

Title: PSI 17/18 - Foundations of Quantum Mechanics - Lecture 13

Date: Feb 14, 2018 10:15 AM

URL: <http://pirsa.org/18020072>

Abstract:

Final comments on Bell's theorem and Contextuality

Bell inequality violations as a resource



Magic is a natural force that can be used to override the usual laws of nature.
-- Harry Potter entry in wikipedia



Magic is a natural force that can be used to override the usual laws of nature.

-- Harry Potter entry in wikipedia

Bell-inequality violations are natural phenomena that can be used to override the usual (classical-like) laws of nature

Quantum Spellcraft

Based on Bell-inequality violation

Reduction in communication complexity

Buhrman, Cleve, van Dam, SIAM J.Comput. 30 1829 (2001)

Brassard, Found. Phys. 33, 1593 (2003)

Device-independent secure key distribution

Barrett, Hardy, Kent, PRL 95, 010503 (2005)

Acin, Gisin, Masanes, PRL. 97, 120405 (2006)

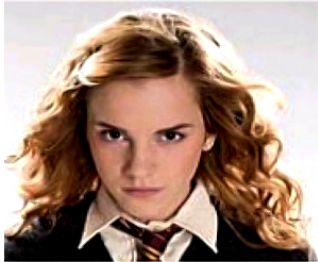
Randomness expansion

Colbeck, Kent, J. Phys. A, 44, 095305 (2011).

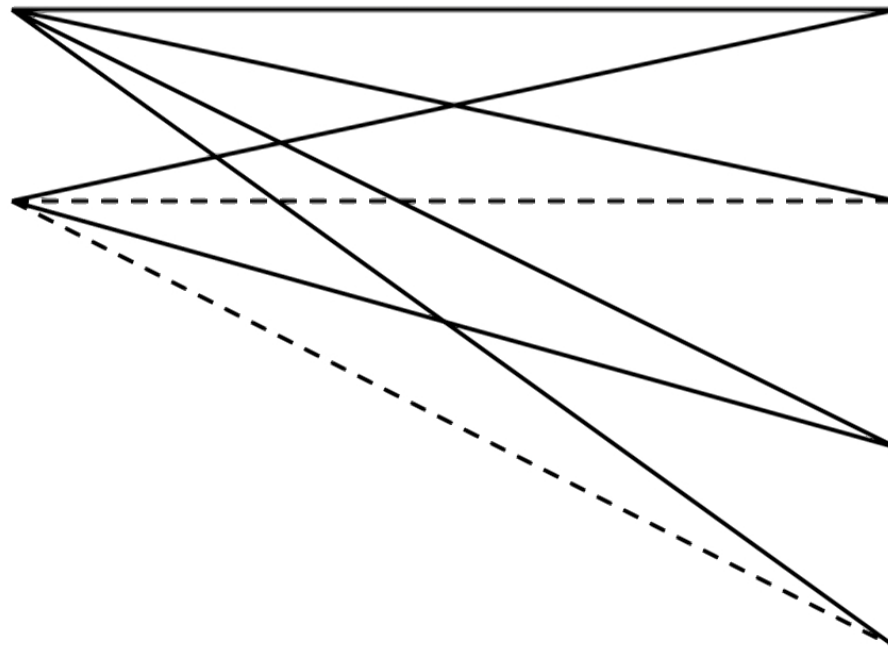
Enhancing zero-error channel capacity

Cubitt, Leung, Matthews, Winter, arXiv:0911.5300

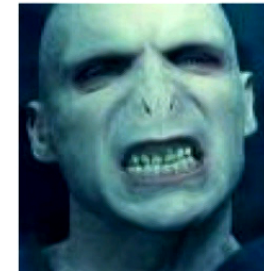
Monogamy of Bell-inequality violating correlations



Alice



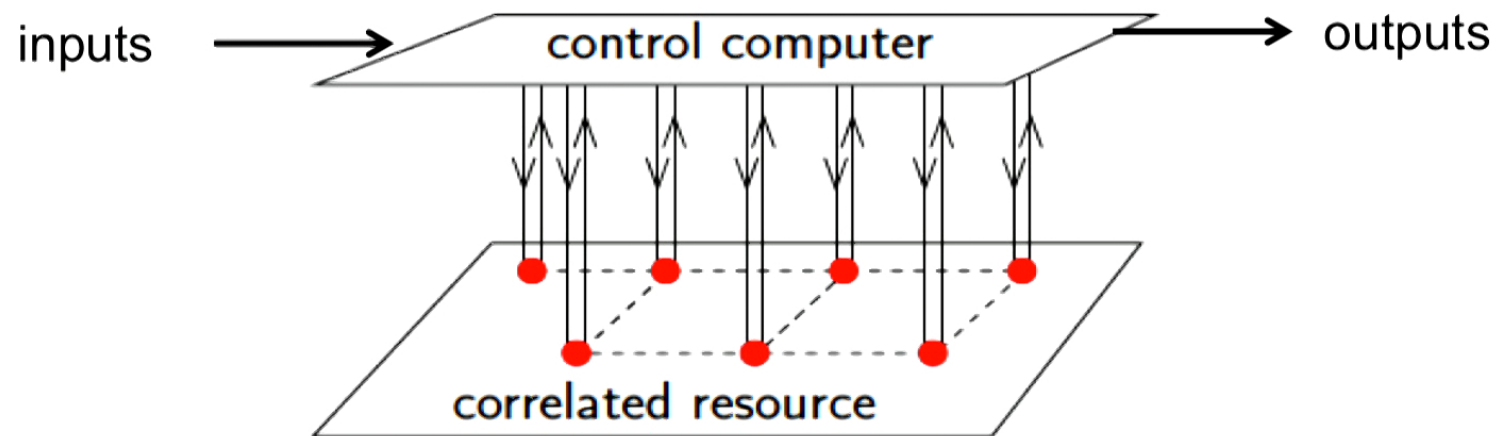
Bob

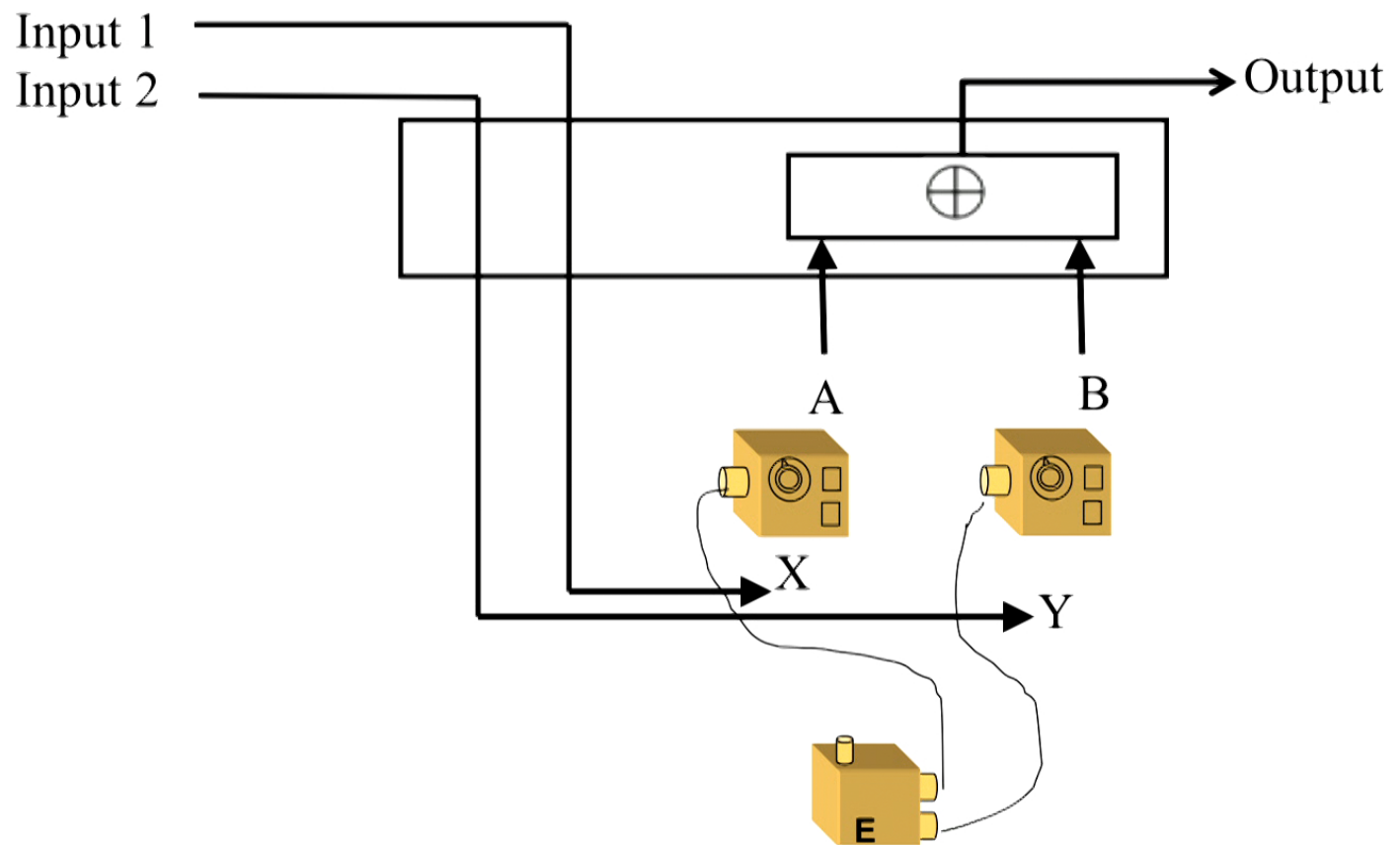


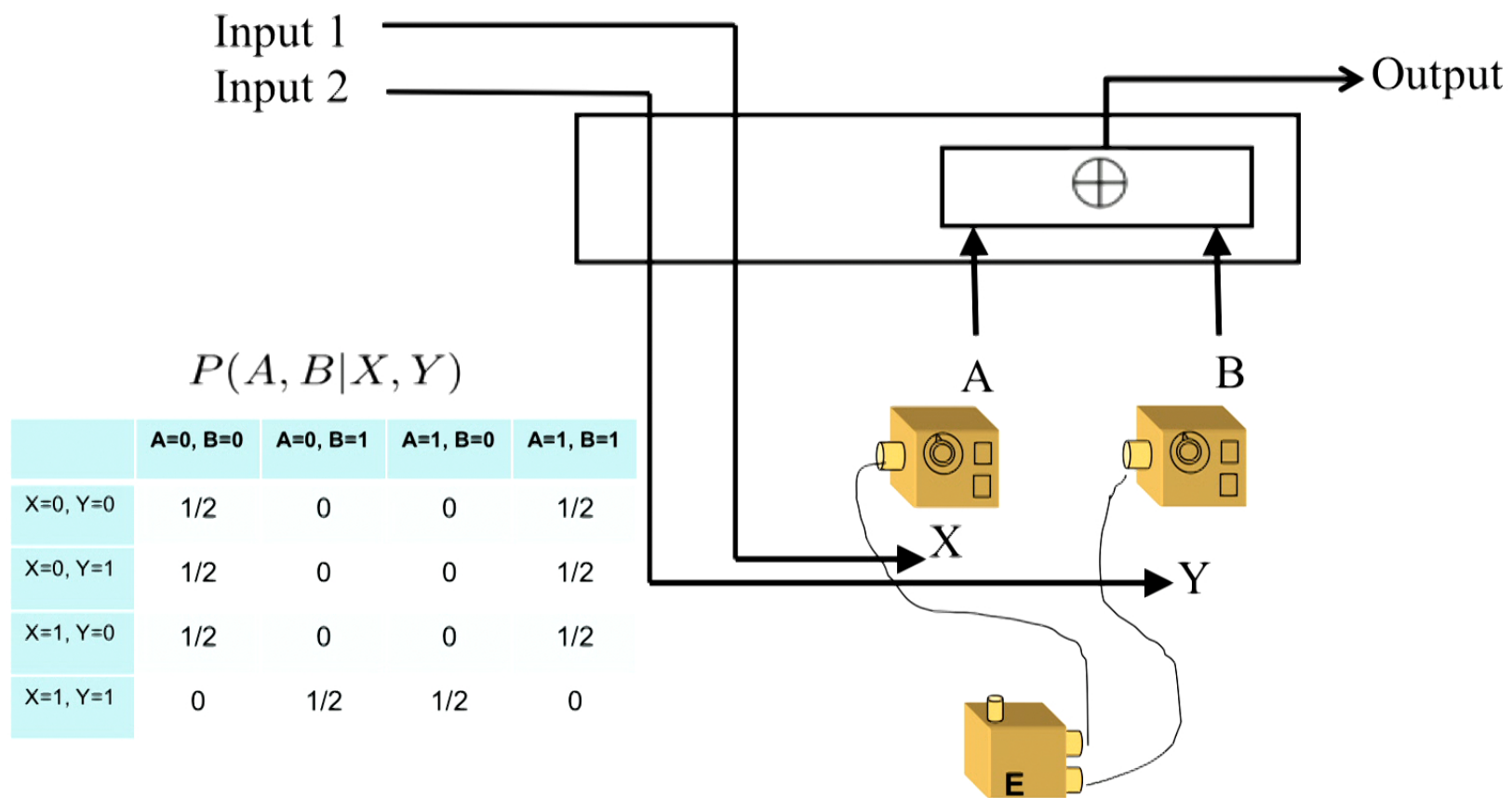
Adversary

Quantum advantages for computation

Anders and Browne. PRL 102, 050502 (2009).







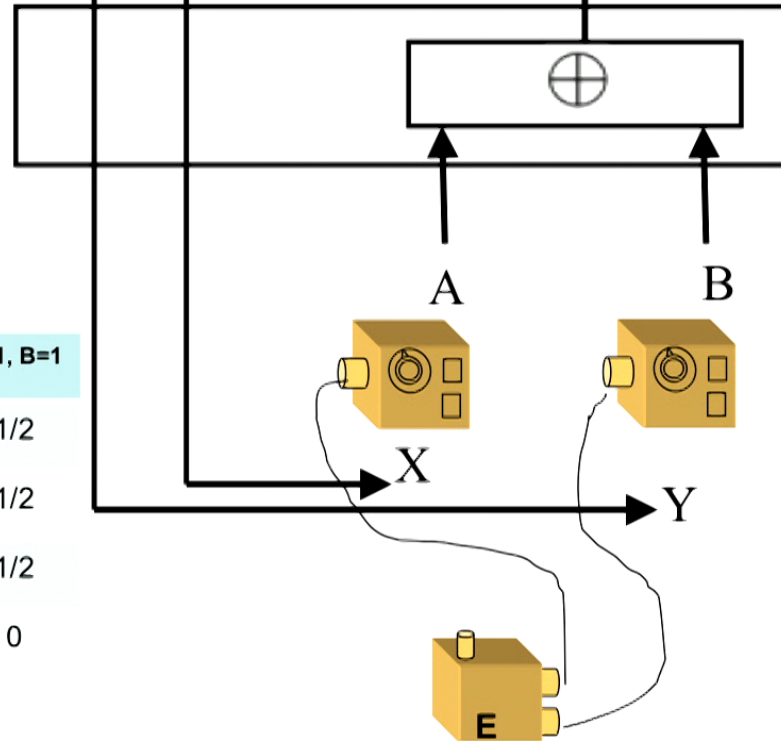
Input 1
Input 2

Output

$P(A, B|X, Y)$

	A=0, B=0	A=0, B=1	A=1, B=0	A=1, B=1
X=0, Y=0	1/2	0	0	1/2
X=0, Y=1	1/2	0	0	1/2
X=1, Y=0	1/2	0	0	1/2
X=1, Y=1	0	1/2	1/2	0

$$A \oplus B = XY$$



Categorizing quantum phenomena

Those arising in a restricted
statistical classical theory

Noncommutativity
Entanglement
Ambiguity of mixtures
EPR Steering
Collapse
Coherent superposition
Teleportation
No cloning

Others...

Weak Nonclassicality

Those not arising in a restricted
statistical classical theory

Bell inequality violations
Contextuality
Computational speed-up
Certain aspects of items on the left
Others...

Strong Nonclassicality

What we want in a notion of nonclassicality

Subject to
direct
experimental
test

Constitutes a
resource

Applicable to a
broad range of
physical
scenarios

What we want in a notion of nonclassicality

Subject to
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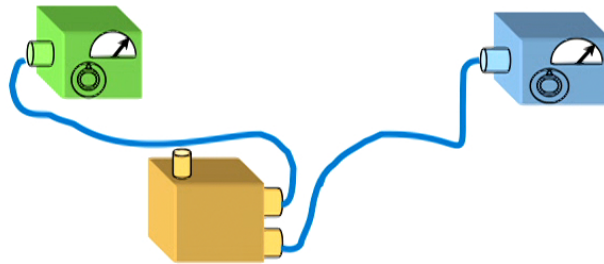
Constitutes a
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Applicable to a
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Failure to admit a
locally causal model



What is needed to witness the failure of local causality



What we want in a notion of nonclassicality

Subject to
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Constitutes a
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Applicable to a
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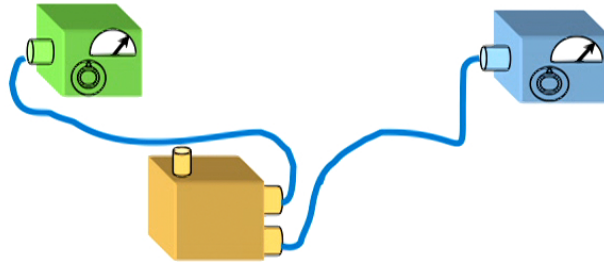
Failure to admit a
locally causal model



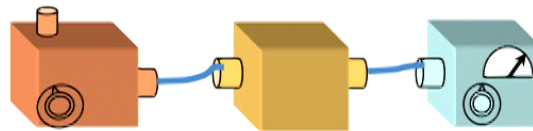
Failure to admit a
noncontextual model



What is needed to witness the failure of local causality



What is needed to witness the failure of noncontextuality



What we want in a notion of nonclassicality

Subject to direct experimental test	Constitutes a resource	Applicable to a broad range of physical scenarios
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Failure to admit a locally causal model



Failure to admit a noncontextual model



What we want in a notion of nonclassicality

Subject to direct experimental test	Constitutes a resource	Applicable to a broad range of physical scenarios
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Failure to admit a locally causal model

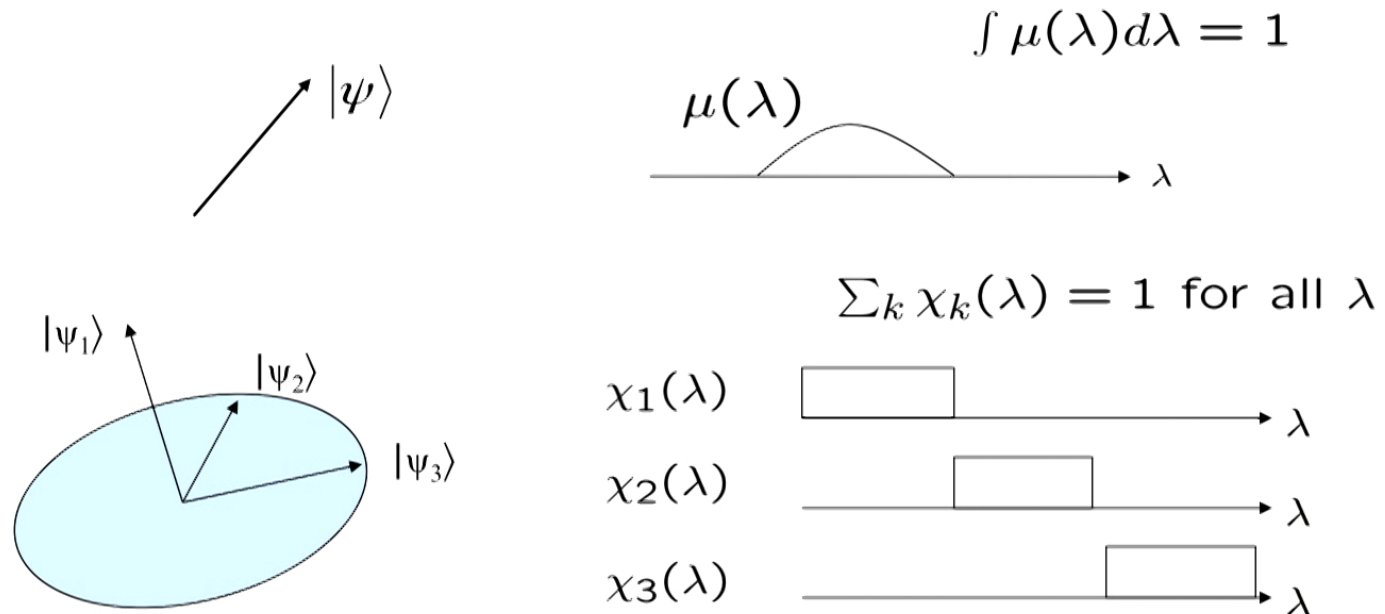


Failure to admit a noncontextual model



The traditional notion of noncontextuality in quantum theory

Outcome-deterministic hidden variable model for pure states and projective measurements

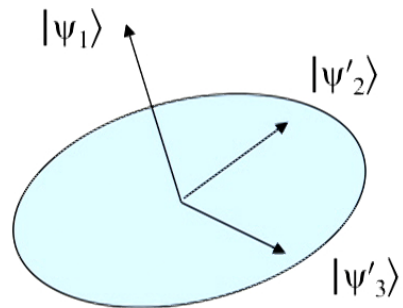
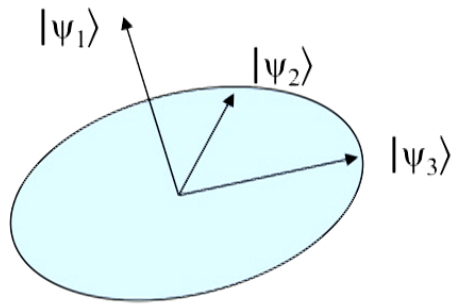


Note: the outcomes are deterministic given λ

$$|\langle\psi|\psi_k\rangle|^2 = \int d\lambda \mu(\lambda) \chi_k(\lambda)$$

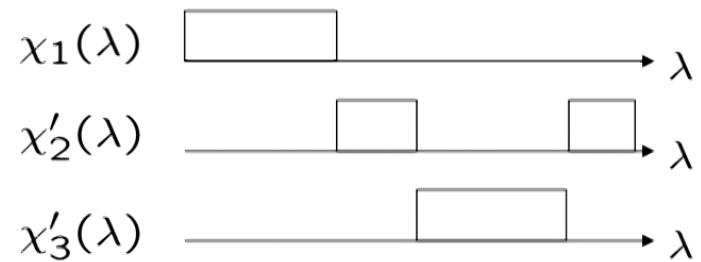
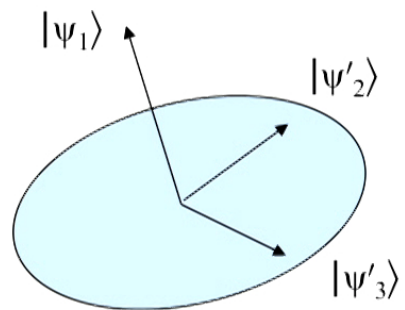
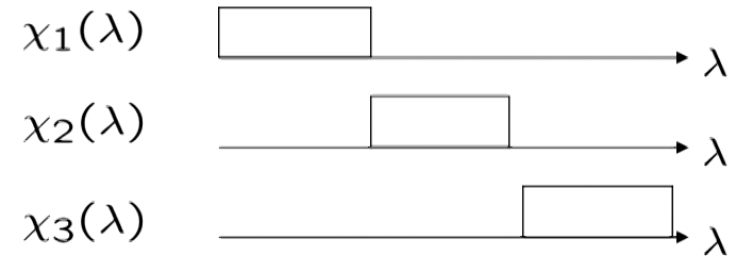
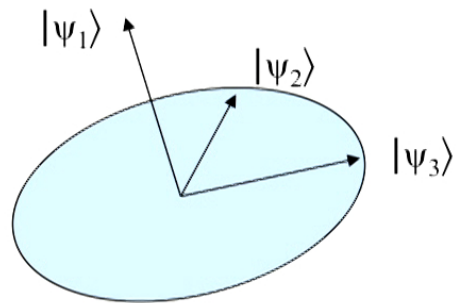
Traditional notion of noncontextuality

A given vector may appear in many different measurements



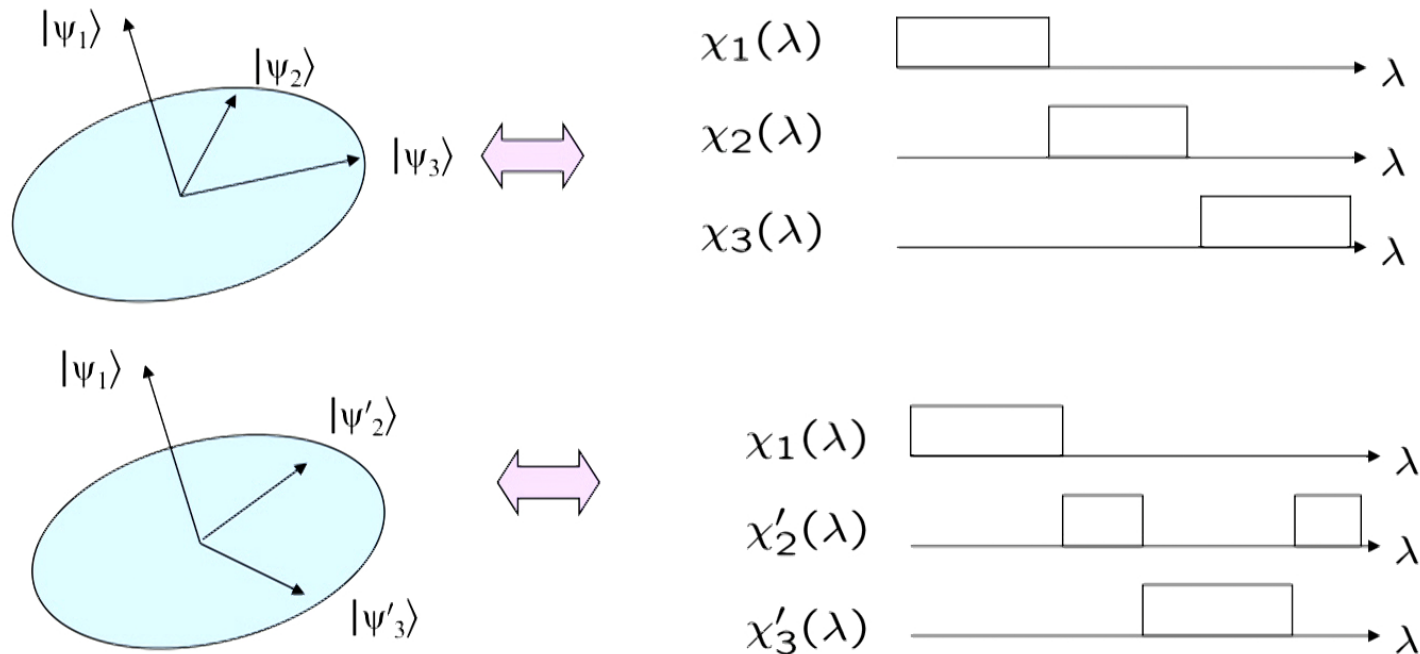
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Traditional notion of noncontextuality

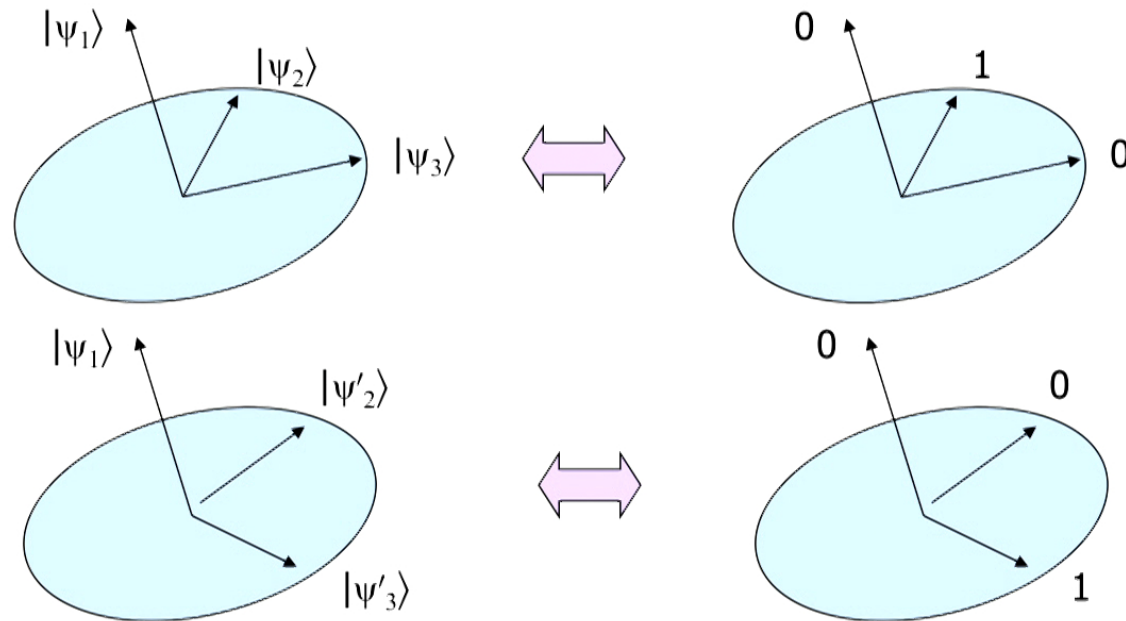
A given vector may appear in many different measurements



The **traditional notion of noncontextuality**:
Every vector is associated with the same $\chi(\lambda)$
regardless of how it is measured (i.e. **the context**)

The **traditional notion of noncontextuality** (take 2):

For every λ , every basis of vectors receives a 0-1 valuation, wherein exactly one element is assigned the value 1 (corresponding to the outcome that would occur for λ), and every vector is assigned the same value regardless of which basis it is considered a part (i.e. **the context**).





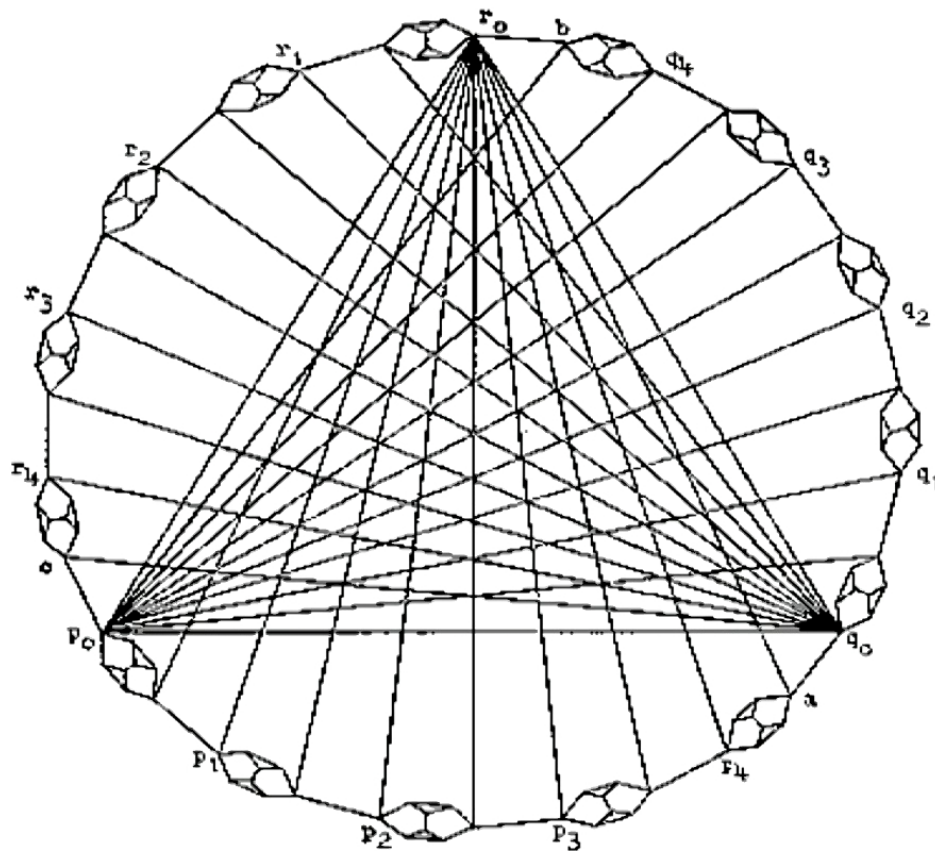
John S. Bell



Ernst Specker (with son) and Simon Kochen

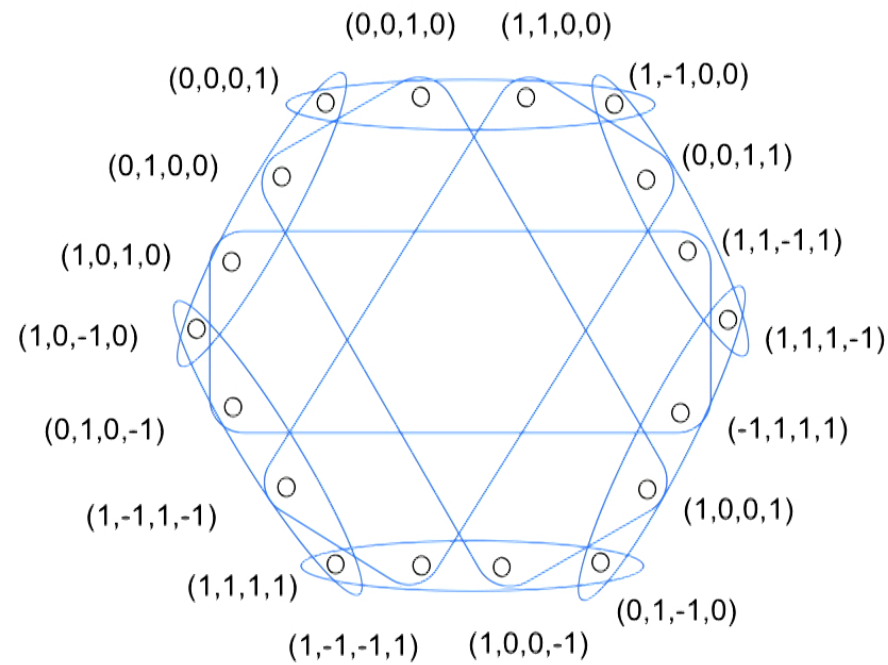
Bell-Kochen-Specker theorem: A traditional noncontextual hidden variable model of quantum theory for Hilbert spaces of dimension 3 or greater is impossible.

Example: Kochen and Specker's original 117 ray proof in 3d

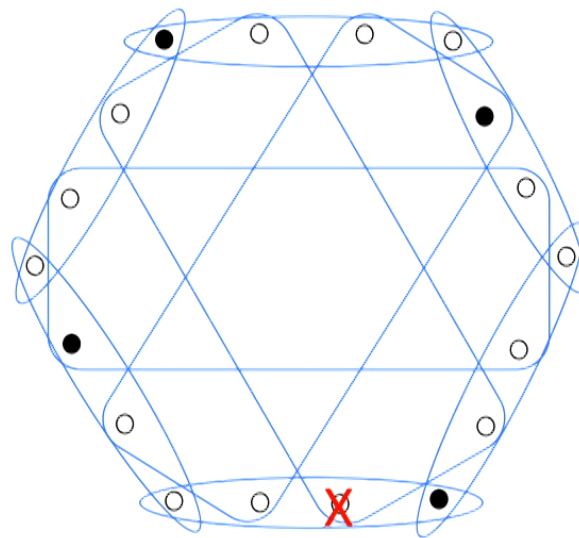


18 ray proof in 4d

Cabello, Estebaranz, Garcia-Alcaine, Phys. Lett. A 212, 183 (1996)



No traditional noncontextual assignments



○ : value 0
● : value 1

If we list all 9 orthogonal quadruples, each ray appears twice in the list

0,0,0,1	0,0,0,1	1,-1,1,-1	1,-1,1,-1	0,0,1,0	1,-1,-1,1	1,1,-1,1	1,1,-1,1	1,1,1,-1
0,0,1,0	0,1,0,0	1,-1,-1,1	1,1,1,1	0,1,0,0	1,1,1,1	1,1,1,-1	-1,1,1,1	-1,1,1,1
1,1,0,0	1,0,1,0	1,1,0,0	1,0,-1,0	1,0,0,1	1,0,0,-1	1,-1,0,0	1,0,1,0	1,0,0,1
1,-1,0,0	1,0,-1,0	0,0,1,1	0,1,0,-1	1,0,0,-1	0,1,-1,0	0,0,1,1	0,1,0,-1	0,1,-1,0

If we list all 9 orthogonal quadruples, each ray appears twice in the list

0,0,0,1	0,0,0,1	1,-1,1,-1	1,-1,1,-1	0,0,1,0	1,-1,-1,1	1,1,-1,1	1,1,-1,1	1,1,1,-1
0,0,1,0	0,1,0,0	1,-1,-1,1	1,1,1,1	0,1,0,0	1,1,1,1	1,1,1,-1	-1,1,1,1	-1,1,1,1
1,1,0,0	1,0,1,0	1,1,0,0	1,0,-1,0	1,0,0,1	1,0,0,-1	1,-1,0,0	1,0,1,0	1,0,0,1
1,-1,0,0	1,0,-1,0	0,0,1,1	0,1,0,-1	1,0,0,-1	0,1,-1,0	0,0,1,1	0,1,0,-1	0,1,-1,0

In each of the 9 quadruples, one ray is assigned 1, the other three 0
Therefore, 9 rays must be assigned 1

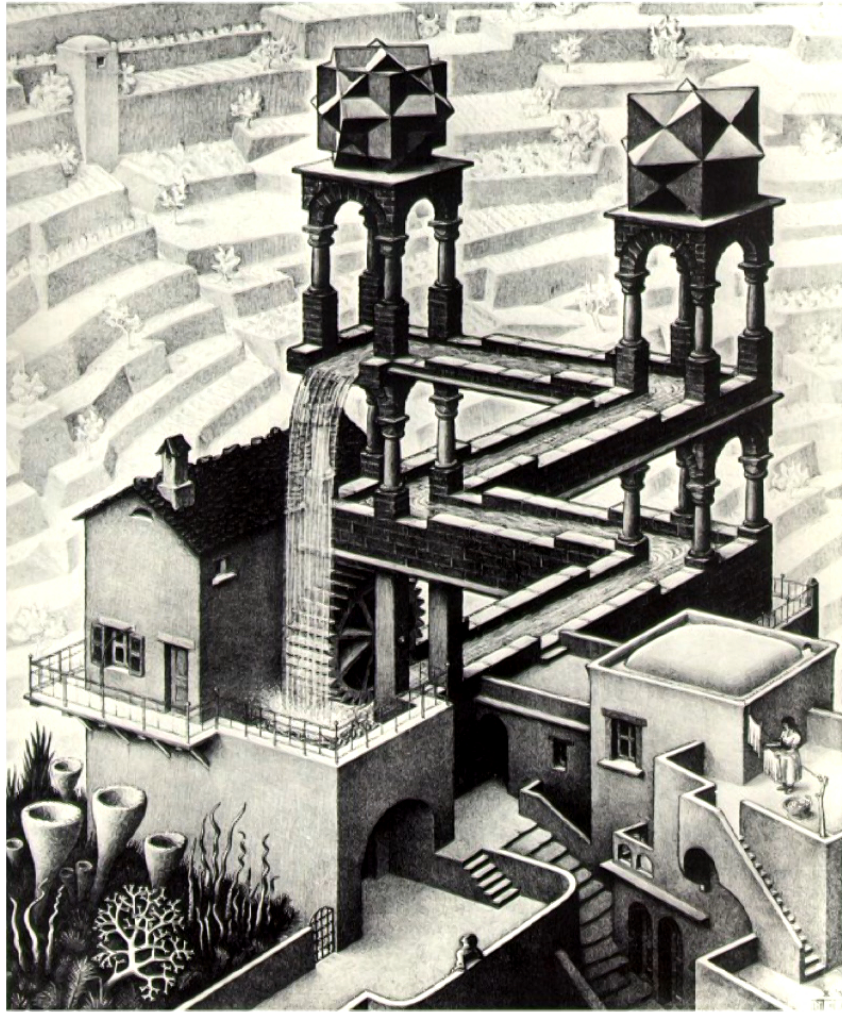
If we list all 9 orthogonal quadruples, each ray appears twice in the list

0,0,0,1	0,0,0,1	1,-1,1,-1	1,-1,1,-1	0,0,1,0	1,-1,-1,1	1,1,-1,1	1,1,-1,1	1,1,1,-1
0,0,1,0	0,1,0,0	1,-1,-1,1	1,1,1,1	0,1,0,0	1,1,1,1	1,1,1,-1	-1,1,1,1	-1,1,1,1
1,1,0,0	1,0,1,0	1,1,0,0	1,0,-1,0	1,0,0,1	1,0,0,-1	1,-1,0,0	1,0,1,0	1,0,0,1
1,-1,0,0	1,0,-1,0	0,0,1,1	0,1,0,-1	1,0,0,-1	0,1,-1,0	0,0,1,1	0,1,0,-1	0,1,-1,0

In each of the 9 quadruples, one ray is assigned 1, the other three 0
Therefore, 9 rays must be assigned 1

But each ray appears twice and so there must be an even number
of rays assigned 1

CONTRADICTION!



The traditional notion of noncontextuality (take 3):

For every λ , every projector Π is assigned a value 0 or 1 regardless of which basis it is a coarse-graining of (i.e. **the context**)

$$v(\Pi) = 0 \text{ or } 1 \quad \text{for all } \Pi$$

Coarse-graining of a measurement implies a coarse-graining of the value (because it is just post-processing)

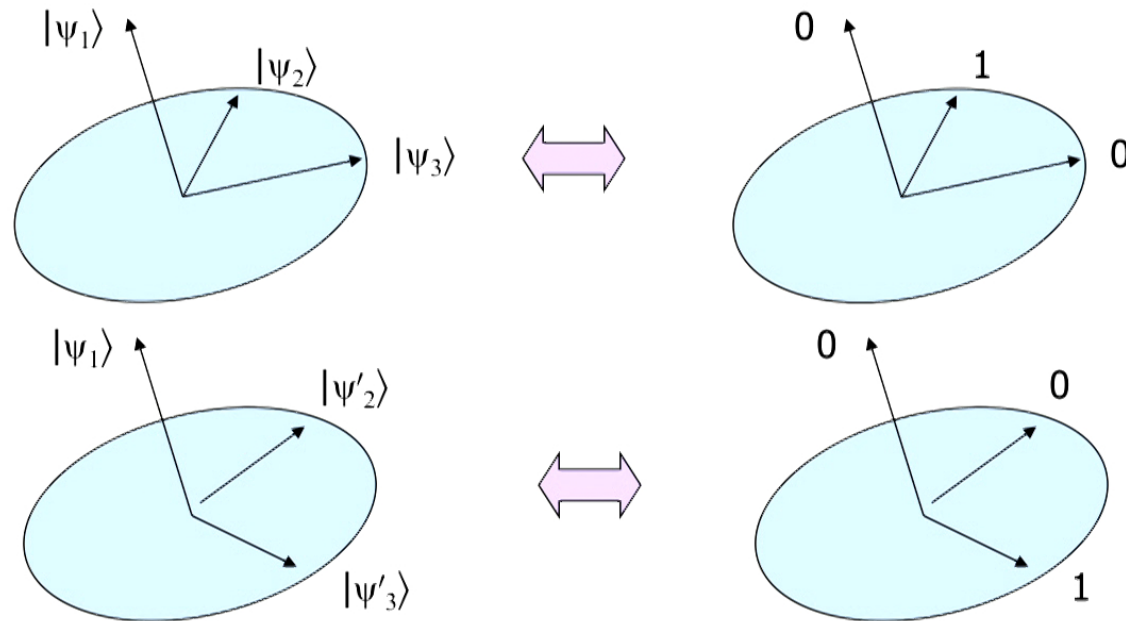
$$v(\sum_k \Pi_k) = \sum_k v(\Pi_k)$$

Every measurement has *some* outcome

$$v(I) = 1$$

The **traditional notion of noncontextuality** (take 2):

For every λ , every basis of vectors receives a 0-1 valuation, wherein exactly one element is assigned the value 1 (corresponding to the outcome that would occur for λ), and every vector is assigned the same value regardless of which basis it is considered a part (i.e. **the context**).



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Every measurement has *some* outcome

$$v(I) = 1$$

The traditional notion of noncontextuality (take 4):

For Hermitian operators A, B, C satisfying

$$[A, B] = 0 \quad [A, C] = 0 \quad [B, C] \neq 0$$

the value assigned to A should be independent of whether it is measured together with B or together with C (i.e. **the context**)

Measure A = measure projectors onto eigenspaces of A, $\{\Pi_a\}$

$$A = \sum_a a \Pi_a \rightarrow v(A) = \sum_a a v(\Pi_a)$$

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$$A = \sum_a a \Pi_a \rightarrow v(A) = \sum_a a v(\Pi_a)$$

Measure A in context of B

= measure projectors onto joint eigenspaces of A and B, $\{\Pi_{ab}\}$

then coarse-grain over B outcome $\Pi_a = \sum_b \Pi_{ab}$

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= measure projectors onto joint eigenspaces of A and B, $\{\Pi_{ab}\}$

then coarse-grain over B outcome $\Pi_a = \sum_b \Pi_{ab}$

Measure A in context of C

= measure projectors onto joint eigenspaces of A and C, $\{\Pi_{ac}\}$

Then coarse-grain over C outcome $\Pi_a = \sum_c \Pi_{ac}$

The traditional notion of noncontextuality (take 4):

For Hermitian operators A, B, C satisfying

$$[A, B] = 0 \quad [A, C] = 0 \quad [B, C] \neq 0$$

the value assigned to A should be independent of whether it is measured together with B or together with C (i.e. **the context**)

Measure A = measure projectors onto eigenspaces of A, $\{\Pi_a\}$

$$A = \sum_a a \Pi_a \rightarrow v(A) = \sum_a a v(\Pi_a)$$

Measure A in context of B

= measure projectors onto joint eigenspaces of A and B, $\{\Pi_{ab}\}$

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Measure A in context of C

= measure projectors onto joint eigenspaces of A and C, $\{\Pi_{ac}\}$

Then coarse-grain over C outcome $\Pi_a = \sum_c \Pi_{ac}$

$v(\Pi_a)$ is independent of context $\rightarrow v(A)$ is independent of context

Functional relationships among commuting Hermitian operators must be respected by their values

$$\begin{array}{l} \text{If } f(L, M, N, \dots) = 0 \\ \text{then } f(v(L), v(M), v(N), \dots) = 0 \end{array}$$

Example: Mermin's magic square proof in 4d

X_1	X_2	X_1X_2
Y_2	Y_1	Y_1Y_2
X_1Y_2	Y_1X_2	Z_1Z_2

I

I

I

I

I

$-I$

$$X_1 X_2 (X_1X_2) = I$$

$$Y_1 Y_2 (Y_1Y_2) = I$$

$$(X_1Y_2) (Y_1X_2) (Z_1Z_2) = I$$

$$X_1 Y_2 (X_1Y_2) = I$$

$$Y_1 X_2 (Y_1X_2) = I$$

$$(X_1X_2) (Y_1Y_2) (Z_1Z_2) = -I$$

Example: Mermin's magic square proof in 4d

X_1	X_2	X_1X_2
Y_2	Y_1	Y_1Y_2
X_1Y_2	Y_1X_2	Z_1Z_2

I

I

I

I

I

$-I$

$$X_1 X_2 (X_1X_2) = I$$

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$$(X_1Y_2) (Y_1X_2) (Z_1Z_2) = I$$

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$$Y_1 X_2 (Y_1X_2) = I$$

$$(X_1X_2) (Y_1Y_2) (Z_1Z_2) = -I$$

$$v(X_1) v(X_2) v(X_1X_2) = 1$$

$$v(Y_1) v(Y_2) v(Y_1Y_2) = 1$$

$$v(X_1Y_2) v(Y_1X_2) v(Z_1Z_2) = 1$$

$$v(X_1) v(Y_2) v(X_1Y_2) = 1$$

$$v(Y_1) v(X_2) v(Y_1X_2) = 1$$

$$v(X_1X_2) v(Y_1Y_2) v(Z_1Z_2) = -1$$

Example: Mermin's magic square proof in 4d

X_1	X_2	X_1X_2	I
Y_2	Y_1	Y_1Y_2	I
X_1Y_2	Y_1X_2	Z_1Z_2	I
I	I	$-I$	

$$v(X_1) v(X_2) v(X_1X_2) = 1$$

$$v(Y_1) v(Y_2) v(Y_1Y_2) = 1$$

$$v(X_1Y_2) v(Y_1X_2) v(Z_1Z_2) = 1$$

$$v(X_1) v(Y_2) v(X_1Y_2) = 1$$

$$v(Y_1) v(X_2) v(Y_1X_2) = 1$$

$$v(X_1X_2) v(Y_1Y_2) v(Z_1Z_2) = -1$$

$$X_1 X_2 (X_1X_2) = I$$

$$Y_1 Y_2 (Y_1Y_2) = I$$

$$(X_1Y_2) (Y_1X_2) (Z_1Z_2) = I$$

$$X_1 Y_2 (X_1Y_2) = I$$

$$Y_1 X_2 (Y_1X_2) = I$$

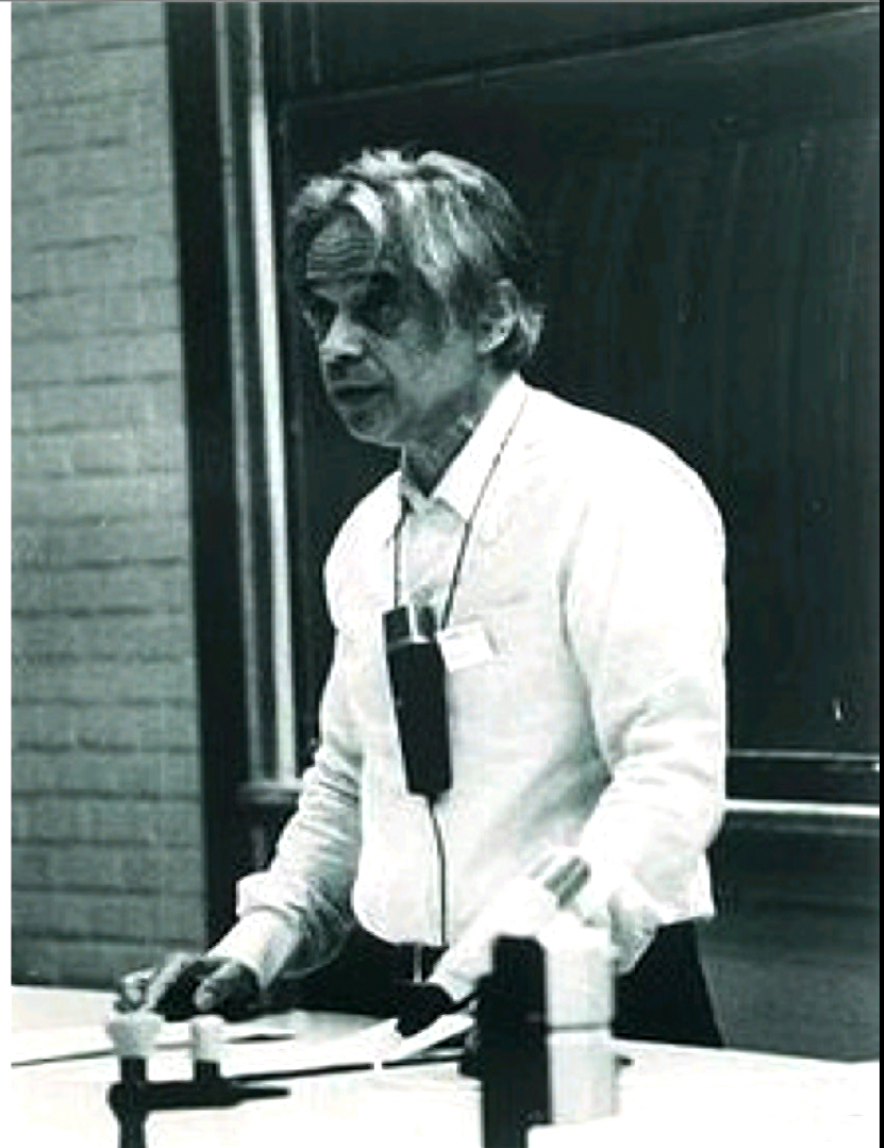
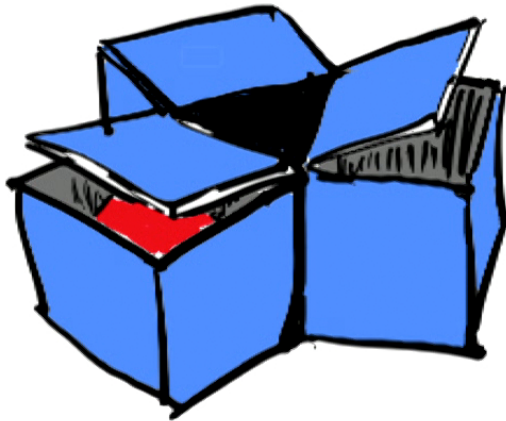
$$(X_1X_2) (Y_1Y_2) (Z_1Z_2) = -I$$

Product of LHSs = +1

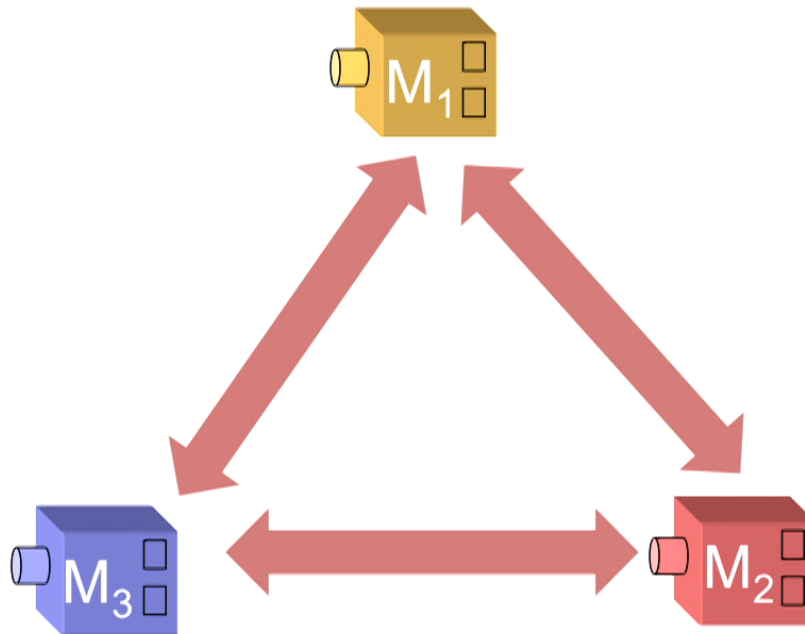
Product of RHSs = -1

CONTRADICTION

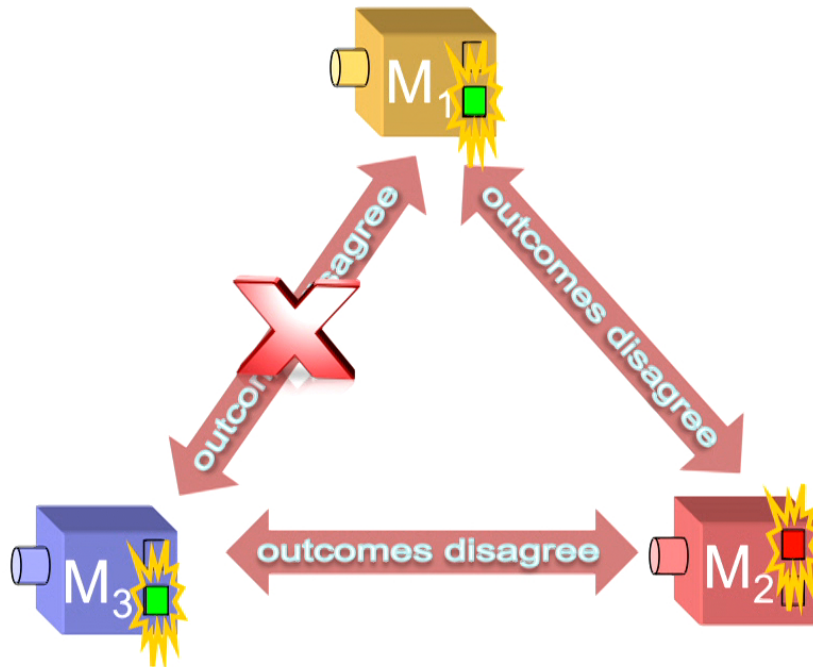
Ernst Specker, "The logic of propositions which are not simultaneously decidable", *Dialectica* 14, 239 (1960).



Specker's example



Specker's example



If the outcomes are fixed **deterministically** by the ontic state and are **independent of the context** in which the measurement is performed, then

$$p(\text{success}) \leq \frac{2}{3}$$

Frustrated Networks

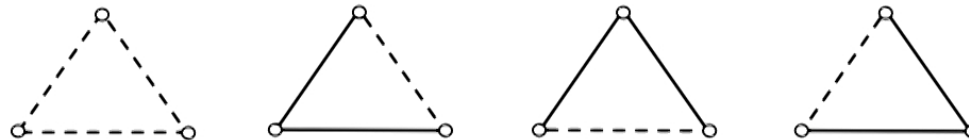
Nodes are binary variables

Edges imply joint measurability

○——○ Outcomes agree

○-----○ Outcomes disagree

Frustration = no valuation satisfying all correlations



Frustrated Networks

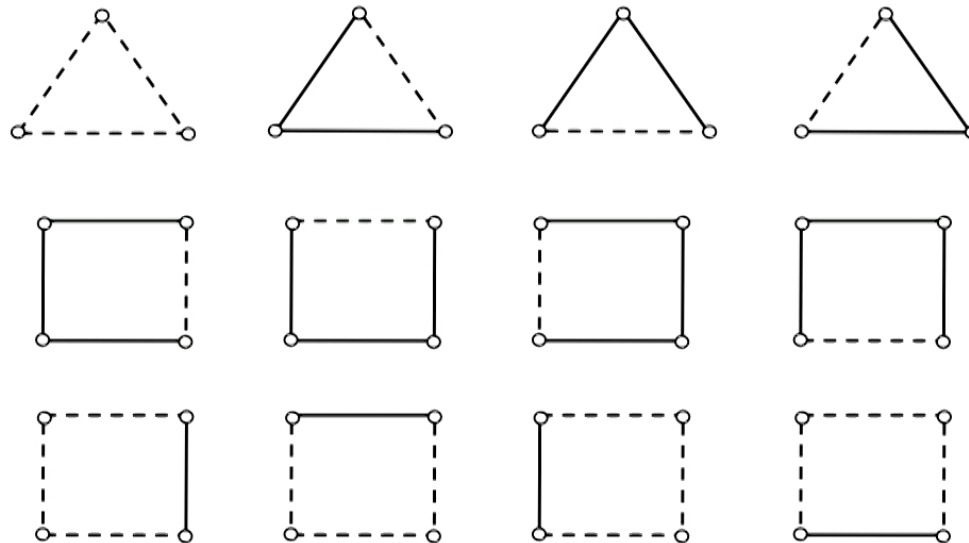
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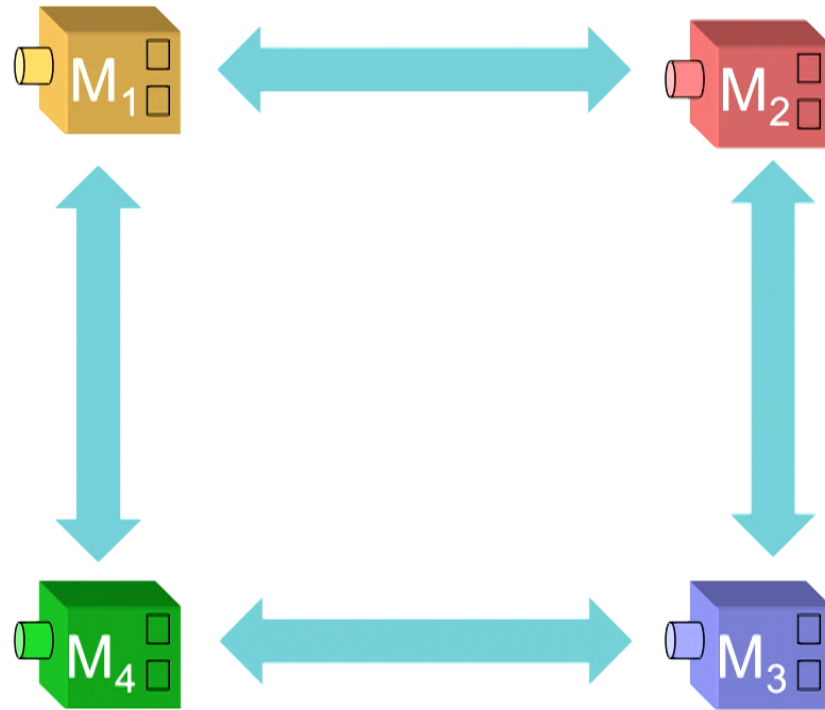
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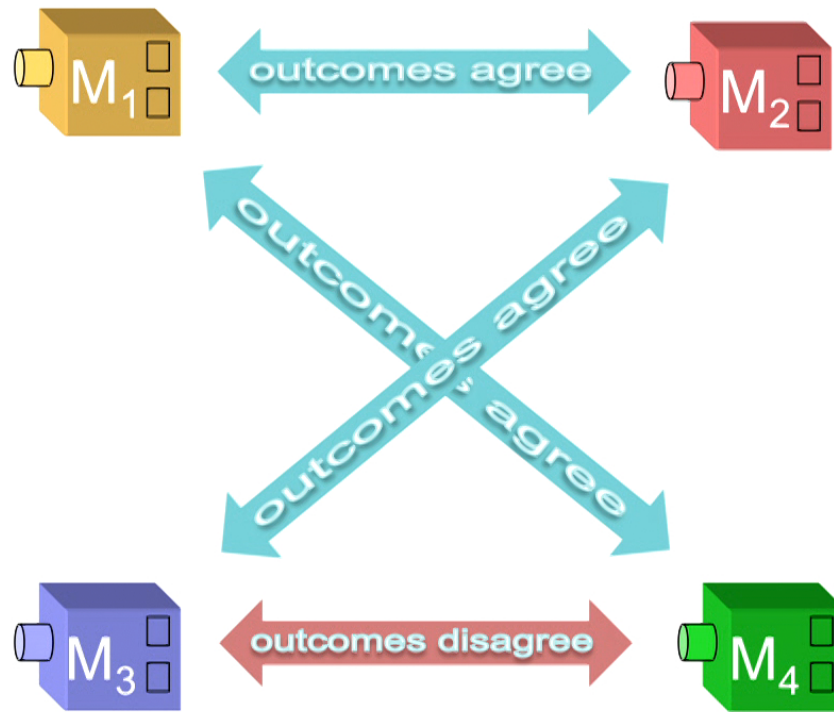
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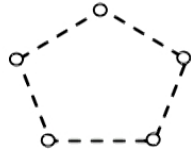
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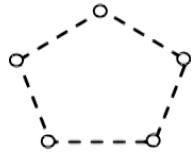
Klyachko, Can, Biniciolu, Shumovsky, PRL 101, 020403 (2008)



In a traditional
noncontextual model

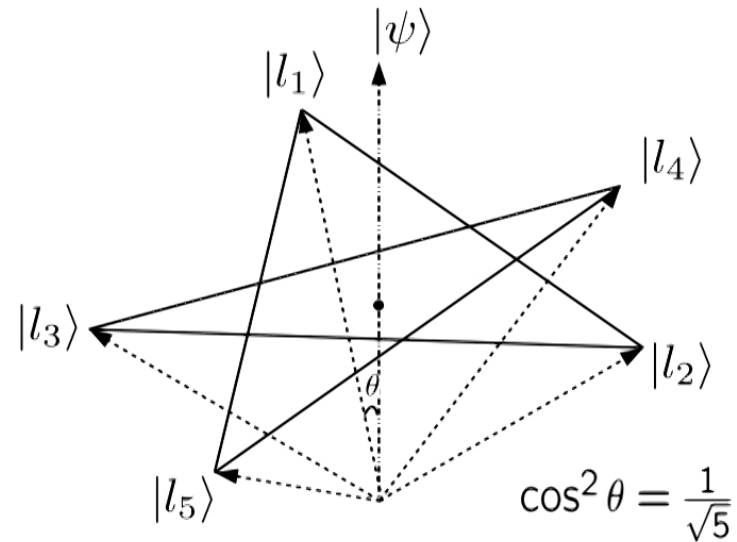
$$p(\text{success}) \leq \frac{4}{5}$$

Klyachko, Can, Biniciolu, Shumovsky, PRL 101, 020403 (2008)



In a traditional
noncontextual model

$$p(\text{success}) \leq \frac{4}{5}$$



Quantum probability of success

$$p(\text{success}) = \frac{2}{\sqrt{5}} \simeq 0.89 > \frac{4}{5}$$