

Title: The Chemical Imagination at Work in VERY Tight Places

Date: Mar 28, 2018 02:00 PM

URL: <http://pirsa.org/18030070>

Abstract: <p>Diamond anvil cells and shock techniques now permit the study of matter under multimegabar (i.e. several hundred GPa) pressures. The properties of matter in this pressure regime differ drastically from those known at 1 atm. Just how different physics and chemistry are at high pressure and the role that a chemical intuition for bonding and structure can have in understanding matter at high pressures will be explored in this lecture. I will discuss in detail an overlapping hierarchy of responses to increased density, consisting of (a) squeezing out van der Waals space (for molecular crystals); (b) increasing coordination; (c) decreasing the bond length of covalent bonds and the size of anions; and (d) an extreme regime of electrons moving off atoms and new modes of correlation. Examples of the startling chemistry and physics that emerge under such extreme conditions will alternate in this account with qualitative chemical ideas about the bonding involved.</p>



# **The Chemical Imagination at Work in Very Tight Places**

*Roald Hoffmann*

*working together with*  
***Neil W. Ashcroft***

**NSF, DOE (EFree)**

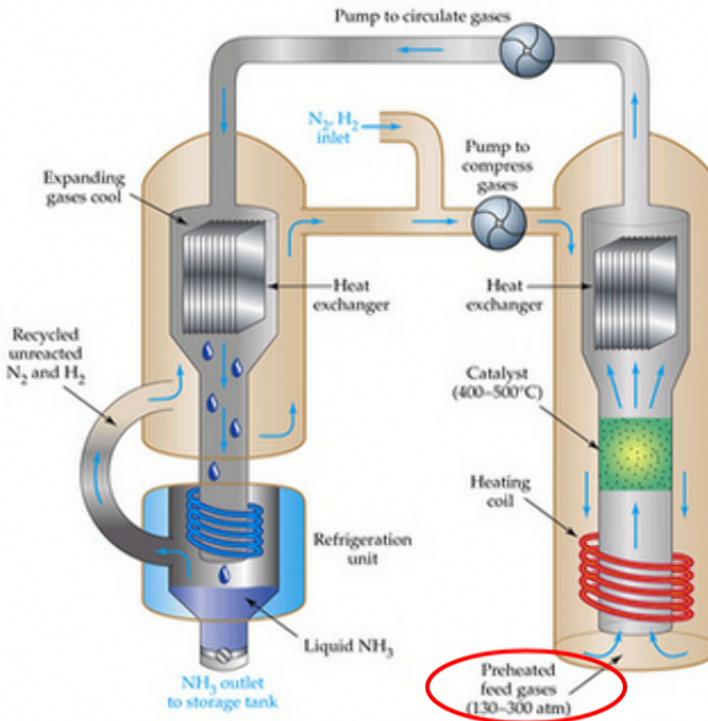
10,000 bar = 9,869 atm = 145,038 lbs/sq.in. = 1 GPa

Pressure at center of earth ~3.5 million atmospheres ~350 Gigapascals

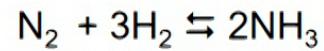
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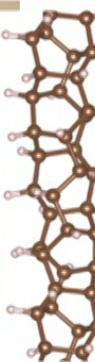
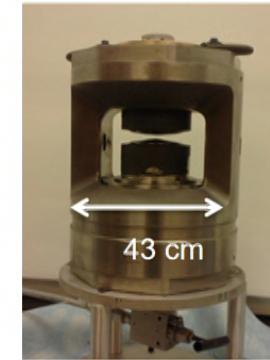
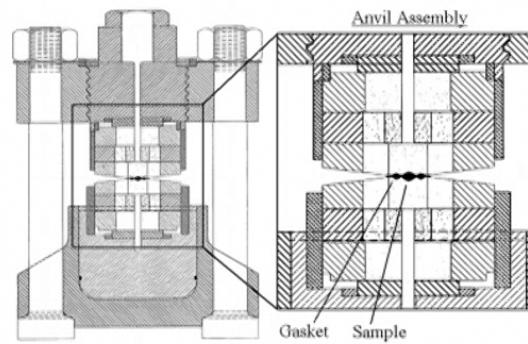
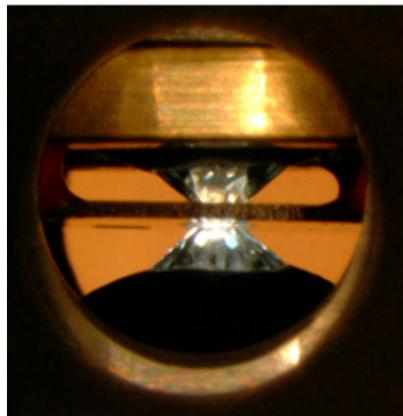
Pressure at center of earth ~3.5 million atmospheres ~350 Gigapascals

On going from 1 atm to 350 Gpa, the **volume** of solid Li or methane or benzene decreases by a factor of ~5, NaCl or H<sub>2</sub>O ~3, Fe, Au ~1.7

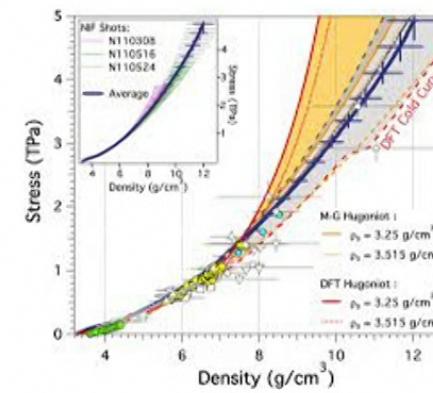
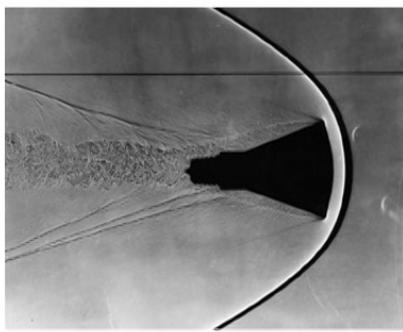


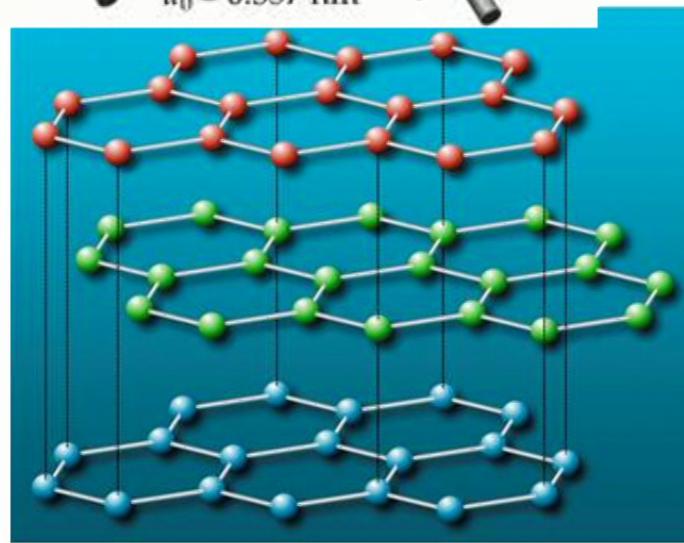
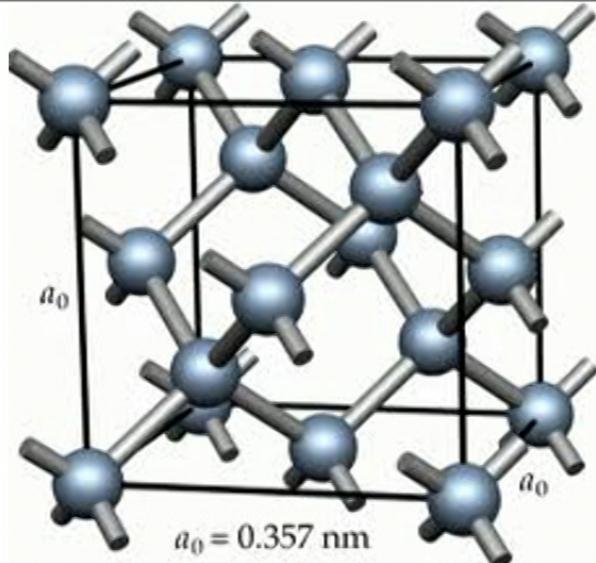
### Haber-Bosch process



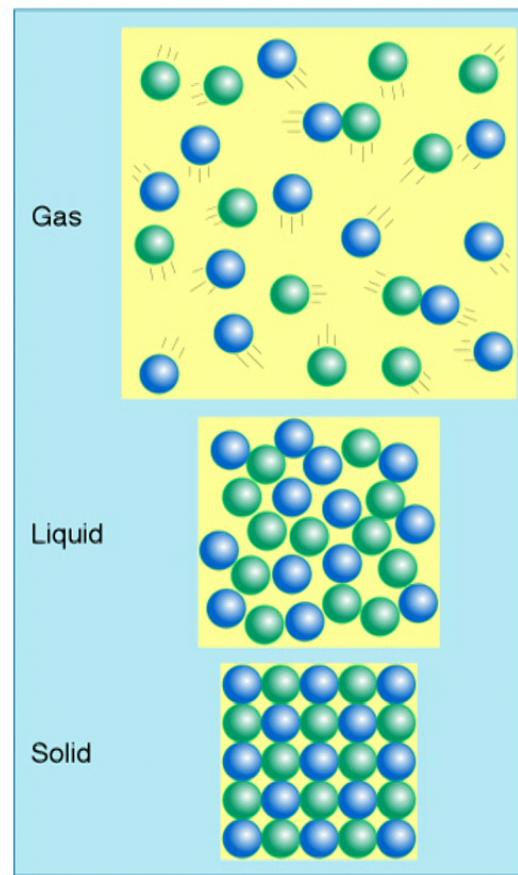


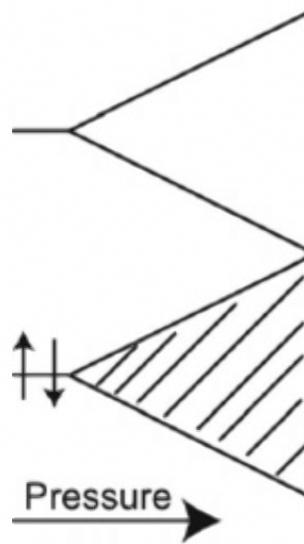
Paris-Edinburgh cell at ORNL





densities of graphite 2.27, diamond 3.51 g/cm<sup>3</sup>

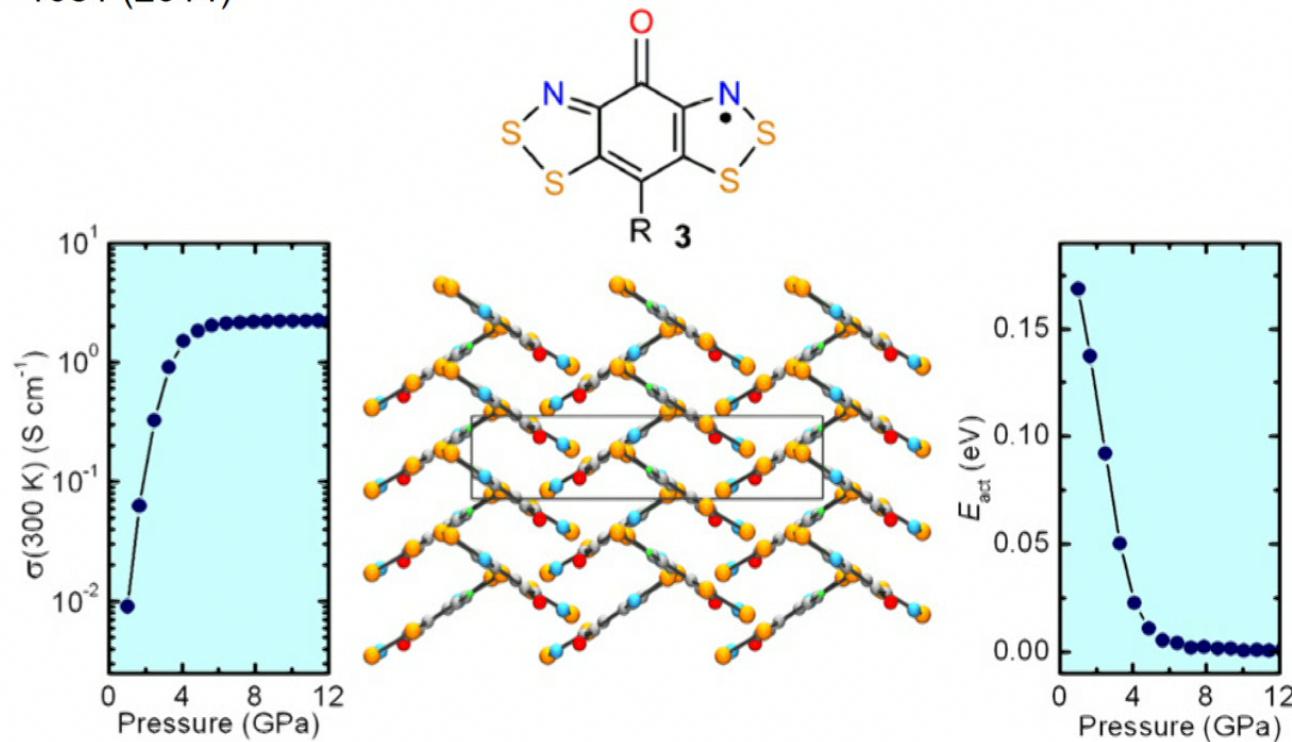




**Virtually all materials become metallic under sufficiently high pressure.**

(but... some – Li, Na -- that are metallic at low P become insulating before eventually returning to metallicity)

J. W. L. Wong, A. Mailman, K. Lekin, S. M. Winter, W. Yong, J. Zhao, S. V. Garimella, J. S. Tse, R. A. Secco, S. Desgreniers, Y. Ohishi, F. Borondics and R. T. Oakley. "Pressure Induced Phase Transitions and Metallization of a Neutral Radical Conductor." *J. Am. Chem. Soc.* **136**, 1070-1081 (2014)

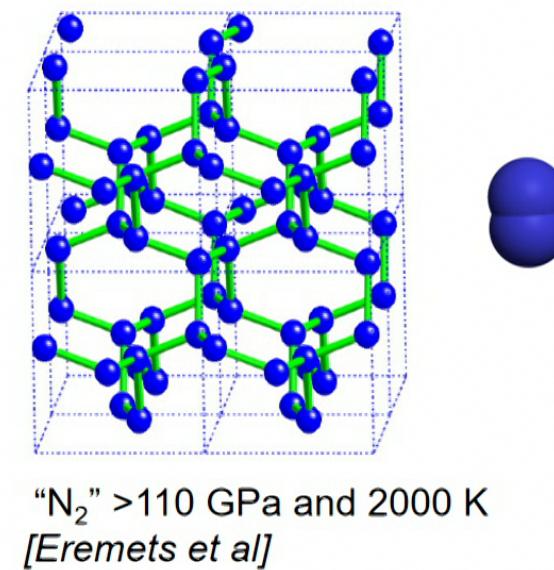
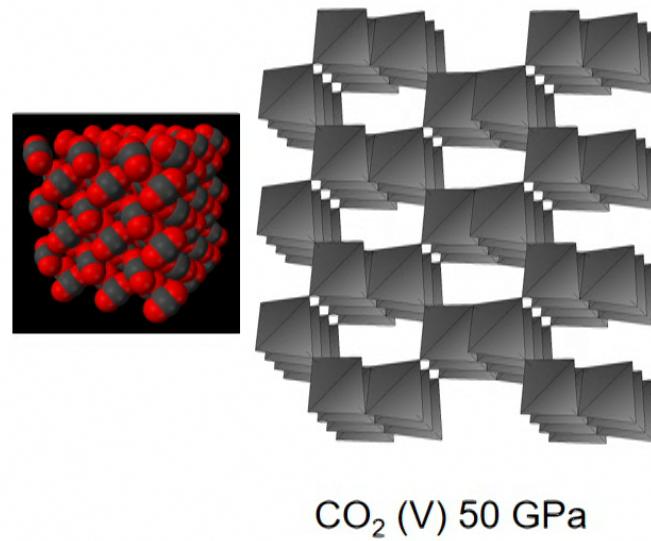


## The extreme landscape of high pressure

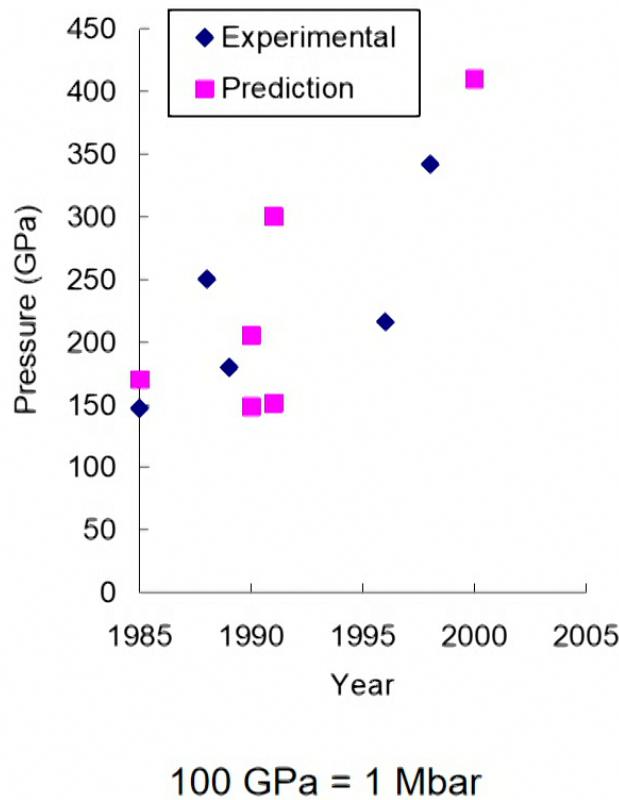
Xe and I<sub>2</sub> can be metallized [*Goettel et al, Ruoff et al*], as can CsI and BaTe [*Eremets et al, Grzybowski et al*]

## The extreme landscape of high pressure

Xe and I<sub>2</sub> can be metallized [Goettel et al, Ruoff et al], as can CsI and BaTe [Eremets et al, Grzybowski et al]



## Metallic hydrogen



[M. I. Eremets & I. A. Troyan](#)

MPI Chemie, Mainz

Nature Materials, 10, 927–931 (2011)

[W.J. Nellis, Arthur L. Ruoff, Isaac F. Silvera](#)

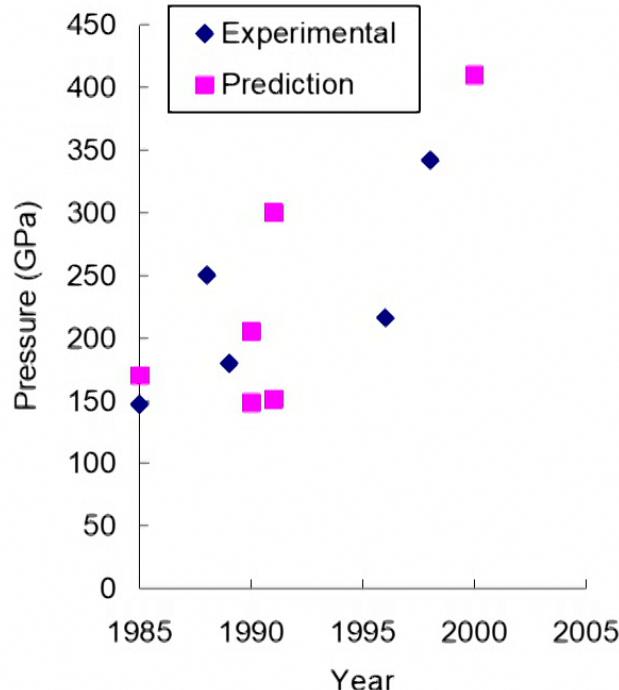
Has Metallic Hydrogen Been Made in a Diamond Anvil Cell?  
arXiv:1201.0407, 2012

Dias, R.; Silvera, I. F. (2016).  
*"Observation of the Wigner-Huntington Transition to Solid Metallic Hydrogen".*  
Science 255, 715 (2017)

[W. J. Nellis](#)

"Dynamic Compression of Materials:  
Metallization of Fluid Hydrogen at High  
Pressures", Rep. Prog. Phys. 69, 1479-  
1580 (2006).

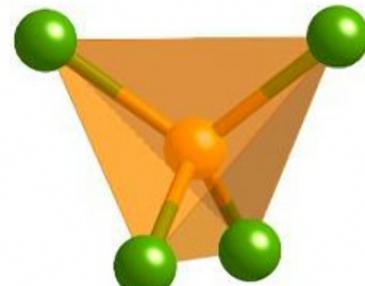
## Metallic hydrogen



100 GPa = 1 Mbar

## Chemical precompression

Neil Ashcroft



$\text{CH}_4$ ,  $\text{SiH}_4$ ,  $\text{GeH}_4$ , ...  
 $\text{LiBH}_4$ ,  $\text{NaAlH}_4$ , ...  
 $\text{K}_2\text{ReH}_9$ ,  $(\text{NH}_4)_2\text{ReH}_9$

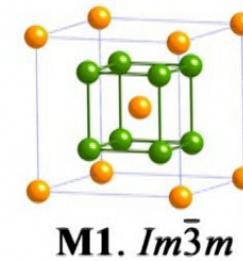
$\text{SiH}_4$  m.p. -185 C, b.p. -112 C

Autoignites 21 C

## Structures and Potential Superconductivity in SiH<sub>4</sub> at High Pressure - *En Route to “Metallic Hydrogen”*

Ji Feng, Wojciech Gochala, Tomasz Jaroń, Roald Hoffmann, Aitor Bergara and N.  
W. Ashcroft

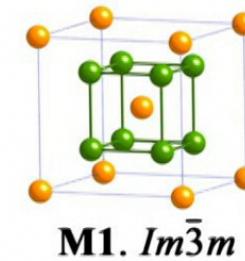
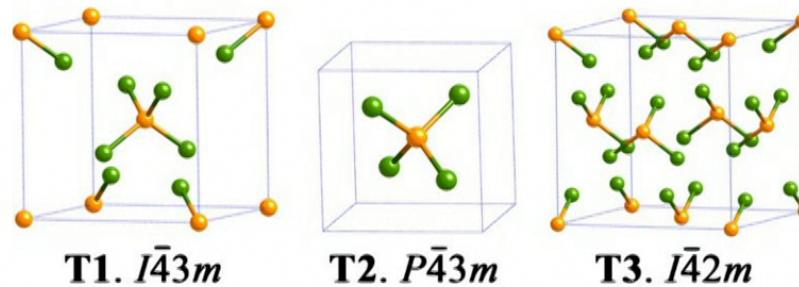
Physical Review Letters, **96**, 017006, (2006).



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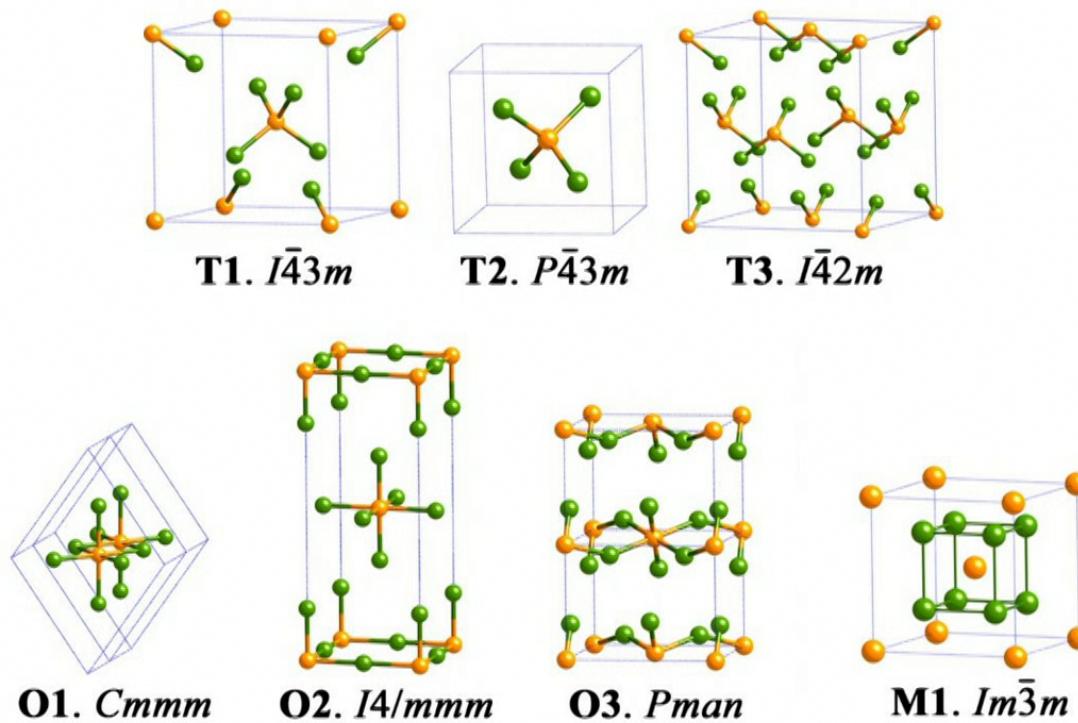
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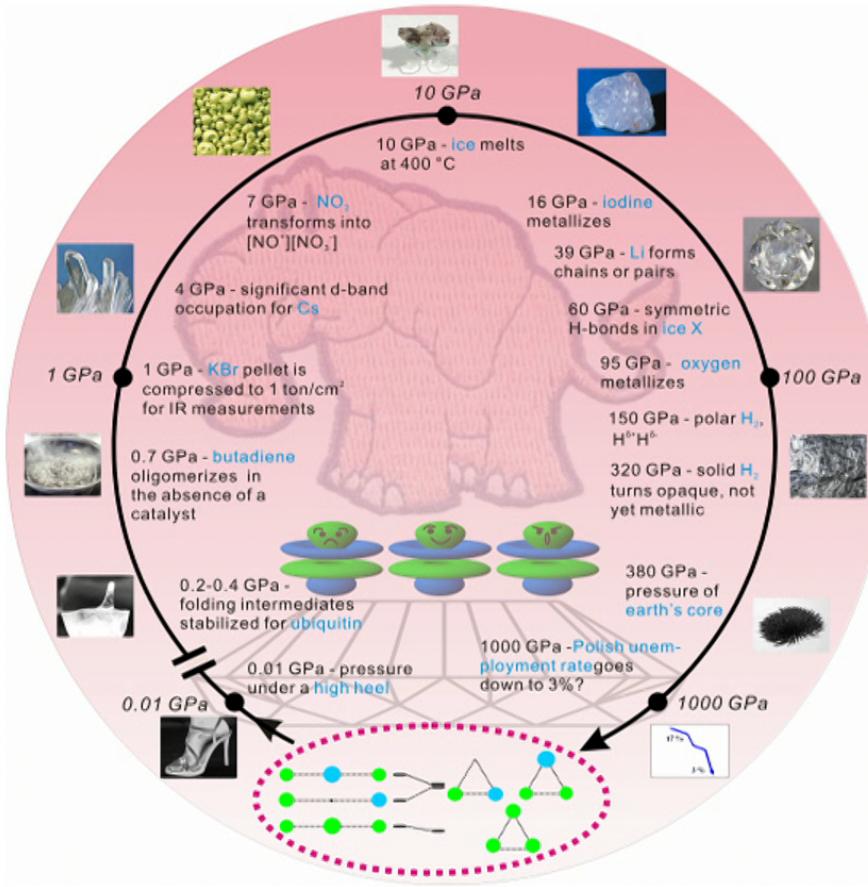
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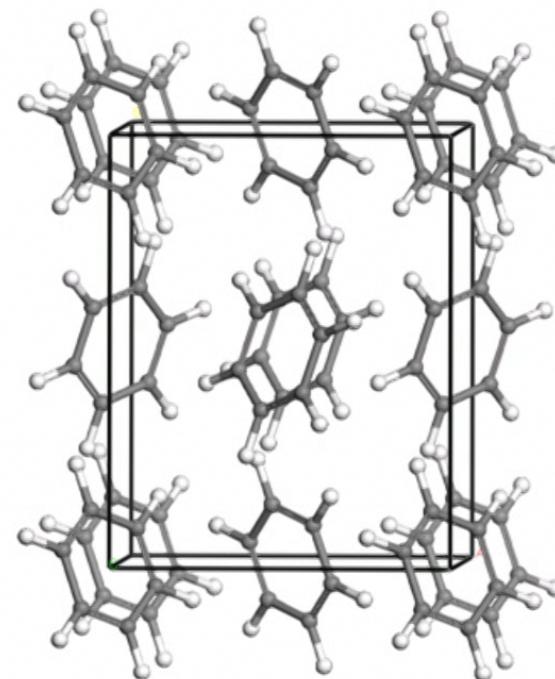
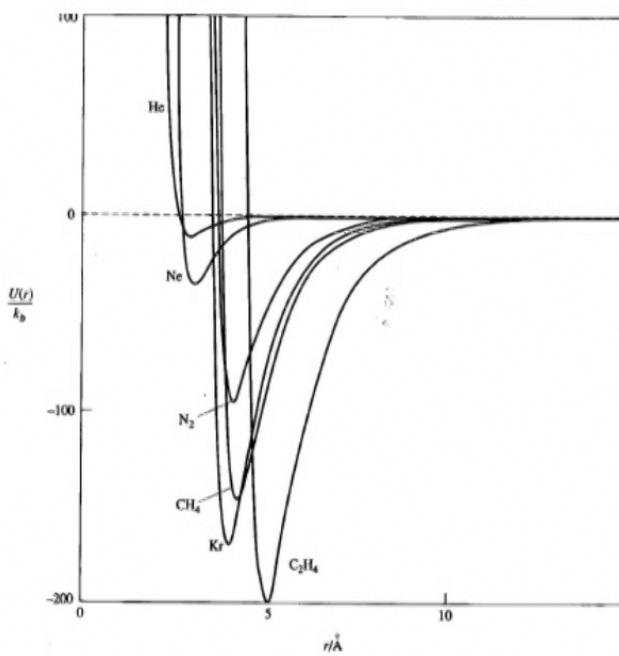
Wojciech Gochala, Roald Hoffmann, Ji Feng, and Neil W. Ashcroft

Angewandte Chemie, **46**(20), 3620-3642 (2007).

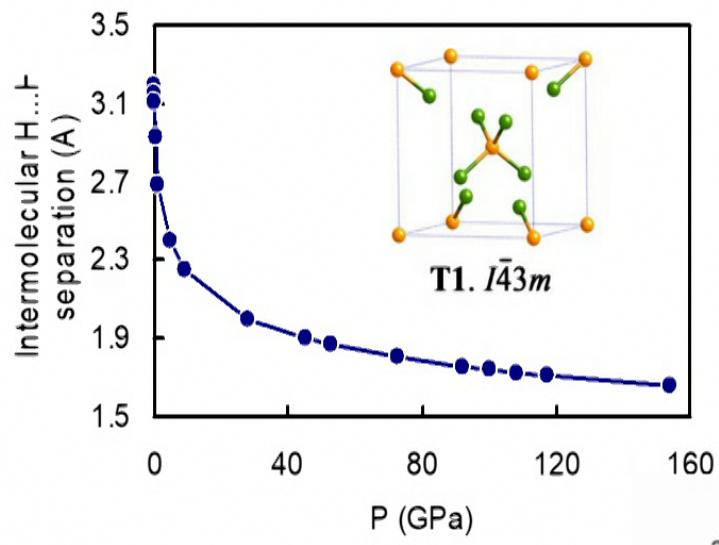


## A Hierarchy of Responses to Pressure in Crystals

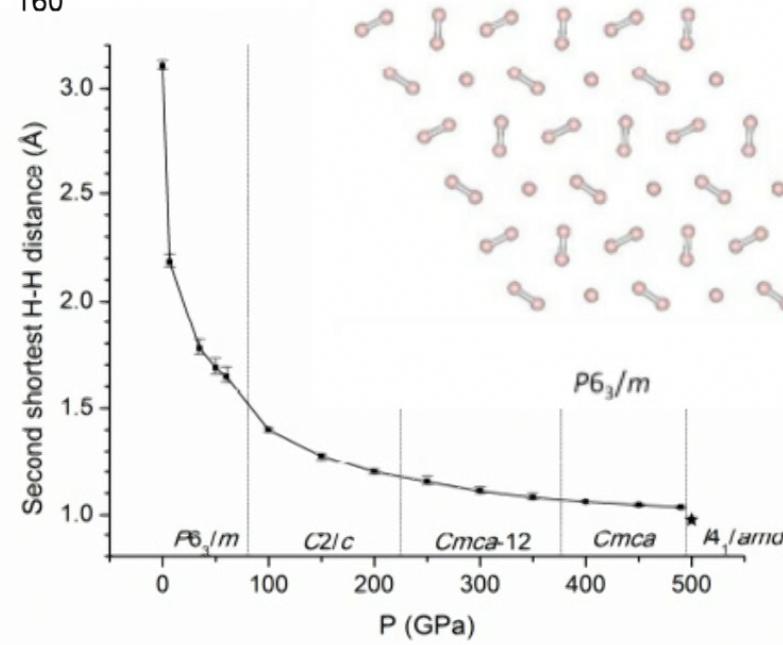
- Penetrating the repulsive region of intermolecular potentials
- Increasing coordination at main group and transition metal elements
- Decreasing the bond length of covalent bonds, and the size of anions
- A new world of electrons moving off their atoms, new modes of correlation



dispersion (= van der Waals) forces -- weak forces between molecules and atoms, attractive at long distance ( $1/r^6$ , due to induced-dipole, induced-dipole interactions), repulsive at short separations

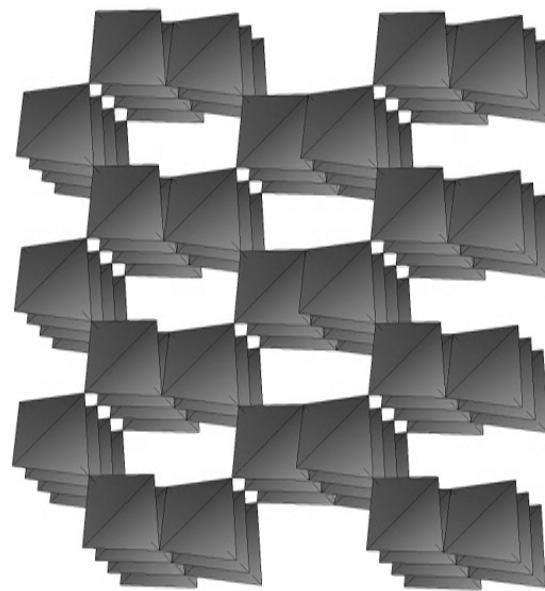
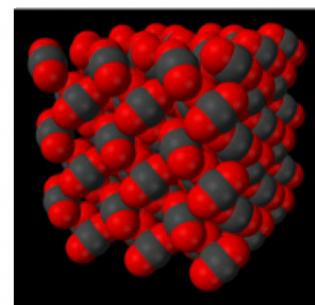


van der Waals space is most easily compressed.



## A Hierarchy of responses to Pressure in Crystals

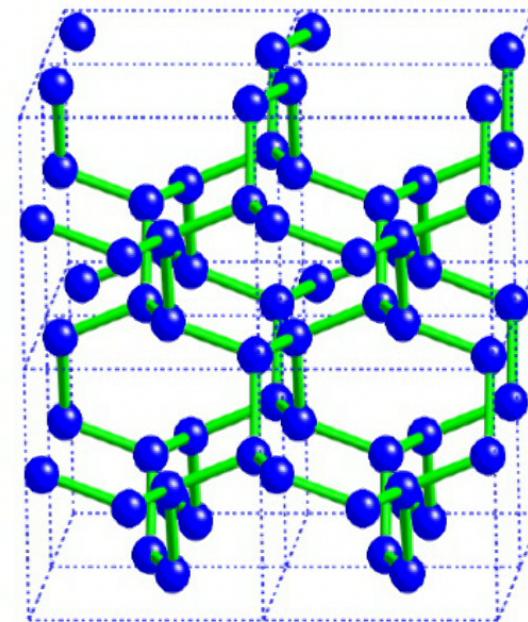
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CO<sub>2</sub> (V) 50 GPa, 1800 K

## Coordination Alchemy

$$\Delta H = \Delta E + P\Delta V$$



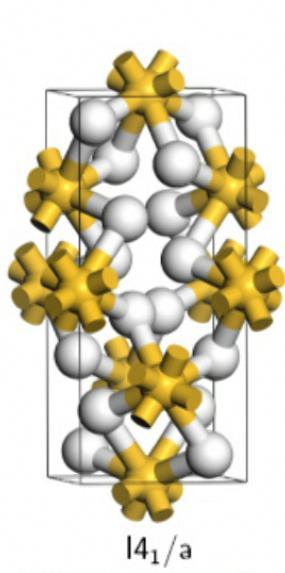
"N<sub>2</sub>" >110 GPa, 2000 K



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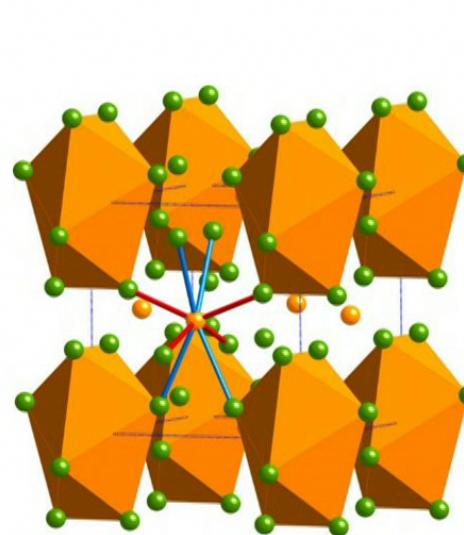




I4<sub>1</sub>/a

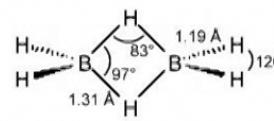
C.J. Pickard, R.J. Needs

each Si atom has 8 nearest neighbor hydrogens, two each bridging to a neighboring Si. The stoichiometry is SiH<sub>8/2</sub>. Each Si-H<sub>2</sub>-Si unit uses 4 electrons (one each from each Si and each H) to bond together in diborane fashion the four centers.

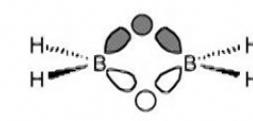


Feng et al

each Si has around it 8 electron deficient bent Si-H-Si three-center bonds



(a)

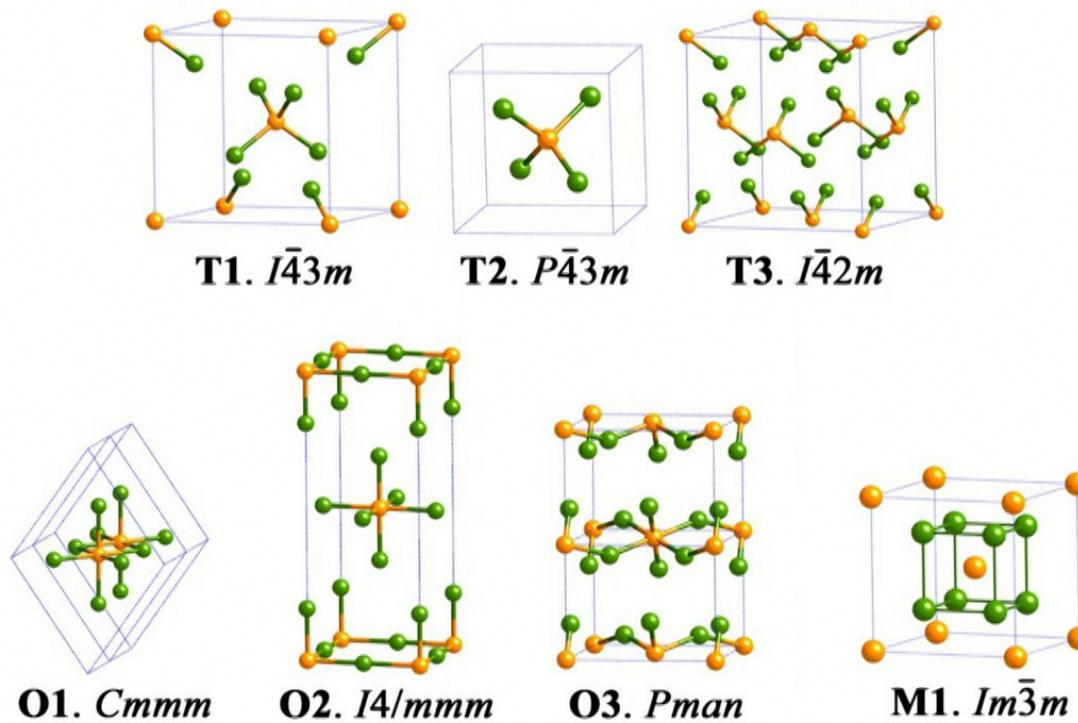


(b)

## Structures and Potential Superconductivity in SiH<sub>4</sub> at High Pressure - En Route to “Metallic Hydrogen”

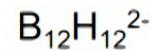
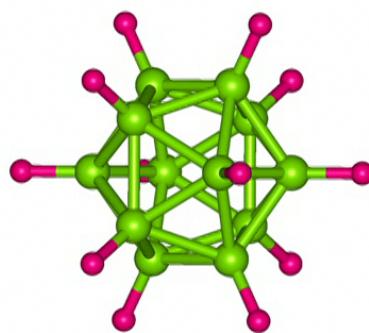
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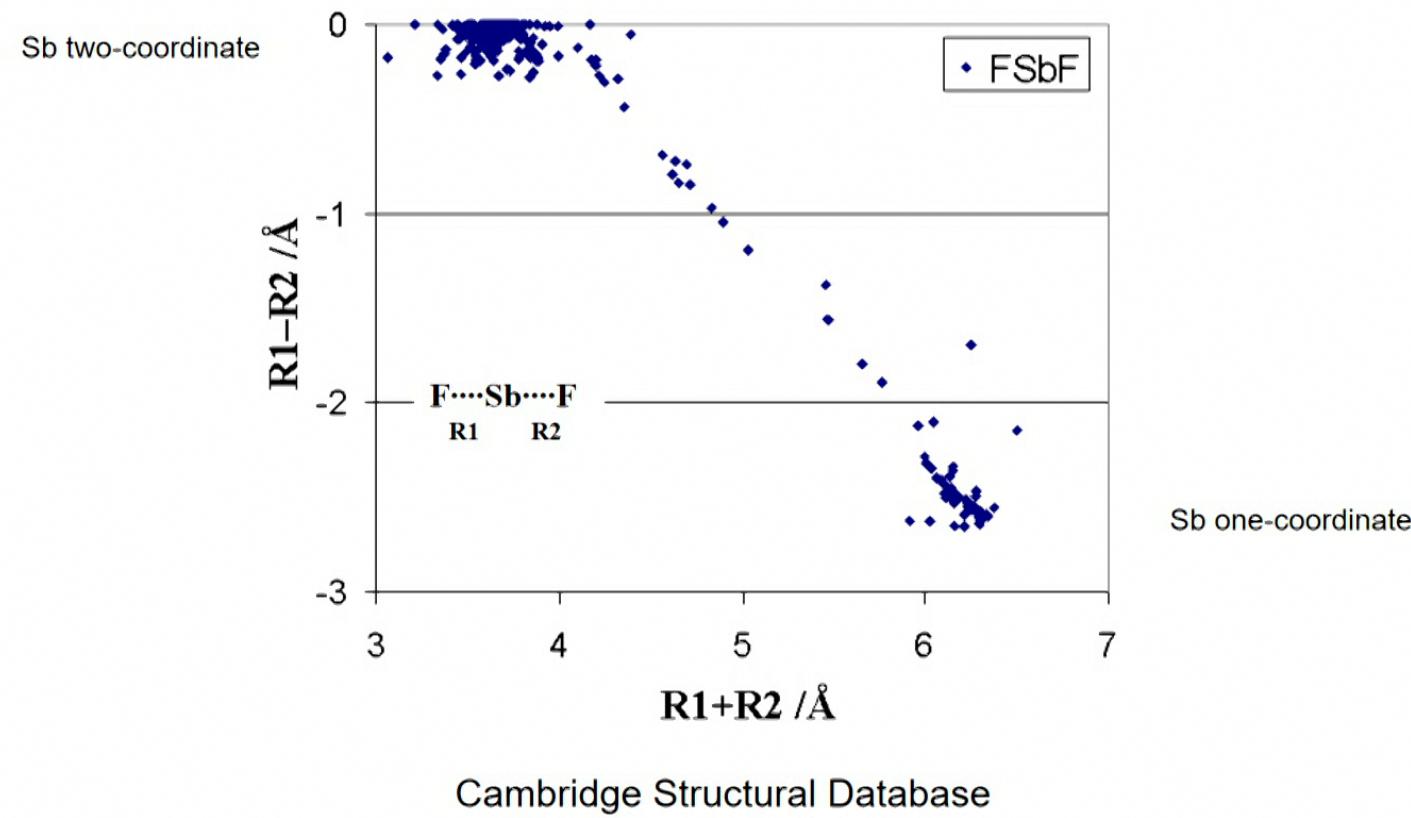


## Increased density with increasing pressure

- higher coordination
- more bonds, full or partial, with same number of electrons
- multicenter bonding, electron-rich or electron-poor



How to use P=1 atm data to understand what will happen at higher pressures



## A Hierarchy of Responses to Pressure in Crystals

- Penetrating the repulsive region of intermolecular potentials
- Increasing coordination at main group and transition metal elements
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# THE PERIODIC TABLE

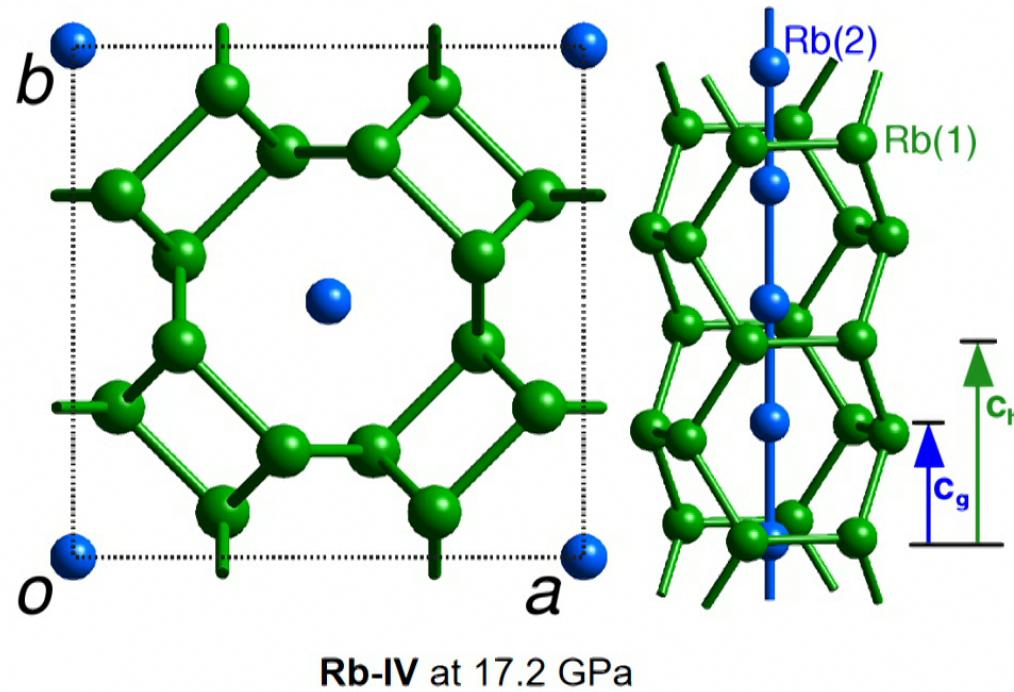
		THE PERIODIC TABLE																						
		1 IA																				18 VIIIA		
		H																				He		
		SYMBOL	ATOMIC NUMBER																			2 4.00 Helium		
		H	1																			18 4.00 Helium		
1		Lithium	3	6.94	4	9.01	Beryllium															He		
2		Li	Be																			2 4.00 Helium		
3		Sodium	Mg																			He		
4		K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As	Se	Br	Kr					
5		Rb	Sr	Y	Zr	Nb	Mo	Tc	Ru	Pd	Ag	Cd	In	Sn	Sb	Te	I	Xe						
6		Cs	Ba	La	Hf	Ta	W	Re	Os	Ir	Pt	Au	Hg	Tl	Pb	Bi	Po	At	Rn					
7		Fr	Ra	Ac	Rf	Db	Sg	Bh	Hs	Mt	Unamed	Unamed	Unamed	Unamed	Unamed	Unamed	Unamed	Unamed	Unamed	Unamed	Unamed			
		ALKALI METALS	ALKALI EARTH METALS		LANTHANIDES	Ce	Pr	Nd	Pm	Sm	Eu	Gd	Tb	Dy	Ho	Er	Tm	Yb	Lu					
		HAYDEN MC NEIL SPECIALTY PRODUCTS		ACTINIDES	Th	Pa	U	Np	Pu	Am	Cm	Bk	Cf	Es	Fm	Md	No	Lr						
		www.hmpublishing.com			58 140.12 Cerium	59 140.91 Praseodymium	60 144.24 Neodymium	61 (145) Promethium	62 150.36 Samarium	63 152.97 Europium	64 157.25 Gadolinium	65 158.93 Terbium	66 162.50 Dysprosium	67 164.93 Holmium	68 167.26 Erbium	69 168.93 Thulium	70 173.04 Ytterbium	71 174.97 Lutetium						
		© Hayden-McNeil Specialty Products			90 232.04 Thorium	91 231.04 Protactinium	92 238.03 Uranium	93 237.05 Neptunium	94 (240) Plutonium	95 243.06 Americium	96 (247) Curium	97 (248) Berkelium	98 (251) Californium	99 252.08 Einsteinium	100 257.10 Fermium	101 (257) Mendelevium	102 259.10 Nobelium	103 262.11 Lawrencium						

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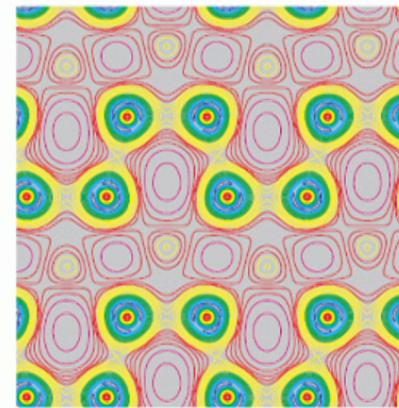
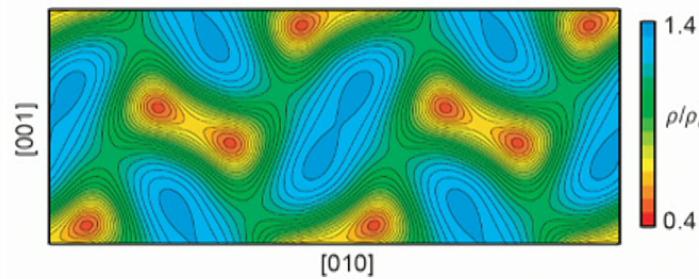
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At 1 atm **Ba** is bcc. At 5.5 GPa **Ba-II** is fcc. At 25GPa **Ba-IV** goes into a structure with 288 atoms in a commensurate unit cell

*R.J. Nelmes, Edinburgh*



## Elemental Li



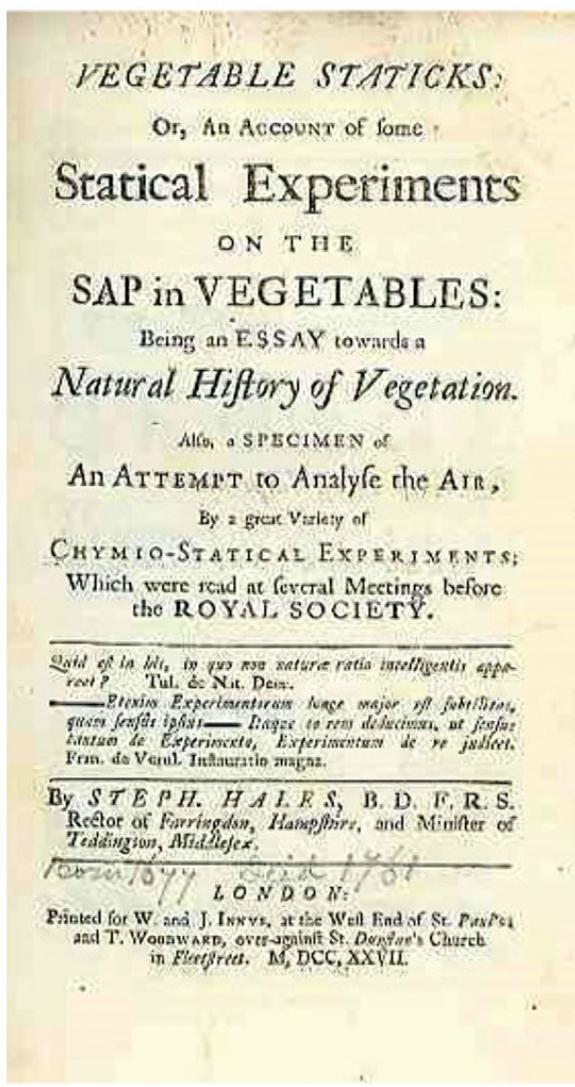
top: Ashcroft group 1999

botttom: Nelmes group 2003

## **close packing is not close enough**

The packing of identical spheres maximizes at ~74% ( $\pi/\sqrt{18}$ )

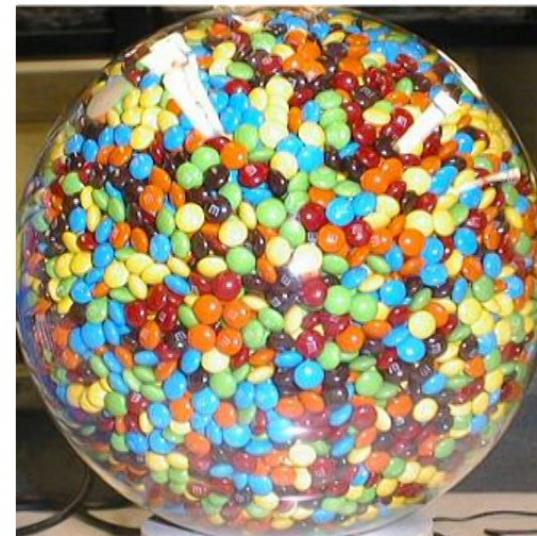
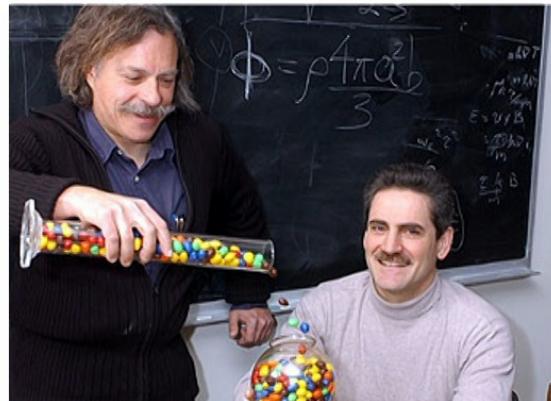




"I compressed several fresh parcels of Pease in the same Pot, with a force equal to 1600, 800, and 400 pounds; in which Experiments, tho' the Pease dilated, yet they did not raise the lever, because what they increased in bulk was, by the great incumbent weight, pressed into the interstices of the Pease, which they adequately filled up, being thereby formed into pretty regular Dodecahedrons."

## close packing is not close enough

The packing of identical spheres maximizes at ~74% ( $\pi/\sqrt{18}$ ). If spheres are deformed, they can pack more efficiently -- *vide* work on ellipsoids and M&M's, and the trivial case of extreme deformation to cubes. Other polyhedra can also pack space fully

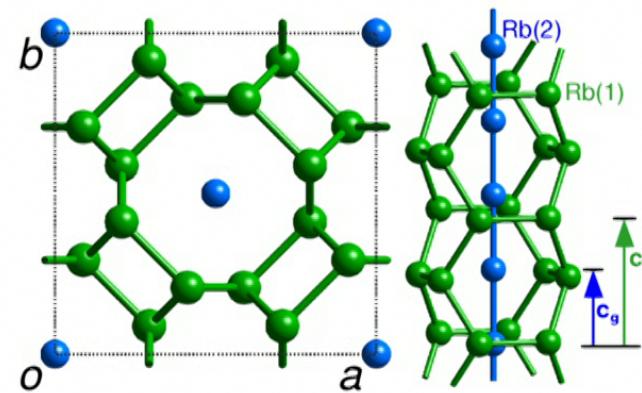


Paul Chaikin NYU

## close packing is not close enough

The packing of identical spheres maximizes at  $\sim 74\%$  ( $\pi/\sqrt{18}$ ). If spheres are deformed, they can pack more efficiently -- *vide* work on ellipsoids and M&M's, and the trivial case of extreme deformation to cubes. Other polyhedra can also pack space fully.

Spheres of different size may pack more efficiently than spheres of identical size..



## Phase Diagram and Structural Diversity of the Densest Binary Sphere Packings

Adam B. Hopkins,<sup>1</sup> Yang Jiao,<sup>2</sup> Frank H. Stillinger,<sup>1</sup> and Salvatore Torquato<sup>1,2,3,4,5</sup>

<sup>1</sup>Department of Chemistry, Princeton University, Princeton, New Jersey 08544, USA

<sup>2</sup>Princeton Institute for the Science and Technology of Materials, Princeton University, Princeton, New Jersey 08544, USA

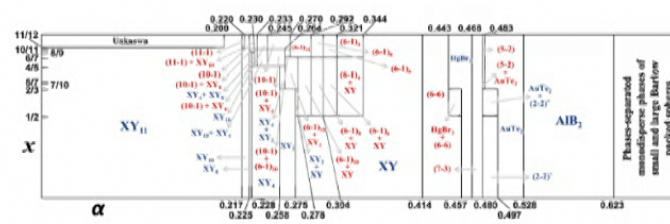
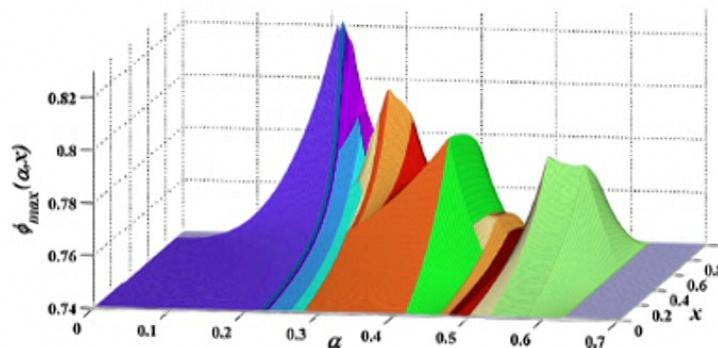
<sup>3</sup>Department of Physics, Princeton University, Princeton, New Jersey 08544, USA

<sup>4</sup>Princeton Center for Theoretical Science, Princeton University, Princeton, New Jersey 08544, USA

<sup>5</sup>Program in Applied and Computational Mathematics, Princeton University, Princeton, New Jersey 08544, USA

(Received 22 June 2011; published 13 September 2011)

The densest binary sphere packings have historically been very difficult to determine. The only rigorously known packings in the  $\alpha$ - $x$  plane of sphere radius ratio  $\alpha$  and relative concentration  $x$  are at the Kepler limit  $\alpha = 1$ , where packings are monodisperse. Utilizing an implementation of the Torquato-Jiao sphere-packing algorithm [S. Torquato and Y. Jiao, *Phys. Rev. E* **82**, 061302 (2010)], we present the most comprehensive determination to date of the phase diagram in  $(\alpha, x)$  for the densest binary sphere packings. Unexpectedly, we find many distinct new densest packings.



## close packing is not close enough

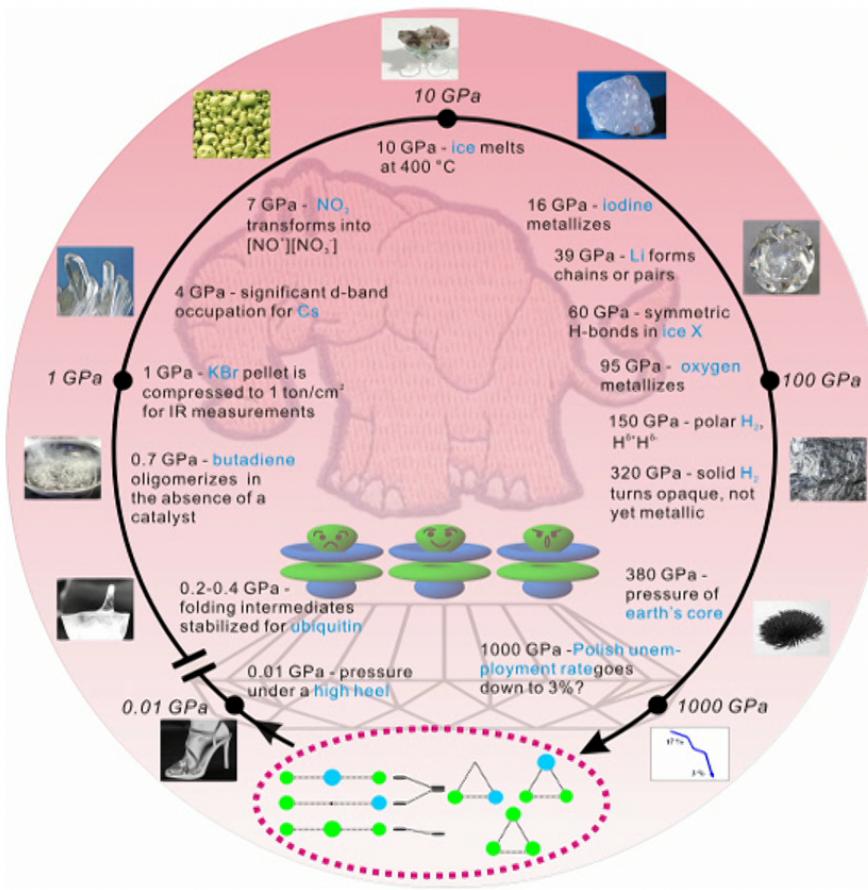
The packing of identical spheres maximizes at ~74% ( $\pi/\sqrt{18}$ ). If spheres are deformed, they can pack more efficiently -- *vide* work on ellipsoids and M&M's, and the trivial case of extreme deformation to cubes. Other polyhedra can also pack space fully.

Spheres of different size may pack more efficiently than spheres of identical size. So.... at some pressure there may be a driving force for a lattice of element E to "electronically disproportionate" to  $(E^{\delta+})_m(E^{\delta-})_n$  sublattices. Such sublattices could be of any dimensionality; there could be more than two sublattices.

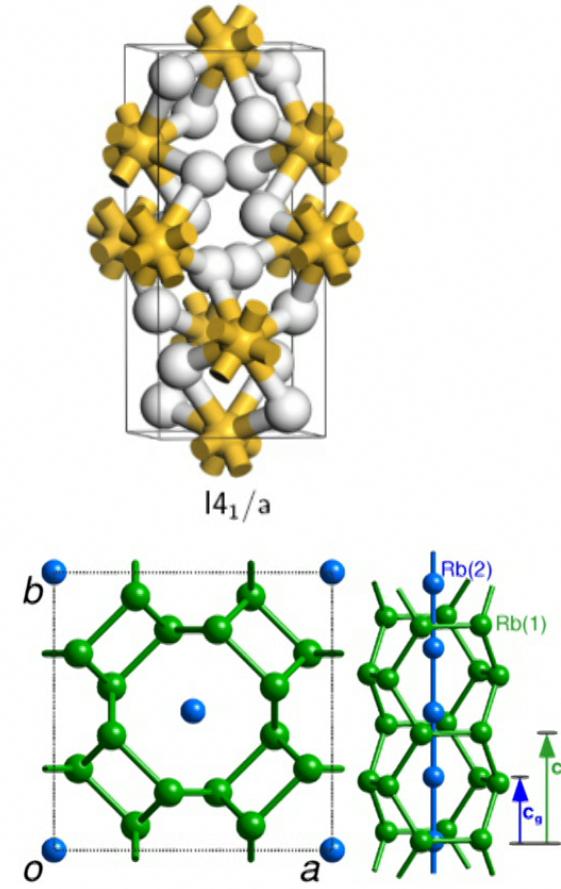
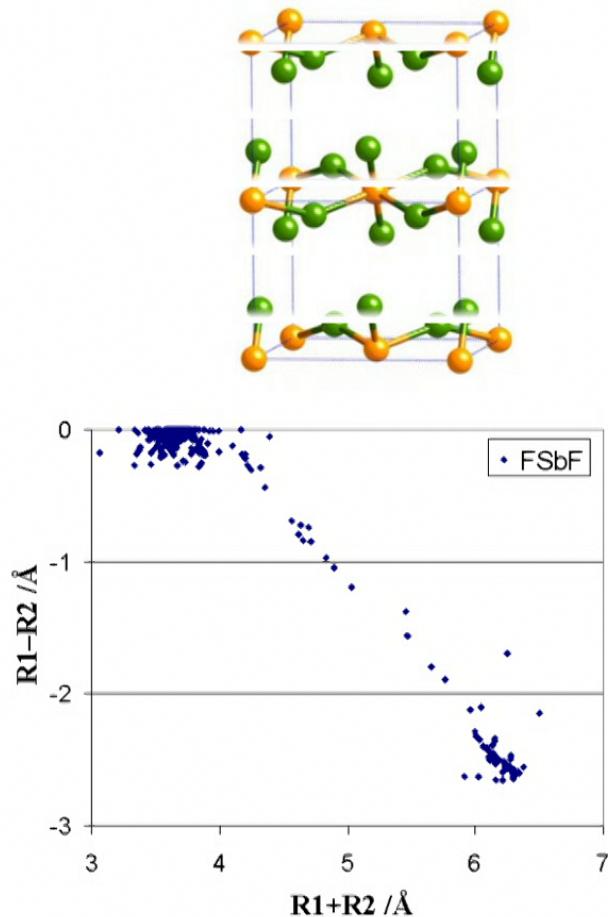
Electron transfer between the sublattices may not be needed to differentiate them. One could imagine a disproportionation to sublattices A, B, C where each is roughly neutral, but characterized by different kinds of bonding – covalent, ionic, or metallic – within each sublattice.

# The Chemical Imagination at Work in *Very Tight Places*

Wojciech Grochala, Roald Hoffmann, Ji Feng, and Neil W. Ashcroft  
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A playground for sure,  
but also another reason,  
a personal and philosophical one....

