Title: PSI 17/18 - Foundations of Quantum Mechanics - Lecture 12

Date: Feb 13, 2018 10:15 AM

URL: http://pirsa.org/18020071

Abstract:

Bell's theorem



John S. Bell (1928-1990)

Pirsa: 18020071 Page 2/46

A pair of two-outcome measurements

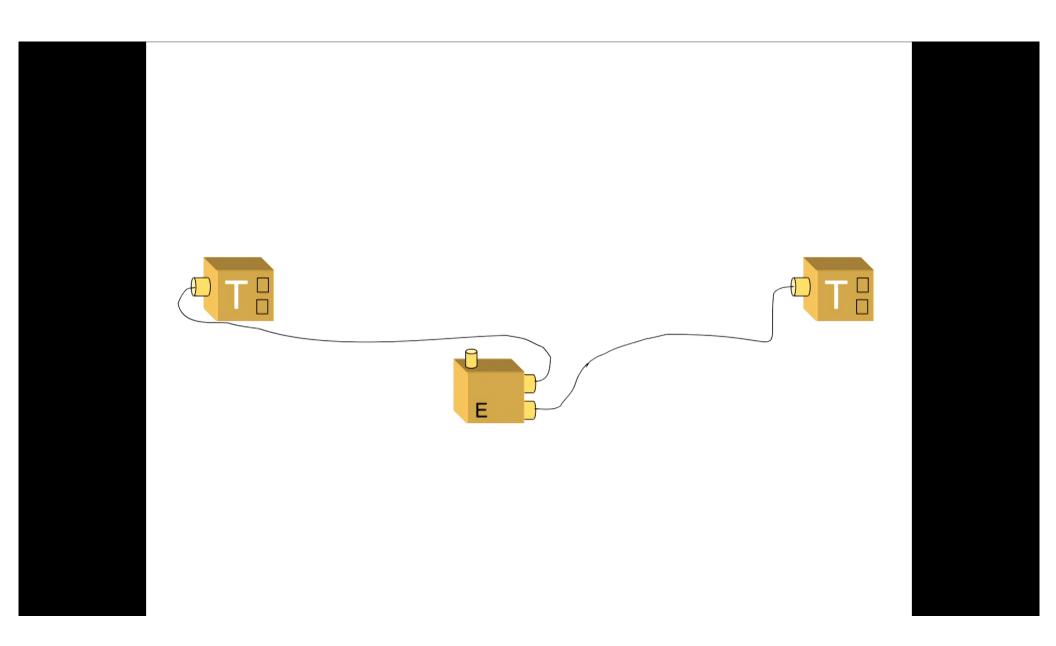








Pirsa: 18020071 Page 3/46



Pirsa: 18020071 Page 4/46

There are two possible measurements, H and T, with two outcomes each: green or red

Suppose which of H or T occurs at each wing is chosen at random

Scenario 1

1. Whenever the same measurement is made on A or and B, the outcomes always agree

H and H or T and T

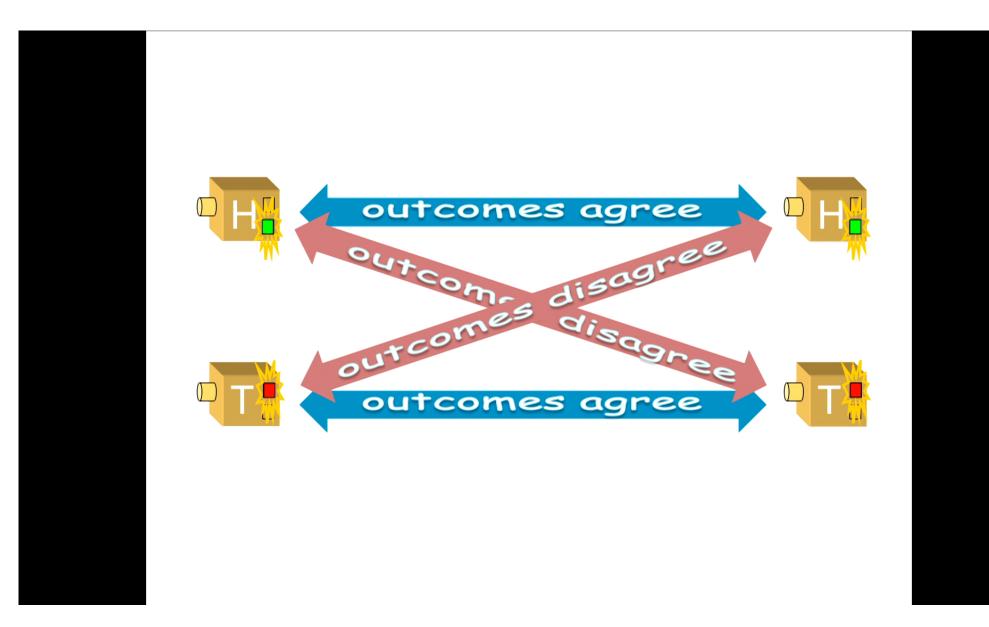
2. Whenever different measurements are made on A and B, the outcomes always disagree

H and T

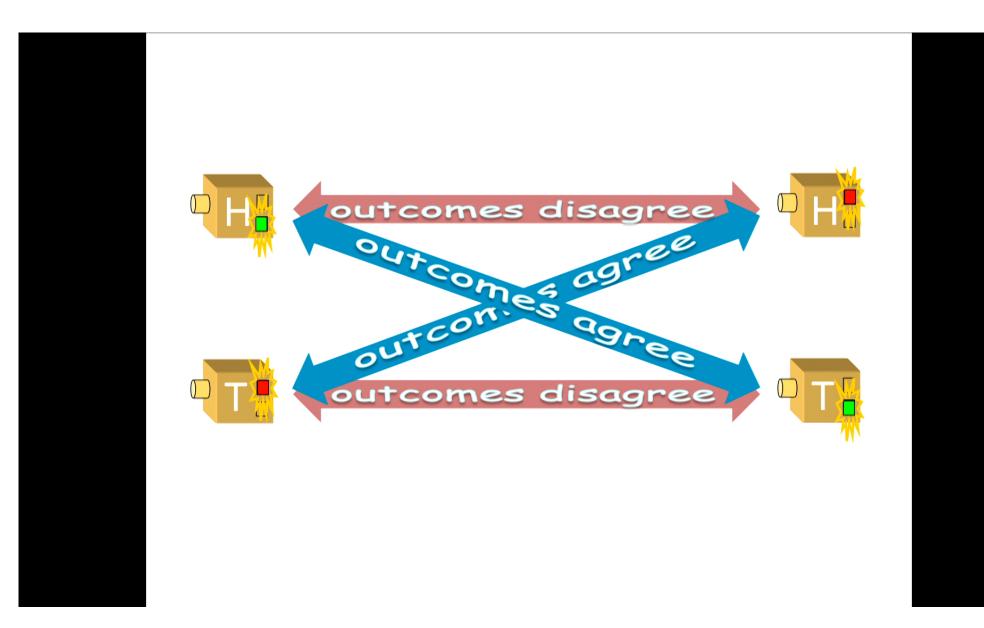
or

T and H

Pirsa: 18020071 Page 5/46



Pirsa: 18020071 Page 6/46



Pirsa: 18020071 Page 7/46

There are two possible measurements, H and T, with two outcomes each: green or red

Suppose which of H or T occurs at each wing is chosen at random

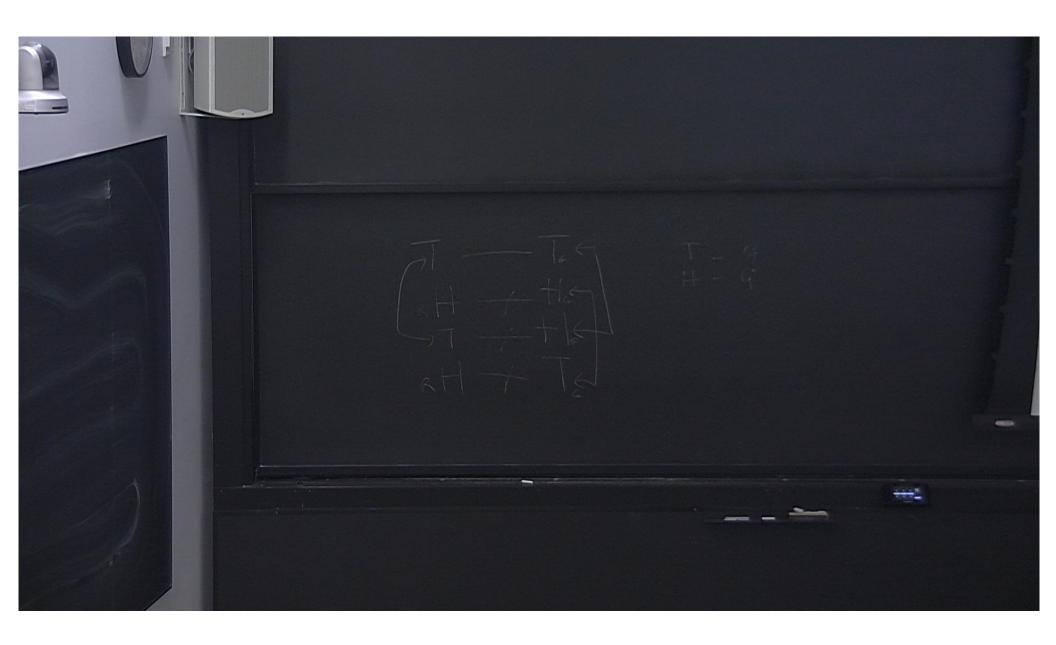
Scenario 3

1. Whenever the measurement T and T T is made on both A and B, the outcomes always disagree

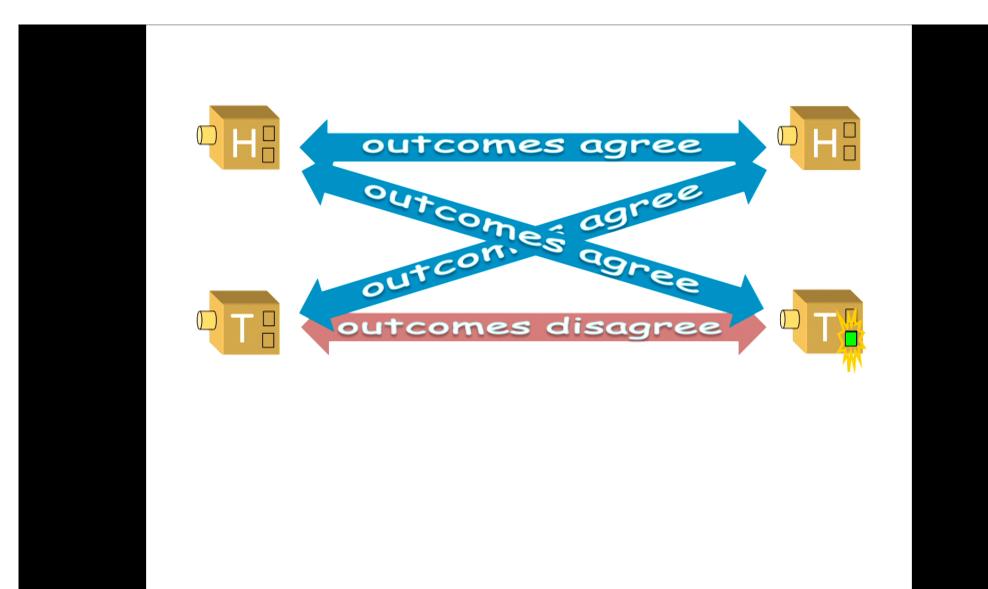
2. Otherwise, the outcomes always agree

Hand H or Hand T or Tand H

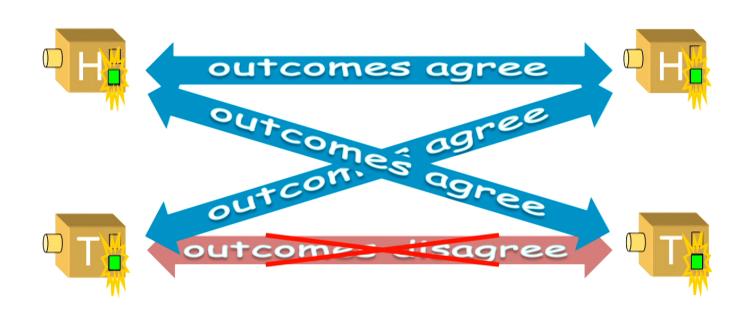
Pirsa: 18020071 Page 8/46



Pirsa: 18020071 Page 9/46

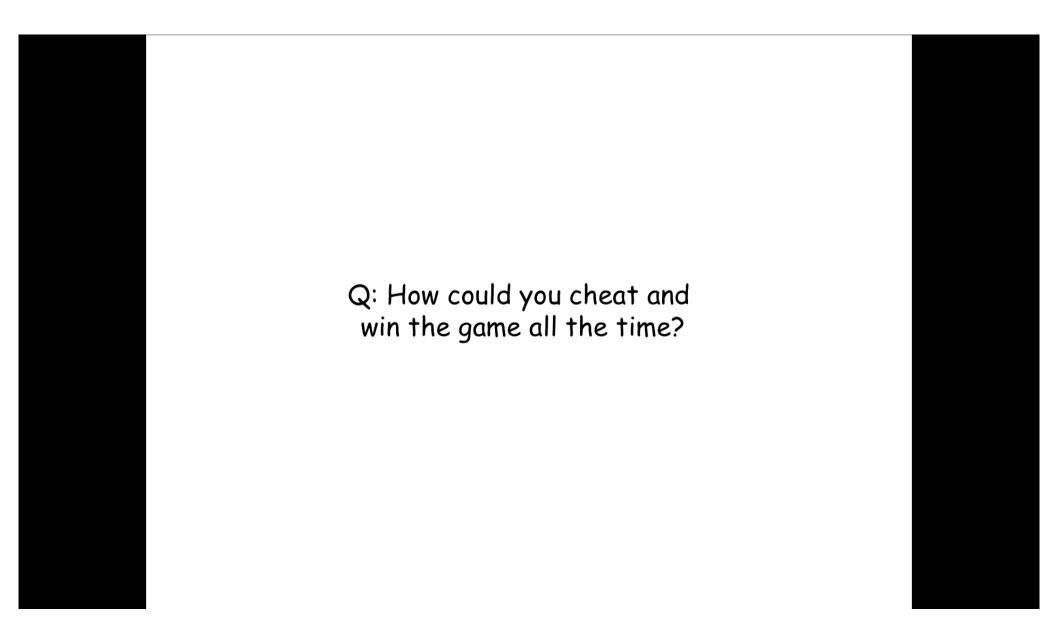


Pirsa: 18020071 Page 10/46

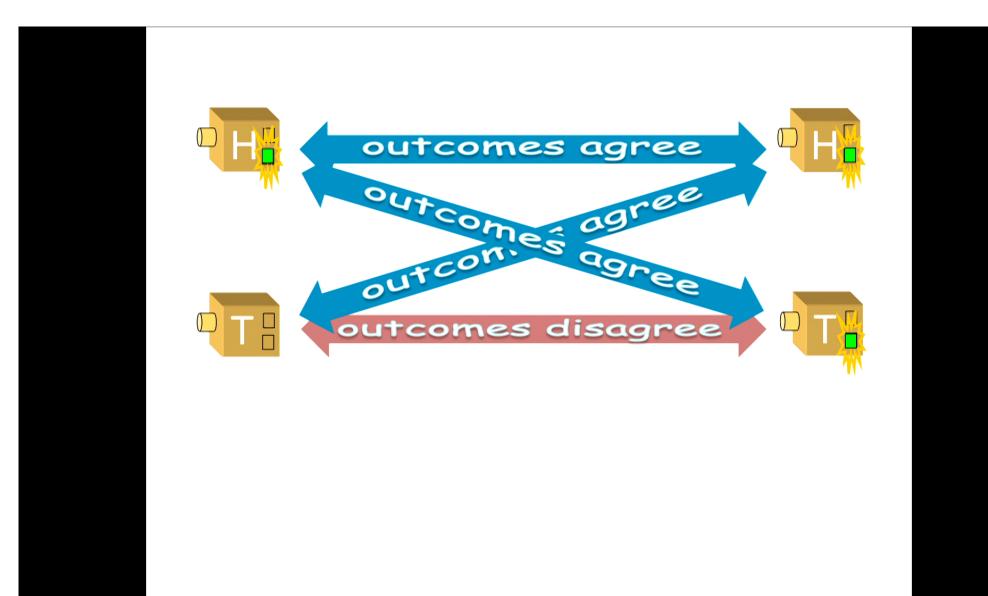


The game can be won at most 75% of the time by local strategies

Pirsa: 18020071 Page 11/46



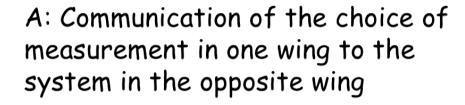
Pirsa: 18020071 Page 12/46



Pirsa: 18020071 Page 13/46

A: Rig the game so that the choices of settings are not random but instead are correlated with the local strategies

Pirsa: 18020071 Page 14/46



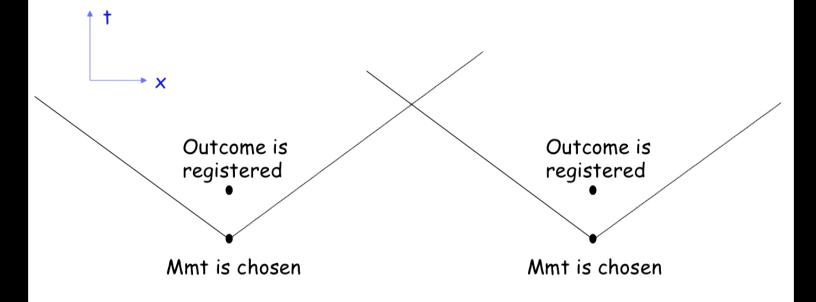
Pirsa: 18020071 Page 15/46

A: Communication of the choice of measurement in one wing to the system in the opposite wing

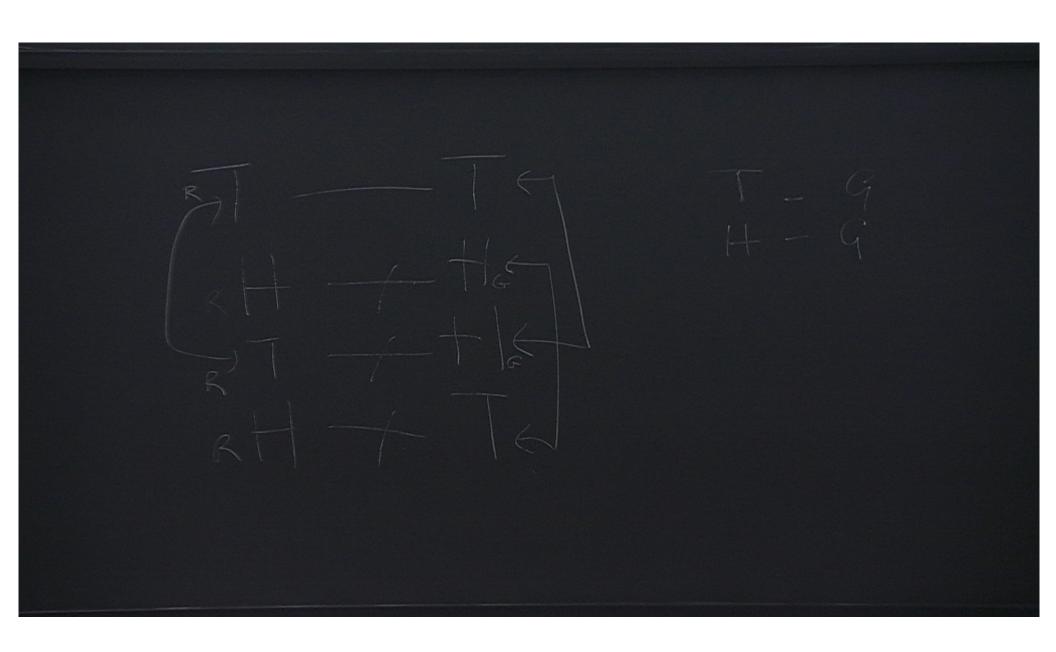
But there's a problem...

Pirsa: 18020071 Page 16/46

Tension with the theory of relativity



Pirsa: 18020071 Page 17/46



Pirsa: 18020071 Page 18/46

If the particles have access to randomness when deciding on their strategy, can it help them to generate the correlations?

No. The degree of correlation can only be the same or less.

If the particles have access to *local* randomness (i.e. independent randomness at the two wings), can it help them to generate the correlations?

No. The degree of correlation can only be the same or less.

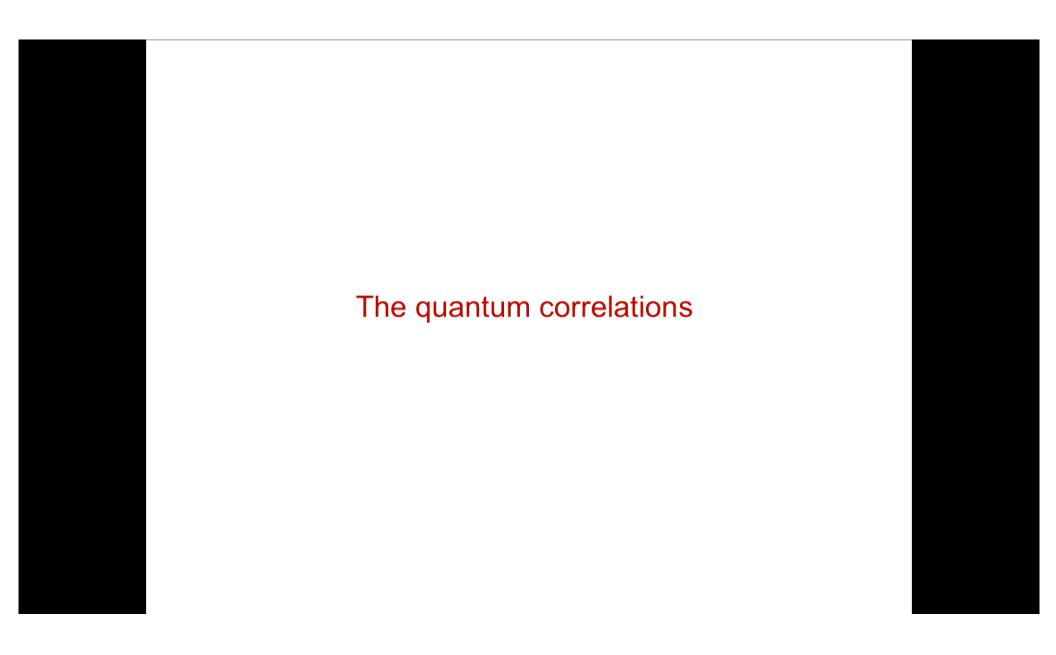
If the detector inefficiencies are sufficiently high (sometimes there is no outcome), can particles obeying local causality simulate the correlations on the detected pairs?

Yes. This is the detector loophole.

Is there a problem if the choice of measurement is made too early?

Yes. This is the locality loophole.

Pirsa: 18020071 Page 19/46



Pirsa: 18020071 Page 20/46

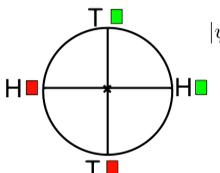
$$p(\text{success}) = \frac{1}{4} [p(\text{agree}|HH) + p(\text{agree}|HT) + p(\text{agree}|TH) + p(\text{disagree}|TT)]$$

Realist theories that are locally causal predict $p(ext{success}) \leq 0.75$ A Bell Inequality

Quantum theory predicts that one can achieve $p(\text{success}) \simeq 0.85$

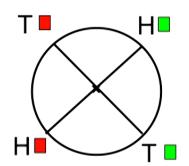
Pirsa: 18020071 Page 21/46

The Bell-inequality violation in quantum theory



$$|\psi\rangle_{AB} = rac{1}{\sqrt{2}}(|0\rangle_A|0\rangle_B + |1\rangle_A|1\rangle_B)$$
 T

$$p(success) = \frac{1}{2} + \frac{1}{2\sqrt{2}}$$
$$\simeq 0.85$$





outcomes agree



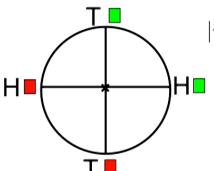


outcomes agree outcomes agree



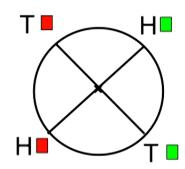
outcomes disagree

The Bell-inequality violation in quantum theory



$$|\psi\rangle_{AB} = \frac{1}{\sqrt{2}}(|0\rangle_A|0\rangle_B + |1\rangle_A|1\rangle_B)$$
 T

$$p(success) = \frac{1}{2} + \frac{1}{2\sqrt{2}}$$
$$\approx 0.85$$

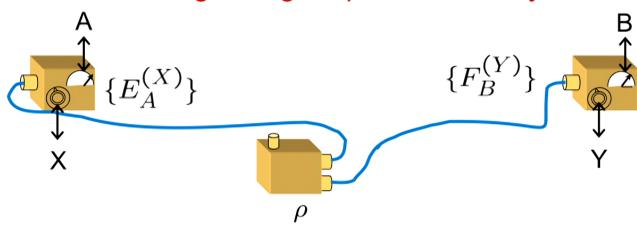


$$_{A}\langle +\hat{n}|\psi\rangle_{AB} = \left[\cos(\theta/2)_{A}\langle 0| + \sin(\theta/2)_{A}\langle 1|\right] \frac{1}{\sqrt{2}} (|0\rangle_{A}|0\rangle_{B} + |1\rangle_{A}|1\rangle_{B})$$
$$= \cos(\theta/2)|0\rangle_{B} + \sin(\theta/2)|1\rangle_{B}$$
$$= |+\hat{n}\rangle_{B}$$

$$|\langle +\hat{n}|_A \langle +\hat{m}|_B |\psi\rangle_{AB}|^2 = |\langle +\hat{m}| +\hat{n}\rangle|^2 = \cos^2(\theta/2)$$

$$p(\text{agree}|HH) = p(\text{agree}|HT) = p(\text{agree}|TH) = p(\text{disagree}|TT)$$
$$= \cos^2(\pi/8) = \frac{1}{2} + \frac{1}{2\sqrt{2}}$$

No signalling in quantum theory



$$p(B|XY) = \sum_{A} p(AB|XY)$$

$$= \sum_{A} \text{Tr}[(E_A^{(X)} \otimes F_B^{(Y)}) \rho]$$

$$= \text{Tr}[(I \otimes F_B^{(Y)}) \rho]$$

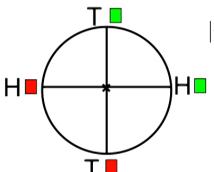
$$= p(B|Y)$$

independent of choice of measurement at left

Note $[E_A^{(X)} \otimes I, I \otimes F_B^{(Y)}] = 0$ because left and right are space-like separated

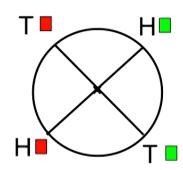
Is the proof robust to imperfection in the state preparation? (the state is mixed rather than pure)

The Bell-inequality violation in quantum theory



$$|\psi\rangle_{AB} = \frac{1}{\sqrt{2}}(|0\rangle_A|0\rangle_B + |1\rangle_A|1\rangle_B)$$
 T

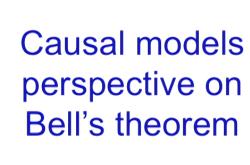
$$p(success) = \frac{1}{2} + \frac{1}{2\sqrt{2}}$$
$$\approx 0.85$$



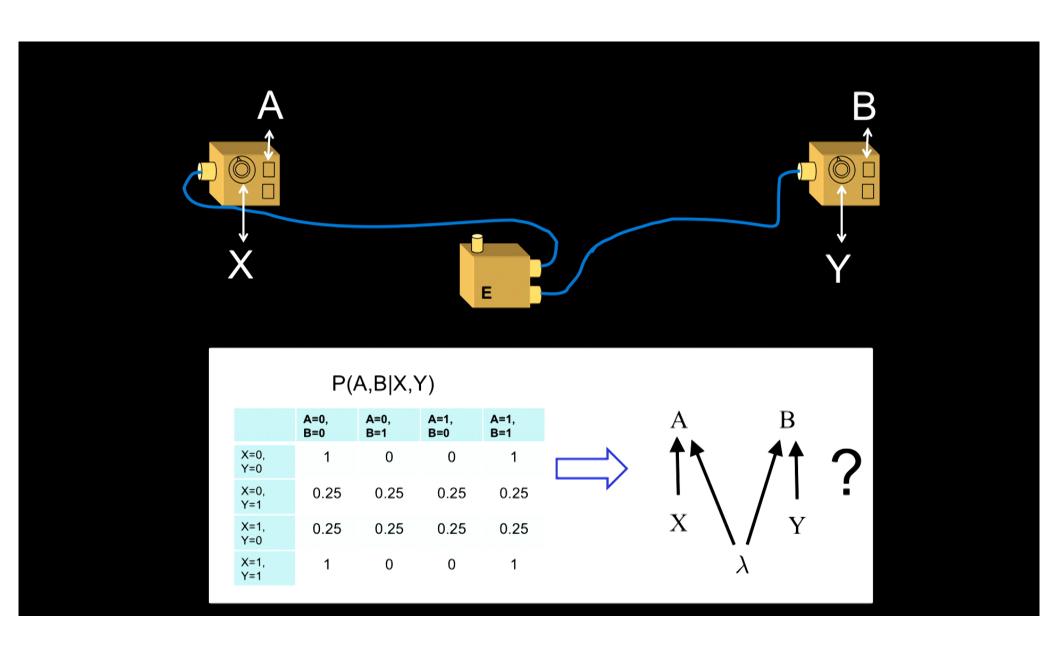
$$A\langle +\hat{n}|\psi\rangle_{AB} = \left[\cos(\theta/2)_A\langle 0| + \sin(\theta/2)_A\langle 1|\right] \frac{1}{\sqrt{2}} (|0\rangle_A|0\rangle_B + |1\rangle_A|1\rangle_B)$$
$$= \cos(\theta/2)|0\rangle_B + \sin(\theta/2)|1\rangle_B$$
$$= |+\hat{n}\rangle_B$$

$$|\langle +\hat{n}|_A \langle +\hat{m}|_B |\psi\rangle_{AB}|^2 = |\langle +\hat{m}| +\hat{n}\rangle|^2 = \cos^2(\theta/2)$$

$$p(\text{agree}|HH) = p(\text{agree}|HT) = p(\text{agree}|TH) = p(\text{disagree}|TT)$$
$$= \cos^2(\pi/8) = \frac{1}{2} + \frac{1}{2\sqrt{2}}$$

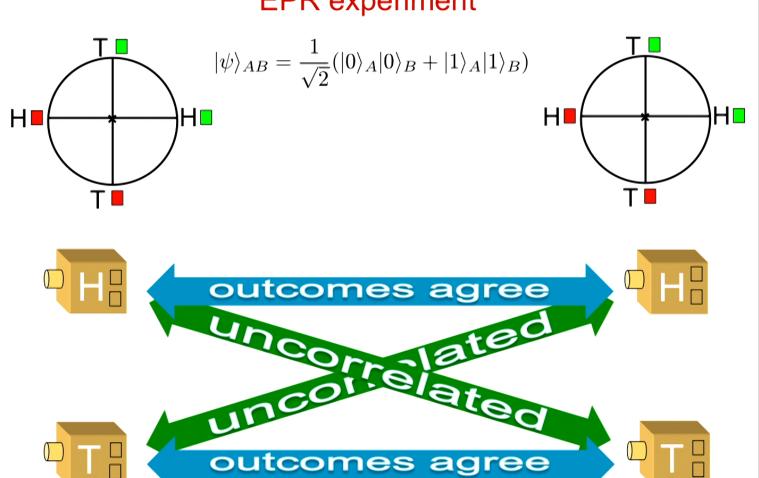


Pirsa: 18020071 Page 27/46

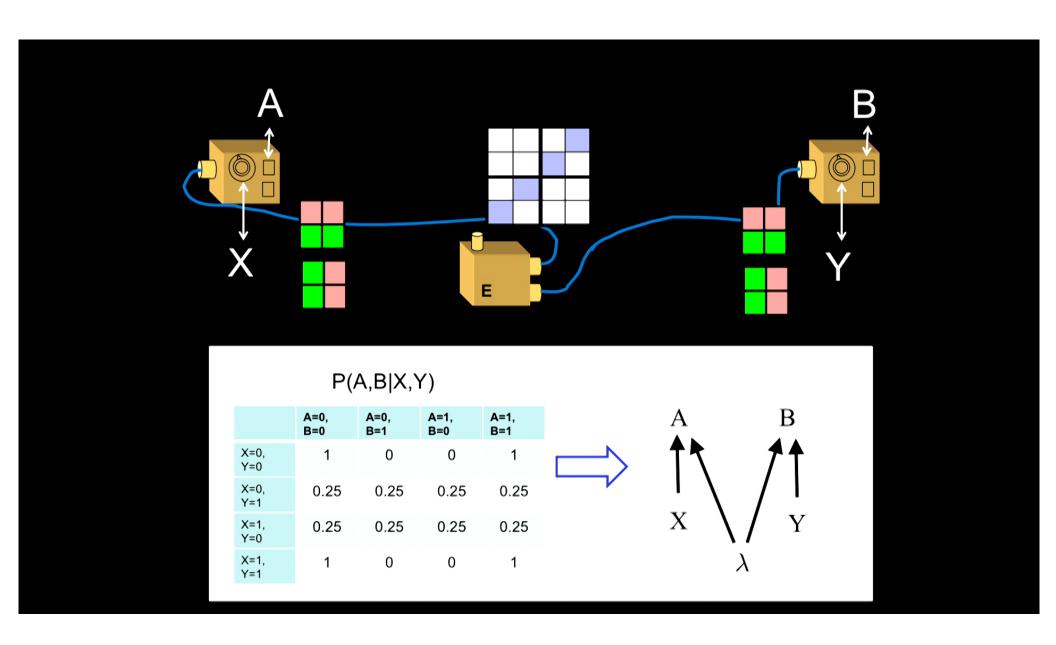


Pirsa: 18020071 Page 28/46

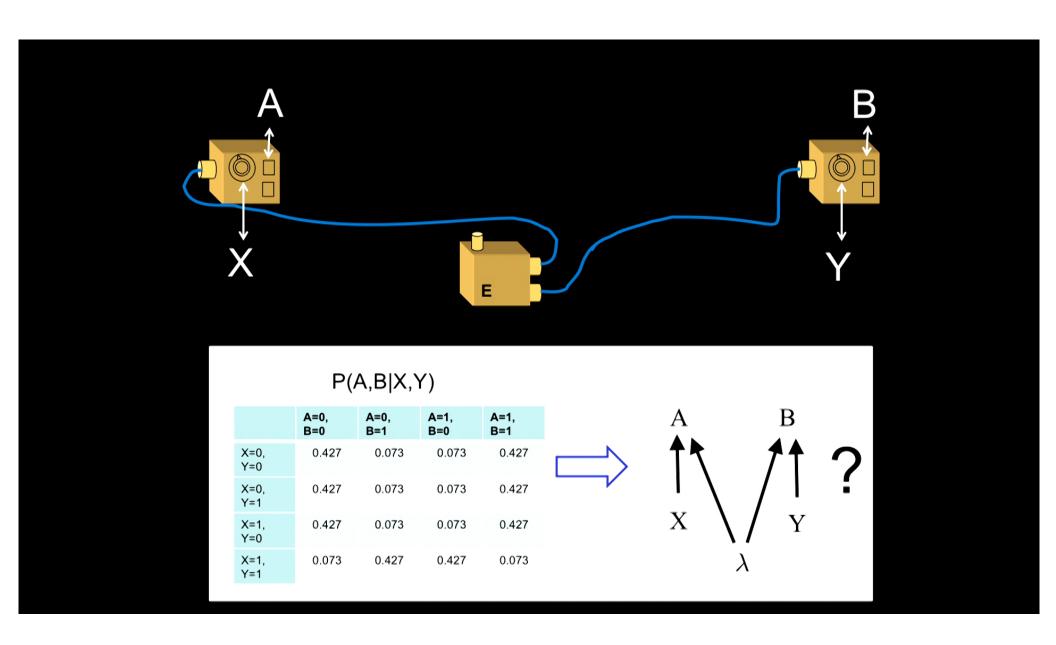
EPR experiment



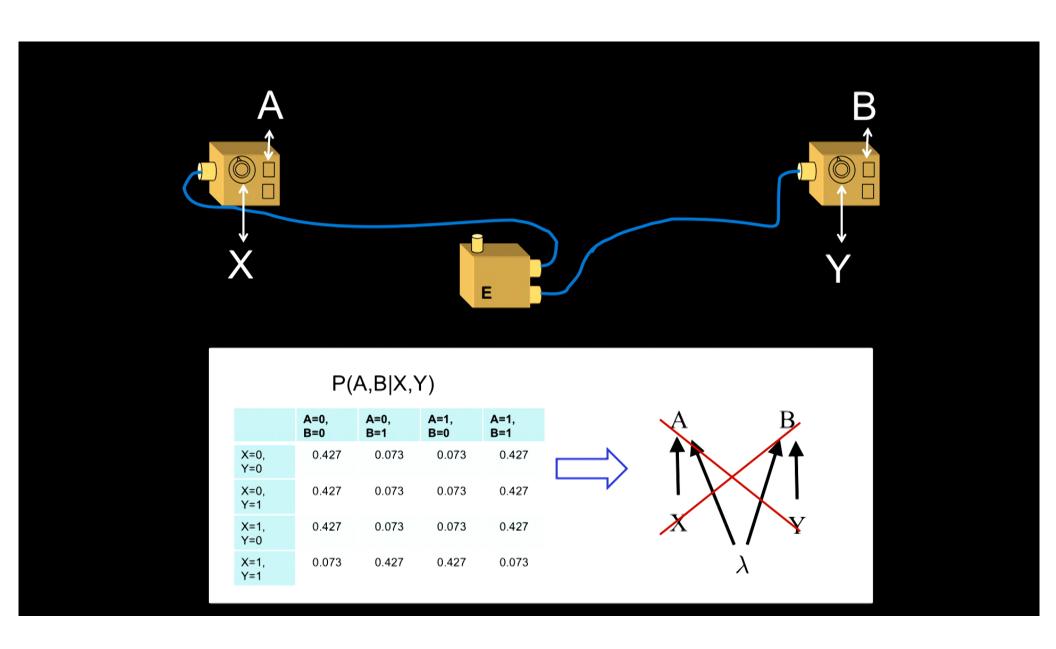
Pirsa: 18020071 Page 29/46



Pirsa: 18020071 Page 30/46



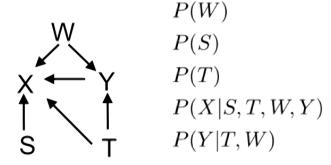
Pirsa: 18020071 Page 31/46



Pirsa: 18020071 Page 32/46

Causal Model

Causal Structure Causal-Statistical Parameters



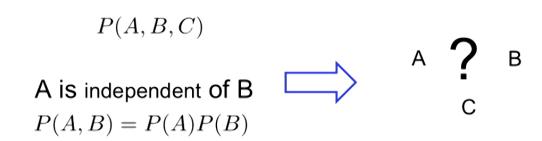
$$P(X, Y, W, S, T) = P(X|S, T, W, Y)P(Y|T, W)P(W)P(S)P(T)$$

Causal inference algorithms seek to solve the inverse problem

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Principle #1 Statistical dependences need to be explained causally

Pirsa: 18020071 Page 34/46



no other independence relations

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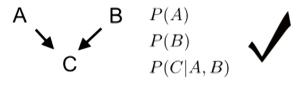


A is independent of B

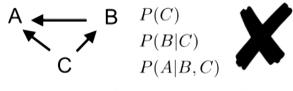


P(A,B) = P(A)P(B)

no other independence relations



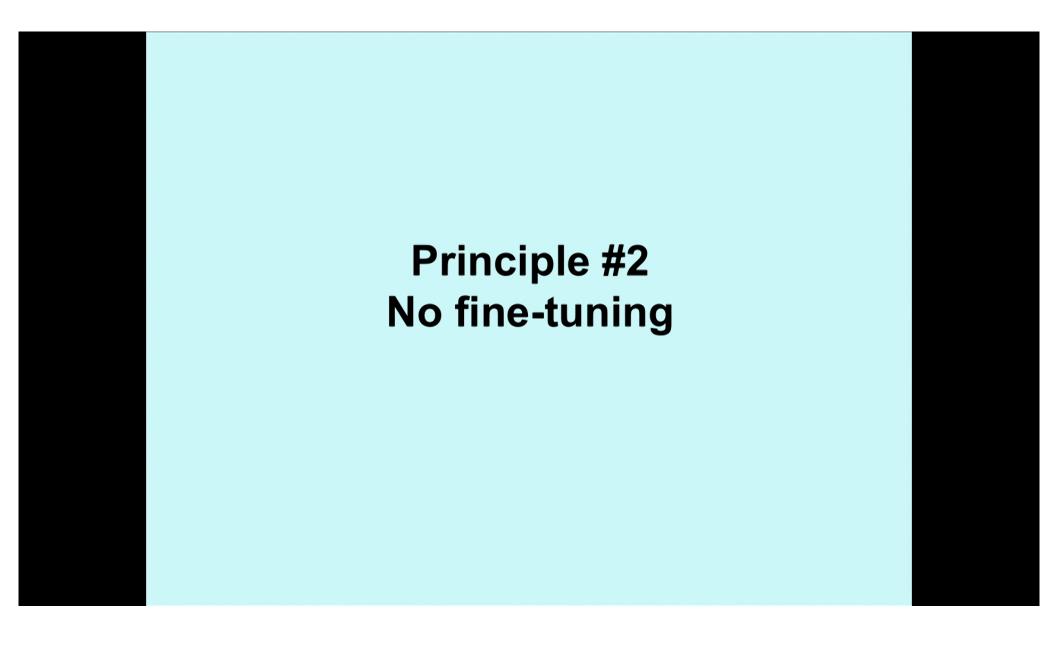
P(A, B, C) = P(C|A, B)P(A)P(B)



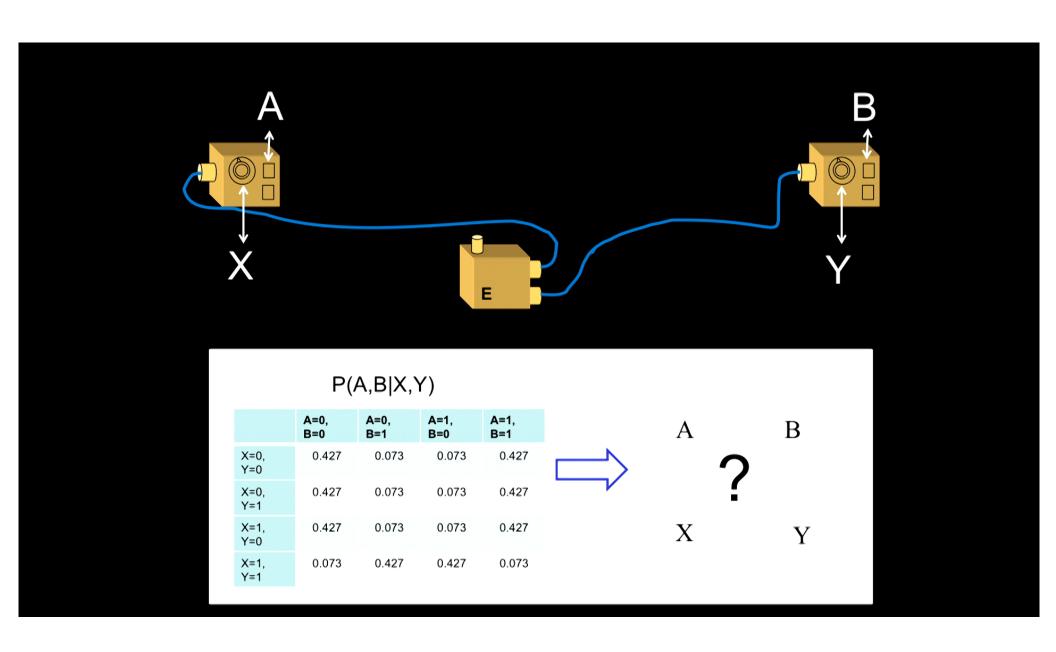
P(A, B, C) = P(A|B, C)P(B|C)P(C)

This model is fine-tuned

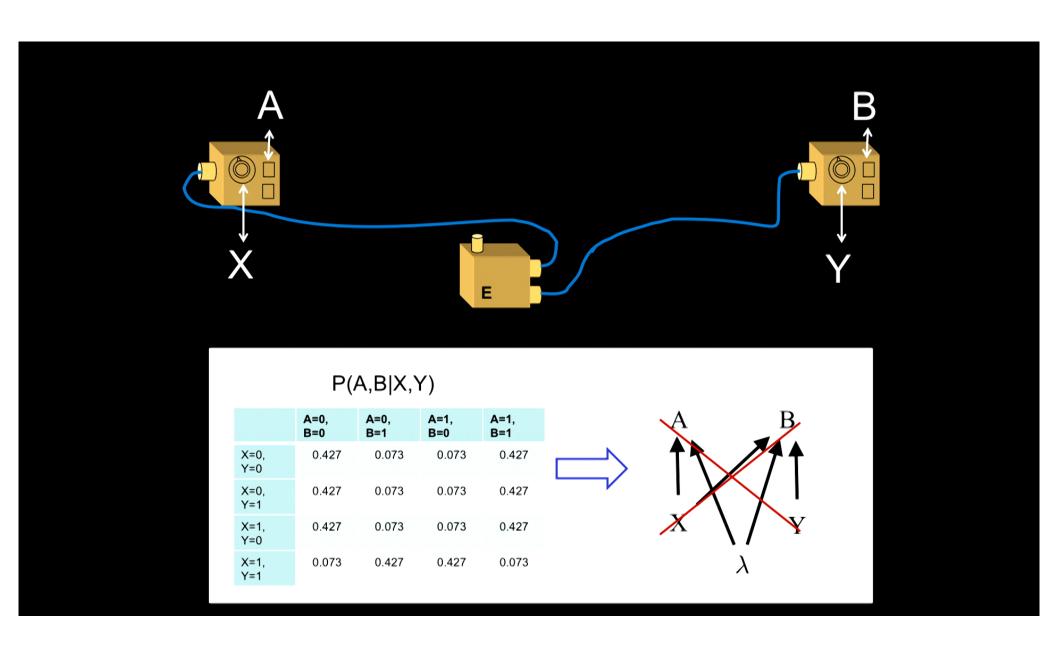
Pirsa: 18020071 Page 36/46

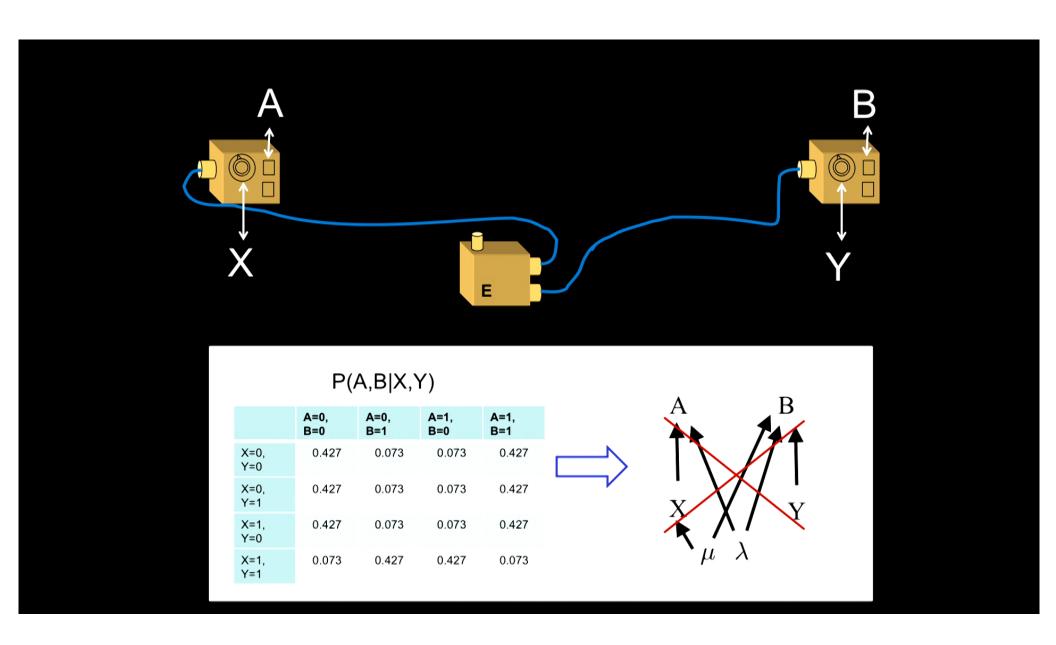


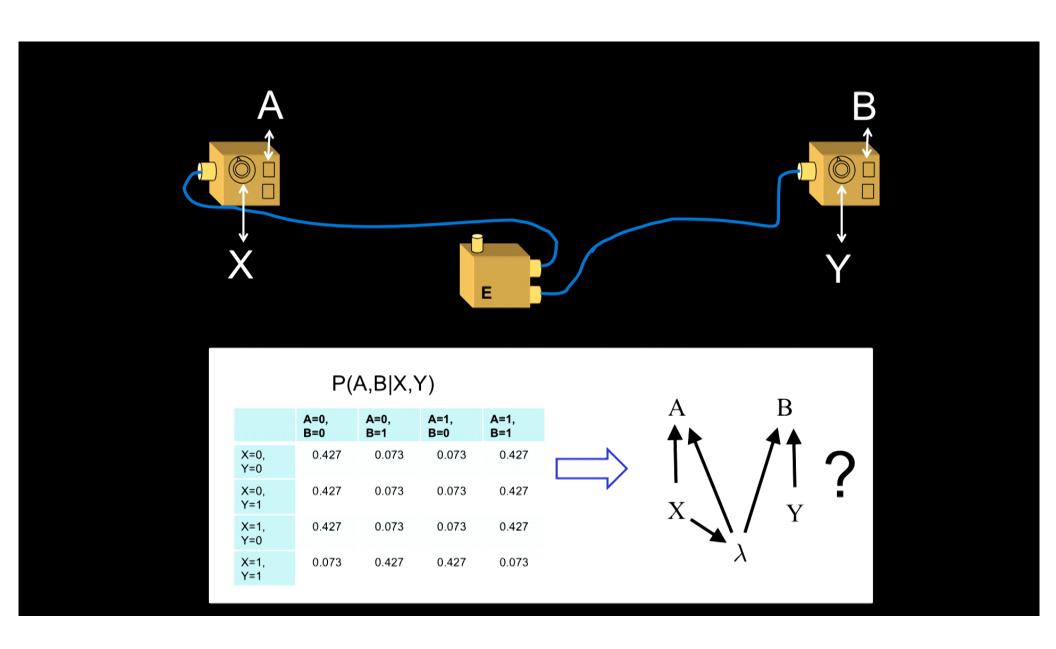
Pirsa: 18020071 Page 37/46



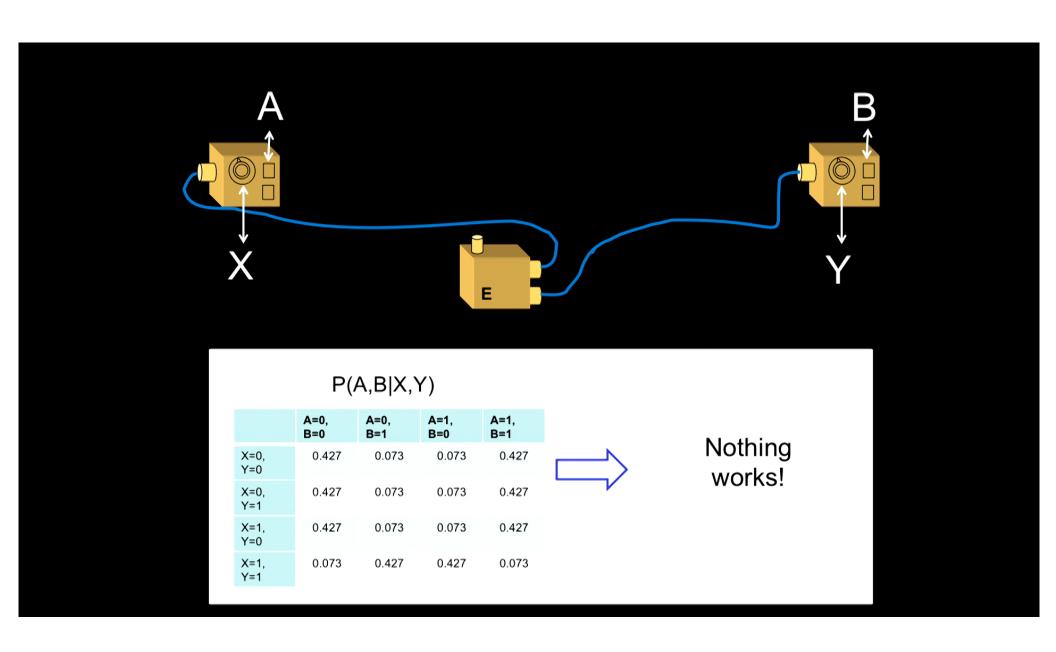
Pirsa: 18020071 Page 38/46







Pirsa: 18020071 Page 41/46





No fine-tuning



Contradiction with quantum theory and experiment

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- Realism
- Locality (no superluminal causal influences)
 - No superdeterminism (free will)
 - · No retrocausation



Contradiction with quantum theory and experiment

Pirsa: 18020071 Page 44/46



No fine-tuning



Contradiction with quantum theory and experiment

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- Realism
- Locality (no superluminal causal influences)
 - No superdeterminism (free will)
 - · No retrocausation



Contradiction with quantum theory and experiment

Pirsa: 18020071 Page 46/46