

Title: PSI 2017/2018 - Quantum Integrable Models - Lecture 8

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Abstract:

So far:

$$\psi_i \partial_{\bar{w}} \psi^i + \bar{\psi}_i \partial_w \bar{\psi}^i + J_j^i \bar{J}_i^j$$

J, \bar{J} currents for $U(n)$ symmetry

chiral + antichiral + $J\bar{J}$
always integrable

σ -models

X a Riemannian manifold

x_i local coordinates

$g^{ij} dx_i dx_j$ metric

\Rightarrow 2d field theory

Fundamental field is a map $\varphi: \mathbb{R}^2 \rightarrow X$

$\varphi = \{\varphi_i\}$

Lagrangian is

$$\int_{\mathbb{R}^2} g^{ij}(\varphi(w, \bar{w})) \partial_w \varphi_i \partial_{\bar{w}} \varphi_j$$

Example If $g^{ij} = \delta^{ij} + x_i x_j$

Then the Lagrangian is

$$\int (\delta^{ij} + \varphi_i \varphi_j) (\partial_w \varphi_i) (\partial_{\bar{w}} \varphi_j)$$

$$g^{ij} = \delta^{ij} + \sum F^{ij, k_1 \dots k_n} x_{k_1} \dots x_{k_n}$$

replace $x_{k_i} \rightarrow \varphi_{k_i}$
in Lagrangian.

integrable

Fundamental field is a map $\varphi: \mathbb{R}^2 \rightarrow X$
 $\varphi = \{\varphi_i\}$
Lagrangian is $\int_{\mathbb{R}^2} g^j(\varphi(u, \bar{u})) \partial_u \varphi_i \partial_{\bar{u}} \varphi_j$

$$g^j = \delta^j + x_i x_j$$

can be

$$\partial_u \varphi_i \partial_{\bar{u}} \varphi_j$$

$$= \delta^j + x_i x_j$$

$$\rightarrow \varphi_i$$

$$\varphi_i = x_i \circ \varphi: \mathbb{R}^2 \rightarrow \mathbb{R}$$

Suppose we can choose local coords x_i, y_i on X

$$u_i = x_i + \sqrt{-1} y_i$$

$$\bar{u}_i = x_i - \sqrt{-1} y_i$$

Such that the metric is

$$g^j du_i d\bar{u}_j \text{ where}$$

$$\bar{g}^j(u, \bar{u}) = g^j(u, \bar{u})$$

$i=1, \dots, n$
 x_i, y_i on X

Lagrangean is \mathbb{R}^2

$$\varphi_i = \alpha_i \circ \varphi : \mathbb{R}^2 \rightarrow \mathbb{R}$$

Suppose we can choose local coords

$$u_i = x_i + \sqrt{-1} y_i$$

$$\bar{u}_i = x_i - \sqrt{-1} y_i$$

Such that the metric is $g^{ij} du_i d\bar{u}_j$ where

$$\bar{g}^{ij}(u, \bar{u}) = g^{ij}(u, \bar{u})$$

local coords x_i, y_i on X $i=1, \dots, n$

X is complex manifold w. Hermitian metric

If in addition,

$$\partial_{\bar{u}_k} g^{ij} - \partial_{\bar{u}_j} g^{ik} = 0$$

X is Kähler

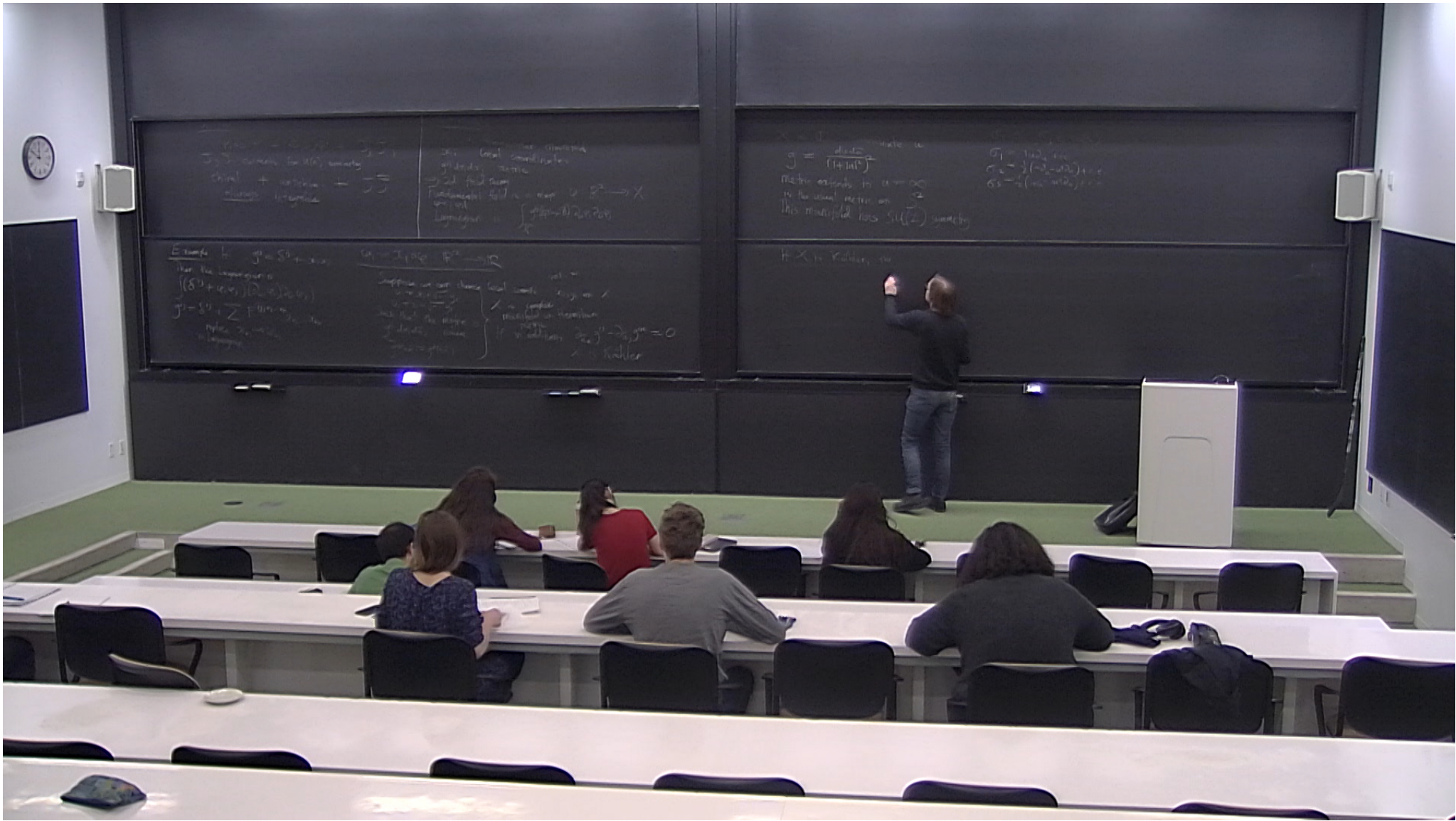
$X = \mathbb{D}$ coordinate u

$$g = \frac{du du}{(1+|u|^2)^2}$$

metric extends to $u = \infty$

is the usual metric on S^2

This manifold has $SU(2)$ symmetry



$$\sigma_1, \sigma_2, \sigma_3 \in \mathfrak{su}(2)$$

$$\sigma_1 = iu\partial_u + \text{c.c.}$$

$$\sigma_2 = \frac{1}{2}(-\partial_u - u^2\partial_u) + \text{c.c.}$$

$$\sigma_3 = \frac{1}{2}(i\partial_u - iu^2\partial_u) + \text{c.c.}$$

If X is Kähler, then
use fields $\sigma_i, \bar{\sigma}_i$
corresponding to u_i, \bar{u}_i coordinates.
Lagrangian is
$$\int \partial_{\bar{w}} \sigma_i \partial_w \bar{\sigma}_j g^{ij}(\sigma, \bar{\sigma})$$

(+ top. term)



$X = \mathbb{C} \cong \mathbb{C}P^1$, with Fubini-Study metric

then σ -model has fields
 $\sigma, \bar{\sigma}$ with Lagrangian

$$\int \frac{\partial_{\bar{w}} \bar{\sigma} \partial_w \sigma}{(1 + \sigma \bar{\sigma})^2}$$

Consider Lag. \nearrow Introduce auxiliary
 fields $\rho^i, \bar{\rho}^i$ of \dim^n and spin
 $(1, -1)$ spins

$$\rho^i \sim dw \quad \bar{\rho}^i \sim d\bar{w}$$

Consider the Lagrangian

$$\int \partial_{\bar{w}} \sigma_i \partial_w \bar{\sigma}_j g^{ij}(\sigma, \bar{\sigma}) - \int \rho^i \bar{\rho}^j g_{ij}(\sigma, \bar{\sigma})$$

$$g_{ij} g^{ik} = \delta_j^k$$

ρ^i : mass, but no kinetic term

Equivalent to previous Lagrangian.

Lagrangian is $\int_{\mathbb{R}^2}$

$$\beta^i = p^i - g^{ij} \partial_{\bar{w}} \bar{\sigma}_j$$

$$\bar{\beta}^i = \bar{p}^i - g^{ij} \partial_w \sigma_j$$

Lagrangian becomes

$$\int \partial_{\bar{w}} \sigma_i \partial_w \bar{\sigma}_j g^{ij} - \int \partial_w \bar{\sigma}_i \partial_{\bar{w}} \sigma_k g^{ik} g_{ij}$$

cancel

$$- \int \beta^i \bar{\beta}^j g_{ij} - \int \beta^i \partial_{\bar{w}} \bar{\sigma}_i - \int \bar{\beta}^i \partial_w \sigma_i$$

Lagrangian is $\int_{\mathbb{R}^2}$

$$\beta^i = p^i - g^{ij} \partial_{\bar{w}} \bar{\sigma}_j$$

$$\bar{\beta}^i = \bar{p}^i - g^{ij} \partial_w \sigma_j$$

Lagrangian becomes

$$\int \partial_{\bar{w}} \sigma_i \partial_w \bar{\sigma}_j g^{ij} - \int \partial_w \bar{\sigma}_i \partial_{\bar{w}} \sigma_k g^{ik} g_{ij}$$

Inverse metric

$$- \int \beta^i \bar{\beta}^j g_{ij}$$

deformation

$$- \int \beta^i \partial_{\bar{w}} \bar{\sigma}_i$$

Chiral
- Only depends on X as a \mathbb{C}^2 manifold

$$- \int \bar{\beta}^i \partial_w \sigma_i$$

Anti-chiral

cancel

$$X = \mathbb{D} \quad g = \frac{du d\bar{u}}{(1+|u|^2)^2}$$

Inverse metric is $g_u (1+|u|^2)^2$

Our Lagrangian is

$$\int \beta \bar{\beta} (1 + \sigma \bar{\sigma})^2 + \int \beta \partial_{\bar{u}} \sigma + \int \bar{\beta} \partial_u \bar{\sigma}$$

$$\sigma_1, \sigma_2, \sigma_3 \in \mathfrak{su}(2)$$

$$\sigma_1 = iu \partial_u + c.c.$$

$$\sigma_2 = \frac{1}{2}(-\partial_u - u^2 \partial_u) + c.c.$$

$$\sigma_3 = \frac{1}{2}(i\partial_u - iu^2 \partial_u) + c.c.$$

$$X = \mathbb{C} \quad g = \frac{du d\bar{u}}{(1+|u|^2)^2}$$

Inverse metric is $g_{ij} (1+|u|^2)^2$

Our Lagrangian is

$$\int \beta \bar{\beta} (1 + \sigma \bar{\sigma})^2 + \int \beta \partial_{\bar{u}} \sigma + \int \bar{\beta} \partial_u \bar{\sigma}$$

Claim $SU(2)$ acts on the free theory $\int \beta \partial_{\bar{u}} \sigma + \int \bar{\beta} \partial_u \bar{\sigma}$, and interaction term is $\mathbb{J} \bar{\mathbb{J}}$

$$\sigma_1, \sigma_2, \sigma_3 \in su(2)$$

$$\sigma_1 = i(u \partial_u + \text{c.c.})$$

$$\sigma_2 = \frac{1}{2} (-\partial_u - u^2 \partial_u) + \text{c.c.}$$

$$\sigma_3 = \frac{1}{2} (i \partial_u - i u^2 \partial_u) + \text{c.c.}$$

What are currents for $\partial_u, u\partial_u, u^2\partial_u$
in the chiral theory $\int \beta \partial_{\bar{u}} \sigma$?

∂_u translate σ : $\sigma \rightarrow \sigma + \epsilon$

If $A_{\bar{u}}$ is a gauge field gauging this symmetry
then $\int \beta \partial_{\bar{u}} \sigma + \int \beta A_{\bar{u}}$

$$J_{\partial_u} = \beta$$

σ - models

X a Riemannian manifold

x_i local coordinates

$g^{ij} dx_i dx_j$ metric

\Rightarrow 2d field theory

Fundamental field is a map $\varphi: \mathbb{R}^2 \rightarrow X$

$\varphi = \{ \varphi_i \}$

Lagrangian is $\int_{\mathbb{R}^2} g^{ij}(\varphi(\omega, \bar{\omega})) \partial_\omega \varphi_i \partial_{\bar{\omega}} \varphi_j$

We can calculate $\sum \gamma_{0i} \gamma_{0i}$ by calculating
the inverse metric:

$$\sigma_1^u \sigma_1^{\bar{u}} + \sigma_2^u \sigma_2^{\bar{u}} + \sigma_3^u \sigma_3^{\bar{u}}$$

$$= \left(-u\bar{u} + \frac{1}{2}u^2 + \frac{1}{2}\bar{u}^2 \right) du d\bar{u}$$

After coordinate change $\bar{u} \rightarrow -\bar{u}$

this becomes $\frac{1}{2}(1+|u|^2)^2 du d\bar{u}$ ✓

We can calculate $\sum \bar{J}_0, \bar{J}_0$ by calculating the inverse metric:

$$\sigma_1^u \sigma_1^{\bar{u}} + \sigma_2^u \sigma_2^{\bar{u}} + \sigma_3^u \sigma_3^{\bar{u}}$$

$$= \left(-u\bar{u} + \frac{1}{2}u^2 + \frac{1}{2}\bar{u}^2 \right) du d\bar{u}$$

After coordinate change $\bar{u} \rightarrow -\bar{u}$
 this becomes $\frac{1}{2}(1+|u|^2)^2 du d\bar{u}$ ✓

Conclusion
 σ -model on S^2
 is free chiral + a.c. + $\bar{J}\bar{J}$
 \Rightarrow integrable.