

Title: PSI 2017/2018 - Quantum Gravity - Lecture 4

Date: Mar 22, 2018 10:15 AM

URL: <http://pirsa.org/18030041>

Abstract:

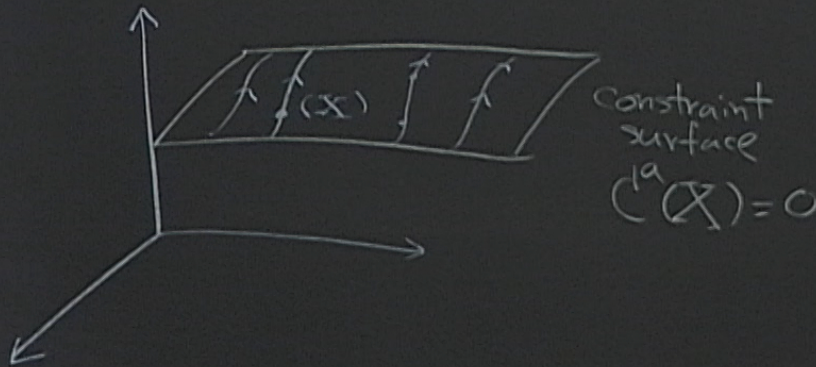
## Kinematical Phase Space.

- constraints  $\{\Phi^\alpha\} = \{1^{\text{st}} \text{ class}\} \cup \{2^{\text{nd}} \text{ class}\}$

$$\{c^a, \Phi^\alpha\} \approx 0 \rightarrow \text{gauge transformation redundancy}$$

- phase space  $\{X^a\}$   $c^a = 0$

$$\frac{dX^a}{d\lambda} = \{X^a, c\}$$



redundancy

given initial conditions  $\rightarrow$  solutions

$\{X, C\}$

reasonable question: "relation" between dynamical variables

$\hookrightarrow$  "gauge fixing"

$$G^a(X) = 0$$

observables.

$$\{F(X), c^a\} = 0$$

$$\{, \} \longrightarrow [, ]$$

$$c \longrightarrow \hat{c} |\psi_{\text{phys}}\rangle = 0$$

$$F \longrightarrow [\hat{F}, \hat{c}] = 0,$$

$$G^a(X) = 0$$

def -  $\{ \neq$  1<sup>st</sup> class -  $\neq$  2<sup>nd</sup> class.

relations

relation "between  
canonical variables

"gauge fixing"  $G^a(X) = 0$

\* kinematical dof - 2 \* 1st class - \* 2nd class.

observables.

$$\{F(X), c^a\} = 0$$

$$\{, \} \longrightarrow [, ]$$

$$c \longrightarrow \hat{c} | \psi_{phys} \rangle = 0$$

$$F \longrightarrow [\hat{F}, \hat{c}] = 0.$$

phase space

$$a) (q, t, p_q, p_t) \longrightarrow \mathcal{H}_{\text{kin}}?$$

$$\bullet \mathcal{H}_{\text{kin}} = \mathcal{L}^2(\mathbb{R}^2)$$

$$\psi(q, t)$$

$$\langle \psi, \phi \rangle =$$

$$\int dq dt \bar{\psi}(q, t) \phi(q, t)$$

$$\bullet \begin{matrix} q \\ t \end{matrix} \longrightarrow \begin{matrix} \hat{q}|\psi\rangle = q|\psi\rangle \\ \hat{t}|\psi\rangle = t|\psi\rangle \end{matrix}$$

$$\begin{matrix} p_q \\ p_t \end{matrix} \longrightarrow$$

$$\begin{matrix} \hat{p}_q = -i\hbar \frac{\partial}{\partial q} \\ \hat{p}_t = -i\hbar \frac{\partial}{\partial t} \end{matrix}$$

phase space  
 $(q, t, p_q, p_t) \longrightarrow \mathcal{H}_{\text{kin}}?$

•  $\mathcal{H}_{\text{kin}} = \mathcal{L}^2(\mathbb{R}^2)$

$\Psi(q, t)$

$\langle \Psi, \Phi \rangle = \int dq dt \overline{\Psi}(q, t) \Phi(q, t)$

•  $\begin{pmatrix} q \\ t \end{pmatrix} \longrightarrow \begin{pmatrix} \hat{q} |\Psi\rangle = q |\Psi\rangle \\ \hat{t} |\Psi\rangle = t |\Psi\rangle \end{pmatrix}$

$\begin{pmatrix} p_q \\ p_t \end{pmatrix} \longrightarrow$

$\hat{p}_q = -i\hbar \frac{\partial}{\partial q}$   
 $\hat{p}_t = -i\hbar \frac{\partial}{\partial t}$

$$b) \quad C = p + \frac{p^2}{2m} \rightarrow \hat{C} = -i\hbar \frac{\partial}{\partial t} - \frac{\hbar^2}{2m} \frac{\partial^2}{\partial q^2}$$

$$\hat{C}|\psi\rangle = 0 \quad \text{Schrödinger eq.}$$



$$b) \quad C = p_t + \frac{p_q^2}{2m} \longrightarrow \hat{C} = -i\hbar \frac{\partial}{\partial t} - \frac{\hbar^2}{2m} \frac{\partial^2}{\partial q^2}$$

$$\hat{C}|\Psi\rangle = 0 \quad \text{Schrödinger eq.}$$

$$a) \quad \hat{C}|\Psi\rangle = 0.$$

$$\Psi(q) = \Psi(q, t=t_0)$$

solution Schrödinger eq.

$$\Psi_{\text{phys}}(q, t) = \exp\left(\frac{-i}{\hbar} \hat{H} t\right) \Psi(q)$$

$$\langle \Psi_{\text{phys}}^{(1)}, \Psi_{\text{phys}}^{(2)} \rangle = \int dt \left( \int dq \overline{\Psi^{(1)}(q)} \Psi^{(2)}(q) \right)$$

→ physical inner product  $\int_{-\infty}^{\infty}$

$$\mathcal{H}_{\text{phys}} = \mathcal{L}^2(\mathbb{R})$$

$$\langle \Psi_{\text{phys}}^{(1)} | \Psi_{\text{phys}}^{(2)} \rangle = \int dq \overline{\Psi^{(1)}(q)} \Psi^{(2)}(q)$$

$$\frac{\hbar^2}{2m} \frac{\partial^2}{\partial q^2}$$

a)  $\hat{C}|\Psi\rangle = 0.$

$$\Psi(q) = \Psi(q, t=t_0)$$

solution Schrödinger eq:

$$\Psi_{\text{phys}}(q, t) = \exp\left(\frac{i}{\hbar} \hat{H} t\right) \Psi(q)$$

d)  $F \rightarrow \hat{F}$

$$\hat{F}_1 = \hat{q} - \frac{\hat{p}_y}{m} (t - t_0), \quad \hat{F}_2 = \hat{p}_y$$

Exercise

$$= \exp\left(\frac{i}{\hbar} \hat{h}(t)\right) \Psi(q)$$

$$\hat{h} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial q^2}$$

→ Physical condition

$$\hat{H}_{\text{phys}} = 0$$

$$\langle \Psi_{\text{phys}}^{(1)} | \Psi_{\text{phys}}^{(2)} \rangle$$

$\hat{P}_q$  Exercise: Check that  $[\hat{F}_i, \hat{C}] = 0$

Hilbert action

$$S = \frac{1}{k} \int R \sqrt{g} d^4x$$

$$k = \frac{16\pi G}{c^3} = 1$$

eq of motion:  $G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 0 \iff R_{\mu\nu} = 0$

Hilbert action

$$S = \frac{1}{k} \int R \sqrt{g} d^3x$$

$$k = \frac{16\pi G}{c^3} = 1$$

eq of motion:  $G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 0 \iff R_{\mu\nu} = 0$

Hilbert action

$$S = \frac{1}{k} \int R \sqrt{g} d^{\hat{n}}x$$

$$k = \frac{16\pi G}{c^3} = 1$$

eq of motion:  $G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 0 \iff R_{\mu\nu} = 0$

if  $n=3$   $R_{\mu\nu} = 0 \implies R_{\mu\nu} \sigma e = 0$

Hilbert action

$$S = \frac{1}{k} \int R \sqrt{g} d^n x$$

$$k = \frac{16\pi G}{c^3} = 1$$

eq of motion:  $G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = 0 \iff R_{\mu\nu} = 0$

if  $n=3$   $R_{\mu\nu} = 0 \implies R_{\mu\nu} \overset{6}{\sqrt{\quad}} = 0$

↓      ↓

3      3

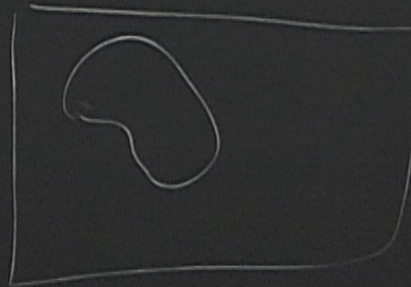


$= 1$   
 $R_{\mu\nu} = 0$

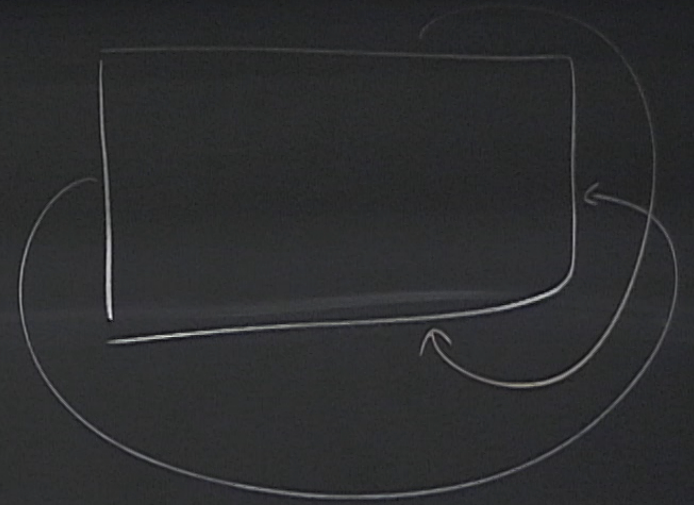
- No local degree of freedom

- Global DOF  $\rightarrow$  topology

Topology characterized by the  
fundamental group



In 3D,  $M = \Sigma \times \mathbb{R}$   
↳ 2d surface  
g: genus

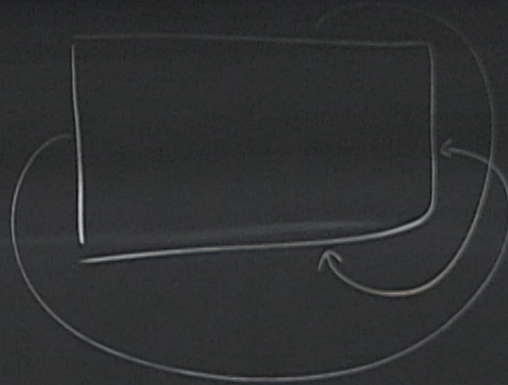


$$d) F \rightarrow \hat{F}$$

$$\hat{F}_1 = \hat{q} - \frac{\hat{p}_q}{m} (t - t_0), \quad \hat{F}_2 = \hat{p}_q$$

Exercise: Check that

In 3D,  $M = \Sigma \times \mathbb{R}$   
 $\hookrightarrow$  2d surface  
 $g$ : genus

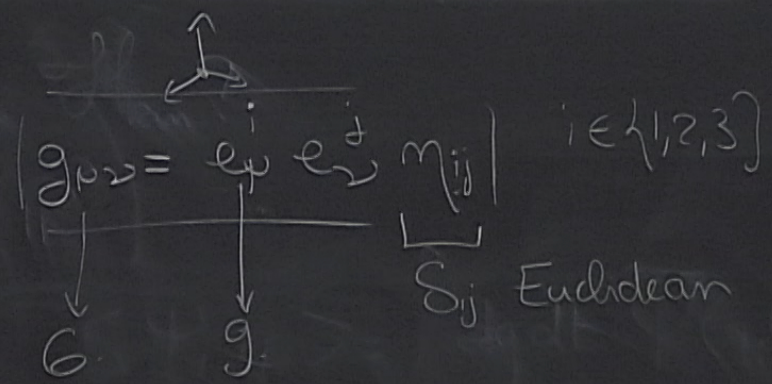


d)  $F \rightarrow \hat{F}$       $F_1 = \hat{q} - \frac{P_1}{m} (t - t_0)$ ,      $F_2 = P_1$

Exercise: Check

First order formulation

$g_{\mu\nu} \rightarrow$



$e_i^{\mu}$ : triad  
 $e_{\mu}^i$ : co-triad

d)  $F \rightarrow \hat{F}$        $\vec{F}_1 = \hat{q} - \frac{P_1}{m} (t - t_0)$ ,       $\vec{F}_2 = P_1$

Exercise: C

First order formulation

$g_{\mu\nu} \rightarrow$

$$g_{\mu\nu} = e_i^{\mu} e_j^{\nu} \eta_{ij} \quad i \in \{1, 2, 3\}$$

$\downarrow$                        $\downarrow$                        $\downarrow$   
 $G$                        $g$                        $S_{ij}$  Euclidean

$e_i^{\mu}$ : triad  
 $e_{\mu}^i$ : co-triad

$$d) F \rightarrow \hat{F} \quad \hat{F}_1 = \hat{q} - \frac{P_1}{m} (t - t_0), \quad \hat{F}_2 = P_1$$

Exercise: Cho

First order formulation

$g_{\mu\nu} \rightarrow$

$$g_{\mu\nu}^{(x)} = e_{\mu}^i e_{\nu}^j \underbrace{m_{ij}}_{\substack{\delta_{ij} \\ \text{Euclidean}}} \quad i \in \{1, 2, 3\}$$

$e_i^{\mu}$ : triad

$e_{\mu}^i$ : co-triad

$$e_{\mu}^i(x) = R_{\mu}^i(x) e_{\mu}^i(x) \quad R \in SO(3)$$

$$e_{\mu}^i e_{\nu}^j \delta^{kl} = \underbrace{(R_{\mu}^k e_{\mu}^i)}_{R^T R = 1} (R_{\nu}^l e_{\nu}^j) \delta^{kl} = g_{\mu\nu}$$

$$d) \mathbf{F} \rightarrow \hat{\mathbf{F}} \quad \mathbf{F}_1 = \hat{\mathbf{q}} - \frac{\mathbf{p}_1}{m} (\hat{t} - t_0), \quad \mathbf{F}_2 = \hat{\mathbf{p}}_1$$

Exercise: C

First order formulation

a) Triads

$g_{\mu\nu} \rightarrow$

$$g_{\mu\nu}(x) = e_{\mu}^i(x) e_{\nu}^j(x) \underbrace{m_{ij}}_{\delta_{ij} \text{ Euclidean}} \quad i \in \{1, 2, 3\}$$

$e_i^{\mu}$ : triad  
 $e_{\mu}^i$ : co-triad

$$e_{\mu}^i(x) = R_{\mu}^i(x) e_{\mu}^i(x) \quad R \in SO(3)$$

$$e_{\mu}^i e_{\nu}^j \delta^{kl} = \underbrace{(R_{\mu}^k e_{\mu}^i)}_{R^T R = 1} (R_{\nu}^l e_{\nu}^j) \delta^{kl} = g_{\mu\nu}$$

use: Check that  $[F, C] = 0$

b) Connection

$\mu, \nu$  - spacetime indices  
 $i, j, k$  - internal indices

connection:

$$\Gamma_{\mu\nu}^{\rho}$$
$$\omega$$

covariant

$$\longrightarrow \nabla_{\mu}$$
$$D_{\mu}$$

- // transport objects with internal indices
- $D_{\mu}$  not  $\rightarrow$  transform well under rotation



$$D_\mu \phi^d = \partial_\mu \phi^d + \omega_{\mu k}^d \phi^k$$

$\omega_{\mu k}^d$ : spin connection

$$D_\mu \phi_j = \partial_\mu \phi_j - \omega_{\mu j}^k \phi_k$$

$$D_\mu (\phi^i \phi_j) = \partial_\mu (\phi^i \phi_j)$$

$$D_\mu \vec{N}^d = \partial_\mu \vec{N}^d - \Gamma_{\mu\nu}^e N_e^d + \omega_{\mu k}^d \phi^k$$

$$D_\mu \phi^j = \partial_\mu \phi^j + \omega_{\mu k}^j \phi^k$$

$\omega_{\mu k}^j$ : spin connection

$$D_\mu \phi_j = \partial_\mu \phi_j - \omega_{\mu j}^k \phi_k$$

$$D_\mu (\phi^i \phi_j) = \partial_\mu (\phi^i \phi_j)$$

$$D_\mu N_{\nu}^d = \partial_\mu N_{\nu}^d - \Gamma_{\mu\nu}^e N_e^d + \omega_{\mu k}^d N_{\nu}^k$$

Compatibility condition

$$D_\mu e^d_\nu = 0$$

$\Rightarrow$  spin connection completely determined by the Levi-Civita connection and triad

$$e^d_\mu D_\nu N^\mu$$

$D_\nu (e^d_\mu \omega^\mu_\nu)$

$$\omega^d_{\mu k} = - e^d_k \nabla_\mu e^d_\nu = e^d_k \Gamma^e_{\mu\nu} e^d_e - e^d_k \partial_\nu e^d_\mu$$

