

Title: PSI 2017/2018 - Relativistic Quantum Information - Lecture 2

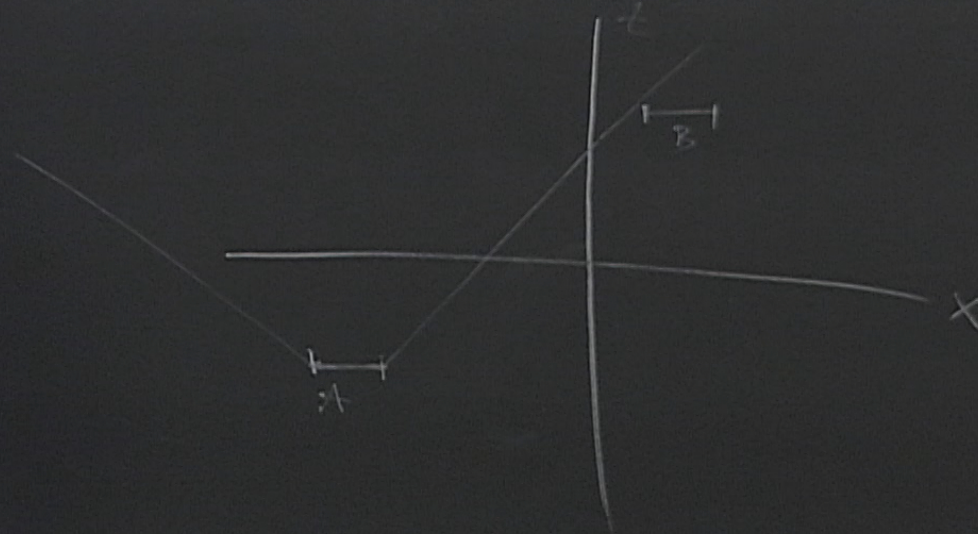
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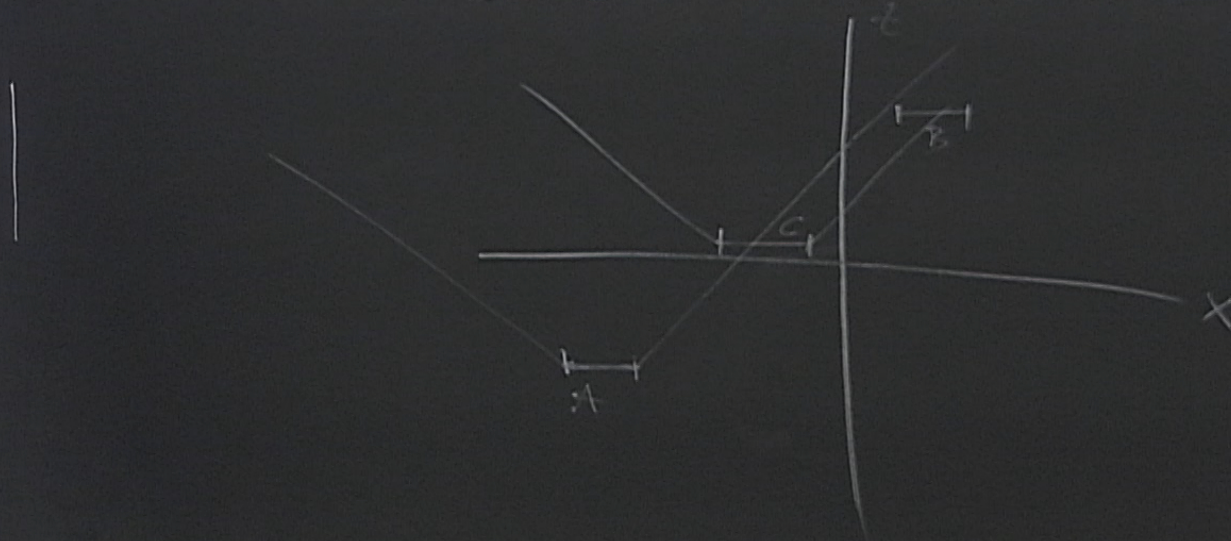
Abstract:

Measuring quantum fields

Measuring quantum fields



Measuring quantum fields



Particle detectors

Particle detectors

What do people do in quantum optics?

Jaynes-Cummings model:

- Two-level "atom" $|e\rangle, |g\rangle$ (energy gap Ω)
- Coupling to a "mode" of the EM field ($\hbar\omega$)

$$\hat{H}_I = \lambda (\hat{\sigma}^+ a e^{i(\Omega-\omega)t} + \hat{\sigma}^- a^\dagger e^{-i(\Omega-\omega)t})$$

$$\sigma^+, \sigma^- \quad \sigma^+ |g\rangle = |e\rangle$$
$$(\sigma^+)^2 = 0 \quad \sigma^- |e\rangle = |g\rangle$$



Particle detectors

What do people do in quantum optics?

Jaynes-Cummings model;

- Two-level "atom" $|e\rangle, |g\rangle$ (energy gap Ω)
- Coupling to a "mode" of the EM field (H.O.)

$$\hat{H}_I = \lambda (\hat{\sigma}^+ \hat{a} e^{i(\Omega - \omega)t} + \hat{\sigma}^- \hat{a}^\dagger e^{-i(\Omega - \omega)t})$$

The light-Matter interaction from first principles

Electron in a Hydrogen-like atom

$$\hat{H} = \frac{\hat{\mathbf{p}}^2}{2m} + V(\hat{\mathbf{x}})$$

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The electron can couple to an EM field

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$$\hat{H} = \frac{(\hat{\vec{p}} - e\vec{A}(t, \hat{\vec{x}}))^2}{2m} + U(t, \hat{\vec{x}}) + V(\hat{\vec{x}})$$

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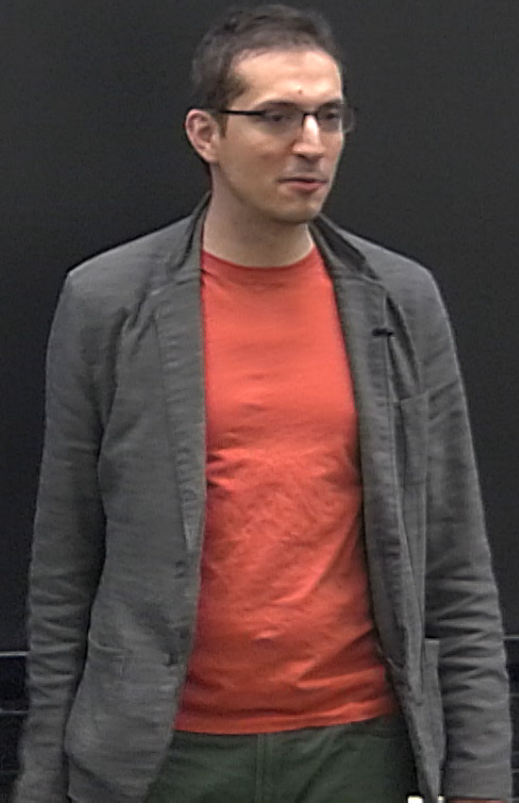
$$\hat{H} = \frac{\hat{\vec{p}}^2}{2m} + eV(\vec{x})$$

The electron can couple to an EM field

$$\hat{H} = \frac{(\hat{\vec{p}} - e\vec{A}(t, \vec{x}))^2}{2m} + eU(t, \vec{x}) + eV(\vec{x})$$

principles

$$\hat{H} = \frac{\hat{P}^2}{2m} + eV(\vec{r})$$



principles

$$\hat{H} = \underbrace{\frac{\hat{\vec{p}}^2}{2m} + eV(\vec{x})}_{\hat{H}_0} - \frac{e}{2m} \left(\hat{\vec{p}} \cdot \vec{A}(t, \vec{x}) + \vec{A}(t, \vec{x}) \cdot \hat{\vec{p}} \right) + eU(t, \vec{x}) + \frac{e^2}{2m} \vec{A}^2(t, \vec{x})$$

principles

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Path of Q.O.

1- Let's neglect $\mathcal{O}(e^2)$

2- Let's work in the Coulomb gauge $[\vec{A}(\vec{x}), \hat{\vec{p}}] = 0 \quad U=0$

Light-matter $H_I = \frac{e}{m} \hat{p} \cdot \vec{A}(t, \hat{x})$

Light-matter $H_{\pm} = \frac{e}{m} \vec{p} \cdot \vec{A}(t, \vec{x})$

Problems: If I pick a different gauge I get different physics

Light-matter $H_{\pm} = \frac{e}{m} \vec{p} \cdot \vec{A}(t, \vec{x})$

Problems: $\left\{ \begin{array}{l} \text{If I pick a different gauge I get different physics} \\ \text{I need to throw } \vec{A}^2 \text{ terms} \end{array} \right.$

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1st solution: Many Göppert-Mayer: Multipole expansion + dipole approximation

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Problems: $\left\{ \begin{array}{l} \text{If I pick a different gauge I get different physics} \\ \text{I need to throw } \vec{A}^2 \text{ terms} \end{array} \right.$

1st solution: Many Göppert-Mayer: Multipole expansion + dipole approximation
(Write H_{\pm} as function of gauge independent variables (\vec{E}, \vec{B})) $H_{\pm} = e \vec{x} \cdot \vec{E}$

physics

perfection + dipole approximation
variables (\vec{E}, \vec{B}) $H_{\text{int}} = e \vec{x} \cdot \vec{E}$ $(\vec{d} \cdot \vec{E})$

Light-matter $H_{\pm} = \frac{e}{m} \hat{\vec{p}} \cdot \vec{A}(t, \hat{\vec{x}})$ ($\vec{p} \cdot \vec{A}$)

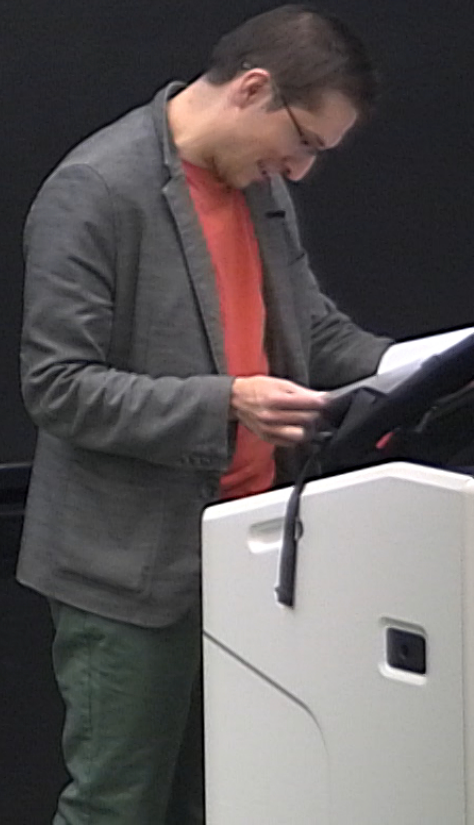
Problems: $\left\{ \begin{array}{l} \text{If I pick a different gauge I get different physics} \\ \text{I need to throw } \vec{A}^2 \text{ terms} \end{array} \right.$

1st solution: Many Göppert-Mayer: Multipole expansion + dipole approximation
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15) solution. Many 001
(Write ψ_{\pm} as function of gauge independent variables (E, l, θ) \neq

Lamb, Scully, Schlicher. When you solve Schrödinger Eq for the Hydrogen atom you picked a gauge!!

Schlicher: When you solve Schrödinger Eq for the Hydrogen atom you pick
- Gauss mismatch between atomic wavefunctions and the Coulomb gauge.

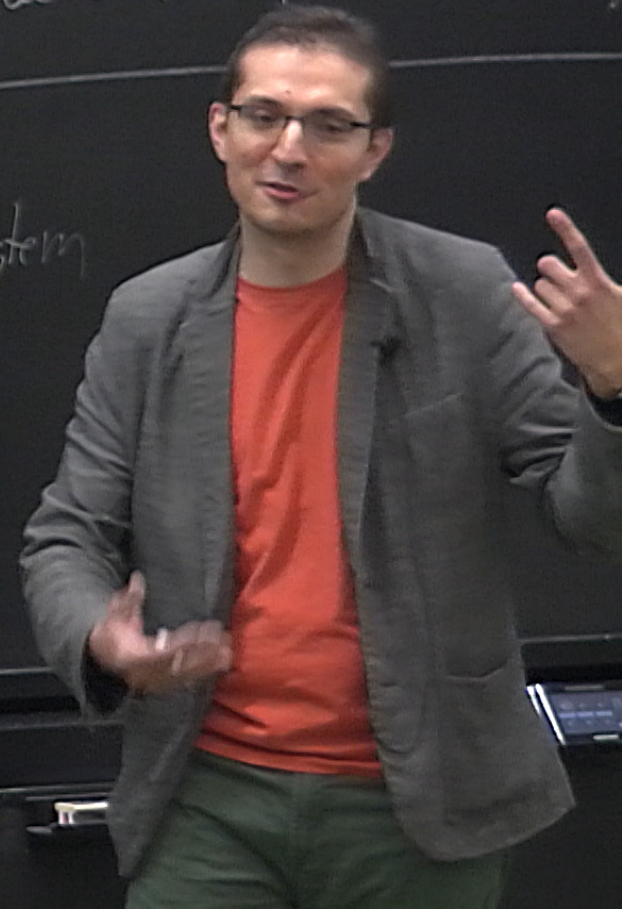


Lamb, Scully, Schlicher: When you solve Schrödinger
1987 - Gauge mismatch between atomic wave

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1987 - Gauge mismatch between atomic wave

What is a "particle detector"

- Localized (in time and space) quantum system
- Coupled to a quantum field
- Easy to measure (clickers quality)



Lamb, Scully, Schlicher: When you solve Schrödinger Eq for
1957 - Gauge mismatch between atomic wavefunctions

What is a "particle detector"?

- Localized (in time and space) quantum system
- Coupled to a quantum field
- Easy to measure (clicking quality) → Not relativistic (first quantized)

When you solve Schrödinger Eq for the Hydrogen atom you
smatch between atomic wavefunctions and the Coulomb gauge.

Quantum system
relativistic (first quantized) } \Rightarrow Hydrogen Atom

$$\hat{H} = \frac{\left(\hat{\vec{p}} - e\vec{A}(t, \hat{\vec{x}})\right)^2}{2m} + eV(t, \hat{\vec{x}}) + eV(\hat{\vec{x}})$$

1 - Let's verify
2 - Let's work in

From dipole coupling to the Uehling-DeWitt model

$$H_{\pm} = e\hat{\vec{x}} \cdot \vec{E}(\hat{\vec{x}})$$

$$\hat{H} = \frac{\left(\hat{\vec{p}} - e\vec{A}(t, \hat{\vec{x}})\right)^2}{2m} + eV(t, \hat{\vec{x}}) + eV(\hat{\vec{x}})$$

1 - Let's recall
2 - Let's work in

From dipole coupling to the Uehling-DeWitt model

$$H_{\pm} = e\hat{\vec{x}} \cdot \vec{E}(\hat{\vec{x}}) = e \sum_{\lambda} \sum_{j} \langle j | \hat{\vec{x}} \cdot \vec{E}(\hat{\vec{x}}) | i \rangle e^{i\Omega_j t} |j\rangle \langle i|$$

can couple to an EM field

$$\frac{e \vec{A}(t, \vec{x})}{2m} + eV(t, \vec{x}) + eV(\vec{x})$$

- 1- Let's neglect $\mathcal{O}(e^2)$
- 2- Let's work in the Coulomb gauge $[\vec{A}(\vec{x}), \vec{P}]$

coupling to the Unruh-DeWitt model

$$\hat{U}(\vec{x}) = e \sum_i \sum_j \langle j | \hat{\vec{x}} \cdot \vec{E}(\vec{x}) | i \rangle e^{i\Omega_j t} |j\rangle \langle i| = \sum_i \sum_j \int d^3x \int d^3x' \langle j | x \rangle \langle x' | \hat{\vec{x}} \cdot \vec{E}(\vec{x}) | x' \rangle \langle x |$$

Let's work in the Coulomb gauge

tt model

$$1 = \sum_i \sum_j \int d^3x \int d^3x' \langle j | x \rangle \langle x | \vec{x} \cdot \vec{E}(x) | x' \rangle \langle x' | i \rangle e^{i\Omega_{ij}t} | j \rangle \langle i |$$

$$H = \frac{1}{2m} (p - eA(t, \mathbf{x}))^2 + eV(t, \mathbf{x}) + e\phi(t, \mathbf{x})$$

2-let

Interaction From dipole coupling to the Unruh-DeWitt

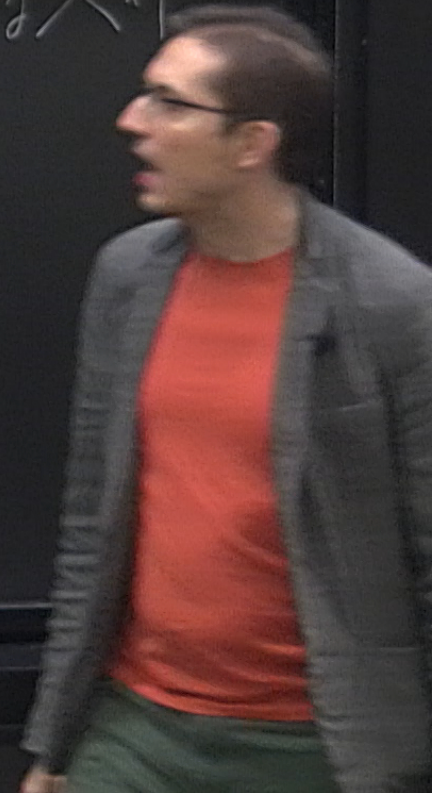
picture

$$H_{\pm} = e \hat{\mathbf{x}} \cdot \vec{E}(\hat{\mathbf{x}}) = e \sum_i \sum_j \langle j | \hat{\mathbf{x}} \cdot \vec{E}(\hat{\mathbf{x}}) | i \rangle e^{i\Omega_j t} |j\rangle \langle i| =$$

Let's work in the Coulomb gauge

tt model

$$1 = \sum_i \sum_j \int d^3x \int d^3x' \langle j|x \rangle \langle x| \vec{x} \cdot \vec{E}(\vec{x}) |x' \rangle \langle x'|i \rangle e^{i\Omega_{ij}t} |j\rangle \langle i|$$



Let's work in the continuum limit

$$\begin{aligned}
 & \text{tt model} \\
 & = \sum_i \sum_j \int d^3x \int d^3x' \overbrace{\langle j|x \rangle}^{\psi_j(x)} \overbrace{\langle x|x' \rangle}^{\vec{x} \cdot \vec{E}(\vec{x})} \overbrace{\langle x'|i \rangle}^{\psi_i(x')} e^{i\Omega_{ij}t} |j\rangle \langle i|
 \end{aligned}$$



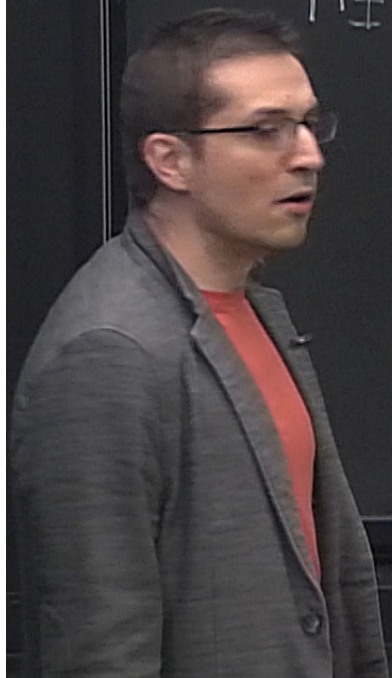
tt model

$$1 = \sum_i \sum_j \int d^3x \int d^3x' \overbrace{\langle j|x \rangle}^{\psi_j(x)} \overbrace{\langle x|x' \rangle}^{\vec{x} \cdot \vec{E}(\vec{x})} \overbrace{\langle x'|i \rangle}^{\psi_i(x')} e^{i\Omega_{ij}t} |j\rangle \langle i|$$

Interaction From dipole coupling to the Uncoupled-DeWitt

Picture D

$$H_{\pm} = e \hat{\vec{x}} \cdot \vec{E}(\hat{\vec{x}}) = e \sum_i \sum_j \langle j | \hat{\vec{x}} \cdot \vec{E}(\hat{\vec{x}}) | i \rangle e^{i\Omega_j t} |j\rangle \langle i| =$$
$$= \sum_i \sum_j \int d^3x \psi_j(\vec{x}) \vec{x} \psi_i(\vec{x}) \cdot \vec{E}(\vec{x}) e^{i\Omega_j t} |j\rangle \langle i|$$



Interaction From dipole coupling to the Uncoupled-Dirac

Picture D

$$H_{\pm} = e \hat{\vec{x}} \cdot \vec{E}(\hat{\vec{x}}) = e \sum_i \sum_j \langle j | \hat{\vec{x}} \cdot \vec{E}(\hat{\vec{x}}) | i \rangle e^{i\Omega_j t} |j\rangle \langle i| =$$
$$= \sum_i \sum_j \int d^3x \psi_j(\vec{x}) \vec{x} \psi_i(\vec{x}) \cdot \vec{E}(\vec{x}) e^{i\Omega_j t} |j\rangle \langle i|$$

\downarrow
 $\propto e^{i\vec{k} \cdot \vec{x}}$

tt model

$$1 = \sum_i \sum_j \int d^3x \int d^3x' \overbrace{\langle j|x \rangle}^{\psi_j(x)} \overbrace{\langle x|\vec{x} \cdot \vec{E}(\vec{x})|x' \rangle}^{\delta^3(x-x')} \overbrace{\langle x'|i \rangle}^{\psi_i(x')} e^{i\Omega_{ij}t} |j\rangle \langle i|$$

tt model

$$1 = \sum_i \sum_j \int d^3x \int d^3x' \overbrace{\langle j|x \rangle}^{\psi_j(x)} \overbrace{\langle x|\vec{x} \cdot \vec{E}(\vec{x})|x' \rangle}^{\delta^{(3)}(\vec{x}-\vec{x}')} \overbrace{\langle x'|i \rangle}^{\psi_i(x')} e^{i\Omega_{ij}t} |j\rangle \langle i|$$

many of these matrix elements are zero
 $i \rightarrow i$ selection rule

tt model

$$1 = \sum_i \sum_j \int d^3x \int d^3x' \overbrace{\langle j|x \rangle}^{\psi_j(x)} \overbrace{\langle x|\vec{x} \cdot \vec{E}(\vec{x})|x' \rangle}^{\delta^{ij}(\vec{x}-\vec{x}')} \overbrace{\langle x'|i \rangle}^{\psi_i(x')} e^{i\Omega_{ij}t} |j\rangle \langle i|$$

many of these matrix elements are zero

$$i \rightarrow \hat{i}$$

Selection rule $\Delta l = 0, \pm 1$
 $l=0 \rightarrow 0$

Interaction From dipole coupling to the Unruh-DeWitt

Picture D

$$H_{\pm} = e \hat{\vec{x}} \cdot \vec{E}(\hat{\vec{x}}) = e \sum_i \sum_j \langle j | \hat{\vec{x}} \cdot \vec{E}(\hat{\vec{x}}) | i \rangle e^{i\Omega_j t} |j\rangle \langle i| =$$

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$\sum_j \psi_j$
 \downarrow
 $e^{i\vec{k}\cdot\vec{x}}$

Interaction From dipole coupling to the Unruh-DeWitt

Picture D

$$H_{\pm} = e \hat{\vec{x}} \cdot \vec{E}(\hat{\vec{x}}) = e \sum_i \sum_j \langle j | \hat{\vec{x}} \cdot \vec{E}(\hat{\vec{x}}) | i \rangle e^{i\Omega_j t} |j\rangle \langle i| = e \sum_i \sum_j \int d^3x \psi_j(\vec{x}) \vec{x} \psi_i(\vec{x}) \cdot \vec{E}(\vec{x}) e^{i\Omega_j t} |j\rangle \langle i|$$

$\sum_j \psi_j$ $\int d^3x \psi_j(\vec{x}) \vec{x} \psi_i(\vec{x}) \cdot \vec{E}(\vec{x}) e^{i\Omega_j t}$

tt model

$$1 = e \sum_i \sum_j \int d^3x \int d^3x' \overbrace{\langle j | x \rangle}^{\psi_j(x)} \overbrace{\langle x | \vec{x} \cdot \vec{E}(\vec{x}) | x' \rangle}^{\delta^{ij}(\vec{x}-\vec{x}')} \overbrace{\langle x' | i \rangle}^{\psi_i(x')} e^{i\Omega_{ij}t} |j\rangle \langle i|$$

many of these matrix elements are zero

$$i \rightarrow i$$

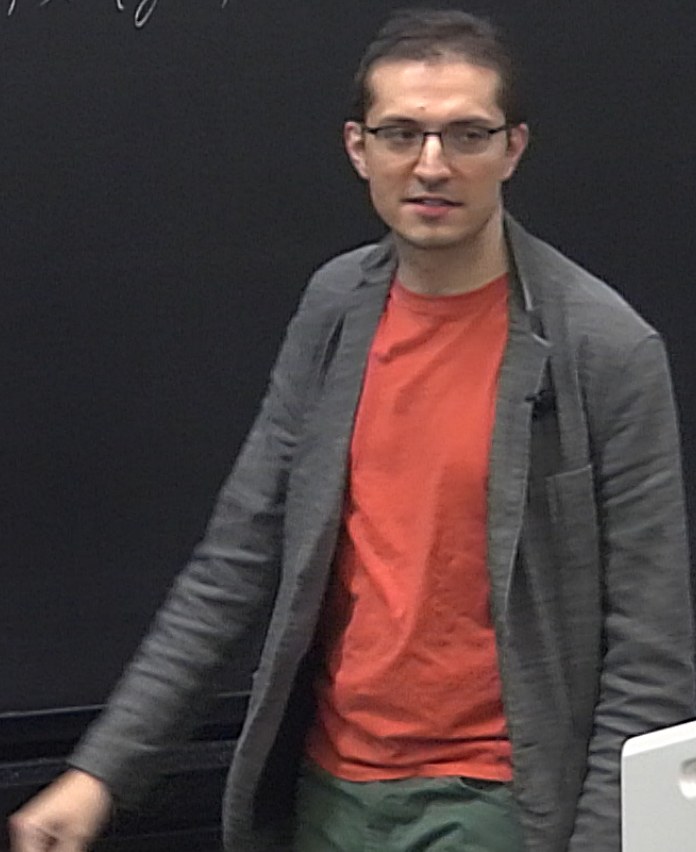
selection rule $\Delta l = 0, \pm 1$
 $l=0 \rightarrow 0$

$$H_T = e \int dx \left(\vec{F}(\vec{x}) e^{i\omega t} \underset{|e\chi g|}{\sigma^+} + \vec{F}^*(\vec{x}) e^{-i\omega t} \underset{|g\chi e|}{\sigma^-} \right) \cdot \vec{E}(\vec{x})$$

 $\vec{F}(\vec{x})$

Smearing vector

$$\vec{F}(\vec{x}) = \psi_e(\vec{x}) \vec{x} \psi_g^*(\vec{x})$$



$$H_{\mp} = e \int dx \left(\vec{F}(\vec{x}) e^{i\omega t} \underset{|e\chi g|}{\sigma^+} + \vec{F}^*(\vec{x}) e^{-i\omega t} \underset{|g\chi e|}{\sigma^-} \right) \cdot \vec{E}(\vec{x}) \quad \vec{F}(\vec{x})$$

$$H_{\mp} = e \int dx \hat{d}(\vec{x}) \cdot \vec{E}(\vec{x})$$

Canonical quantization of $E(\vec{x})$

Smearing vector

$$\vec{F}(\vec{x}) = \psi_e(\vec{x}) \vec{x} \psi_g^*(\vec{x})$$

$$\text{of } E(\vec{x}) \rightarrow \hat{E}(\vec{x})$$

$$e^{i\vec{k}\cdot\vec{x}} \left(\sigma^+ + \frac{\vec{k}^*}{|\vec{k}|} e^{-i\vec{k}\cdot\vec{x}} \sigma^- \right) \cdot \vec{E}(\vec{x})$$

Smearing vector

$$\vec{F}(\vec{x}) = \psi_e(\vec{x}) \vec{x} \psi_g^*(\vec{x})$$

Canonical quantization of $E(\vec{x}) \rightarrow \hat{E}(\vec{x})$

$$\vec{E}(\vec{x})$$

$$\hat{E}(t, \vec{x}) = \sum_s \int \frac{d^3k}{(2\pi)^3} \sqrt{\frac{|\vec{k}|}{2}} \left[-i \hat{a}_{\vec{k},s} \vec{E}(\vec{k},s) e^{i\vec{k}\cdot\vec{x}} + i \hat{a}_{\vec{k},s}^\dagger \vec{E}^*(\vec{k},s) e^{-i\vec{k}\cdot\vec{x}} \right]$$

$$H_{\text{F}} = \int d^3x e^{i\vec{k}\cdot\vec{x}} \left(\vec{F}(\vec{x}) e^{i\omega t} \underset{|e\chi g|}{\sigma^+} + \vec{F}^*(\vec{x}) e^{-i\omega t} \underset{|g\chi e|}{\sigma^-} \right) \cdot \vec{E}(\vec{x})$$

$$H_{\text{F}} = \int d^3x \hat{\vec{d}}(\vec{x}) \cdot \vec{E}(\vec{x})$$

Canonical quantization of $E(\vec{x})$

$$\hat{\vec{E}}(t, \vec{x}) = \sum_{\vec{s}} \int \frac{d^3k}{\sqrt{(2\pi)^3}} \sqrt{\frac{|\vec{k}|}{2}} \left[-i \hat{a}_{\vec{k}, \vec{s}} \vec{E}(\vec{k}, \vec{s}) e^{i\vec{k}\cdot\vec{x}} \right]$$

$$H_{\text{F}} = \int dx e^{i\vec{k}\cdot\vec{x}} \left(\vec{F}(\vec{x}) e^{i\omega t} \underset{|e^{\vec{k}\cdot\vec{x}}|}{\sigma^+} + \vec{F}^*(\vec{x}) e^{-i\omega t} \underset{|e^{\vec{k}\cdot\vec{x}}|}{\sigma^-} \right) \cdot \vec{E}(\vec{x})$$

$$H_{\text{F}} = \int dx \hat{\vec{d}}(\vec{x}) \cdot \vec{E}(\vec{x})$$

Canonical quantization of $E(\vec{x})$

$$\hat{\vec{E}}(t, \vec{x}) = \sum_{\vec{s}} \int \frac{d^3k}{\sqrt{(2\pi)^3}} \sqrt{\frac{|\vec{k}|}{2}} \left[-i \hat{a}_{\vec{k}, \vec{s}} \vec{E}(\vec{k}, \vec{s}) e^{i\vec{k}\cdot\vec{x}} \right]$$

$$(\sigma^+ + \sigma^-) (a + a^\dagger)$$

$$H_{\text{F}} = \int dx e^{i\vec{k}\cdot\vec{x}} \left(\vec{F}(\vec{x}) e^{i\omega t} \underset{|e\chi g|}{\sigma^+} + \vec{F}(\vec{x})^* e^{-i\omega t} \underset{|e\chi e|}{\sigma^-} \right) \cdot \vec{E}(\vec{x})$$

$$H_{\text{F}} = \int dx \hat{\vec{d}}(\vec{x}) \cdot \vec{E}(\vec{x})$$

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$$(\sigma^+ + \sigma^-) (a + a^\dagger)$$

$$\sigma^+ a^\dagger \quad \sigma^- a$$