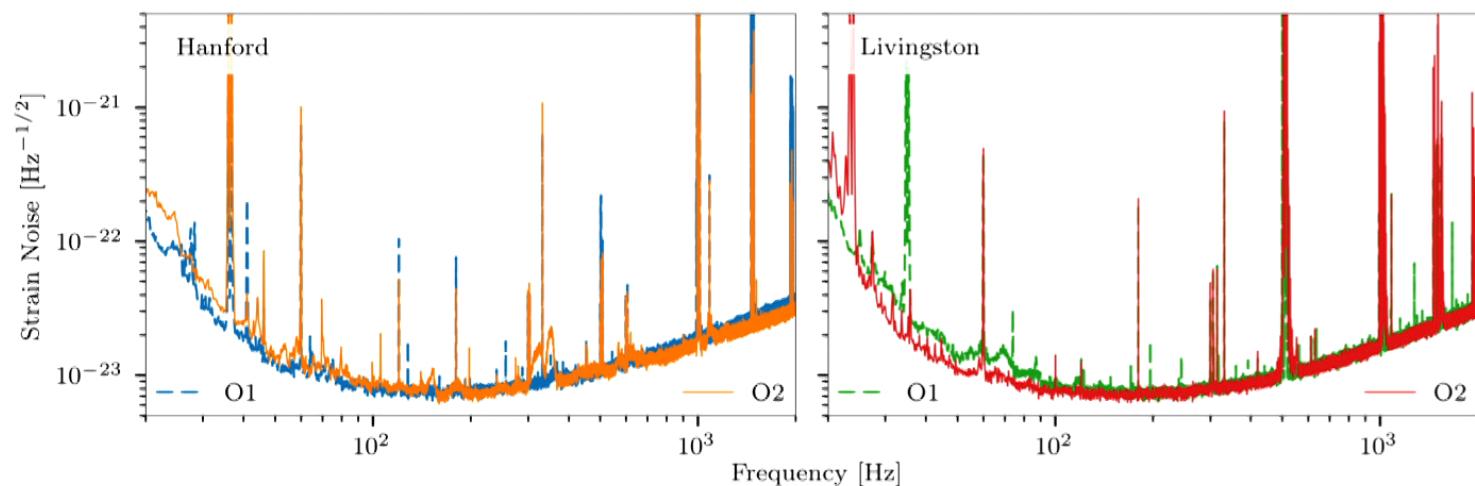
- ❑ Mitigate some non-fundamental noise sources seen in O1
- ❑ Raise laser powers to reduce fundamental noise and demonstrate mitigation of parametric instability (PI) associated with high power

Mishap at Livingston derailed high-power plans for 2016, but other noise mitigation paid off well

Hanford learned to cope well with PIs at higher power, but encountered other technical problems at higher power and had to back off (for O2)

The O2 Run

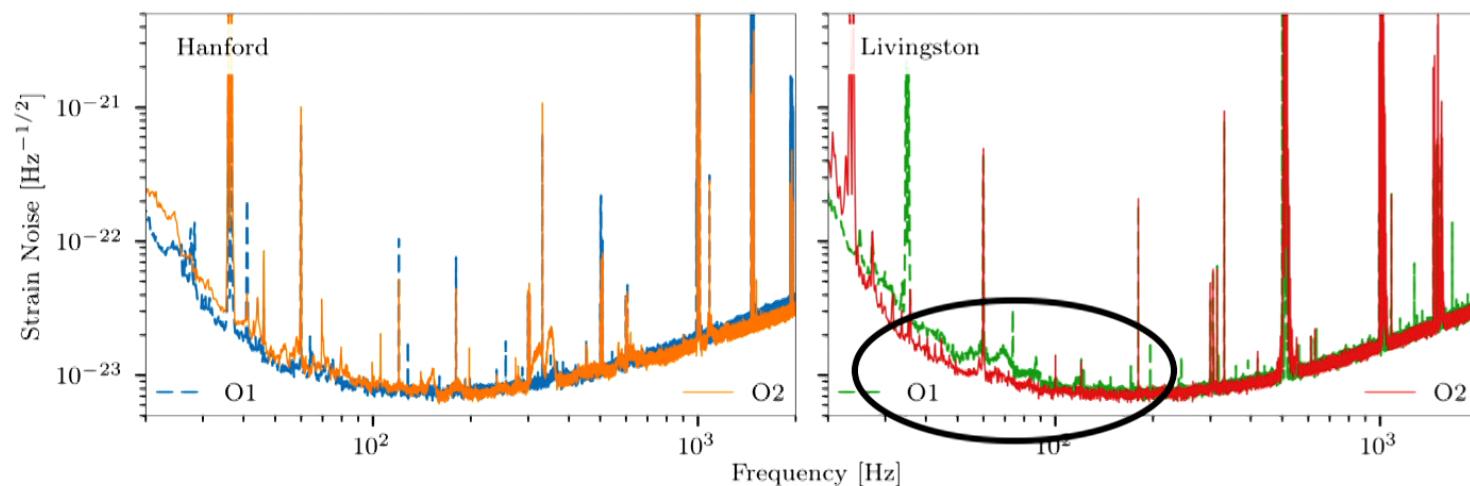
- Hanford less sensitive 😞
- Livingston more sensitive in O2 than in O1 😊



B.P. Abbott *et al*, PRL 118, 221101 (2017)

The O2 Run

- Hanford less sensitive 😞
- Livingston more sensitive in O2 than in O1 😊

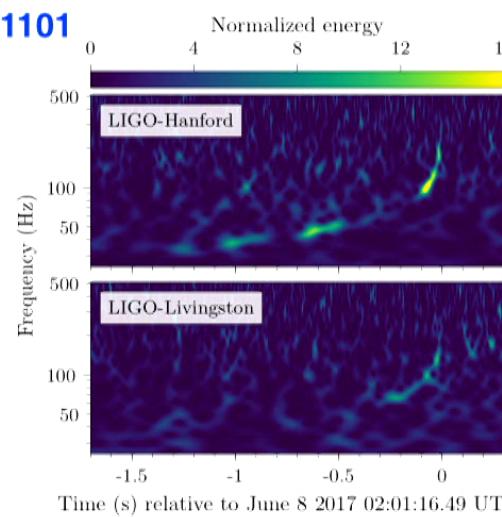
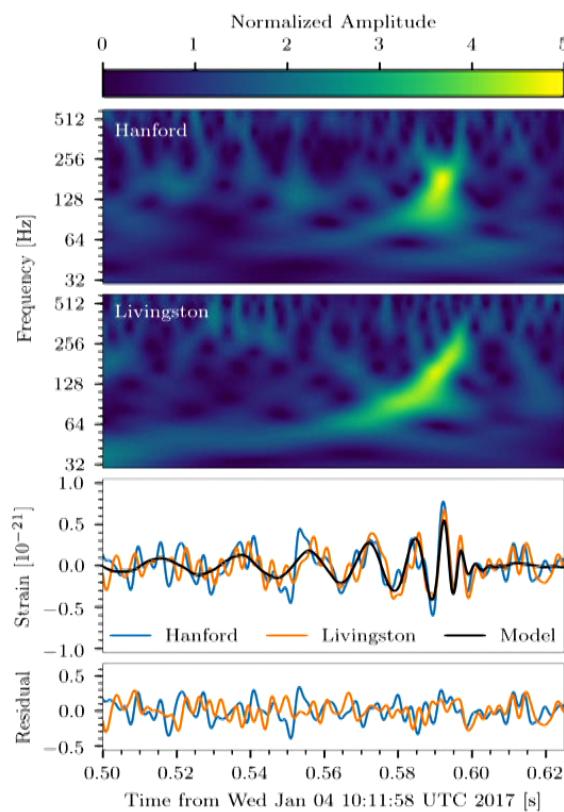


B.P. Abbott *et al*, PRL 118, 221101 (2017)

K. Riles - Overview of CW Searches

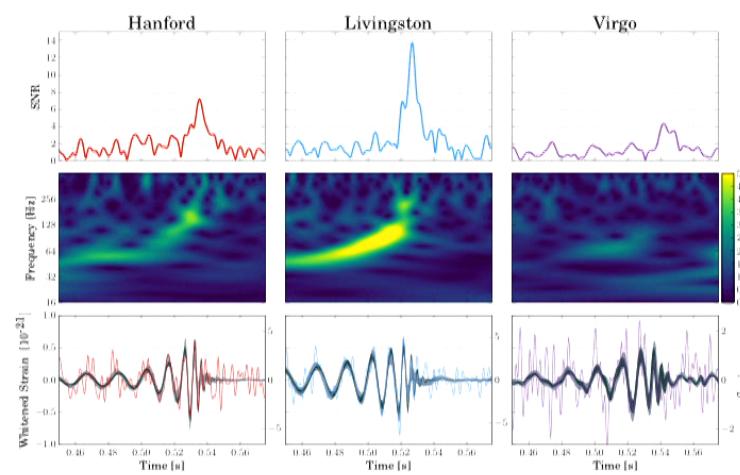
This band helpful for
detecting binary black holes
– and young pulsars
– and QCD axions?

B.P. Abbott et al., PRL 118 (2017) 221101



GW170608

B.P. Abbott et al.,
ApJL 851 (2017) L35



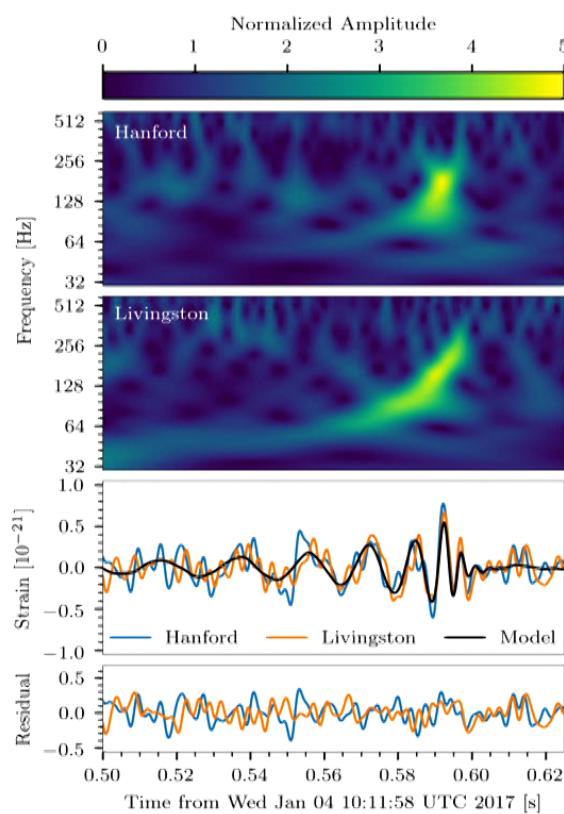
B.P. Abbott et al.,
PRL 119 (2017) 141101

GW170814

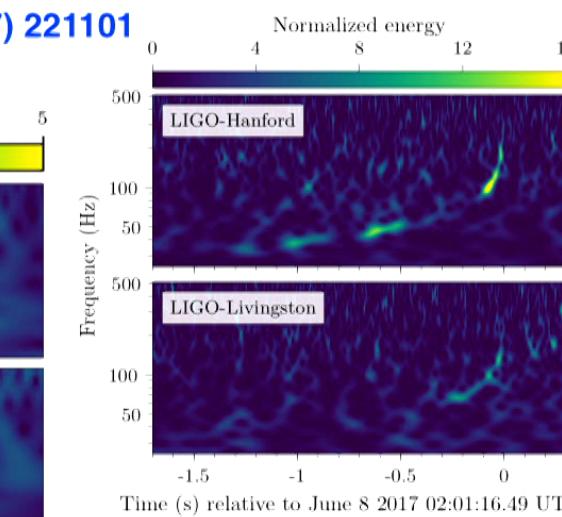
8

More black hole
mergers

B.P. Abbott et al., PRL 118 (2017) 221101

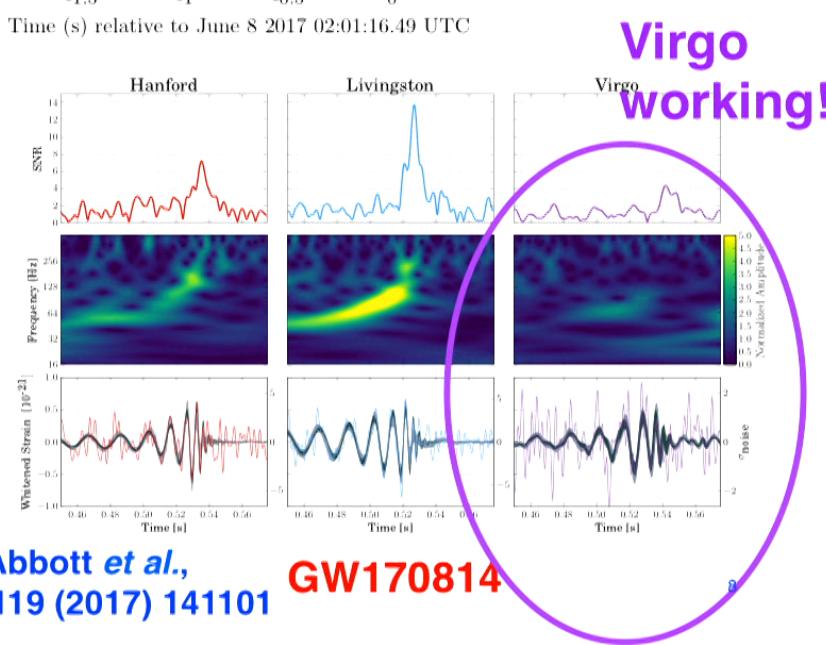


More black hole
mergers



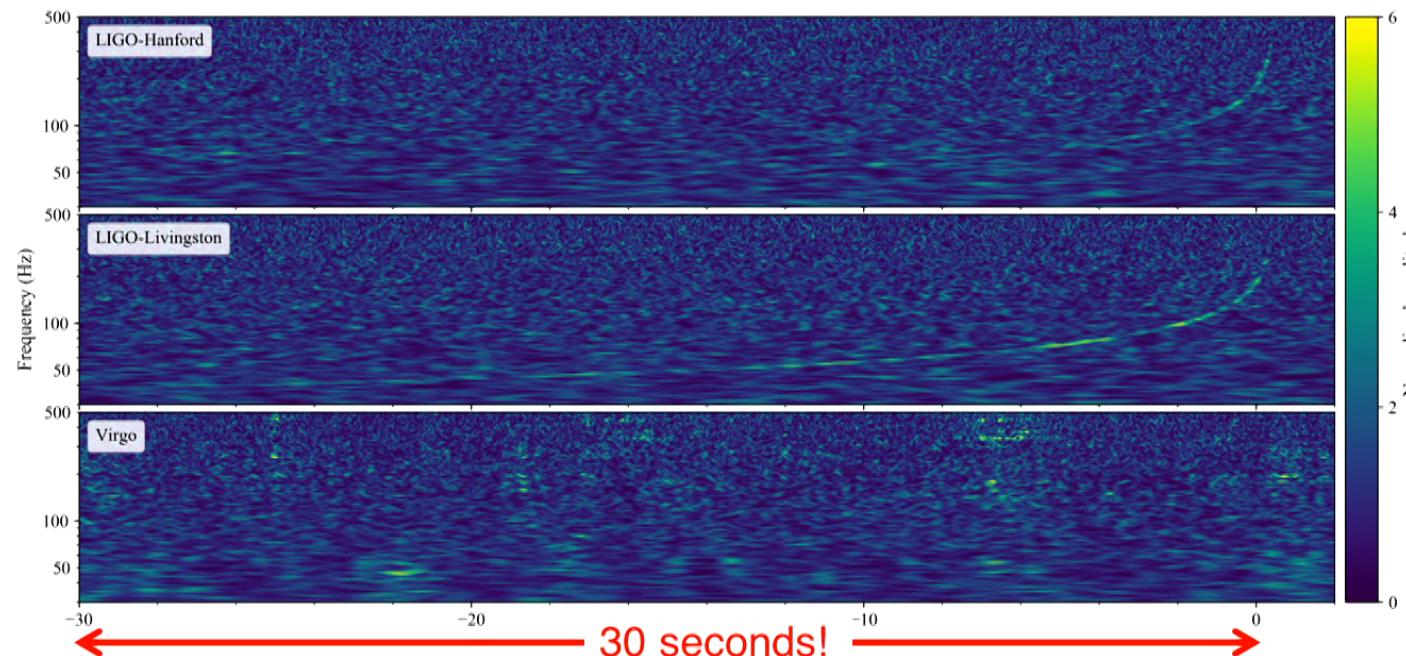
GW170608

B.P. Abbott et al.,
ApJL 851 (2017) L35



Not to mention a binary neutron star merger...

GW170817



B.P. Abbott *et al.*,
PRL 120 (2017) 161101



CAASTRO

9

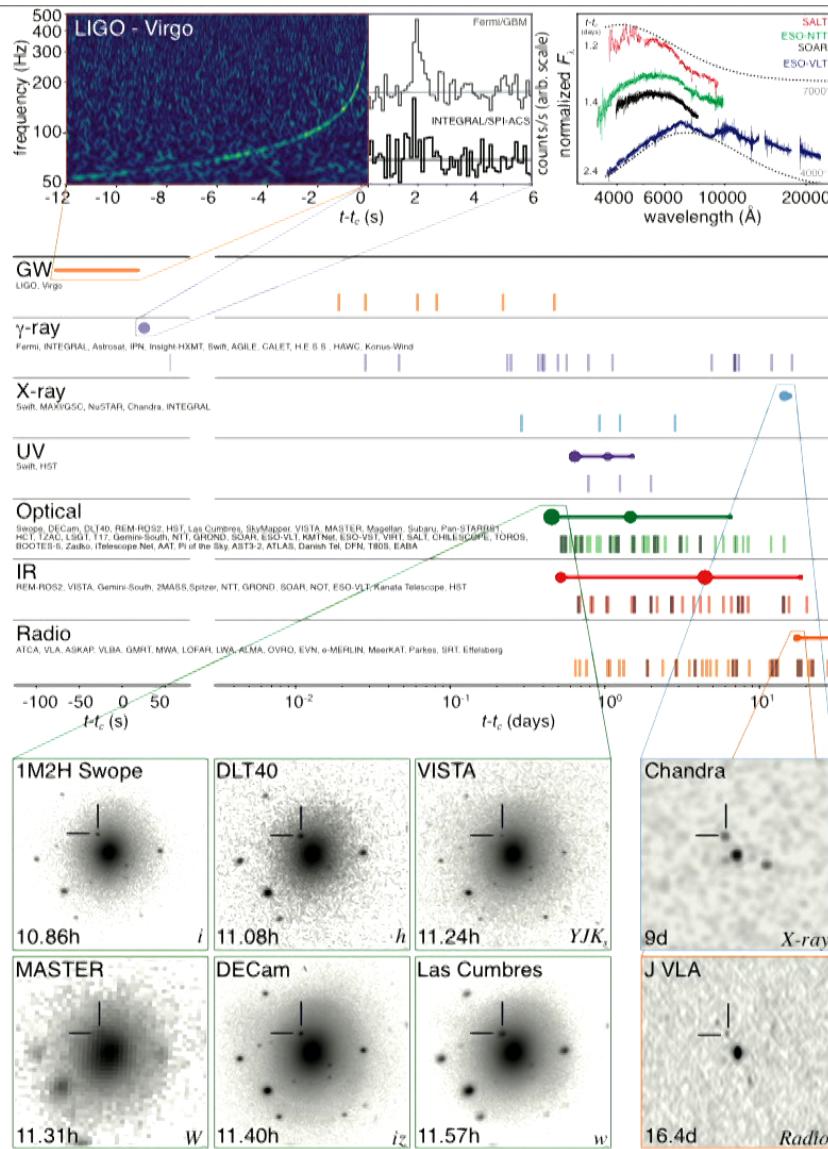
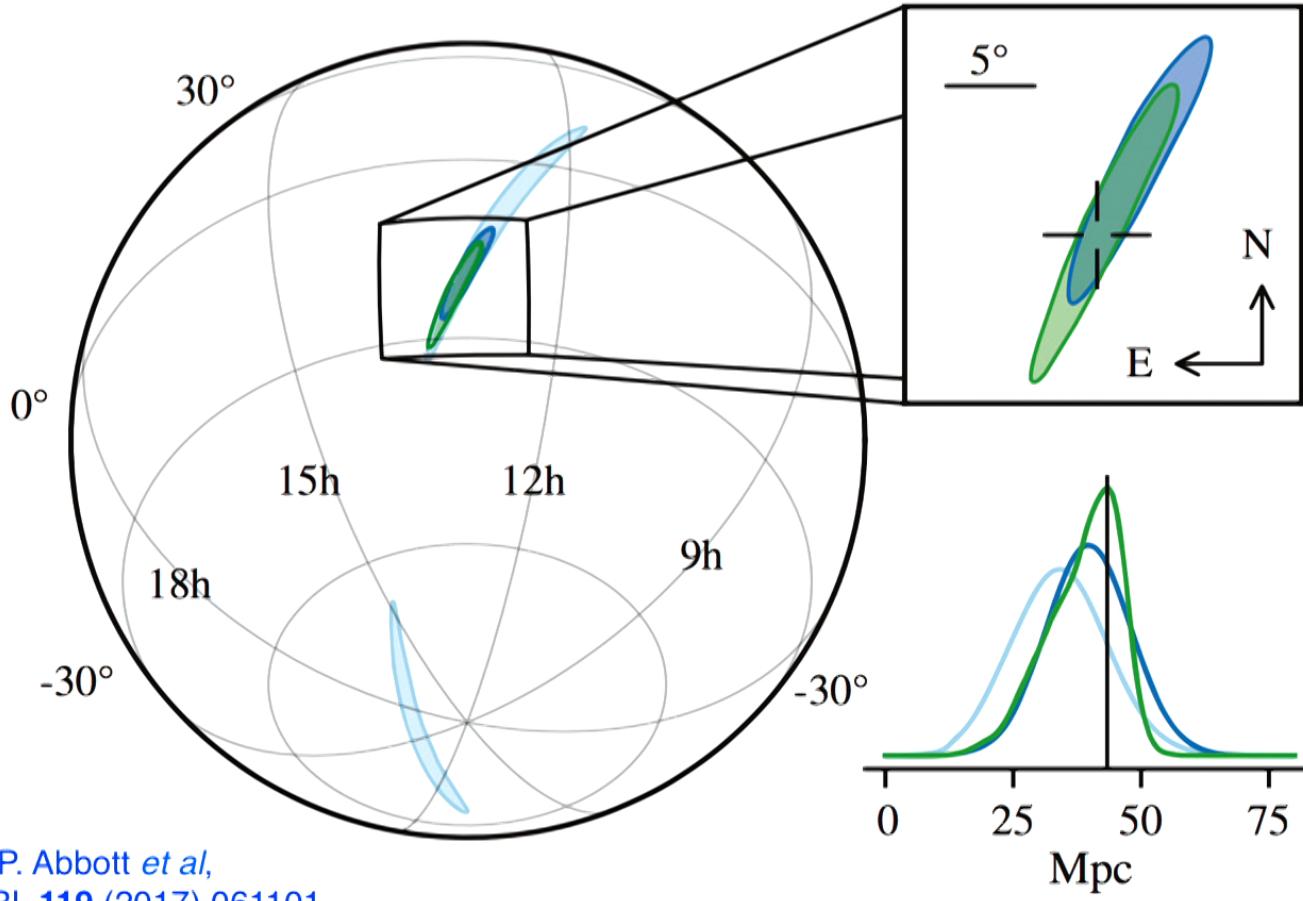


Figure from

“Multi-Messenger Observations of a Binary Neutron Star Merger”

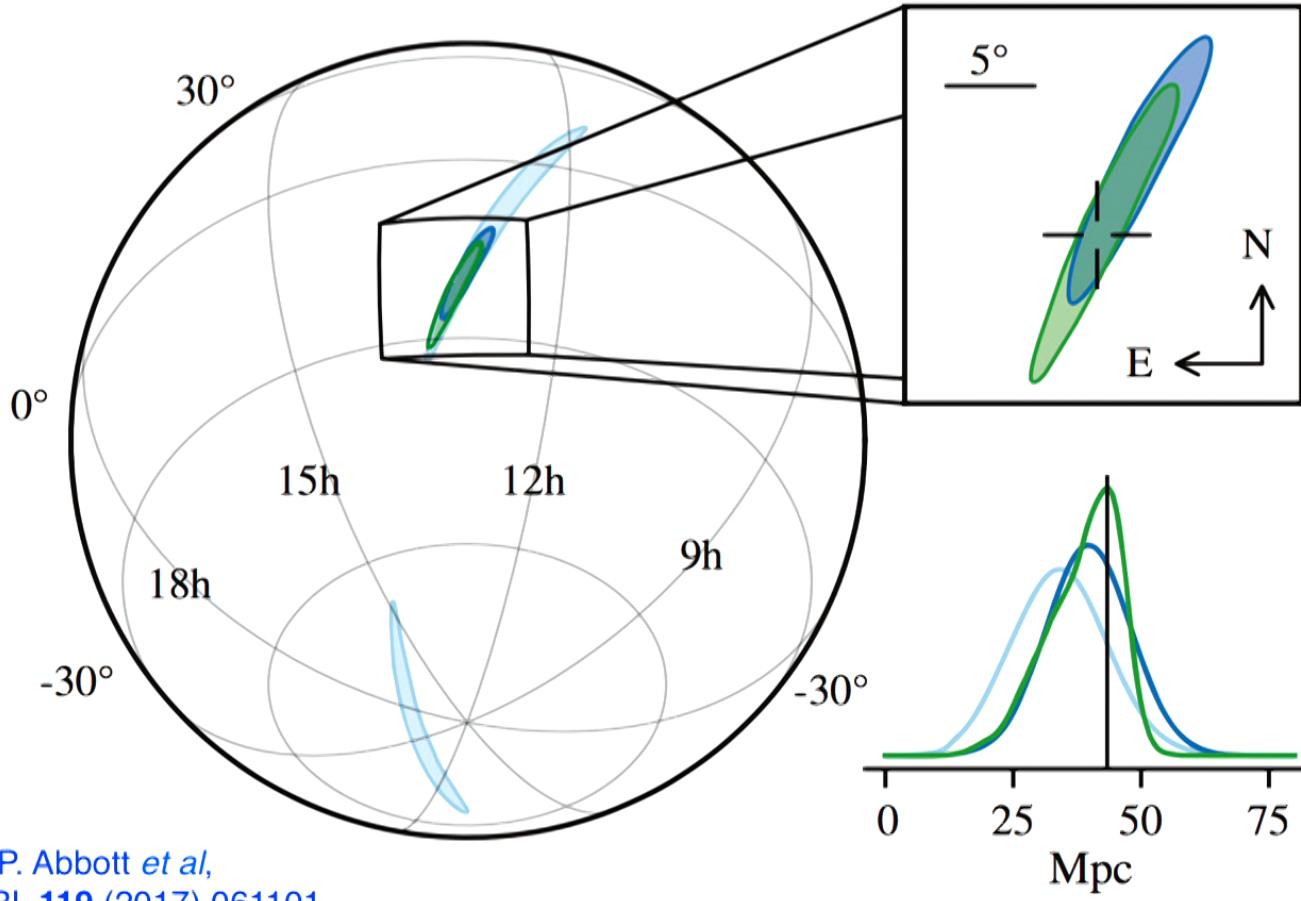
*B. Abbott et al.,
ApJL 848 (2017) L12*

59-page “letter” (!)
More than 3000 authors,
~70 collaborations



B.P. Abbott *et al.*,
PRL 119 (2017) 061101

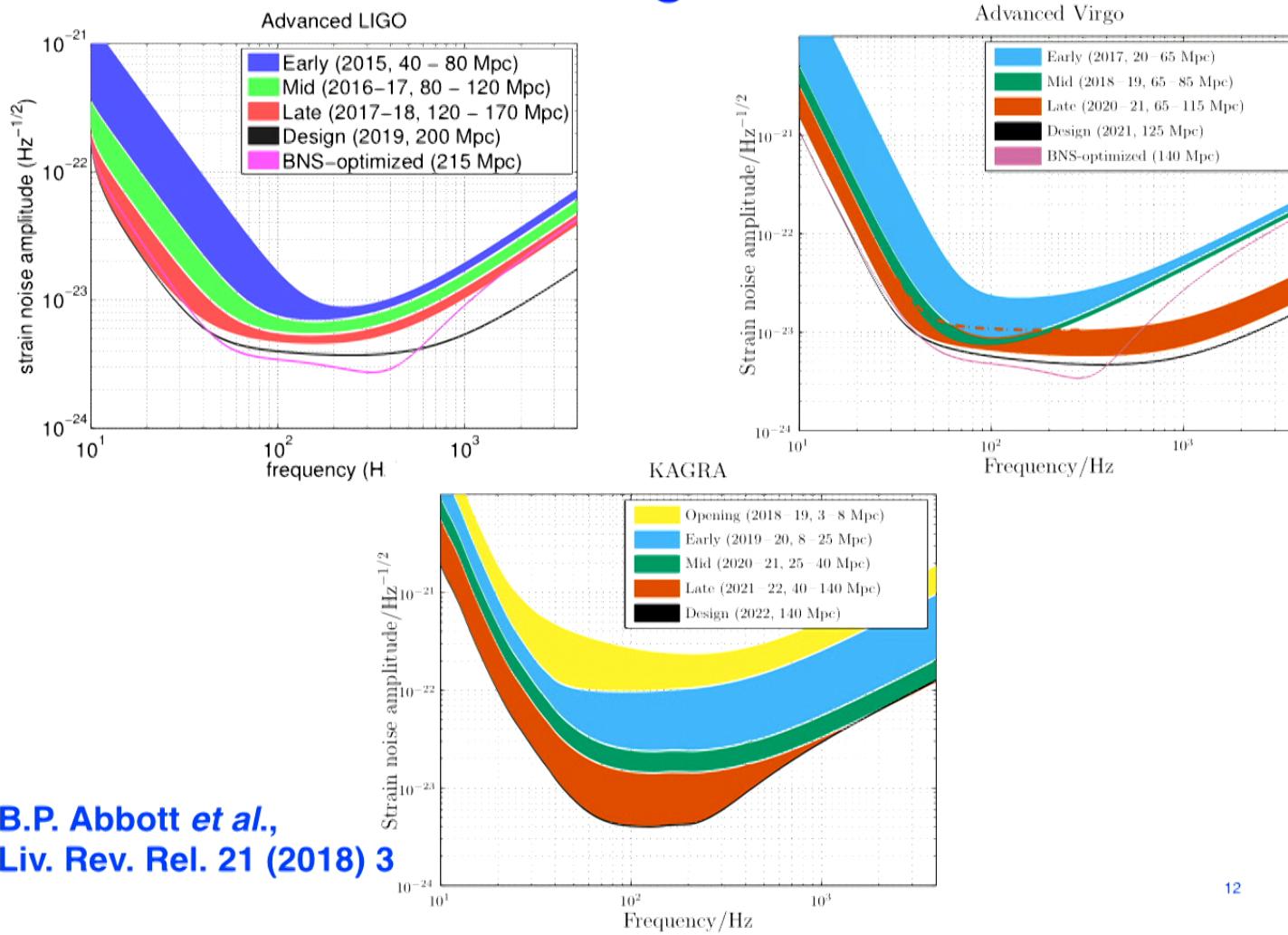
11



**A CW detection could give far better localization – sub-arcsec
(Aperture of detector is Earth's orbit for a 1-year observation)**

¹¹

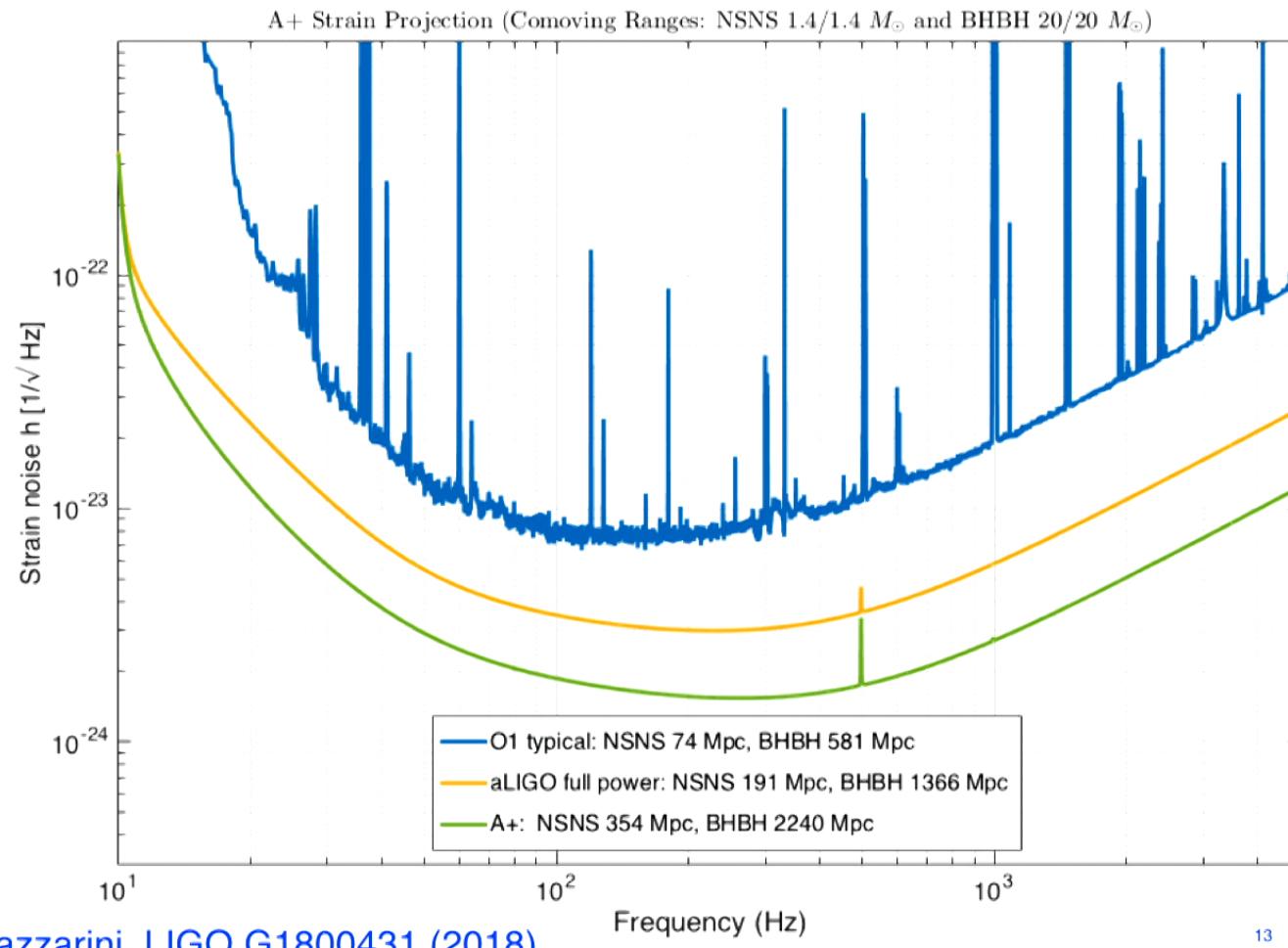
Looking Ahead



B.P. Abbott *et al.*,
Liv. Rev. Rel. 21 (2018) 3

12

Looking Ahead



A. Lazzarini, LIGO G1800431 (2018)

13

CW Search Assumptions

The LIGO and Virgo joint CW search group fields a diverse suite of programs to cover various source scenarios

Diversity applies to source types:

- Isolated neutron stars
- Neutron stars in binary systems
- Neutron stars undergoing accretion
- Newborn neutron stars (including BNS post-merger remnants)
- Glitching neutron stars

CW Search Assumptions

The LIGO and Virgo joint CW search group fields a diverse suite of programs to cover various source scenarios

Diversity applies to source types:

- Isolated neutron stars
- Neutron stars in binary systems
- Neutron stars undergoing accretion
- Newborn neutron stars (including BNS post-merger remnants)
- Glitching neutron stars
- **Black hole super-radiance**

Diversity applies to algorithms:

- Matched filters
- Time domain heterodynes
- Semi-coherent power sums
- “Loose coherence” power sums
- Hough transforms
- Viterbi dynamical programming
- Sideband summing in orbital systems
- Double-Fourier Spectra

14

CW Search Assumptions – Emission Mechanisms

- Radiation generated by quadrupolar mass movements:

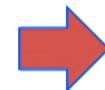
$$h_{\mu\nu} = \frac{2G}{rc^4} \frac{d^2}{dt^2} [I_{\mu\nu}]$$

($I_{\mu\nu}$ = quadrupole tensor, r = source distance)

CW Search Assumptions – Emission Mechanisms

- Radiation generated by quadrupolar mass movements:

$$h_{\mu\nu} = \frac{2G}{rc^4} \frac{d^2}{dt^2} [I_{\mu\nu}]$$

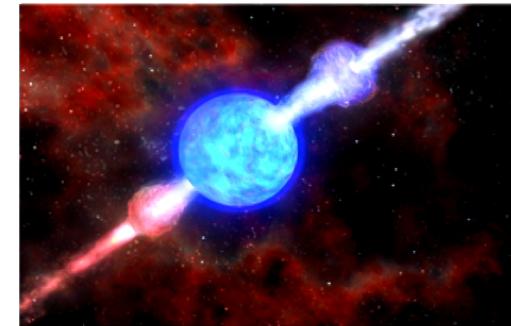


No GW from axisymmetric object rotating about symmetry axis

($I_{\mu\nu}$ = quadrupole tensor, r = source distance)

- Spinning neutron star with equatorial ellipticity $\varepsilon_{\text{equat}}$

$$\varepsilon_{\text{equat}} = \frac{|I_{xx} - I_{yy}|}{I_{zz}}$$

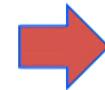


Courtesy: U. Liverpool

CW Search Assumptions – Emission Mechanisms

- Radiation generated by quadrupolar mass movements:

$$h_{\mu\nu} = \frac{2G}{rc^4} \frac{d^2}{dt^2} [I_{\mu\nu}]$$

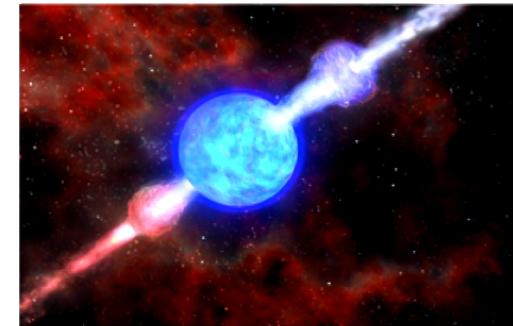


No GW from axisymmetric object rotating about symmetry axis

($I_{\mu\nu}$ = quadrupole tensor, r = source distance)

- Spinning neutron star with equatorial ellipticity ϵ_{equat}

$$\epsilon_{\text{equat}} = \frac{|I_{xx} - I_{yy}|}{I_{zz}}$$



Courtesy: U. Liverpool

gives a strain amplitude h ($f_{\text{GW}} = 2 \cdot f_{\text{Rot}}$):

$$h = 1.1 \times 10^{-24} \left[\frac{kpc}{r} \right] \left[\frac{f_{\text{GW}}}{\text{kHz}} \right]^2 \left[\frac{\epsilon}{10^{-6}} \right] \left[\frac{I_{zz}}{10^{38} \text{ kg} \cdot \text{m}^2} \right]$$

15

Gravitational CW mechanisms

- Equatorial ellipticity (e.g., – mm-high “bulge”):

$$h \propto \epsilon_{\text{equat}} \quad \text{with} \quad f_{GW} = 2f_{\text{rot}}$$

- Poloidal ellipticity (natural) + wobble angle (precessing star):

$$h \propto \epsilon_{\text{poloidal}} \times \theta_{\text{wobble}} \quad \text{with} \quad f_{GW} = f_{\text{rot}} \pm f_{\text{precess}}$$

→ Precession due to different \mathbf{L} and $\boldsymbol{\Omega}$ axes

Gravitational CW mechanisms

- Equatorial ellipticity (e.g., – mm-high “bulge”):

$$h \propto \epsilon_{\text{equat}} \quad \text{with} \quad f_{GW} = 2f_{\text{rot}}$$

- Poloidal ellipticity (natural) + wobble angle (precessing star):

$$h \propto \epsilon_{\text{poloidal}} \times \theta_{\text{wobble}} \quad \text{with} \quad f_{GW} = f_{\text{rot}} \pm f_{\text{precess}}$$

→ Precession due to different L and Ω axes

- Two-component (crust+superfluid) → $f_{GW} = f_{\text{rot}}$ and $2f_{\text{rot}}$

Gravitational CW mechanisms

- Equatorial ellipticity (e.g., – mm-high “bulge”):

$$h \propto \epsilon_{\text{equat}} \quad \text{with} \quad f_{GW} = 2f_{\text{rot}}$$

- Poloidal ellipticity (natural) + wobble angle (precessing star):

$$h \propto \epsilon_{\text{poloidal}} \times \theta_{\text{wobble}} \quad \text{with} \quad f_{GW} = f_{\text{rot}} \pm f_{\text{precess}}$$

→ Precession due to different L and Ω axes

- Two-component (crust+superfluid) → $f_{GW} = f_{\text{rot}}$ and $2f_{\text{rot}}$

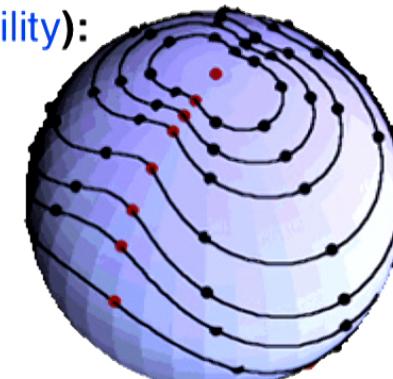
- r modes (rotational oscillations – CFS-driven instability):

N. Andersson, ApJ 502 (1998) 708

S. Chandrasekhar, PRL 24 (1970) 611

J. Friedman, B.F. Schutz, ApJ 221 (1978) 937

$$h \propto \alpha_{\text{r-mode}} \quad \text{with} \quad f_{GW} \equiv \frac{4}{3} f_{\text{rot}}$$



C. Hanna & B.J. Owen¹⁶

Gravitational CW mechanisms

- Equatorial ellipticity (e.g., – mm-high “bulge”):

$$h \propto \epsilon_{\text{equat}} \quad \text{with} \quad f_{GW} = 2f_{\text{rot}}$$

- Poloidal ellipticity (natural) + wobble angle (precessing star):

$$h \propto \epsilon_{\text{poloidal}} \times \theta_{\text{wobble}} \quad \text{with} \quad f_{GW} = f_{\text{rot}} \pm f_{\text{precess}}$$

→ Precession due to different L and Ω axes

- Two-component (crust+superfluid) → $f_{GW} = f_{\text{rot}}$ and $2f_{\text{rot}}$

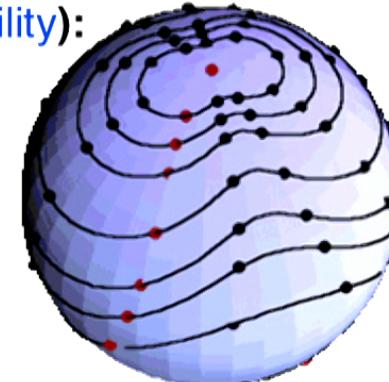
- r modes (rotational oscillations – CFS-driven instability):

N. Andersson, ApJ 502 (1998) 708

S. Chandrasekhar, PRL 24 (1970) 611

J. Friedman, B.F. Schutz, ApJ 221 (1978) 937

$$h \propto \alpha_{\text{r-mode}} \quad \text{with} \quad f_{GW} \equiv \frac{4}{3} f_{\text{rot}}$$



C. Hanna & B.J. Owen¹⁶

Gravitational CW mechanisms

- Equatorial ellipticity (e.g., – mm-high “bulge”):

$$h \propto \epsilon_{\text{equat}} \quad \text{with} \quad f_{GW} = 2f_{\text{rot}}$$

- Poloidal ellipticity (natural) + wobble angle (precessing star):

$$h \propto \epsilon_{\text{poloidal}} \times \theta_{\text{wobble}} \quad \text{with} \quad f_{GW} = f_{\text{rot}} \pm f_{\text{precess}}$$

→ Precession due to different L and Ω axes

- Two-component (crust+superfluid) → $f_{GW} = f_{\text{rot}}$ and $2f_{\text{rot}}$

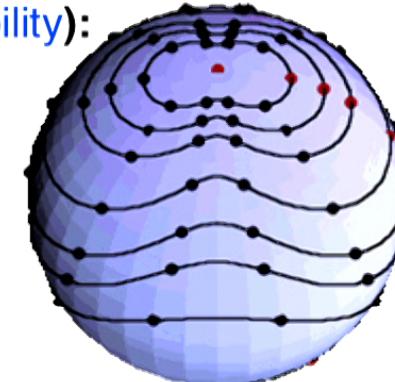
- r modes (rotational oscillations – CFS-driven instability):

N. Andersson, ApJ 502 (1998) 708

S. Chandrasekhar, PRL 24 (1970) 611

J. Friedman, B.F. Schutz, ApJ 221 (1978) 937

$$h \propto \alpha_{\text{r-mode}} \quad \text{with} \quad f_{GW} \equiv \frac{4}{3} f_{\text{rot}}$$



C. Hanna & B.J. Owen

Gravitational CW mechanisms

Assumption we (LSC, Virgo) have usually made to date:

Bulge is best bet for detection

→ Look for GW emission at twice the EM frequency

e.g., look for Crab Pulsar (29.6 Hz) at 59.3 Hz
(troublesome frequency in North America!)

Gravitational CW mechanisms

Assumption we (LSC, Virgo) have usually made to date:

Bulge is best bet for detection

→ Look for GW emission at twice the EM frequency

e.g., look for Crab Pulsar (29.6 Hz) at 59.3 Hz
(troublesome frequency in North America!)

What is allowed for ϵ_{equat} ?

Old maximum (?) $\approx 5 \times 10^{-7}$ [$\sigma/10^{-2}$] (“ordinary” neutron star)
with σ = breaking strain of crust

G. Ushomirsky, C. Cutler & L. Bildsten, MNRAS 319 (2000) 902

Gravitational CW mechanisms

Assumption we (LSC, Virgo) have usually made to date:

Bulge is best bet for detection

→ Look for GW emission at twice the EM frequency

e.g., look for Crab Pulsar (29.6 Hz) at 59.3 Hz
(troublesome frequency in North America!)

What is allowed for ϵ_{equat} ?

Old maximum (?) $\approx 5 \times 10^{-7}$ $[\sigma/10^{-2}]$ (“ordinary” neutron star)
with σ = breaking strain of crust

G. Ushomirsky, C. Cutler & L. Bildsten, MNRAS 319 (2000) 902

More recent finding: $\sigma \approx 10^{-1}$ supported by detailed numerical simulation
C.J. Horowitz & K. Kadau, PRL 102, (2009) 191102

Gravitational CW mechanisms

Assumption we (LSC, Virgo) have usually made to date:

Bulge is best bet for detection

→ Look for GW emission at twice the EM frequency

e.g., look for Crab Pulsar (29.6 Hz) at 59.3 Hz
(troublesome frequency in North America!)

What is allowed for ϵ_{equat} ?

Old maximum (?) $\approx 5 \times 10^{-7}$ [$\sigma/10^{-2}$] (“ordinary” neutron star)
with σ = breaking strain of crust

G. Ushomirsky, C. Cutler & L. Bildsten, MNRAS 319 (2000) 902

More recent finding: $\sigma \approx 10^{-1}$ supported by detailed numerical simulation
C.J. Horowitz & K. Kadau, PRL 102, (2009) 191102

Recent re-evaluation: $\epsilon_{\text{equat}} < 10^{-5}$

N.K. Johnson-McDaniel & B.J. Owen, PRD 88 (2013) 044004

17

CW Search Assumptions – Promising sources

Three broad categories of searches have dominated analysis:

→ See Sylvia Zhu's talk for details on 1st two categories

Targeted searches for known pulsars using radio / X-ray / γ-ray ephemerides

→ Exact phase tracking over O(years) – low trials factor

→ Variation (“narrowband”) allows for EM/GW $\Delta f \sim O(10^{-3}) f_{EM}$

CW Search Assumptions – Promising sources

Three broad categories of searches have dominated analysis:

→ See Sylvia Zhu's talk for details on 1st two categories

Targeted searches for known pulsars using radio / X-ray / γ-ray ephemerides

→ Exact phase tracking over O(years) – low trials factor

→ Variation (“narrowband”) allows for EM/GW $\Delta f \sim O(10^{-3}) f_{EM}$

Directed searches for known sources / locations

(unknown / poorly known frequencies)

- Isolated and binary sources treated separately
- Fully coherent searches over days/weeks
- Semi-coherent searches over full data runs

CW Search Assumptions – Promising sources

Three broad categories of searches have dominated analysis:

→ See Sylvia Zhu's talk for details on 1st two categories

Targeted searches for known pulsars using radio / X-ray / γ-ray ephemerides

→ Exact phase tracking over O(years) – low trials factor

→ Variation (“narrowband”) allows for EM/GW $\Delta f \sim O(10^{-3}) f_{EM}$

Directed searches for known sources / locations

(unknown / poorly known frequencies)

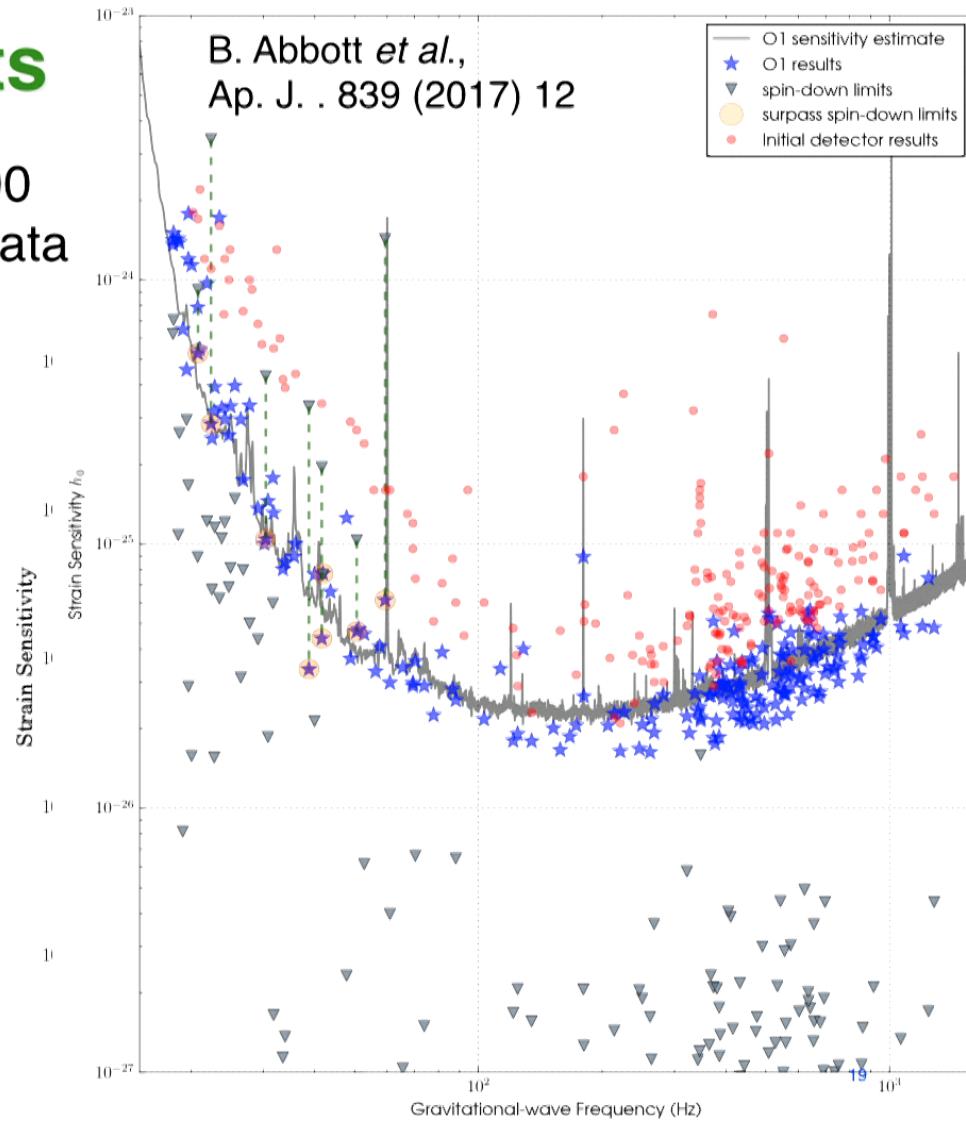
- Isolated and binary sources treated separately
- Fully coherent searches over days/weeks
- Semi-coherent searches over full data runs

All-sky searches for unknown sources

- Isolated and binary sources treated separately
(binary esp. challenging)
- Semi-coherent searches over full data runs

Recent results

Targeted search for 200 known pulsars in O1 data

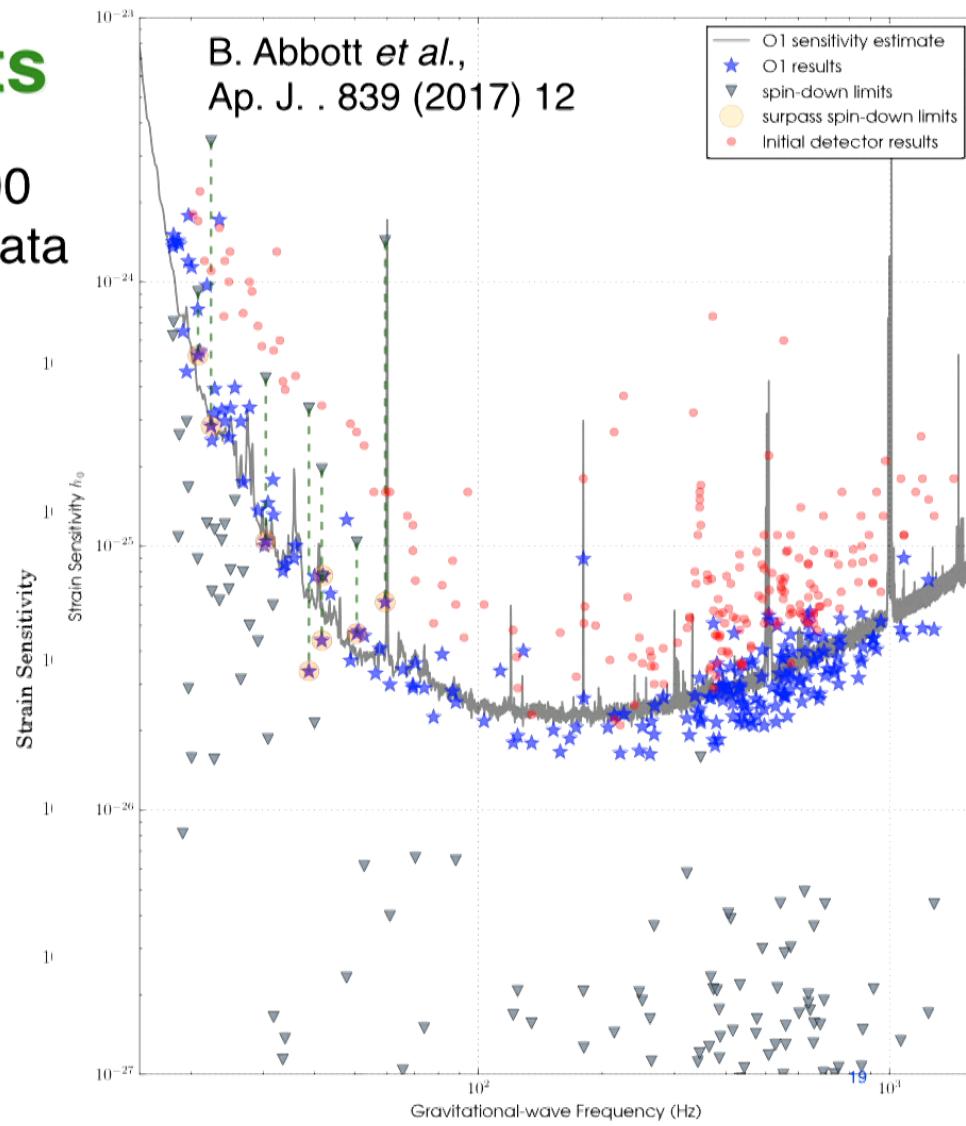


Recent results

Targeted search for 200 known pulsars in O1 data

Lowest (best) upper limit on strain:

$$h_0 < 1.6 \times 10^{-26}$$



Recent results

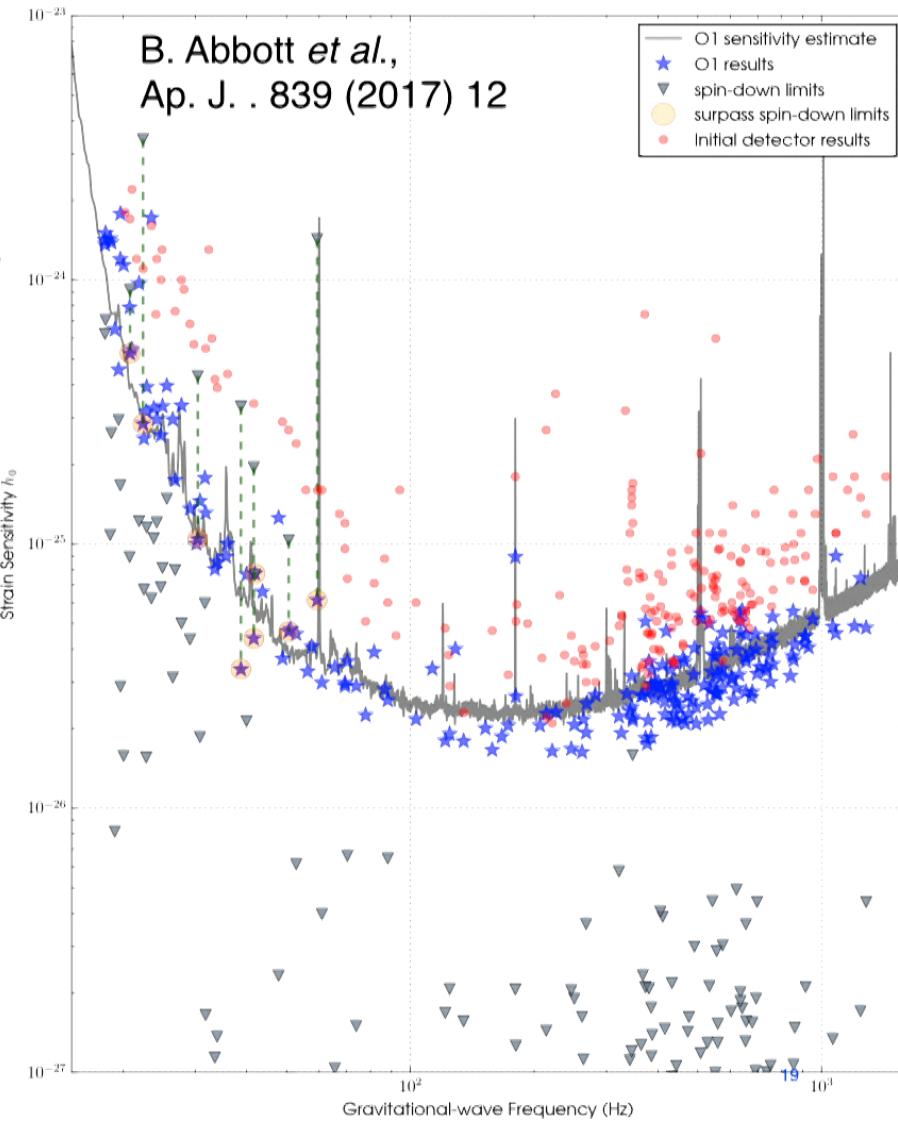
Targeted search for 200 known pulsars in O1 data

Lowest (best) upper limit on strain:

$$h_0 < 1.6 \times 10^{-26}$$

Lowest (best) upper limit on ellipticity:

$$\varepsilon < 1.3 \times 10^{-8}$$



Recent results

Targeted search for 200 known pulsars in O1 data

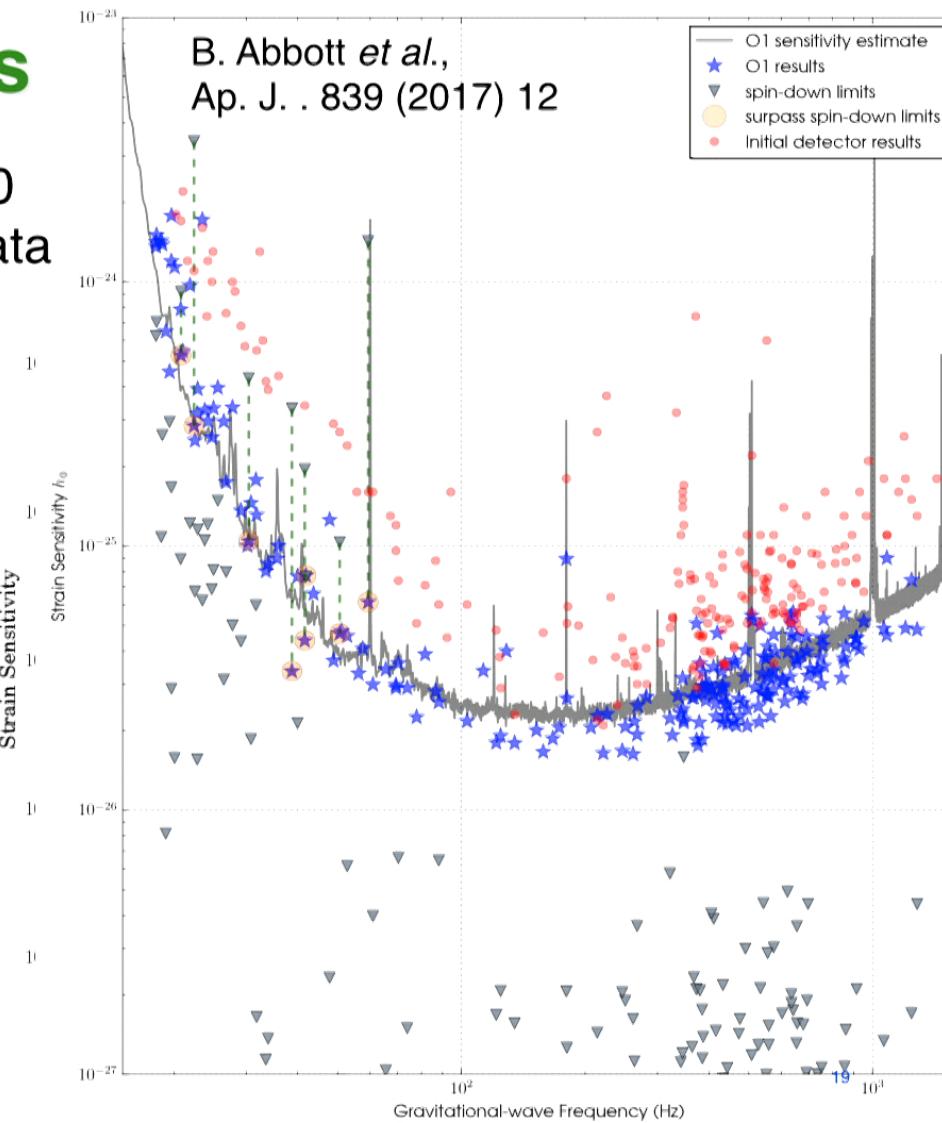
Lowest (best) upper limit on strain:

$$h_0 < 1.6 \times 10^{-26}$$

Lowest (best) upper limit on ellipticity:

$$\varepsilon < 1.3 \times 10^{-8}$$

*Crab limit at 0.2% of total energy loss
(beats “spindown limit”)*



Sensitivity Depth

Different search algorithms use different fractions of available data from a run and have different scalings of sensitivity with observation time [e.g., $(T_{\text{OBS}})^{1/2}$, $(T_{\text{OBS}})^{1/4}$]

Sensitivity Depth

Examples: (current algorithms)

- Targeted search over 1 year: ~500
- “Narrowband” search over 4 months: ~100-150
- Directed isolated search over 10 days: ~30
- Semi-coherent directed isolated search over 1 year: ~60
- Cross-correlation directed Sco X-1 search over 4 months: ~80
- All-sky isolated search over 4 months: ~25
- All-sky binary search over 1 year: ~3

These depths affect how we choose to focus resources on searches
(computational, human)

Computing challenges – All-sky Searches

Serious technical difficulty: Doppler frequency shifts

- Frequency modulation from earth's rotation ($v/c \sim 10^{-6}$)
 - Frequency modulation from earth's orbital motion ($v/c \sim 10^{-4}$)
- Coherent integration of 1 year gives frequency resolution of 30 nHz
→ 1 kHz source spread over 6 million bins in ordinary FFT!

Computing challenges – All-sky Searches

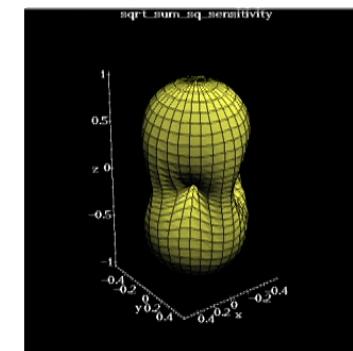
Serious technical difficulty: Doppler frequency shifts

- Frequency modulation from earth's rotation ($v/c \sim 10^{-6}$)
- Frequency modulation from earth's orbital motion ($v/c \sim 10^{-4}$)
- Coherent integration of 1 year gives frequency resolution of 30 nHz
- 1 kHz source spread over 6 million bins in ordinary FFT!

Additional, related complications:

**Daily amplitude modulation of antenna pattern
(polarization dependent)**

Still have frequency derivatives to address...



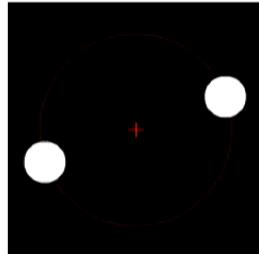
Computing challenges – All-sky Searches

Serious technical difficulty: Doppler frequency shifts

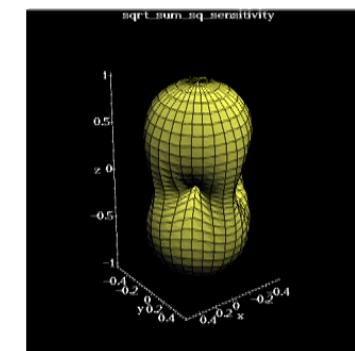
- Frequency modulation from earth's rotation ($v/c \sim 10^{-6}$)
- Frequency modulation from earth's orbital motion ($v/c \sim 10^{-4}$)
- Coherent integration of 1 year gives frequency resolution of 30 nHz
- 1 kHz source spread over 6 million bins in ordinary FFT!

Additional, related complications:

Daily amplitude modulation of antenna pattern
(polarization dependent)
Still have frequency derivatives to address...



Orbital motion of sources in binary systems



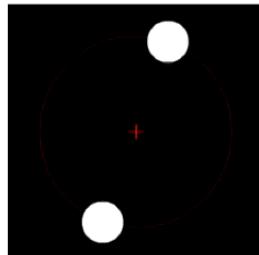
Computing challenges – All-sky Searches

Serious technical difficulty: Doppler frequency shifts

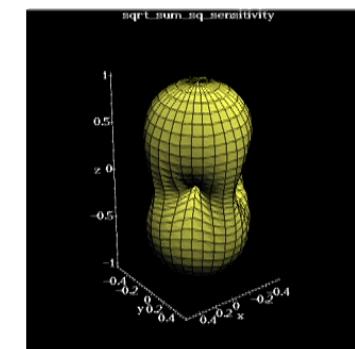
- Frequency modulation from earth's rotation ($v/c \sim 10^{-6}$)
- Frequency modulation from earth's orbital motion ($v/c \sim 10^{-4}$)
- Coherent integration of 1 year gives frequency resolution of 30 nHz
- 1 kHz source spread over 6 million bins in ordinary FFT!

Additional, related complications:

Daily amplitude modulation of antenna pattern
(polarization dependent)
Still have frequency derivatives to address...



Orbital motion of sources in binary systems



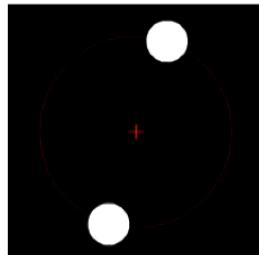
Computing challenges – All-sky Searches

Serious technical difficulty: Doppler frequency shifts

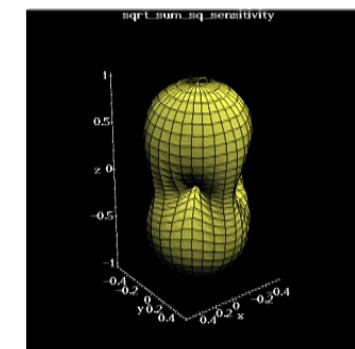
- Frequency modulation from earth's rotation ($v/c \sim 10^{-6}$)
- Frequency modulation from earth's orbital motion ($v/c \sim 10^{-4}$)
- Coherent integration of 1 year gives frequency resolution of 30 nHz
- 1 kHz source spread over 6 million bins in ordinary FFT!

Additional, related complications:

Daily amplitude modulation of antenna pattern
(polarization dependent)
Still have frequency derivatives to address...



Orbital motion of sources in binary systems



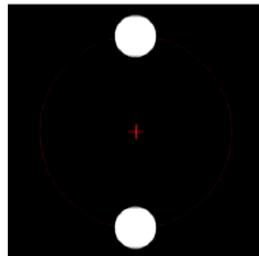
Computing challenges – All-sky Searches

Serious technical difficulty: Doppler frequency shifts

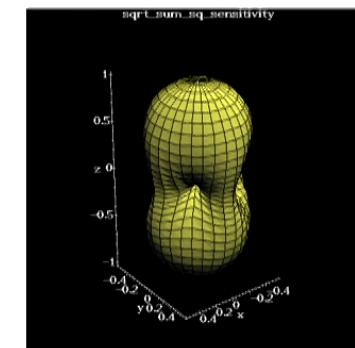
- Frequency modulation from earth's rotation ($v/c \sim 10^{-6}$)
- Frequency modulation from earth's orbital motion ($v/c \sim 10^{-4}$)
- Coherent integration of 1 year gives frequency resolution of 30 nHz
- 1 kHz source spread over 6 million bins in ordinary FFT!

Additional, related complications:

Daily amplitude modulation of antenna pattern
(polarization dependent)
Still have frequency derivatives to address...



Orbital motion of sources in binary systems



Computing challenges – All-sky Searches

Modulations / drifts complicate analysis enormously:

- Simple Fourier transform inadequate
- Every sky direction requires different demodulation
- →Number of distinct sky location scales like $f^2 T_{\text{COH}}^2$
- →Computing cost scales as $(T_{\text{COH}})^6$ without 2nd frequency derivative!
- All-sky survey at full sensitivity not possible (at this time)

Computing challenges – All-sky Searches

Modulations / drifts complicate analysis enormously:

- Simple Fourier transform inadequate
- Every sky direction requires different demodulation
- → Number of distinct sky location scales like $f^2 T_{\text{COH}}^2$
- → Computing cost scales as $(T_{\text{COH}})^6$ without 2nd frequency derivative!
- All-sky survey at full sensitivity not possible (at this time)

Tradeoffs to cope with astronomical computing costs:

- Restrict $T_{\text{COH}} <$ several days for all-sky search
 - Compute demodulated spectra (e.g., “F-Statistic”)
 - Exploit coincidence among different time segments
 - Sensitivity scales like $(T_{\text{COH}})^{1/2}$

Computing challenges – All-sky Searches

Modulations / drifts complicate analysis enormously:

- Simple Fourier transform inadequate
- Every sky direction requires different demodulation
- → Number of distinct sky location scales like $f^2 T_{\text{COH}}^2$
- → Computing cost scales as $(T_{\text{COH}})^6$ without 2nd frequency derivative!
- All-sky survey at full sensitivity not possible (at this time)

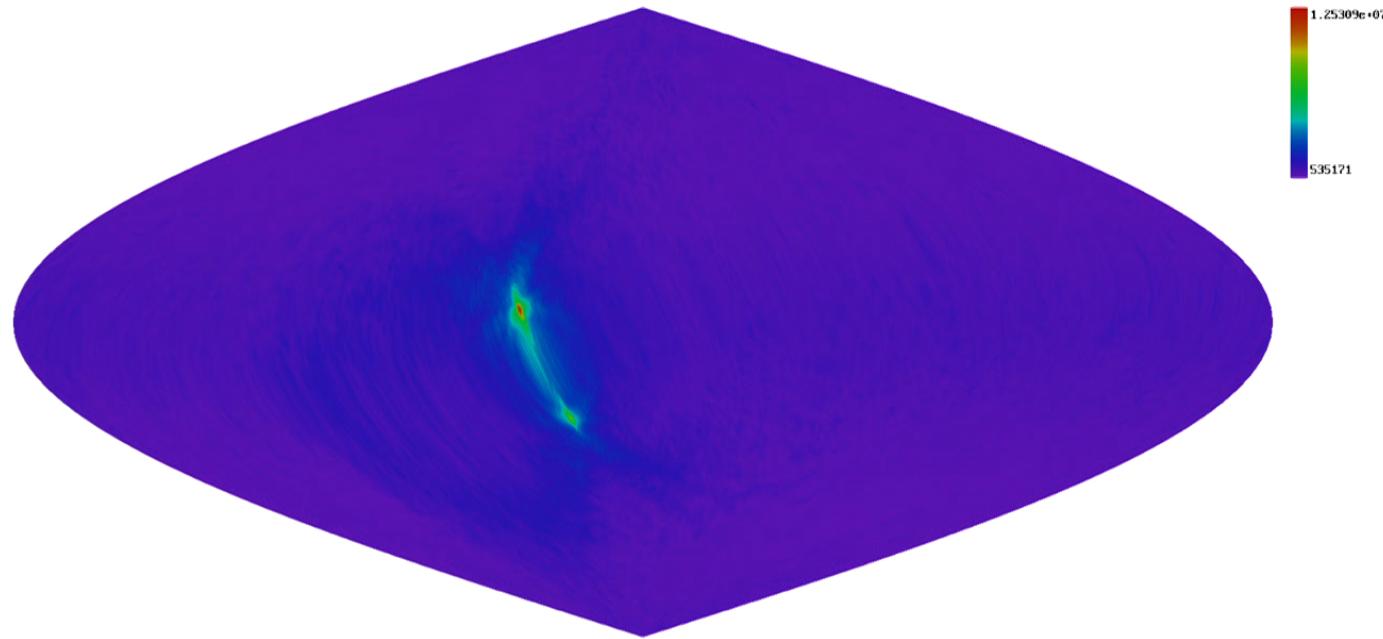
Tradeoffs to cope with astronomical computing costs:

- Restrict $T_{\text{COH}} <$ several days for all-sky search
 - Compute demodulated spectra (e.g., “F-Statistic”)
 - Exploit coincidence among different time segments
 - Sensitivity scales like $(T_{\text{COH}})^{1/2}$
- Semi-coherent stacking ($N_{\text{stack}} = T_{\text{OBS-RUN}} / T_{\text{COH}}$)
 - Raw / weighted spectral powers ($T_{\text{COH}} \sim 0.5\text{-}2.0$ hours)
 - Thresholded powers – Hough transforms
 - Demodulated spectra – F-Statistic
 - Sensitivity improves only as $(N_{\text{stack}})^{1/4}$

23

But three substantial benefits from Doppler modulations:

- Reality of signal confirmed by need for corrections
- Corrections give precise direction of source
- Single interferometer can make definitive discovery

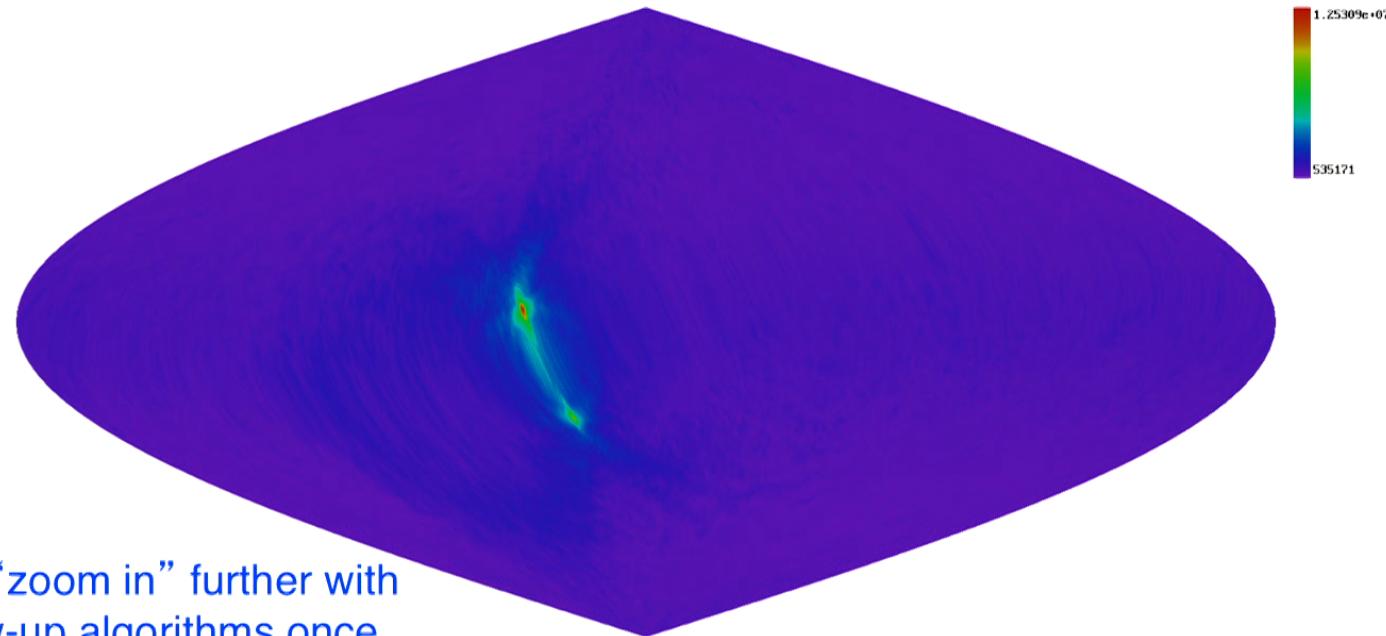


Example sky map of strain power for signal injection
(semi-coherent search)

24

But three substantial benefits from Doppler modulations:

- Reality of signal confirmed by need for corrections
- Corrections give precise direction of source
- Single interferometer can make definitive discovery



Can “zoom in” further with follow-up algorithms once we lock on to source

V. Dergachev, PRD 85 (2012) 062003
M. Shaltev & R. Prix, PRD 87 (2013) 084057
A. Singh *et al.*, PRD 96 (2017) 082003

Example sky map of strain power for signal injection (semi-coherent search)
²⁴

All-sky Searches for Unknown Isolated Stars

Search methods used in Initial LIGO data:

- Coherent F -Statistic ($T_{COH} = 10$ hours)
[P. Jaranowski, A. Krolak & B. Schutz, PRD 58 (1998) 063001;
B. Abbott *et al.*, PRD 76 (2007) 082001]
- Stacked power spectra (“stack-slide” & variants) ($T_{coh} = 0.5\text{-}2.0$ hours]
 - Stack-slide [P. Brady *et al.*, PRD 57 (1998) 2101; P. Brady & T. Creighton, PRD 61 (2000) 082001; B. Abbott *et al.*, PRD 77 (2008) 022001]
 - Sky Hough transform [A.M. Sintes & B. Krishnan, JPCS 32 (2006) 206]

All-sky Searches for Unknown Isolated Stars

Search methods used in Initial LIGO data:

- Coherent F -Statistic ($T_{\text{COH}} = 10$ hours)
[P. Jaranowski, A. Krolak & B. Schutz, PRD 58 (1998) 063001;
B. Abbott *et al.*, PRD 76 (2007) 082001]
- Stacked power spectra (“stack-slide” & variants) ($T_{\text{coh}} = 0.5\text{-}2.0$ hours]
 - Stack-slide [P. Brady *et al.*, PRD 57 (1998) 2101; P. Brady & T. Creighton, PRD 61 (2000) 082001; B. Abbott *et al.*, PRD 77 (2008) 022001]
 - Sky Hough transform [A.M. Sintes & B. Krishnan, JPCS 32 (2006) 206]
 - PowerFlux / Loose coherence [B. Abbott *et al.*, PRD 77 (2008) 022001;
V. Dergachev, CQG 27 (2010) 205017]
 - Frequency Hough transform [P. Astone *et al.*, PRD 90 (2014) 042002]
- Multi-segment F -Statistic ($T_{\text{COH}} \sim 1\text{-few days}$)

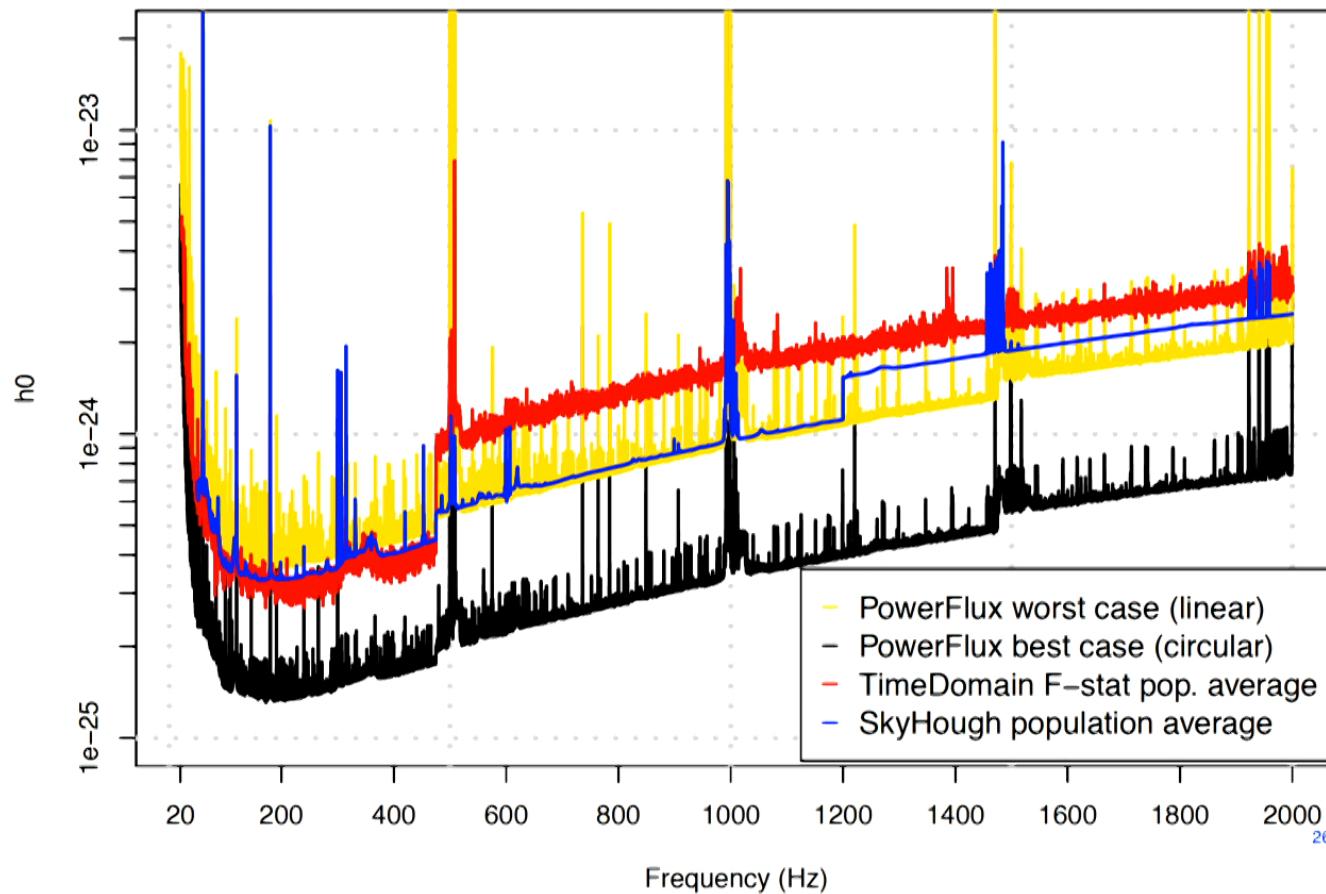
All-sky Searches for Unknown Isolated Stars

Search methods used in Initial LIGO data:

- Coherent *F*-Statistic ($T_{\text{COH}} = 10$ hours)
[P. Jaranowski, A. Krolak & B. Schutz, PRD 58 (1998) 063001;
B. Abbott *et al.*, PRD 76 (2007) 082001]
- Stacked power spectra (“stack-slide” & variants) ($T_{\text{coh}} = 0.5\text{-}2.0$ hours]
 - Stack-slide [P. Brady *et al.*, PRD 57 (1998) 2101; P. Brady & T. Creighton, PRD 61 (2000) 082001; B. Abbott *et al.*, PRD 77 (2008) 022001]
 - Sky Hough transform [A.M. Sintes & B. Krishnan, JPCS 32 (2006) 206]
 - PowerFlux / Loose coherence [B. Abbott *et al.*, PRD 77 (2008) 022001;
V. Dergachev, CQG 27 (2010) 205017]
 - Frequency Hough transform [P. Astone *et al.*, PRD 90 (2014) 042002]
- Multi-segment *F*-Statistic ($T_{\text{COH}} \sim 1\text{-few days}$)
 - Coincident *F*-Statistic
[P. Astone *et al.*, PRD 82 (2010) 022005; J. Aasi *et al.*, CQG 31 (2014) 165014]
 - Stacked *F*-Statistic (Einstein@Home)
[R. Prix, PRD 75 (2007) 023004; H.J. Pletsch & B. Allen, PRL 103 (2009) 181102;
K. Wette & R. Prix, PRD 88 (2013) 123005; D. Keitel, PRD 93 (2016) 084024;
M. Shaltev *et al.*, PRD 89 (2014) 124030; M.A. Papa *et al.*, PRD 94 (2016) 122006]

Recent results – All-sky search

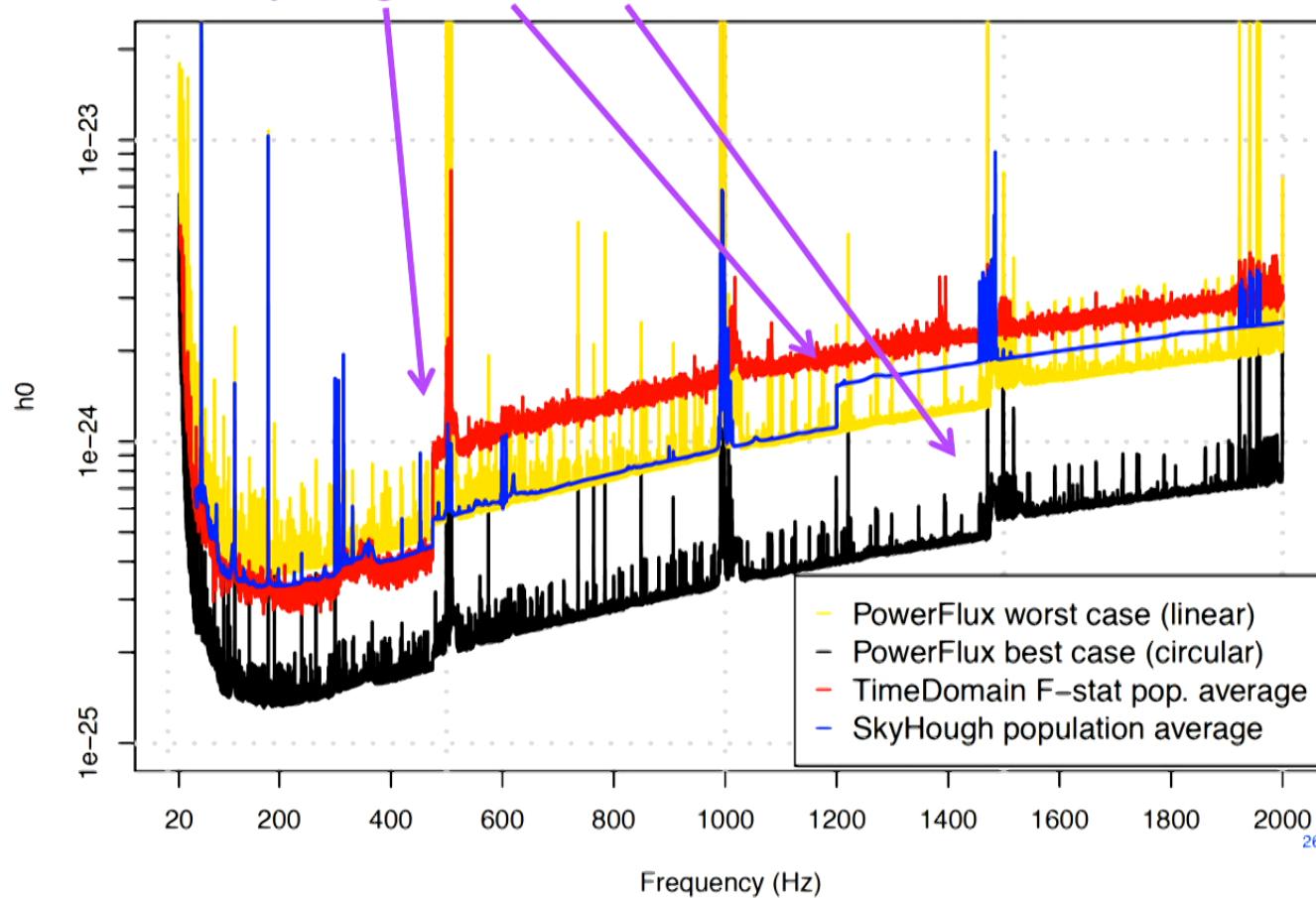
B. Abbott *et al.*,
arXiv:1802.05241, Feb 2018



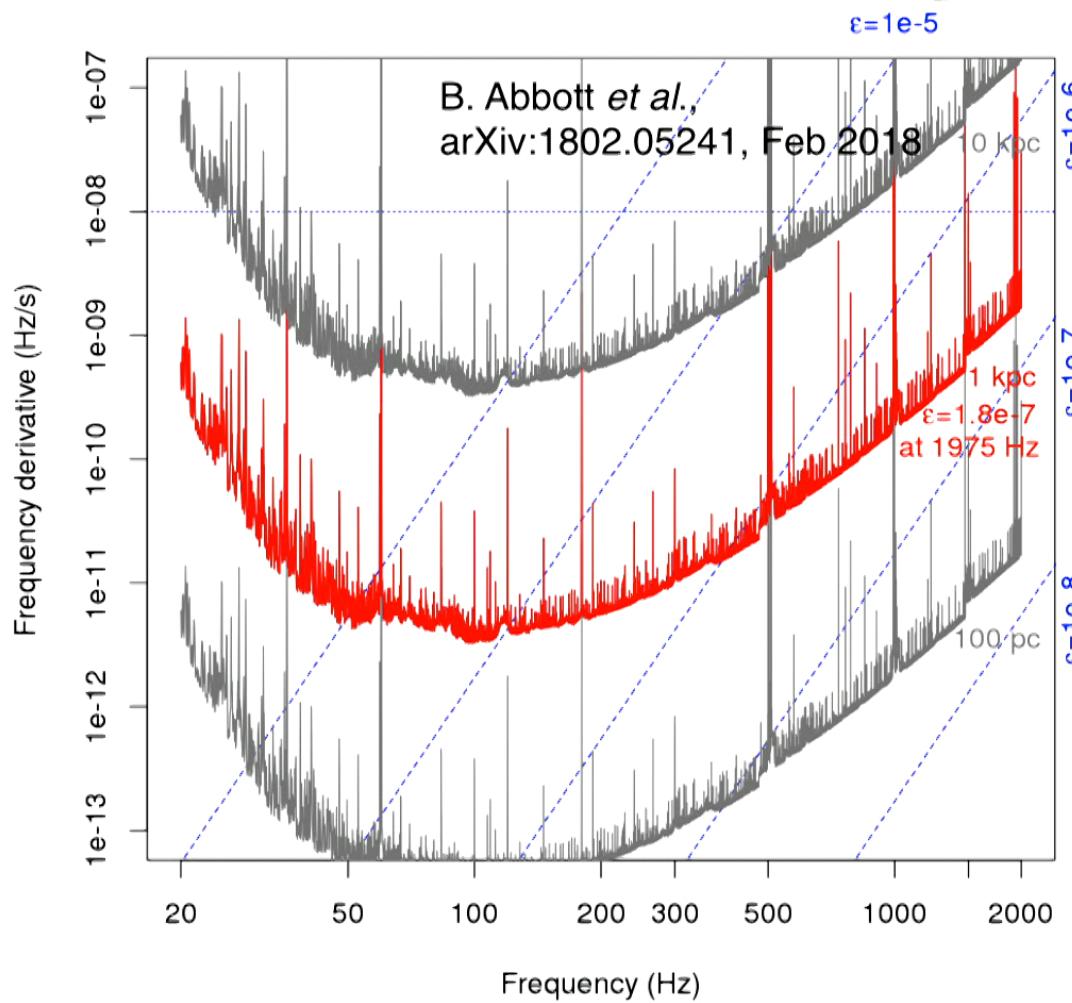
Recent results – All-sky search

Sensitivity steps come from
computing tradeoffs

B. Abbott *et al.*,
arXiv:1802.05241, Feb 2018



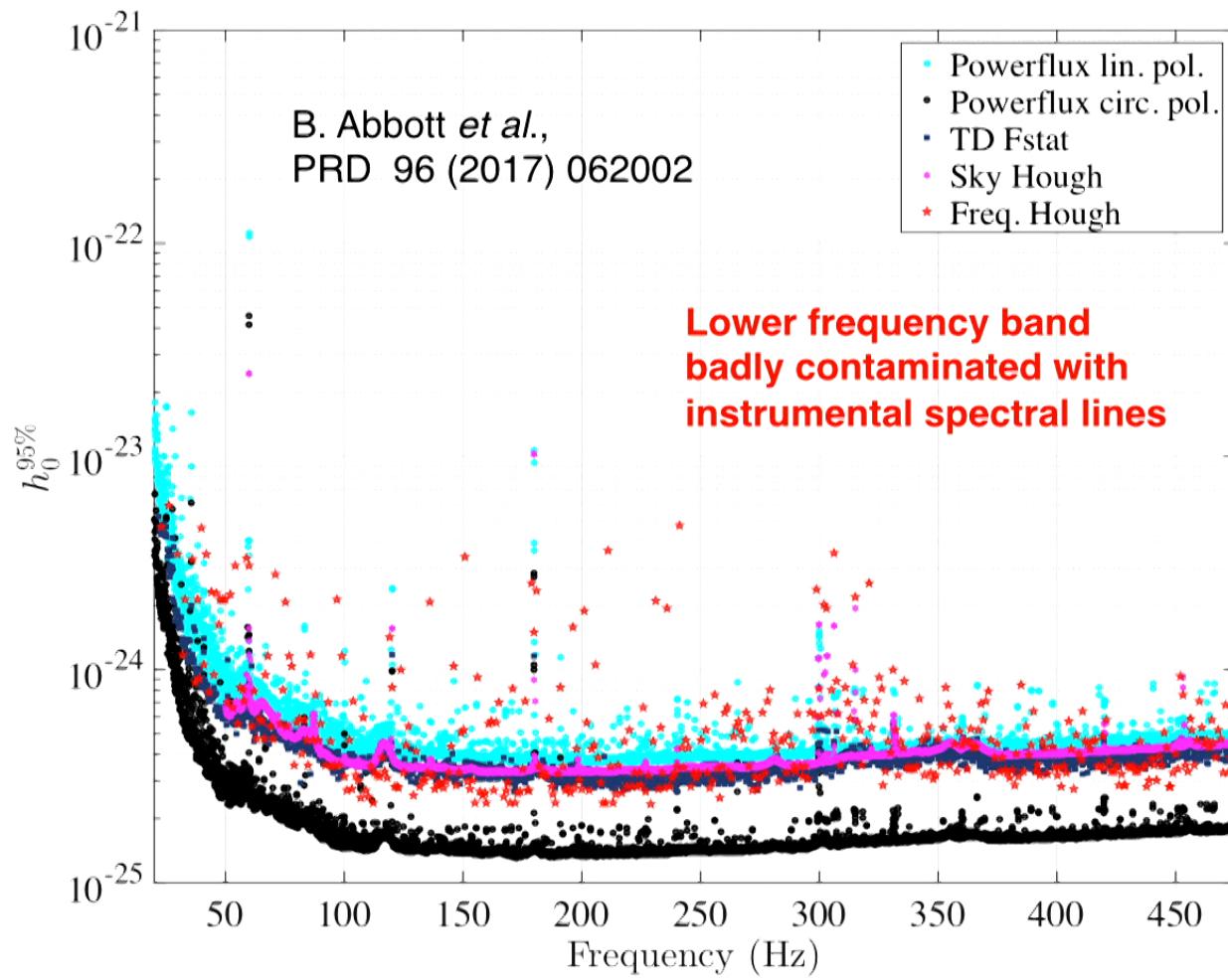
Recent results – All-sky search



Astrophysical reach for most favorable inclination
→ Circular Polarization

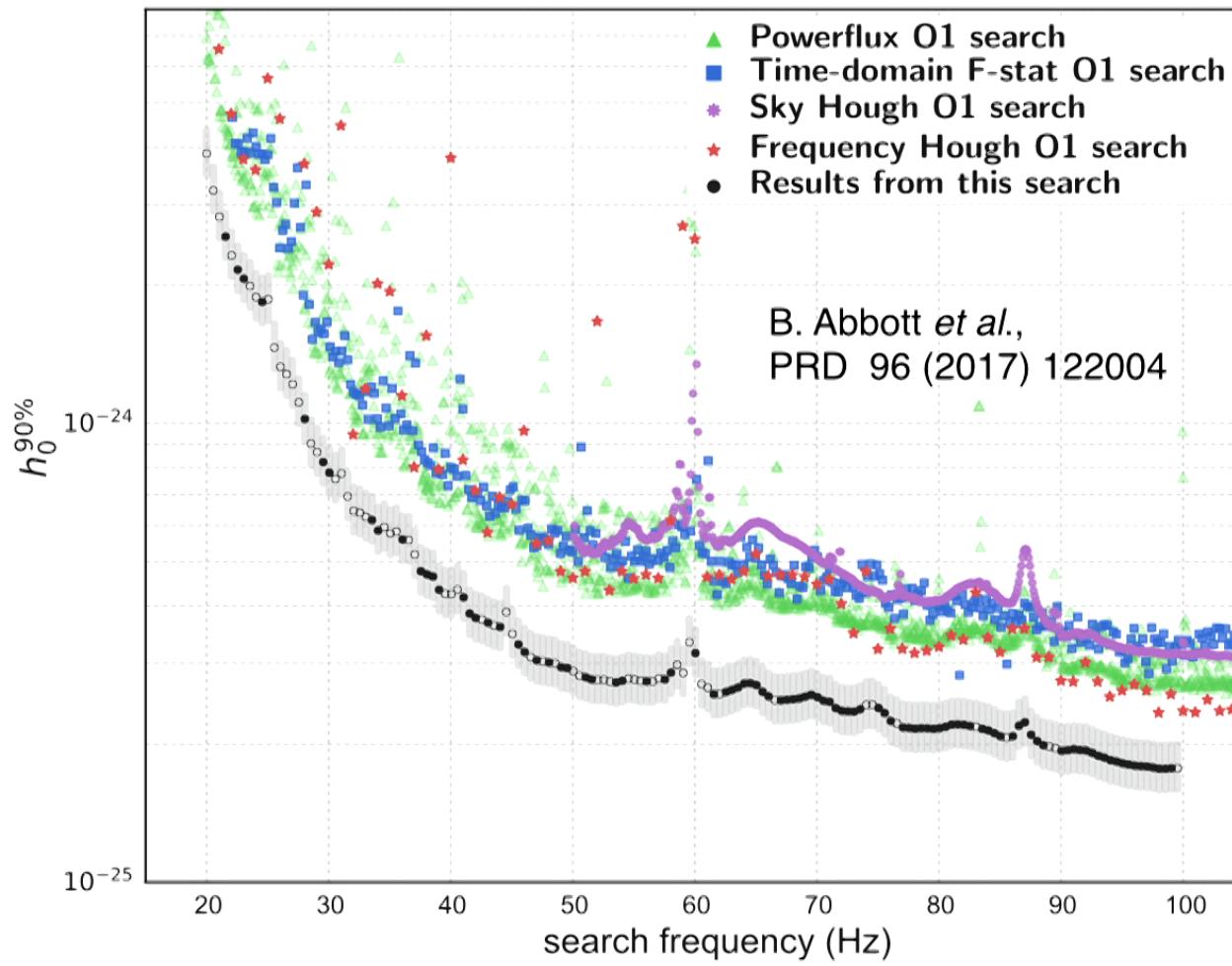
27

Recent results – All-sky search



28

Recent results – All-sky search



29

All-sky Searches for Unknown Binary Stars

Search method used in Initial LIGO data:

Double Fourier spectra (“TwoSpect”) – similar to semicoherent PowerFlux
[E. Goetz & K. Riles, CQG 28 (2011) 215006;
J. Aasi *et al.*, PRD 90 (2014) 062010]

→ Sensitivity tradeoff to cover enormous parameter space is severe

All-sky Searches for Unknown Binary Stars

Search method used in Initial LIGO data:

Double Fourier spectra (“TwoSpect”) – similar to semicoherent PowerFlux
[E. Goetz & K. Riles, CQG 28 (2011) 215006;
J. Aasi *et al.*, PRD 90 (2014) 062010]

→ Sensitivity tradeoff to cover enormous parameter space is severe

Other algorithms on near horizon or under development

- Radiometer method with sidereal-folded data
[E. Thrane *et al.*, PRD 91 (2017) 124012]

All-sky Searches for Unknown Binary Stars

Search method used in Initial LIGO data:

Double Fourier spectra (“TwoSpect”) – similar to semicoherent PowerFlux
[E. Goetz & K. Riles, CQG 28 (2011) 215006;
J. Aasi *et al.*, PRD 90 (2014) 062010]

→ Sensitivity tradeoff to cover enormous parameter space is severe

Other algorithms on near horizon or under development

- Radiometer method with sidereal-folded data
[E. Thrane *et al.*, PRD 91 (2017) 124012]
- “Polynomial” method using short- T_{COH} filters and coincidence
[S. van der Putten *et al.*, JPCS 228 (2010) 012005]

All-sky Searches for Unknown Binary Stars

Search method used in Initial LIGO data:

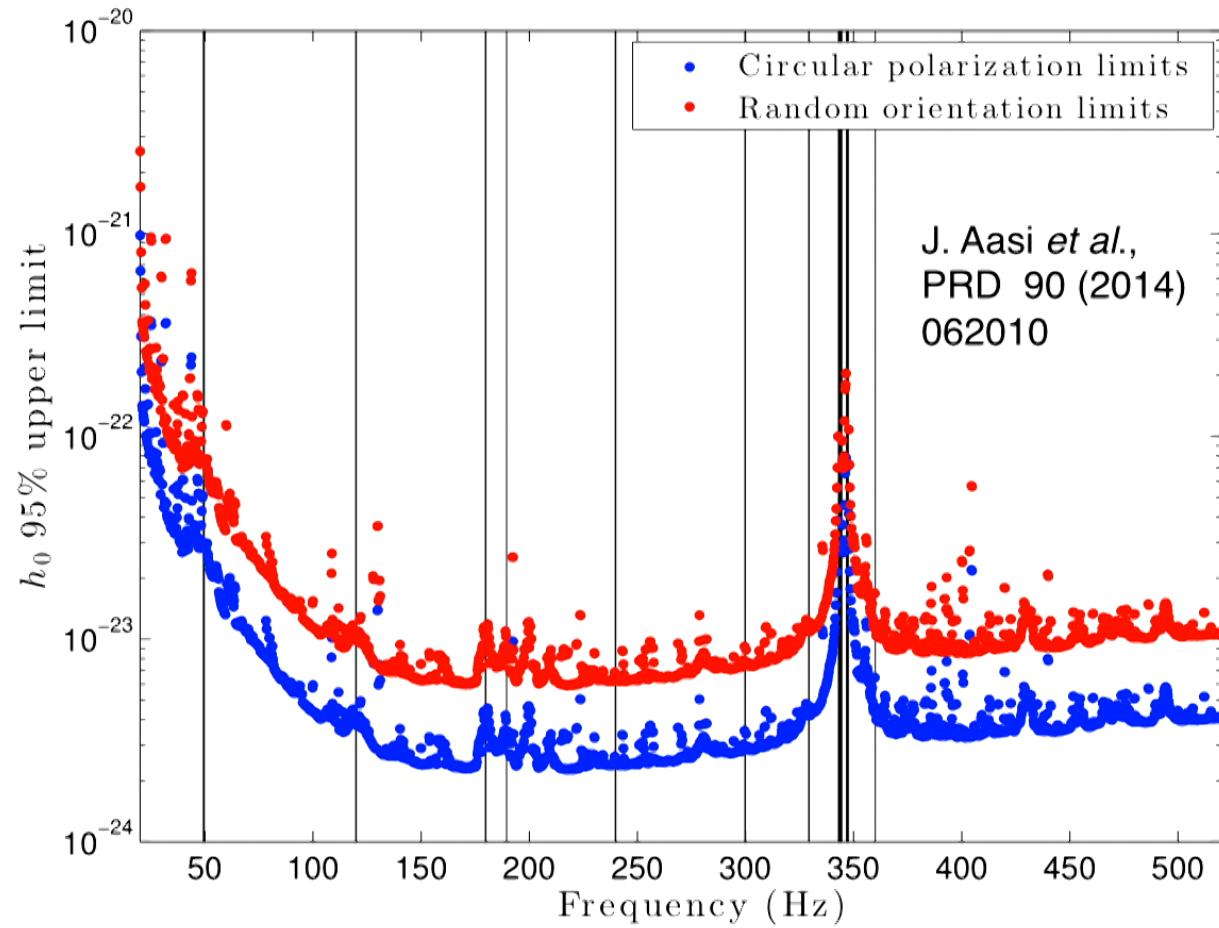
Double Fourier spectra (“TwoSpect”) – similar to semicoherent PowerFlux
[E. Goetz & K. Riles, CQG 28 (2011) 215006;
J. Aasi *et al.*, PRD 90 (2014) 062010]

→ Sensitivity tradeoff to cover enormous parameter space is severe

Other algorithms on near horizon or under development

- Radiometer method with sidereal-folded data
[E. Thrane *et al.*, PRD 91 (2017) 124012]
- “Polynomial” method using short- T_{COH} filters and coincidence
[S. van der Putten *et al.*, JPCS 228 (2010) 012005]
- Autocorrelation of spectrograms
[A. Vicere & M. Yvert, CQG 33 (2016) 165006]

Initial LIGO results – All-sky binary search



31

What changes when searching for superradiance?

Different models have (somewhat) different consequences:

- String axion

[Arvanitaki et al., PRD 81 (2010) 123530;
Arvanitaki & Dubrovsky, 83 (2011) 044026;
Kodama & Yoshino, PTEP 2014 (2014) 043E02]

- QCD axion

[Arvanitaki, Baryakhtar & Huang, PRD 91 (2015) 084011 = ABH]

What changes when searching for superradiance?

Different models have (somewhat) different consequences:

- String axion

[Arvanitaki et al., PRD 81 (2010) 123530;
Arvanitaki & Dubrovsky, 83 (2011) 044026;
Kodama & Yoshino, PTEP 2014 (2014) 043E02]

- QCD axion

[Arvanitaki, Baryakhtar & Huang, PRD 91 (2015) 084011 = ABH]

See also R. Brito et al., PRL 119 (2017) 131101;

R. Brito et al., PRD 96, 064050 (2017)

→ **Numerical calculations and consideration of extragalactic sources predict more optimistic GW strengths / rates than earlier analytic approximations**

What changes when searching for superradiance?

Different models have (somewhat) different consequences:

- String axion

[Arvanitaki et al., PRD 81 (2010) 123530;
Arvanitaki & Dubrovsky, 83 (2011) 044026;
Kodama & Yoshino, PTEP 2014 (2014) 043E02]

- QCD axion

[Arvanitaki, Baryakhtar & Huang, PRD 91 (2015) 084011 = ABH]

See also R. Brito et al., PRL 119 (2017) 131101;

R. Brito et al., PRD 96, 064050 (2017)

→ *Numerical calculations and consideration of extragalactic sources predict more optimistic GW strengths / rates than earlier analytic approximations*

Will focus here on QCD scenario (weak self-interactions)

K. Riles - Overview of CW Searches

32

What changes when searching for superradiance?

Level transition emission

[$1 \times 10^{-11} - 6 \times 10^{-11}$ eV/c² in nominal LIGO/Virgo CW 10-2000 Hz band for ABH-favored $6g \rightarrow 5g$ transition and $M_{BH} = 10M_{\odot}$]

→ Signal frequency determined by axion mass, but also by self-interaction and varying (black hole mass)²
(still must deal with Doppler modulations)

- Spinup parameter space small but non-zero

$$\frac{df}{dt} \approx 10^{-11} \frac{\text{Hz}}{\text{s}} \left(\frac{f}{90 \text{ Hz}} \right) \left(\frac{M}{10M_{\text{Sun}}} \right) \left(\frac{10^{17} \text{ GeV}}{f_a} \right) \left(\frac{5 \text{ yr}}{T_{\text{Signal}}} \right)$$

- Rare events (ABH) – once per stellar lifetime for ~ 10 years

What changes when searching for superradiance?

Annihilation emission – Splendid long-lived source!

[$2 \times 10^{-14} - 4 \times 10^{-12}$ eV/c² in nominal LIGO/Virgo CW 10-2000 Hz band]

→ Signal frequency determined by fixed axion mass,
but with small quadratic corrections in $(M/10M_\odot)^2$
(still must deal with Doppler modulations)

- Spinup parameter space nearly vanishes
→ Especially helpful for long observation times
- Stochastic GW searches could see superposition of multiple, incoherent sources in narrow band
(detection of closest source likely first?)

Some Questions

- Could we get significantly better sensitivity by narrowing search parameter space (smaller magnitude of frequency derivative)?

Some Questions

- Could we get significantly better sensitivity by narrowing search parameter space (smaller magnitude of frequency derivative)?
- OTOH, could we be missing superradiance altogether in BEC clouds with strong self-interaction and unpredictable frequency evolution? (stochastic GW searches serve as backstop)

Some Questions

- Could we get significantly better sensitivity by narrowing search parameter space (**smaller magnitude of frequency derivative?**)?
- OTOH, could we be missing superradiance altogether in BEC clouds with strong self-interaction and unpredictable frequency evolution? (**stochastic GW searches serve as backstop**)
- Is there theoretical consensus on QCD rates / strengths? (**analytical and numerical agree well now?**)
- Are unknown isolated black holes or known black hole binaries better bets? (**all-sky comes “for free”; searches for specific binaries impose non-trivial computing and analysis costs**)

Some Questions

- Could we get significantly better sensitivity by narrowing search parameter space (smaller magnitude of frequency derivative)?
- OTOH, could we be missing superradiance altogether in BEC clouds with strong self-interaction and unpredictable frequency evolution? (stochastic GW searches serve as backstop)
- Is there theoretical consensus on QCD rates / strengths? (analytical and numerical agree well now?)
- Are unknown isolated black holes or known black hole binaries better bets? (all-sky comes “for free”; searches for specific binaries impose non-trivial computing and analysis costs)
- Which binaries are most promising? (Ideal: close & medium-spin)

Some Questions

- Just how narrow is the signal band for annihilation?

$$f_{\text{GW}} = \frac{2\mu_a c^2}{h} \left(1 - \frac{\alpha^2}{2n^2}\right) \Rightarrow \frac{\Delta f_{\text{GW}}}{f_{\text{GW}}} = \frac{\Delta(\alpha^2)}{2n^2}$$

$$\alpha = \frac{GM_{\text{BH}}\mu_a}{\hbar c} = (0.075) \left(\frac{M_{\text{BH}}}{10M_{\text{Sun}}} \right) \left(\frac{\mu_a}{10^{-12} \text{ eV/c}^2} \right)$$

→ If band extremely narrow, must be careful about source confusion in following up all-sky outliers (we wish!)

- Should CW searches go above 290 Hz? ($\mu_a > 6 \times 10^{-13} \text{ eV/c}^2$)
[excluded for QCD axions; what about stringy axions?]
- If so, should CW searches go above 2000 Hz?
[very costly for all-sky searches]

Summary

CW searches are diverse in both targets and algorithms
→ Trying to keep our eyes wide open

Summary

CW searches are diverse in both targets and algorithms

→ Trying to keep our eyes wide open

But we do make choices on where to spend computing resources and analyst time

Have mainly assumed galactic neutron star is prime source and that BH superradiance comes for free in all-sky searches – at some sensitivity

Summary

CW searches are diverse in both targets and algorithms

→ Trying to keep our eyes wide open

But we do make choices on where to spend computing resources and analyst time

Have mainly assumed galactic neutron star is prime source and that BH superradiance comes for free in all-sky searches – at some sensitivity

If superradiance gives stable frequencies, we should be safe

But if boson self-interactions give significant spin wandering, we could miss the signal with current CW searches

(but still catch extra-galactic signal from stochastic signal?)

→ How much should we worry about this?

Should we focus on particular black hole binaries?

Guidance is welcome!

37