

Title: PSI 17/18 - Foundations of Quantum Mechanics - Lecture 11

Date: Feb 12, 2018 10:15 AM

URL: <http://pirsa.org/18020070>

Abstract:

Some ψ -epistemic ontological models

Classical statistical theory
+
fundamental restriction on statistical distributions
↓
A large part of quantum theory

Classical theory

Mechanics

Statistical theory for the classical theory

Liouville mechanics

Epistemically restricted statistical theory for the classical theory

Restricted Liouville mechanics
= Gaussian quantum mechanics

Classical theory

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Bits

Statistical theory for the classical theory

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Statistical theory of bits

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Restricted statistical theory of bits
 \simeq Stabilizer theory for qubits

Classical theory

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Trits

Optics

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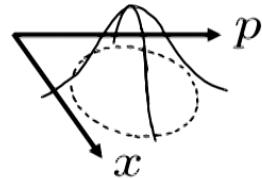
Restricted statistical theory of trits
= Stabilizer theory for qutrits

Restricted statistical optics
= linear quantum optics

Epistemically restricted Liouville mechanics
= Gaussian Quantum Mechanics

Liouville mechanics

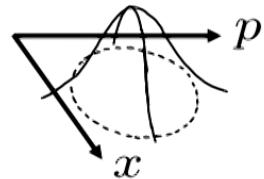
$$\mu(x, p)$$



What is a good epistemic restriction to apply?
-- look to quantum mechanics

Liouville mechanics

$$\mu(x, p)$$



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Quantum mechanics

Uncertainty principle:

$$\Delta^2 x \Delta^2 p - C_{x,p}^2 \geq (\hbar/2)^2$$

where

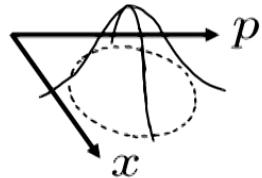
$$\Delta^2 x \equiv \langle \hat{x}^2 \rangle - \langle \hat{x} \rangle^2$$

$$C_{x,p} \equiv \frac{1}{2} \langle \hat{x} \hat{p} + \hat{p} \hat{x} \rangle - \langle \hat{x} \rangle \langle \hat{p} \rangle$$

$$\langle \hat{A} \rangle \equiv \text{Tr}(\hat{A} \hat{\rho})$$

Liouville mechanics

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Liouville mechanics with an epistemic restriction

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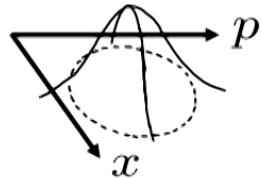
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Liouville mechanics with an epistemic restriction

Assume:

The classical uncertainty principle (for a single particle in 1D):

The only Liouville distributions that can be prepared are those that satisfy

$$\Delta^2 x \Delta^2 p - C_{x,p}^2 \geq (\hbar/2)^2$$

and that have maximal entropy for a given set of second-order moments.

Liouville mechanics with an epistemic restriction

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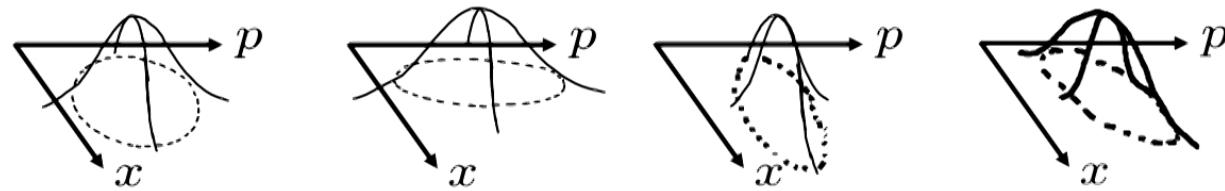
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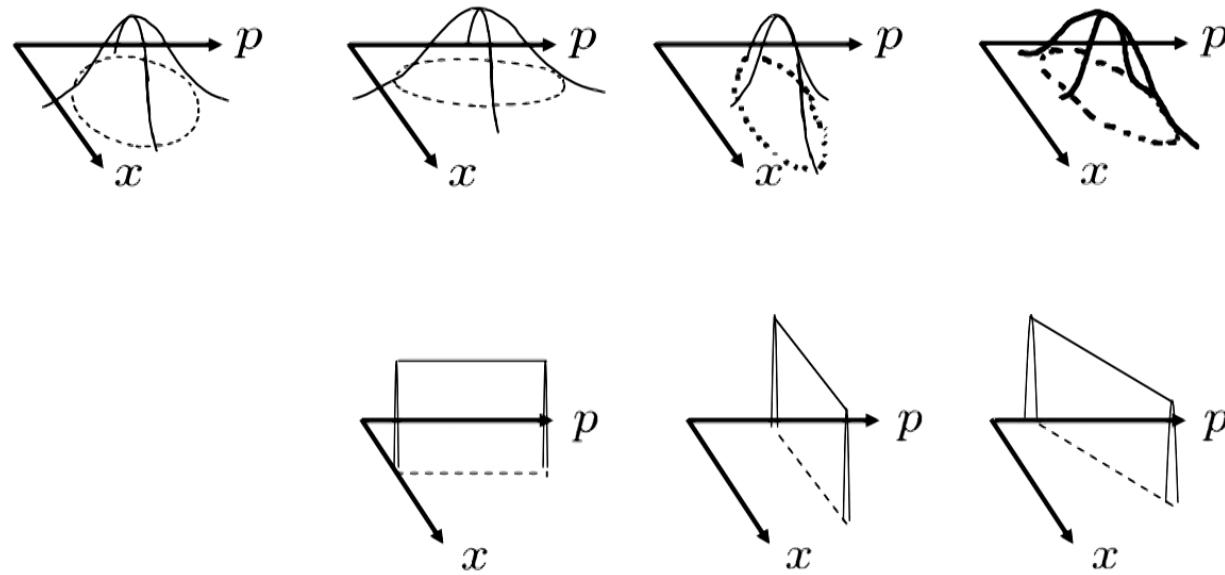
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Among $\mu(x,p)$ with a given set of second-order moments, Gaussian distributions maximize the entropy

Valid pure epistemic states for one canonical system



Valid pure epistemic states for one canonical system



The Wigner representation

Phase point operators

$$A(x, p) = \frac{1}{2\pi\hbar} \int e^{ipy/\hbar} \left| x + \frac{1}{2}y \right\rangle \left\langle x - \frac{1}{2}y \right| dy.$$

Wigner representation of an operator O

$$W_{\hat{O}}(x, p) = \text{Tr}[\hat{O}\hat{A}(x, p)]$$

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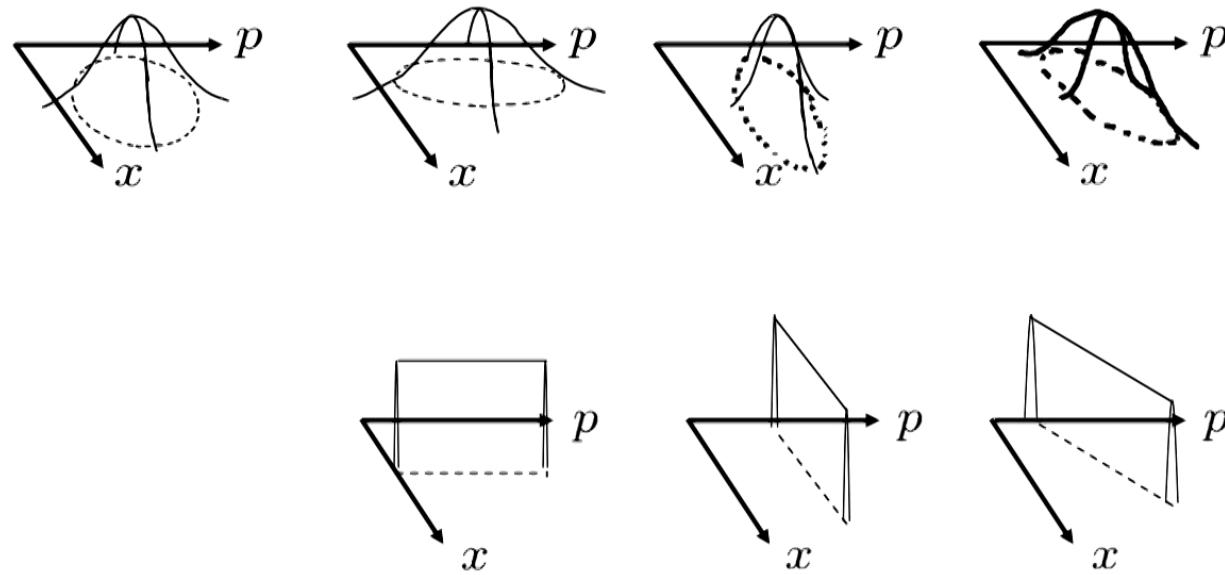
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Wigner representation of a map \mathcal{E}

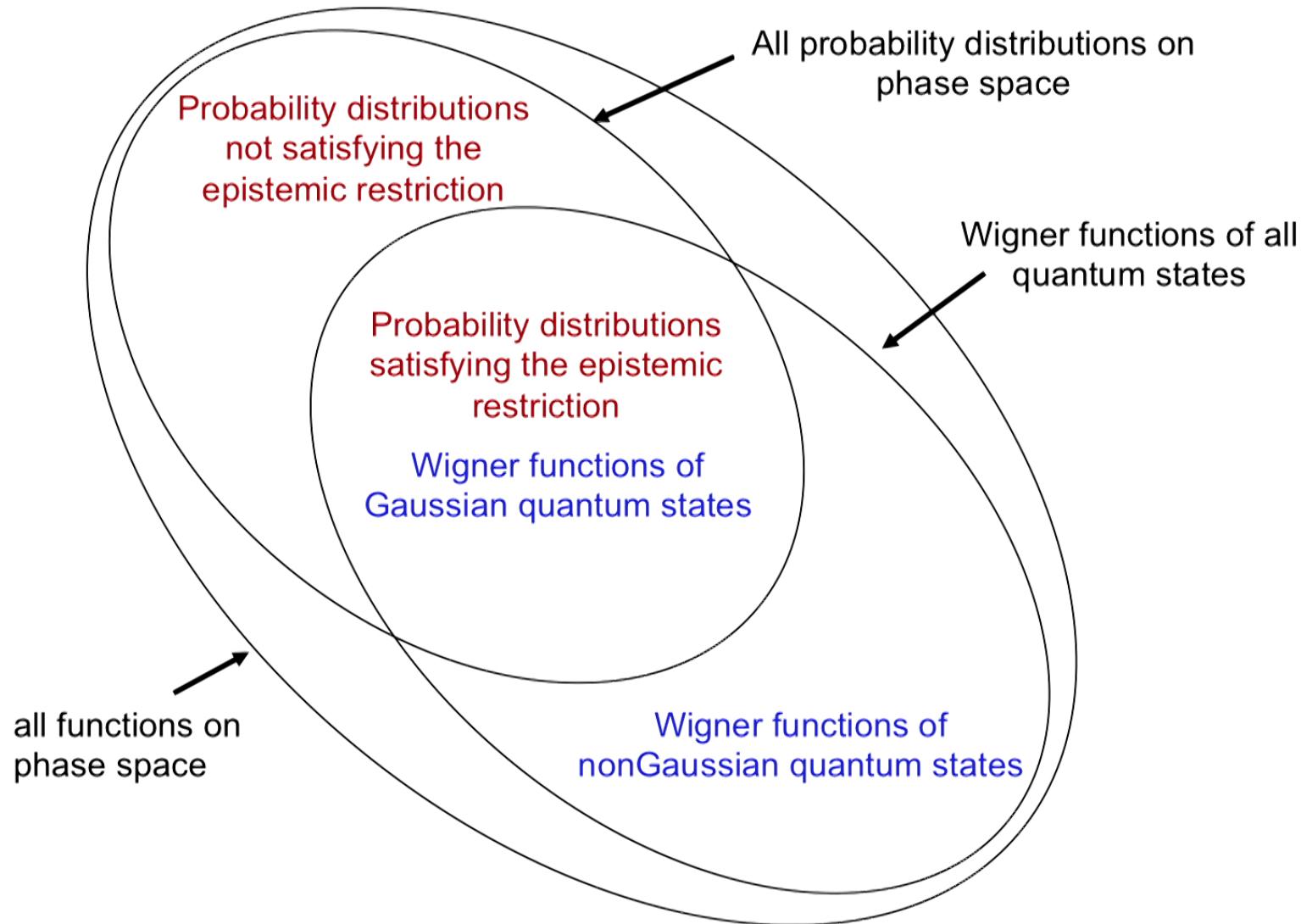
$$W_{\mathcal{E}}(x', p' | x, p) = \text{Tr}[\hat{A}(x', p') \mathcal{E}(\hat{A}(x, p))]$$

$$\text{and } W_{\mathcal{E}(\rho)}(x, p) = \int dx' dp' W_{\mathcal{E}}(x, p | x', p') W_{\hat{\rho}}(x', p')$$

Valid pure epistemic states for one canonical system



These correspond to the Wigner representations of the pure squeezed states in quantum mechanics

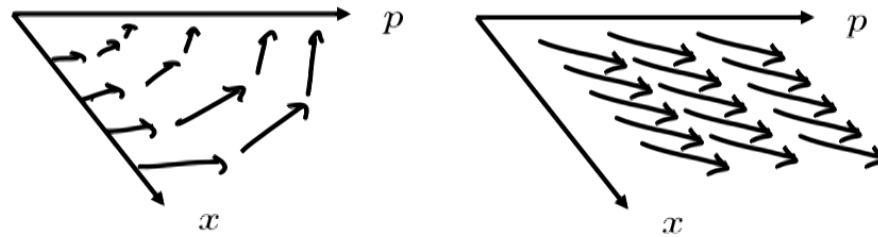


Valid deterministic transformations

The group of canonical transformations with quadratic Hamiltonian

Only canonical transformations preserve the uncertainty principle

Only quadratic Hamiltonians preserve the gaussianity

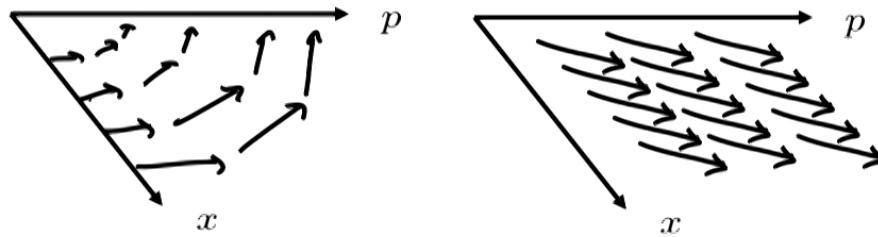


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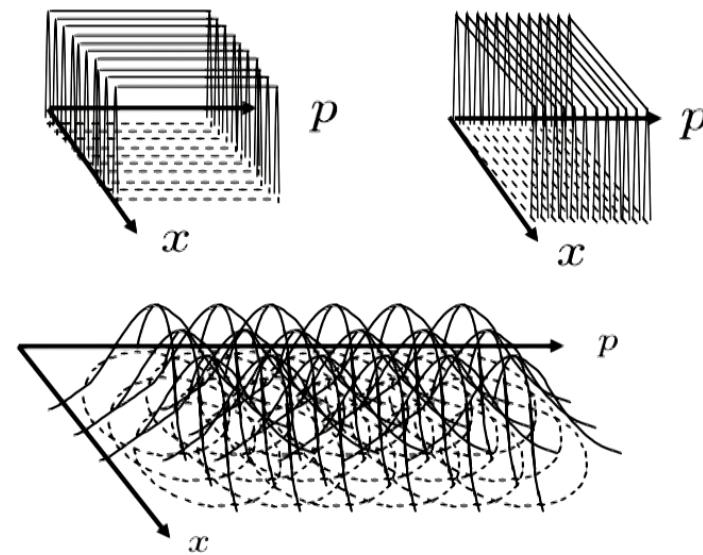


→ These correspond to the Wigner representations of the unitaries associated with these Hamiltonians

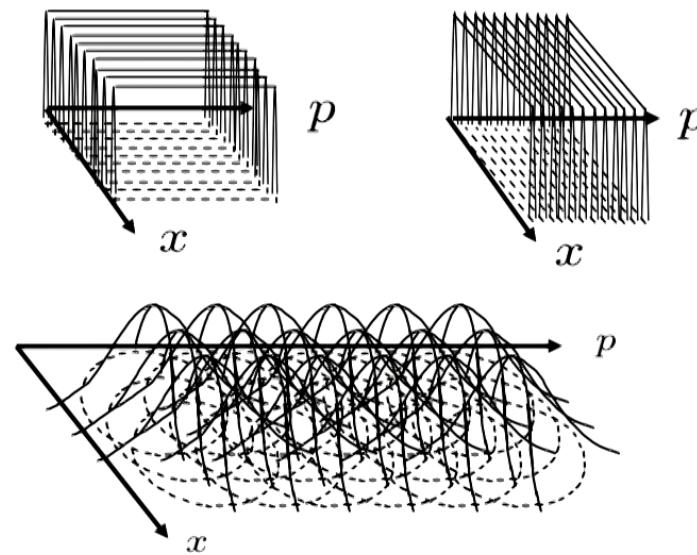
$$\text{e.g. } D_{x_0, p_0} = e^{-ix_0 \hat{P} - ip_0 \hat{X}}$$

$$\rho' = D_{x_0, p_0} \rho D_{x_0, p_0}^\dagger \quad \rightarrow \quad W_{\rho'}(x, p) = W_\rho(x + x_0, p + p_0)$$

Valid measurements



Valid measurements



→ These also correspond to the Wigner representations of quantum measurements

Extension to multiple systems

Use a generalization of the uncertainty principle for multiple systems
or:

Allow all products of valid epistemic states

Allow canonical transformations with quadratic Hamiltonians on these

E.g.

$$\begin{array}{ll} x_A \rightarrow x_A - x_B & x_B \rightarrow x_A + x_B \\ p_A \rightarrow p_A - p_B & p_B \rightarrow p_A + p_B \end{array}$$

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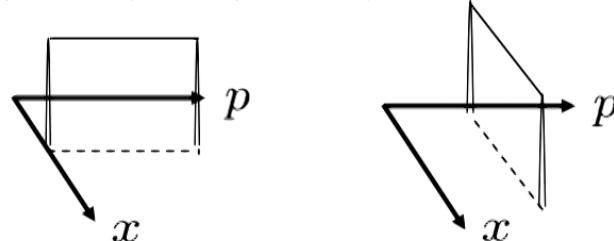
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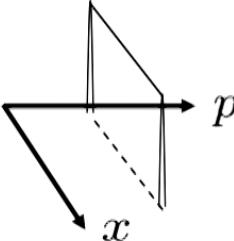
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know X_A and P_B

$$\mu(x_A, p_A) \propto \delta(x_A - a)$$



$$\mu(x_B, p_B) \propto \delta(p_B - b)$$



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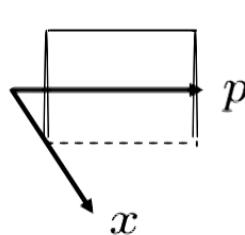
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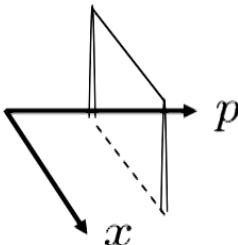


know $X_A - X_B$ and $P_A + P_B$

$$\mu(x_A, p_A, x_B, p_B)$$

$$\propto \delta(x_A - x_B - a)\delta(p_A + p_B - b)$$

corresponds to EPR state



Q: How can one characterize the set of variables that can be jointly known?

A: They commute relative to the Poisson bracket!

Configuration space: $\mathbb{R} \ni x$

Phase space: $\Omega \equiv \mathbb{R}^2 \ni (x, p) \equiv m$

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A canonically conjugate pair $[F, G] = 1$

e.g. $\{X_1, P_1\}, \{X_2, P_2\}$, and $\{X_1 + X_2, P_1 + P_2\}$

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Canonical transformations preserve the Poisson bracket

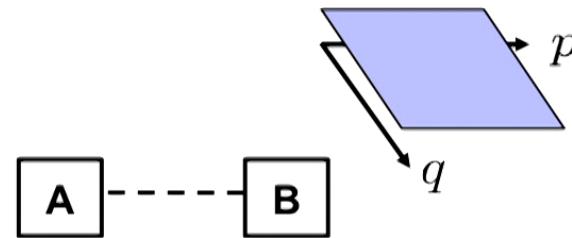
Gaussian Quantum Mechanics

- = States with Gaussian Wigner rep'ns (a.k.a. squeezed states)
- + Measurements with Gaussian Wigner rep'ns
- + Transformations that preserve Gaussianity of states

Epistemically restricted Liouville mechanics

- = Gaussian quantum mechanics in the Wigner representation

EPR effect in Epistemically Restricted Liouville mechanics

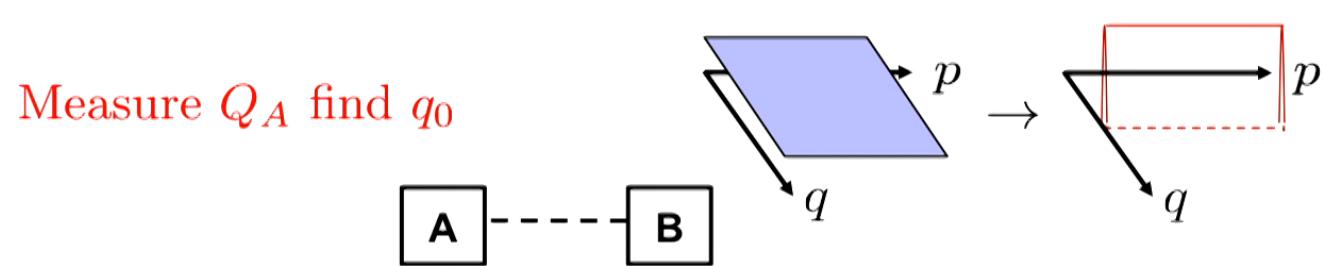


$$P_{\text{EPR}}(q_A, p_A, q_B, p_B) \propto \delta(q_A - q_B)\delta(p_A + p_B)$$

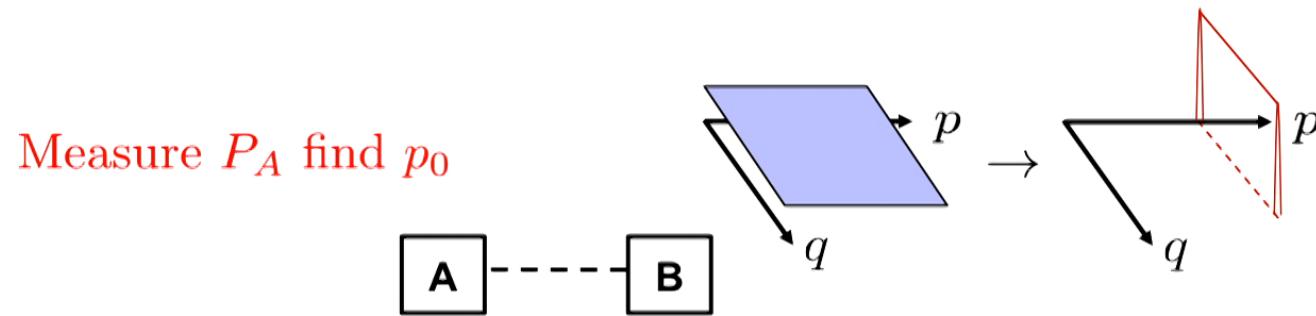
$$Q_B - Q_A = 0$$

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EPR effect in Epistemically Restricted Liouville mechanics

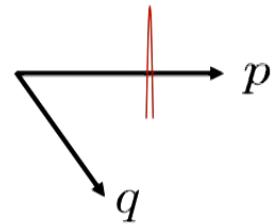
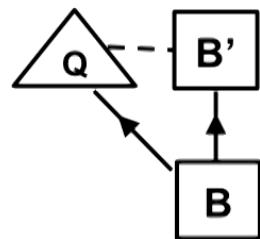


EPR effect in Epistemically Restricted Liouville mechanics

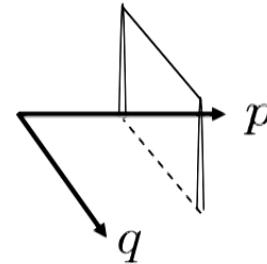


Collapse Rule in Epistemically Restricted Liouville mechanics

Measure Q_B find q_0

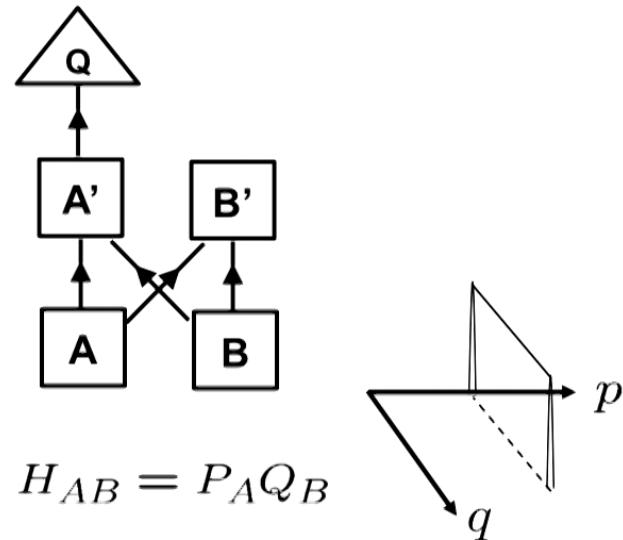


But this would violate the epistemic restriction!



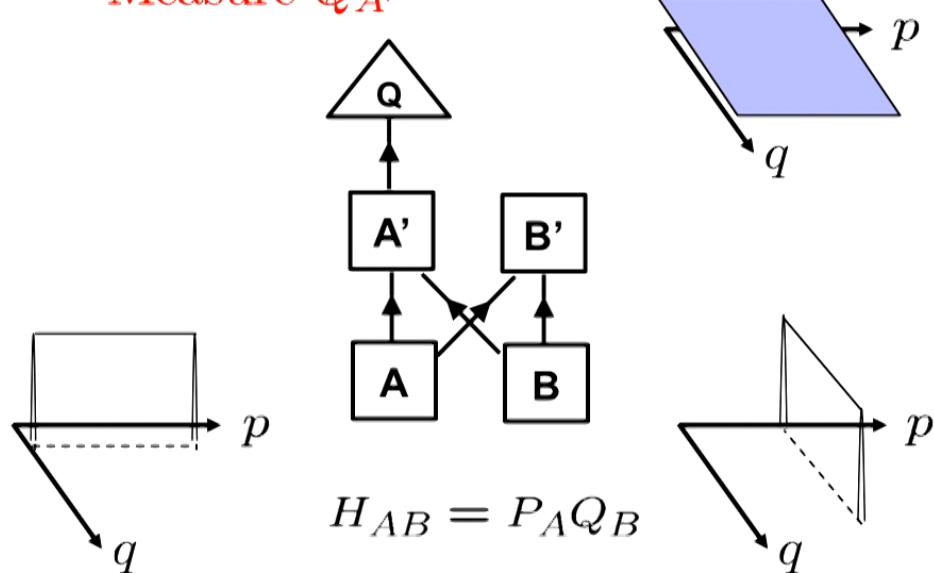
Collapse Rule in Epistemically Restricted Liouville mechanics

Measure $Q_{A'}$



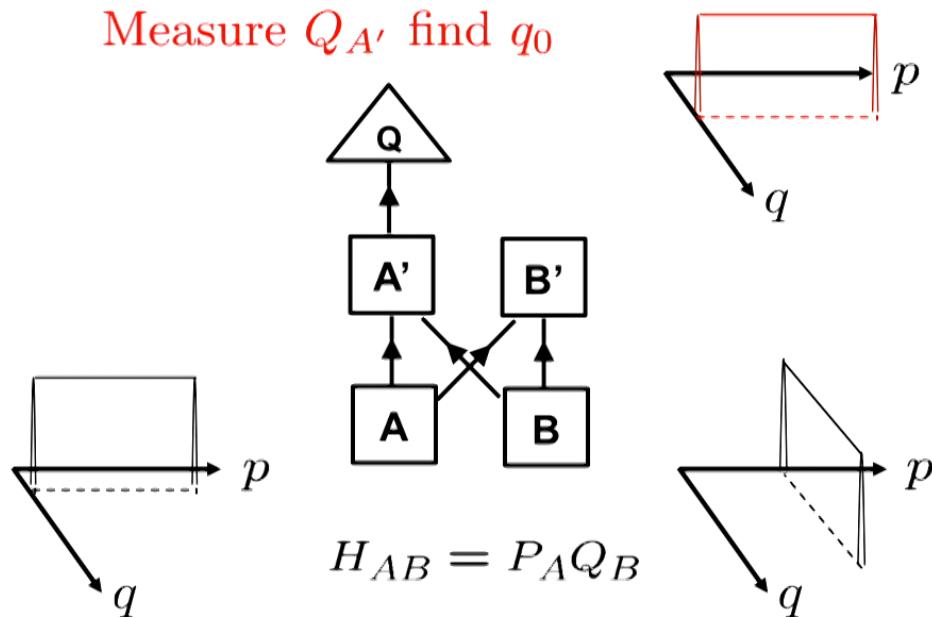
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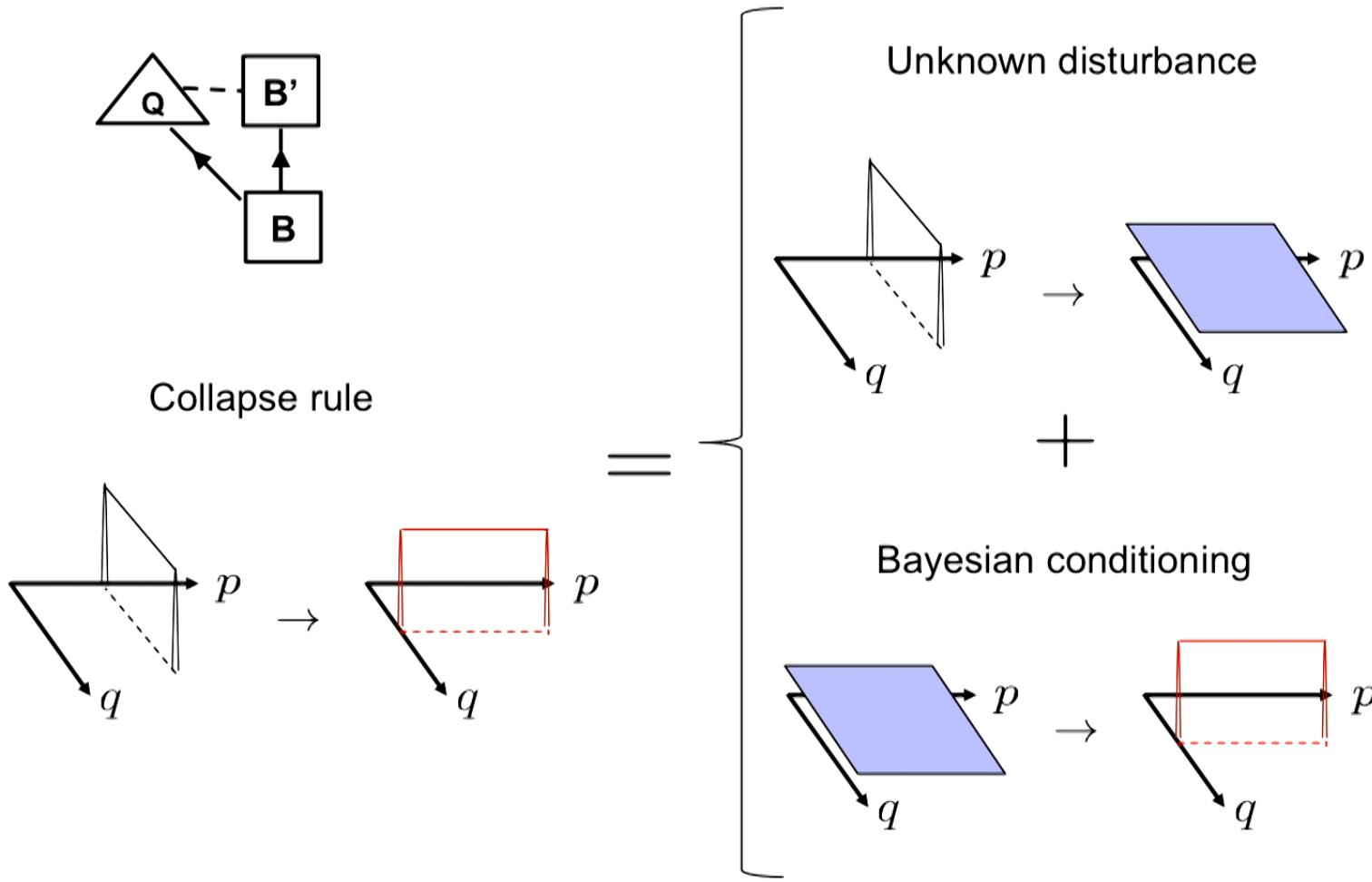


Collapse Rule in Epistemically Restricted Liouville mechanics

Measure $Q_{A'}$ find q_0



Collapse Rule in Epistemically Restricted Liouville mechanics



Epistemically restricted statistical theory of bits

Recall: particle mechanics

Configuration space: $\mathbb{R}^n \ni (x_1, x_2, \dots, x_n)$

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“bit mechanics” $\mathbb{Z}_2 = \{0, 1\}$

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$$\begin{aligned}[F, G](m) \equiv \sum_{i=1}^n & (F[m + e_{x_i}] - F[m])(G[m + e_{p_i}] - G[m]) \\ & - (F[m + e_{p_i}] - F[m])(G[m + e_{q_i}] - G[m])\end{aligned}$$

The epistemic restriction:

An observer can only have knowledge of the values of a set of canonical variables that commute relative to the Poisson bracket and is maximally ignorant otherwise.

A single bit

| | | |
|-----|---|-----|
| X | 1 | |
| 0 | | |
| | 0 | 1 |
| | | P |

Canonical variables

$$aX + bP \quad a, b \in \mathbb{Z}_2 \quad \text{Addition is mod2}$$

$$X, \quad P, \quad X + P$$

A single bit

| | | |
|-----|---|-----|
| X | 1 | |
| 0 | | |
| | 0 | 1 |
| | | P |

Canonical variables

$$aX + bP \quad a, b \in \mathbb{Z}_2 \quad \text{Addition is mod2}$$

$$X, \quad P, \quad X + P$$

Statistical distributions

X known

| | | |
|-----|---|-----|
| X | 1 | |
| 0 | | |
| | 0 | 1 |
| | | P |

P known

| | | |
|-----|---|-----|
| X | 1 | |
| 0 | | |
| | 0 | 1 |
| | | P |

$X + P$ known

| | | |
|-----|---|-----|
| X | 1 | |
| 0 | | |
| | 0 | 1 |
| | | P |

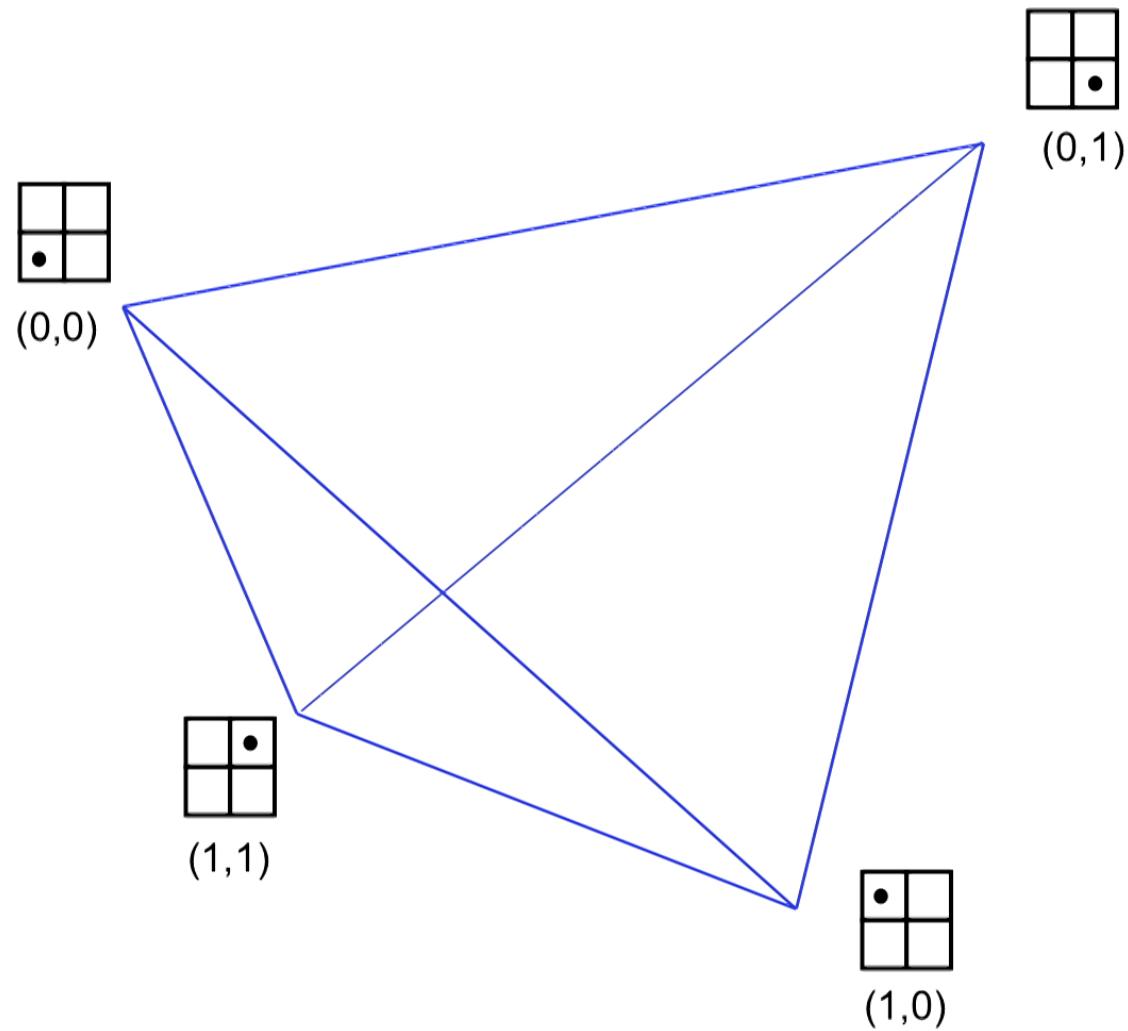
Nothing known

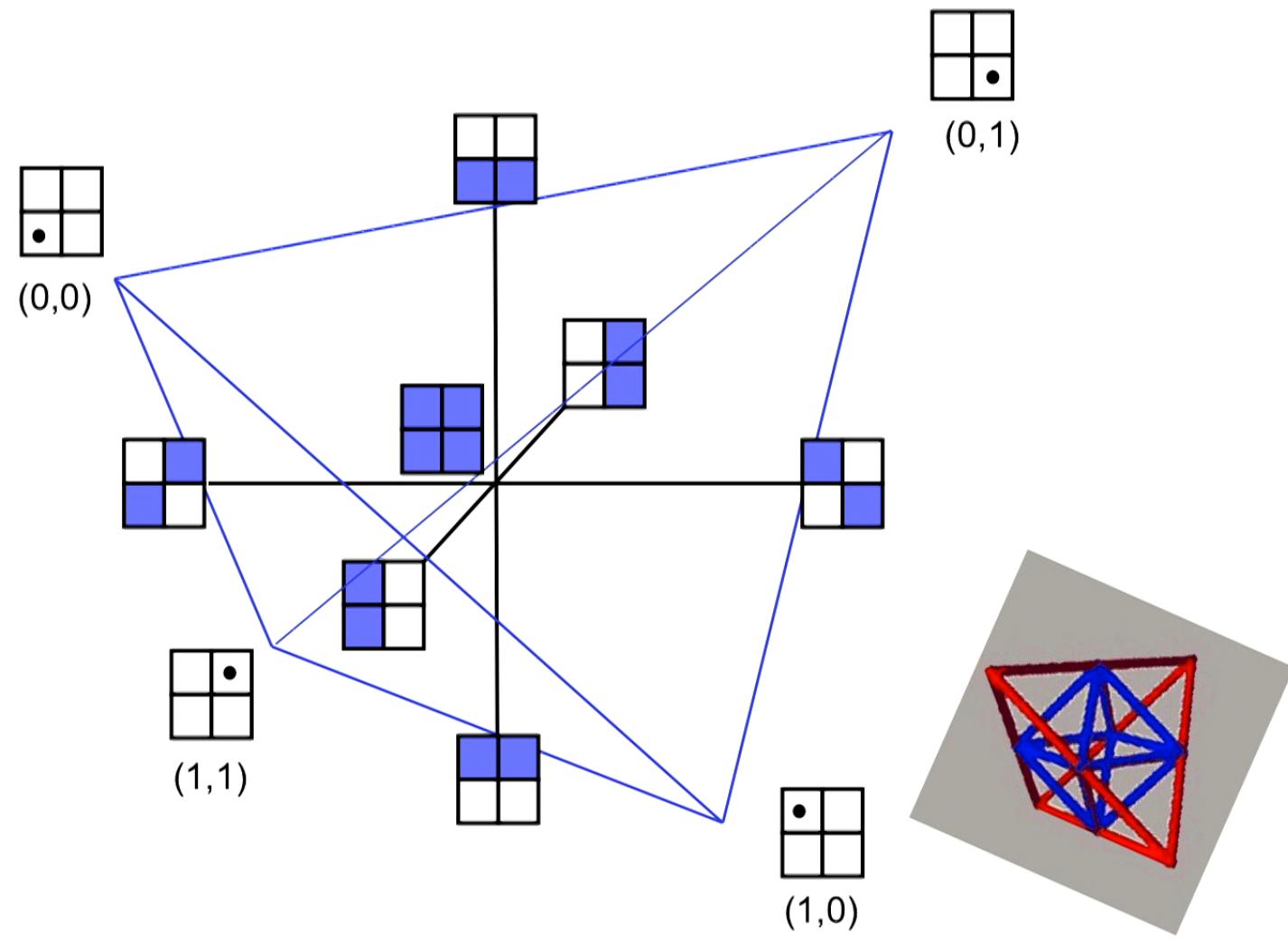
| | | |
|-----|---|-----|
| X | 1 | |
| 0 | | |
| | 0 | 1 |
| | | P |

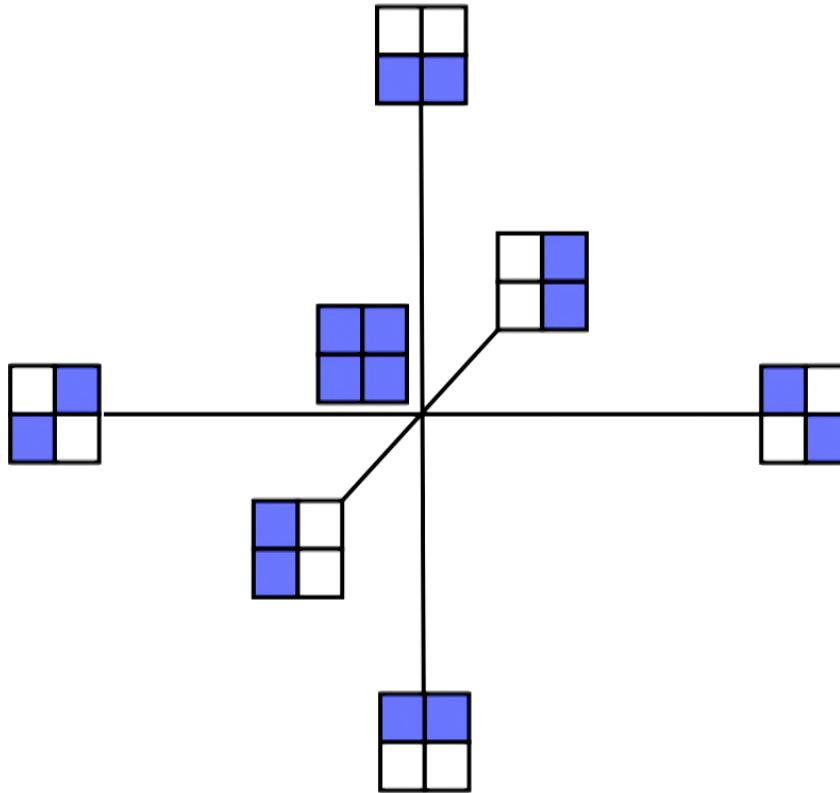
Convex combination

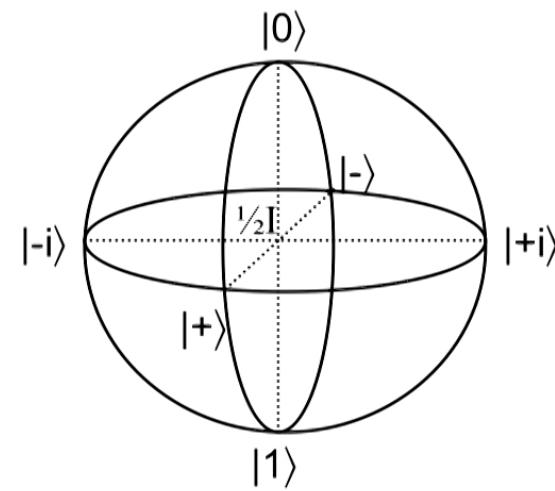
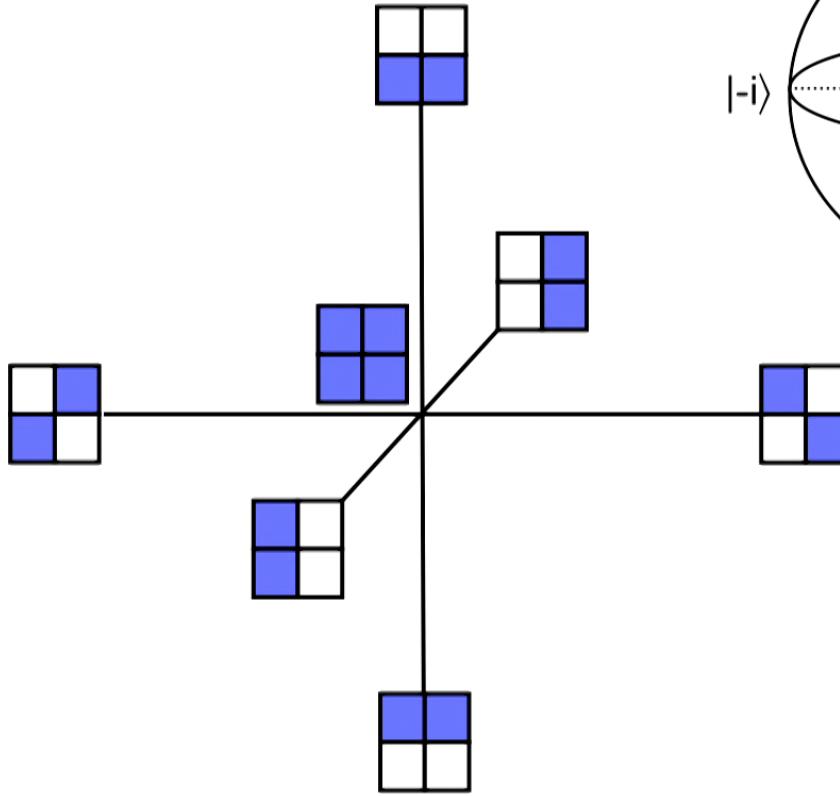
$$\begin{array}{c} \text{Diagram of a } 2 \times 2 \text{ matrix with all entries blue.} \\ = \quad \begin{array}{c} \text{Diagram of a } 2 \times 2 \text{ matrix with top-left and bottom-right entries blue, others white.} \\ +_{\text{cx}} \quad \begin{array}{c} \text{Diagram of a } 2 \times 2 \text{ matrix with top-right and bottom-left entries blue, others white.} \\ \end{array} \end{array} \\ \begin{array}{c} \text{Diagram of a } 2 \times 2 \text{ matrix with top-left and bottom-right entries blue, others white.} \\ +_{\text{cx}} \quad \begin{array}{c} \text{Diagram of a } 2 \times 2 \text{ matrix with top-right and bottom-left entries blue, others white.} \\ \end{array} \end{array} \\ \begin{array}{c} \text{Diagram of a } 2 \times 2 \text{ matrix with top-left and bottom-right entries blue, others white.} \\ +_{\text{cx}} \quad \begin{array}{c} \text{Diagram of a } 2 \times 2 \text{ matrix with top-right and bottom-left entries blue, others white.} \\ \end{array} \end{array} \end{array}$$

$$\begin{aligned} \frac{1}{2}I &= \frac{1}{2}|0\rangle\langle 0| + \frac{1}{2}|1\rangle\langle 1| \\ &= \frac{1}{2}|+\rangle\langle +| + \frac{1}{2}|-\rangle\langle -| \\ &= \frac{1}{2}|+i\rangle\langle +i| + \frac{1}{2}|-i\rangle\langle -i| \end{aligned}$$









Valid reproducible measurements:

Any **commuting** set of canonical variables

X

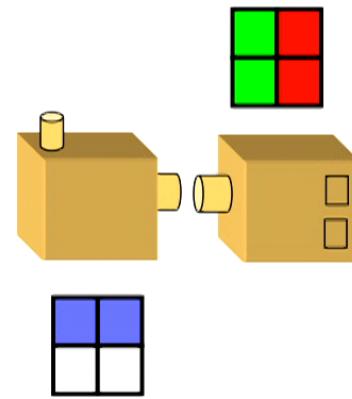
| | | |
|-----|-----|--|
| X | 1 | |
| 0 | | |
| | P | |

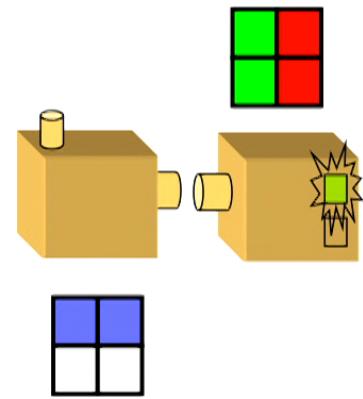
P

| | | |
|-----|-----|--|
| X | 1 | |
| 0 | | |
| | P | |

$X + P$

| | | |
|-----|-----|--|
| X | 1 | |
| 0 | | |
| | P | |





$\frac{1}{2}$ of the time

Valid reversible transformations:

The discrete canonical transformations

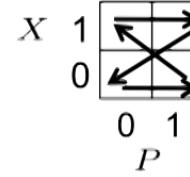
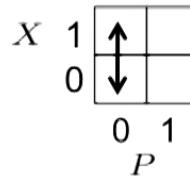
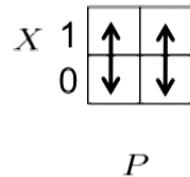
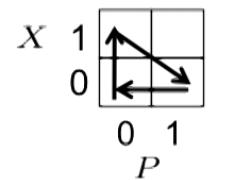
(those that preserve the discrete Poisson bracket)

Valid reversible transformations:

The discrete canonical transformations

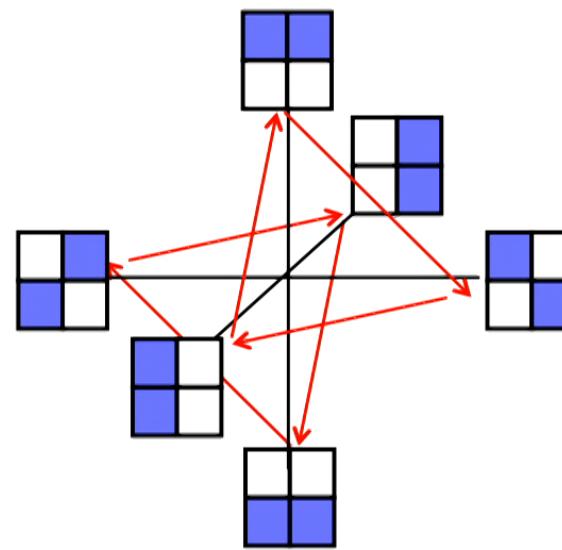
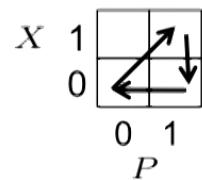
(those that preserve the discrete Poisson bracket)

→ All 24 permutations of the four ontic states



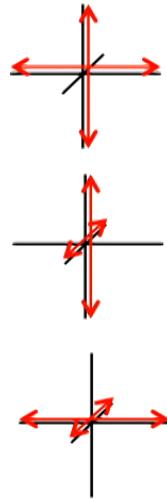
etcetera

A 3-cycle

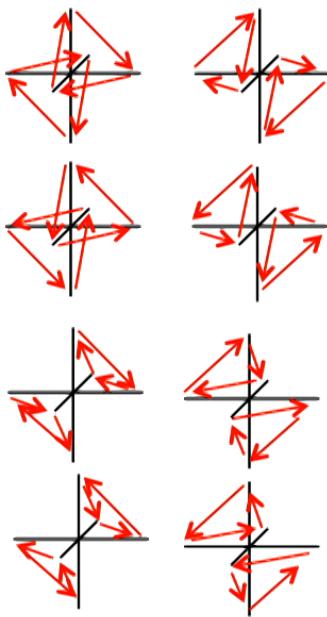


Reversible transformations:

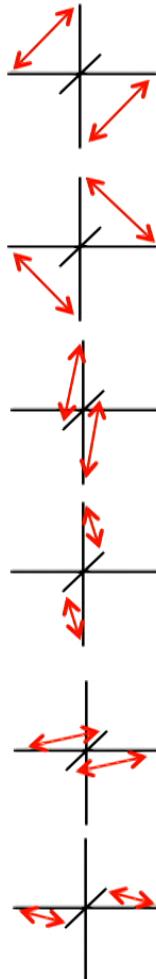
Pairs of 2-cycles



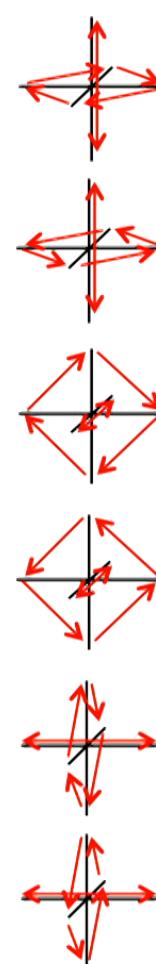
3-cycles



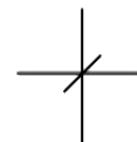
2-cycles



4-cycles

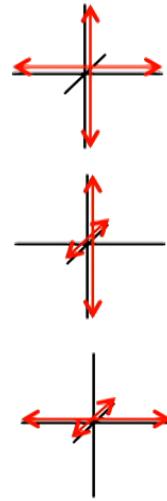


identity

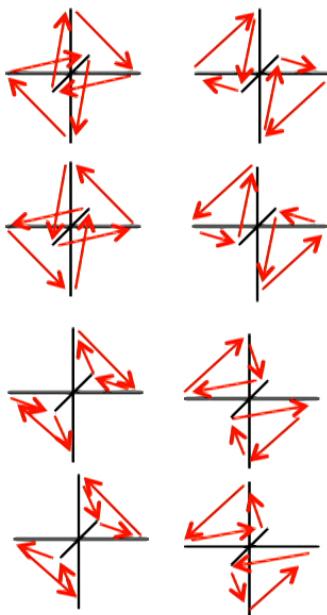


Reversible transformations:

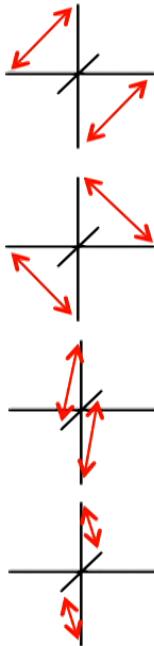
Pairs of 2-cycles



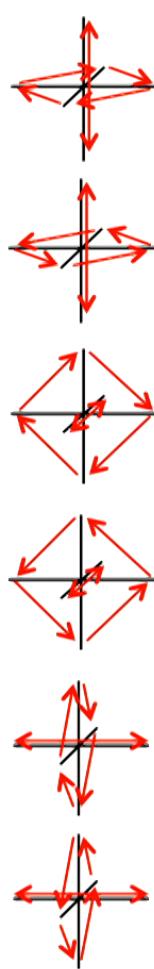
3-cycles



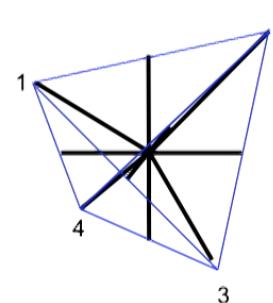
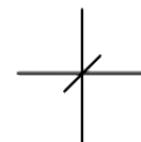
2-cycles



4-cycles



identity



Symmetries of the tetrahedron under rotations and reflections

Coherent superposition

Find U such that $|1\rangle = U|0\rangle$

Find all U' such that $(U')^2 = U$

$U'|0\rangle$ = superpos'n of $|0\rangle$ and $|1\rangle$.

$R_y(\pi)$

$R_x(\pi)$

$R_y(-\pi/2)$

$R_y(\pi/2)$

$R_x(\pi/2)$

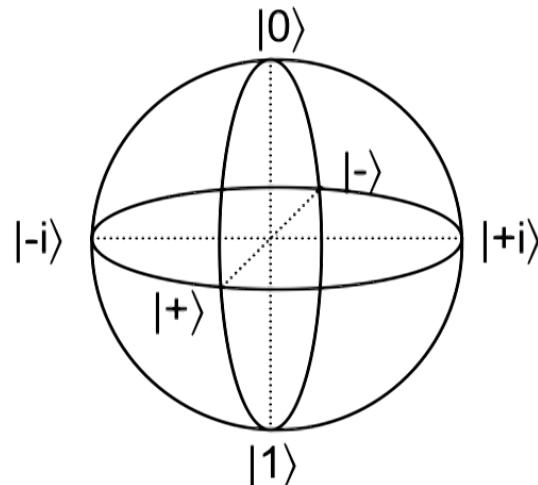
$R_x(-\pi/2)$

$|+\rangle$

$|-\rangle$

$|+i\rangle$

$| -i \rangle$

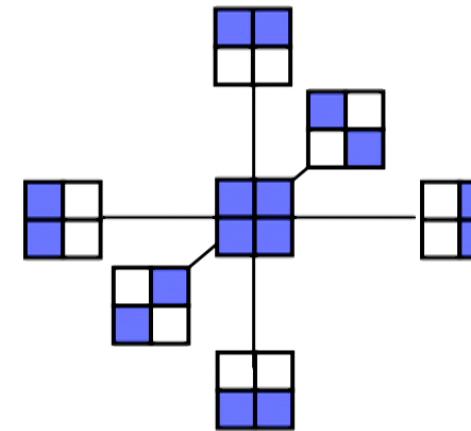


Find P such that

$$P'(\begin{array}{|c|c|}\hline \textcolor{blue}{\square} & \textcolor{blue}{\square} \\ \hline \textcolor{blue}{\square} & \textcolor{blue}{\square} \\ \hline \end{array}) = P(\begin{array}{|c|c|}\hline \textcolor{blue}{\square} & \textcolor{blue}{\square} \\ \hline \textcolor{blue}{\square} & \textcolor{blue}{\square} \\ \hline \end{array})$$

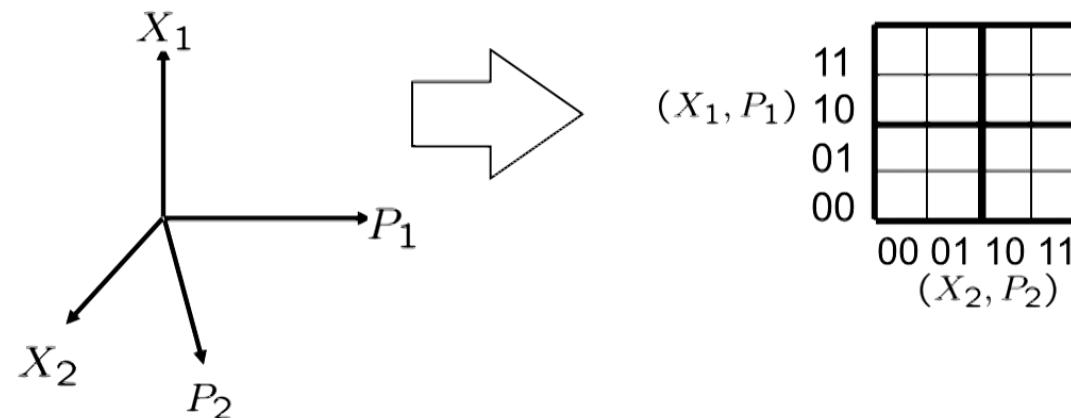
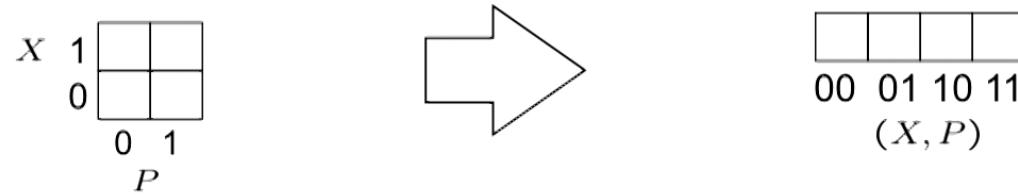
Find all P' such that $(P')^2 = P$

$P'(\begin{array}{|c|c|}\hline \textcolor{blue}{\square} & \textcolor{blue}{\square} \\ \hline \textcolor{blue}{\square} & \textcolor{blue}{\square} \\ \hline \end{array})$ = superpos'n of $\begin{array}{|c|c|}\hline \textcolor{blue}{\square} & \textcolor{blue}{\square} \\ \hline \textcolor{blue}{\square} & \textcolor{blue}{\square} \\ \hline \end{array}$ and $\begin{array}{|c|c|}\hline \textcolor{blue}{\square} & \textcolor{blue}{\square} \\ \hline \textcolor{blue}{\square} & \textcolor{blue}{\square} \\ \hline \end{array}$



A pair of bits

Canonical variables $a_1 X_1 + b_1 P_1 + a_2 X_2 + b_2 P_2 \quad a_1, b_1, a_2, b_2 \in \mathbb{Z}_2$



1 variable known

X_1 known

| | | | | |
|----|----|----|--------------|----|
| | | | | |
| 11 | | | | |
| 10 | | | | |
| 01 | | | | |
| 00 | | | | |
| | 00 | 01 | 10 | 11 |
| | | | (X_2, P_2) | |

P_2 known

| | | | | |
|----|--------------|----|----|----|
| | | | | |
| 11 | | | | |
| 10 | | | | |
| 01 | | | | |
| 00 | | | | |
| | 00 | 01 | 10 | 11 |
| | (X_2, P_2) | | | |

2 variables known

X_1 and P_2 known

1 variable known

$X_1 + X_2$ known

| | | | |
|----------------|----|-------------|----------------|
| | 11 | | |
| | 10 | | |
| (X_1, P_1) | 01 | | |
| | 00 | 00 01 10 11 | |
| | | | (X_2, P_2) |

$P_1 + P_2$ known

| | | | |
|----------------|----|-------------|----------------|
| | 11 | | |
| | 10 | | |
| (X_1, P_1) | 01 | | |
| | 00 | 00 01 10 11 | |
| | | | (X_2, P_2) |

2 variables known

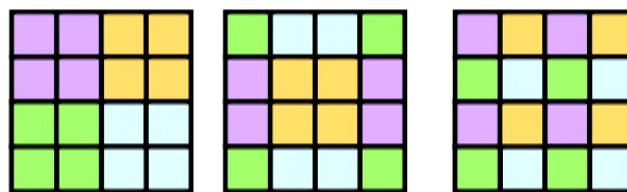
$X_1 + X_2$ and $P_1 + P_2$ known

| | | | |
|----------------|----|-------------|----------------|
| | 11 | | |
| | 10 | | |
| (X_1, P_1) | 01 | | |
| | 00 | 00 01 10 11 | |
| | | | (X_2, P_2) |

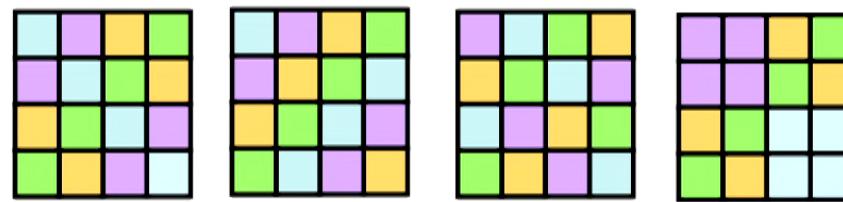
$$\frac{1}{\sqrt{2}}(|0\rangle \otimes |0\rangle + |1\rangle \otimes |1\rangle)$$

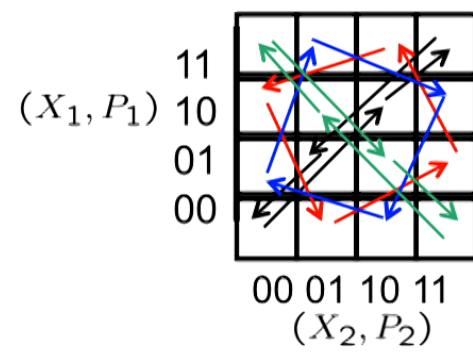
Maximally informative reproducible measurements

Product basis Measurements



Entangled basis Measurements





The way to understand EPR steering and the collapse rule are precisely analogous to how it is done in epistemically restricted Liouville mechanics

$$|\phi^\pm\rangle = \frac{1}{\sqrt{2}}(|0\rangle|0\rangle \pm |1\rangle|1\rangle)$$

$$|\psi^\pm\rangle = \frac{1}{\sqrt{2}}(|0\rangle|1\rangle \pm |1\rangle|0\rangle)$$

Teleportation

I, X, Y, Z

$\{|\psi^-\rangle, |\psi^+\rangle, |\phi^-\rangle, |\phi^+\rangle\}$

