Title: PSI 17/18 - Foundations of Quantum Mechanics - Lecture 5

Date: Feb 02, 2018 10:15 AM

URL: http://pirsa.org/18020064

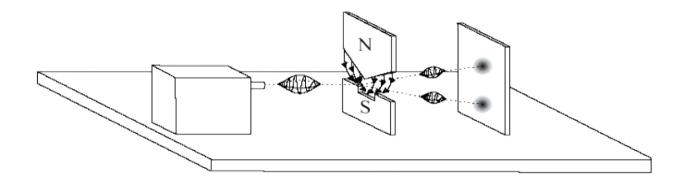
Abstract:

Pirsa: 18020064 Page 1/62

Dynamical Collapse Theories

Pirsa: 18020064 Page 2/62

The quantum measurement problem



$$\begin{array}{c} (a|\uparrow\rangle+b|\downarrow\rangle)|\text{ "ready"}\,\rangle\\ \\ \to a|\uparrow\rangle|\text{ "up"}\,\rangle+b|\downarrow\rangle|\text{ "down"}\,\rangle \end{array}$$

Pirsa: 18020064 Page 3/62

Inconsistencies of the orthodox interpretation

By the collapse postulate (applied to the system)

By unitary evolution postulate (applied to isolated system that includes the apparatus)

Indeterministic and discontinuous evolution

Deterministic and continuous evolution

Determinate properties

Indeterminate properties

Pirsa: 18020064 Page 4/62

Inconsistencies of the orthodox interpretation

By the collapse postulate (applied to the system)

By unitary evolution postulate (applied to isolated system that includes the apparatus)

Indeterministic and discontinuous evolution

Deterministic and continuous evolution

Determinate properties

Indeterminate properties

Posit a new nonunitary dynamical evolution law (either nonlinear or indeterministic or both)

Macroscopic systems obey collapse postulate dynamics to good approximation

Microscopic systems obey unitary dynamics to good approximation

Pirsa: 18020064 Page 5/62

Responses to the measurement problem

- 1. Deny realism
 - Purely operational account of quantum theory
- 2. Deny the universality of unitary dynamics
 - Dynamical collapse theories
- 3. Deny that ψ is a complete representation of reality
 - Hidden variable models
 - Models of reality beyond hidden variables?
- 4. Deny indeterminism and discontinuity, except as subjective illusions
 - Everett's relative state interpretation, or "many worlds"

Pirsa: 18020064 Page 6/62

The goal:

$$(a|\uparrow\rangle+b|\downarrow\rangle)|$$
 "ready" \rangle \rightarrow $|\uparrow\rangle|$ "up" \rangle with probability $|a|^2$ \rightarrow $|\downarrow\rangle|$ "down" \rangle with probability $|b|^2$

The preferred decomposition issue

Into what states do collapses occur?

Pirsa: 18020064 Page 7/62

The goal:

The preferred decomposition issue

Into what states do collapses occur?

$$\begin{aligned} \left(a|+\rangle+b|-\rangle\right) |O_{0}\rangle & \longrightarrow a|+\rangle |O_{+}\rangle+b|-\rangle |O_{-}\rangle \\ &= \left(\frac{a|+\rangle+b|-\rangle}{\sqrt{2}}\right) \left(\frac{|O_{+}\rangle+|O_{-}\rangle}{\sqrt{2}}\right) \\ &+ \left(\frac{a|+\rangle-b|-\rangle}{\sqrt{2}}\right) \left(\frac{|O_{+}\rangle-|O_{-}\rangle}{\sqrt{2}}\right) \end{aligned}$$

Pirsa: 18020064

The goal:

$$(a|\uparrow\rangle+b|\downarrow\rangle)|$$
 "ready" \rangle \rightarrow $|\uparrow\rangle|$ "up" \rangle with probability $|a|^2$ \rightarrow $|\downarrow\rangle|$ "down" \rangle with probability $|b|^2$

The preferred decomposition issue

Into what states do collapses occur?

A: Choose a collapse mechanism that tends to localize in space

Pirsa: 18020064 Page 9/62

The goal:

$$(a|\uparrow\rangle+b|\downarrow\rangle)|$$
 "ready" \rangle \rightarrow $|\uparrow\rangle|$ "up" \rangle with probability $|a|^2$ \rightarrow $|\downarrow\rangle|$ "down" \rangle with probability $|b|^2$

The preferred decomposition issue

Into what states do collapses occur?

A: Choose a collapse mechanism that tends to localize in space

The trigger issue

When and how do collapses occur?

Pirsa: 18020064 Page 10/62

The goal:

$$(a|\uparrow\rangle+b|\downarrow\rangle)|$$
 "ready" \rangle \rightarrow $|\uparrow\rangle|$ "up" \rangle with probability $|a|^2$ \rightarrow $|\downarrow\rangle|$ "down" \rangle with probability $|b|^2$

The preferred decomposition issue

Into what states do collapses occur?

A: Choose a collapse mechanism that tends to localize in space

The trigger issue

When and how do collapses occur?

A: Postulate collapse as a new primitive physical mechanism

Pirsa: 18020064 Page 11/62

The Ghirardi-Rimini-Weber model

At most times:

$$i\hbar rac{\partial}{\partial t} |\psi(t)
angle = H |\psi(t)
angle$$
 Schrödinger's equation

Every τ/N time interval on average

$$|\psi(t+dt)\rangle = \frac{1}{\sqrt{p(\mathbf{q}_k)}}\sqrt{E^{(k)}(\mathbf{q}_k)}|\psi(t)\rangle$$
 "Collapse"

where
$$E^{(k)}(\mathbf{q}_k) = \int d\mathbf{r}_k \ K \exp(-\frac{(\mathbf{r}_k - \mathbf{q}_k)^2}{\sigma^2}) |\mathbf{r}_k\rangle \langle \mathbf{r}_k|$$

$$p(\mathbf{q}_k) = \langle \psi(t) | E^{(k)}(\mathbf{q}_k) | \psi(t) \rangle$$

k is chosen uniformly at random

 \mathbf{q}_k is chosen by sampling from $p(\mathbf{q}_k)$

Two new fundamental constants:

 $au pprox 10^{15}
m s pprox 100 \ million \ years$ mean time between collapses for one particle $\sigma pprox 10^{-7}
m m pprox size \ of \ large \ molecule$ Localization width

Pirsa: 18020064 Page 12/62

The Ghirardi-Rimini-Weber model

At most times:

$$i\hbar rac{\partial}{\partial t} |\psi(t)
angle = H |\psi(t)
angle$$
 Schrödinger's equation

Every τ/N time interval on average

$$|\psi(t+dt)\rangle = \frac{1}{\sqrt{p(\mathbf{q}_k)}}\sqrt{E^{(k)}(\mathbf{q}_k)}|\psi(t)\rangle$$
 "Collapse"

where
$$E^{(k)}(\mathbf{q}_k) = \int d\mathbf{r}_k \ K \exp(-\frac{(\mathbf{r}_k - \mathbf{q}_k)^2}{\sigma^2}) |\mathbf{r}_k\rangle \langle \mathbf{r}_k|$$

$$p(\mathbf{q}_k) = \langle \psi(t) | E^{(k)}(\mathbf{q}_k) | \psi(t) \rangle$$

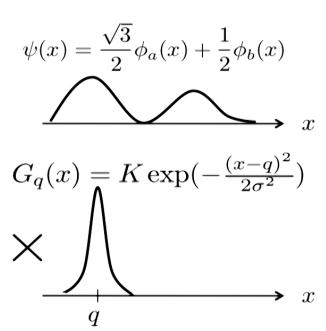
k is chosen uniformly at random

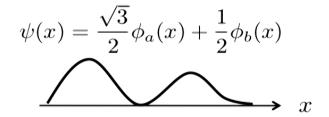
 \mathbf{q}_k is chosen by sampling from $p(\mathbf{q}_k)$

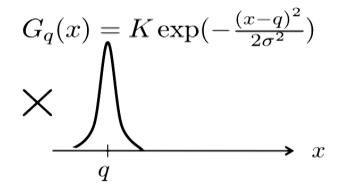
Two new fundamental constants:

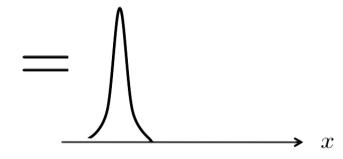
 $au pprox 10^{15}
m s pprox 100 \ million \ years$ mean time between collapses for one particle $\sigma pprox 10^{-7}
m m pprox size \ of \ large \ molecule$ Localization width

Pirsa: 18020064 Page 13/62



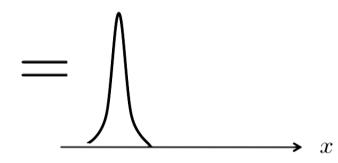






$$\psi(x) = \frac{\sqrt{3}}{2}\phi_a(x) + \frac{1}{2}\phi_b(x)$$

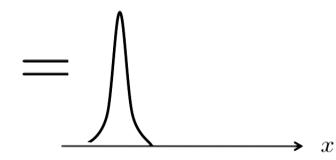
$$K \exp(-\frac{(x-q)^2}{2\sigma^2})$$



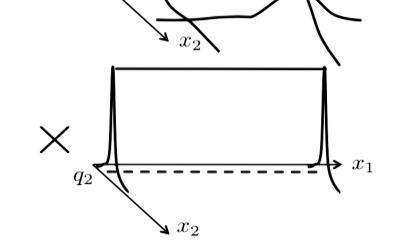
$$\psi(x_1, x_2) = \frac{\sqrt{3}}{2} \phi_a(x_1) \chi_a(x_2) + \frac{1}{2} \phi_b(x_1) \chi_b(x_2)$$

$$\psi(x) = \frac{\sqrt{3}}{2}\phi_a(x) + \frac{1}{2}\phi_b(x)$$

$$K \exp(-\frac{(x-q)^2}{2\sigma^2})$$

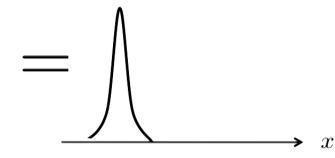


$$\psi(x_1, x_2) = \frac{\sqrt{3}}{2} \phi_a(x_1) \chi_a(x_2) + \frac{1}{2} \phi_b(x_1) \chi_b(x_2)$$

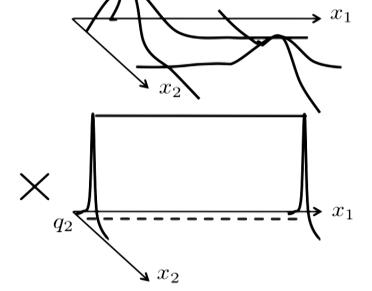


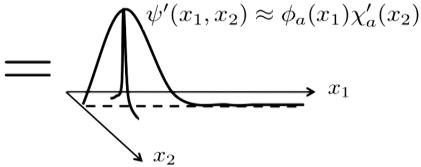
$$\psi(x) = \frac{\sqrt{3}}{2}\phi_a(x) + \frac{1}{2}\phi_b(x)$$

$$G_q(x) = K \exp(-\frac{(x-q)^2}{2\sigma^2})$$



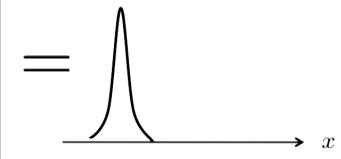
$$\psi(x_1, x_2) = \frac{\sqrt{3}}{2} \phi_a(x_1) \chi_a(x_2) + \frac{1}{2} \phi_b(x_1) \chi_b(x_2)$$



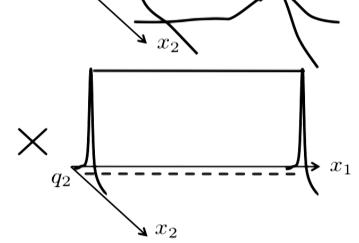


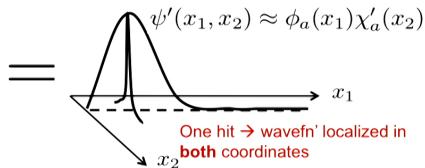
$$\psi(x) = \frac{\sqrt{3}}{2}\phi_a(x) + \frac{1}{2}\phi_b(x)$$

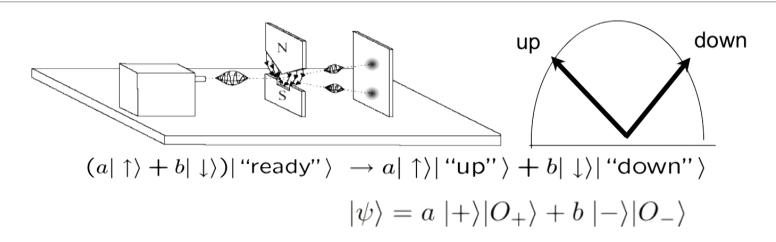
$$K \exp(-\frac{(x-q)^2}{2\sigma^2})$$



$$\psi(x_1, x_2) = \frac{\sqrt{3}}{2} \phi_a(x_1) \chi_a(x_2) + \frac{1}{2} \phi_b(x_1) \chi_b(x_2)$$

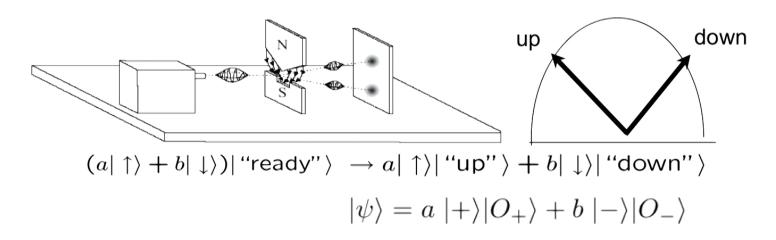






Suppose that for macroscopic # of components k, $|\langle O_+|\mathbf{r}_k\rangle|^2|\langle O_-|\mathbf{r}_k\rangle|^2\approx 0$ One particle is hit \rightarrow wavefn' is localized in all coordinates

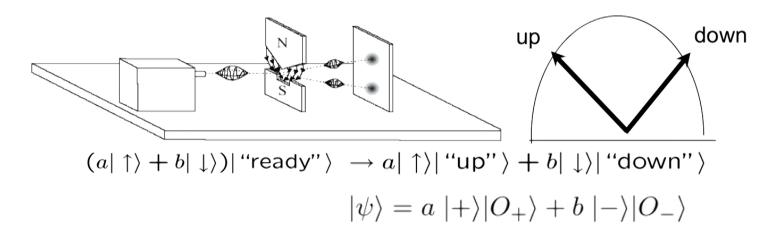
Pirsa: 18020064 Page 20/62



Suppose that for macroscopic # of components k, $|\langle O_+|{f r}_k\rangle|^2|\langle O_-|{f r}_k\rangle|^2pprox 0$

One particle is hit → wavefn' is localized in all coordinates

$$\begin{split} |\psi\rangle &\mapsto |\psi_+\rangle \approx \ |+\rangle |O'_+\rangle &\quad \text{with probability} &\quad |a|^2 \\ |\psi\rangle &\mapsto |\psi_-\rangle \approx \ |-\rangle |O'_-\rangle &\quad \text{with probability} &\quad |b|^2 \end{split}$$



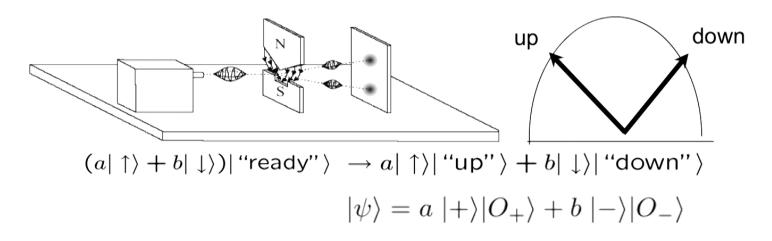
Suppose that for macroscopic # of components k, $|\langle O_+|{\bf r}_k\rangle|^2|\langle O_-|{\bf r}_k\rangle|^2pprox 0$

One particle is hit \rightarrow wavefn' is localized in all coordinates

$$|\psi\rangle\mapsto|\psi_{+}\ranglepprox |+\rangle|O'_{+}\rangle$$
 with probability $|a|^2$ $|\psi\rangle\mapsto|\psi_{-}\ranglepprox |-\rangle|O'_{-}\rangle$ with probability $|b|^2$

If apparatus contains $\approx 10^{20}$ particles

This happens every
$$\frac{10^{15}\mathrm{S}}{10^{20}}pprox10^{-5}\mathrm{s}$$



Suppose that for macroscopic # of components k, $|\langle O_+|{f r}_k\rangle|^2|\langle O_-|{f r}_k\rangle|^2pprox 0$

One particle is hit → wavefn' is localized in all coordinates

$$|\psi\rangle\mapsto|\psi_{+}\ranglepprox\ |+\rangle|O'_{+}\rangle$$
 with probability $|a|^2$ $|\psi\rangle\mapsto|\psi_{-}\ranglepprox\ |-\rangle|O'_{-}\rangle$ with probability $|b|^2$

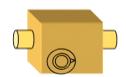
If apparatus contains $\approx 10^{20}$ particles

This happens every
$$\frac{10^{15} \mathrm{S}}{10^{20}} pprox 10^{-5} \mathrm{s}$$

The apparatus gets determinate properties And we recover the collapse postulate

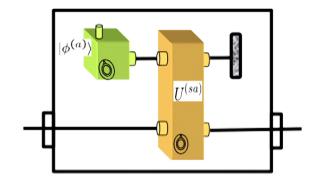
Any anomolous decoherence is also consistent with unitary coupling to novel degrees of freedom

$$\rho \to \sum_{k} \int \mathrm{d}q_k \ \sqrt{E^{(k)}(\mathbf{q}_k)} \ \rho \ \sqrt{E^{(k)}(\mathbf{q}_k)}$$



Trace-preserving completely positive linear map (CP map)

 \mathcal{T}

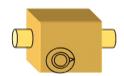


Stinespring dilation

$$|\phi^{(a)}\rangle, U^{(sa)}$$

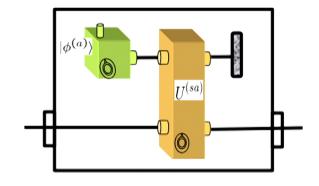
Any anomolous decoherence is also consistent with unitary coupling to novel degrees of freedom

$$\rho \to \sum_{k} \int \mathrm{d}q_k \ \sqrt{E^{(k)}(\mathbf{q}_k)} \ \rho \ \sqrt{E^{(k)}(\mathbf{q}_k)}$$



Trace-preserving completely positive linear map (CP map)

 \mathcal{T}



Stinespring dilation

$$|\phi^{(a)}\rangle, U^{(sa)}$$

The "smoking gun" experimental signature of dynamical collapse does not rule decisively in its favour

How localized is localized enough? The "tails" problem.

$$\psi(x_1, x_2) = \frac{\sqrt{3}}{2} \phi_a(x_1) \chi_a(x_2) + \frac{1}{2} \phi_b(x_1) \chi_b(x_2)$$

$$\times q_2$$

$$x_1$$

$$x_2$$

$$\psi'(x_1, x_2) \approx \phi_a(x_1) \chi'_a(x_2)$$

$$x_1$$

Failure of energy conservation

$$\psi(x_1, x_2) = \frac{\sqrt{3}}{2} \phi_a(x_1) \chi_a(x_2) + \frac{1}{2} \phi_b(x_1) \chi_b(x_2)$$

$$\times q_2$$

$$x_1$$

$$x_2$$

$$\psi'(x_1, x_2) \approx \phi_a(x_1) \chi'_a(x_2)$$

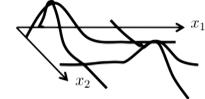
$$x_1$$

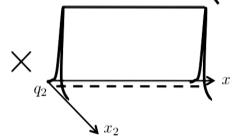
$$x_2$$

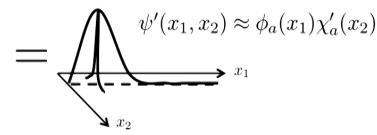
$$x_3$$

The fact that a hit on one system leads to an approximate localization of a distant system is a failure of local causality and of Lorentz invariance

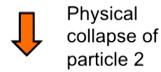
$$\psi(x_1, x_2) = \frac{\sqrt{3}}{2} \phi_a(x_1) \chi_a(x_2) + \frac{1}{2} \phi_b(x_1) \chi_b(x_2)$$

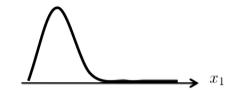












Pirsa: 18020064

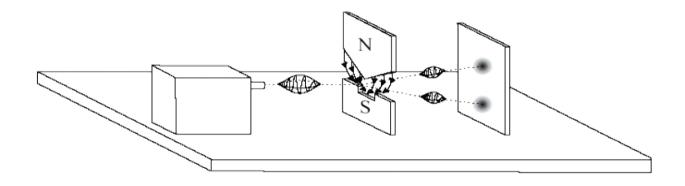
The Everett interpretation "Many Worlds"



Hugh Everett, III (1930-1982)

Pirsa: 18020064 Page 29/62

The quantum measurement problem



$$\begin{array}{c} (a|\uparrow\rangle+b|\downarrow\rangle)|\text{ "ready"}\,\rangle\\ \\ \to a|\uparrow\rangle|\text{ "up"}\,\rangle+b|\downarrow\rangle|\text{ "down"}\,\rangle \end{array}$$

Pirsa: 18020064 Page 30/62

Inconsistencies of the orthodox interpretation

By the collapse postulate (applied to the system)

By unitary evolution postulate (applied to isolated system that includes the apparatus)

Indeterministic and discontinuous evolution

Determinate properties

Deterministic and continuous evolution

Indeterminate properties



Subjective illusion

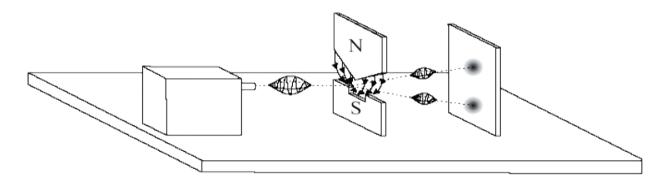
Pirsa: 18020064 Page 31/62

Responses to the measurement problem

- 1. Deny realism
 - Purely operational account of quantum theory
- 2. Deny the universality of unitary dynamics
 - Dynamical collapse theories
- 3. Deny that ψ is a complete representation of reality
 - Hidden variable models
 - Models of reality beyond hidden variables?
- 4. Deny indeterminism and discontinuity, except as subjective illusions
 - Everett's relative state interpretation, or "many worlds"

Pirsa: 18020064 Page 32/62

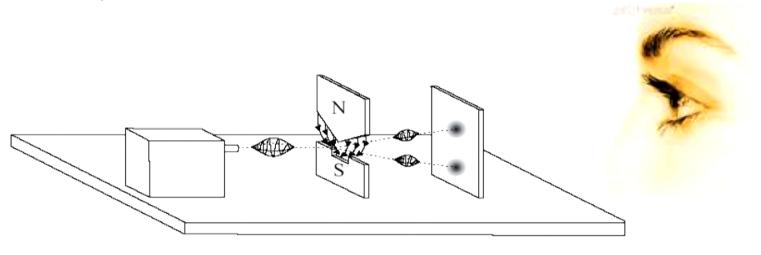
Quantum measurement



$$\begin{array}{c} (a|\uparrow\rangle + b|\downarrow\rangle)|\text{ "ready"}\,\rangle \\ \rightarrow a|\uparrow\rangle|\text{ "up"}\,\rangle + b|\downarrow\rangle|\text{ "down"}\,\rangle \end{array}$$

Pirsa: 18020064 Page 33/62

Quantum measurement with observer

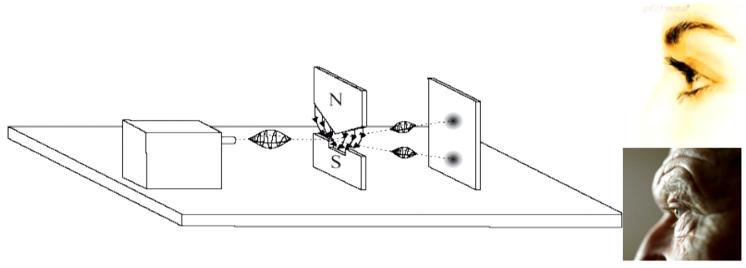


 $(a|\uparrow\rangle + b|\downarrow\rangle)|$ "ready" $\rangle|$ "ready to observe" \rangle

 $\rightarrow a|\uparrow\rangle|$ "up" $\rangle|$ "observe up" $\rangle+b|\downarrow\rangle|$ "down" $\rangle|$ "observe down" $\rangle|$

Pirsa: 18020064 Page 34/62

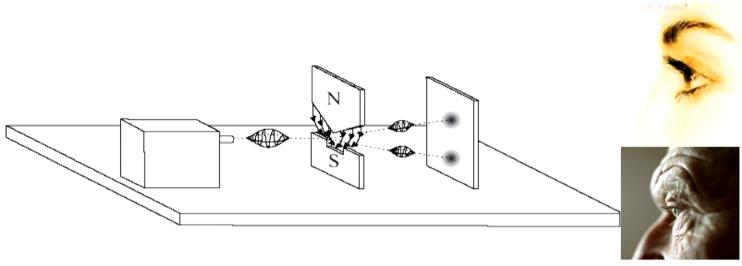
Quantum measurement with many observers



 $(a|\uparrow\rangle+b|\downarrow\rangle)|$ "ready" $\rangle|$ "ready to observe" $\rangle|$ "ready to observe" $\rangle|$

Pirsa: 18020064 Page 35/62

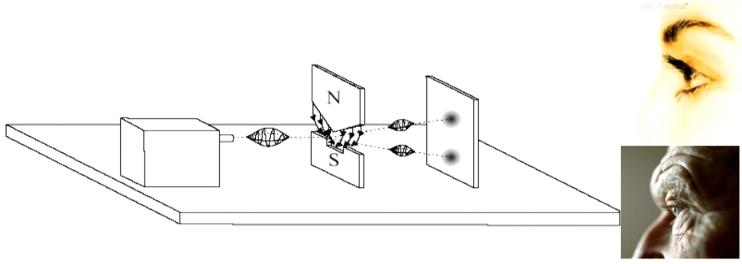
Quantum measurement with many observers



 $(a|\uparrow\rangle+b|\downarrow\rangle)|$ "ready" $\rangle|$ "ready to observe" $\rangle|$ "ready to observe" $\rangle|$ $\rightarrow (a|\uparrow\rangle|$ "up" $\rangle|$ "observe up" $\rangle+a|\downarrow\rangle|$ "down" $\rangle|$ "observe down" $\rangle|$ $\otimes|$ "ready to observe" $\rangle|$

Pirsa: 18020064 Page 36/62

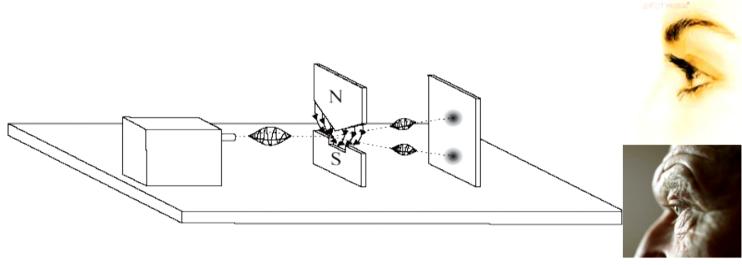
Quantum measurement with many observers



```
(a|\uparrow\rangle+b|\downarrow\rangle)| "ready" \rangle| "ready to observe" \rangle| "ready to observe" \rangle
\to (a|\uparrow\rangle| "up" \rangle| "observe up" \rangle+a|\downarrow\rangle| "down" \rangle| "observe down" \rangle)
\otimes| "ready to observe" \rangle
\to a|\uparrow\rangle| "up" \rangle| "observe up" \rangle| "observe up" \rangle
+b|\downarrow\rangle| "down" \rangle| "observe down" \rangle| "observe down" \rangle
```

Pirsa: 18020064 Page 37/62

Quantum measurement with many observers

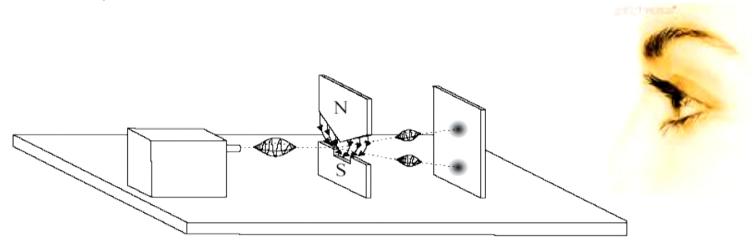


```
(a|\uparrow\rangle+b|\downarrow\rangle)| "ready"\rangle| "ready to observe"\rangle| "ready to observe"\rangle| \rightarrow (a|\uparrow\rangle| "up"\rangle| "observe up"\rangle+a|\downarrow\rangle| "down"\rangle| "observe down"\rangle| \otimes| "ready to observe"\rangle| \rightarrow a|\uparrow\rangle| "up"\rangle| "observe up"\rangle| "observe up"\rangle| "observe down"\rangle| "observe down"\rangle|
```

There is intersubjective agreement about the outcome

Pirsa: 18020064 Page 38/62

Quantum measurement with observer



$$(a|\uparrow\rangle+b|\downarrow\rangle)|$$
 "ready" $\rangle|$ "ready to observe" \rangle

$$\rightarrow a|\uparrow\rangle|$$
 "up" $\rangle|$ "observe up" $\rangle + b|\downarrow\rangle|$ "down" $\rangle|$ "observe down" $\rangle|$

again rewrite as

$$(a|+\rangle+b|-\rangle)|O_0\rangle \rightarrow a|+\rangle|O_+\rangle+b|-\rangle|O_-\rangle$$

Pirsa: 18020064 Page 39/62

$$(a|+\rangle+b|-\rangle)|O_{0}\rangle \implies a|+\rangle|O_{+}\rangle+b|-\rangle|O_{-}\rangle$$

$$= \left(\frac{a|+\rangle+b|-\rangle}{\sqrt{2}}\right)\left(\frac{|O_{+}\rangle+|O_{-}\rangle}{\sqrt{2}}\right)$$

$$+ \left(\frac{a|+\rangle-b|-\rangle}{\sqrt{2}}\right)\left(\frac{|O_{+}\rangle-|O_{-}\rangle}{\sqrt{2}}\right)$$

Adrian Kent: Everett needs a basis selection rule

Pirsa: 18020064 Page 40/62

$$(a|+\rangle|O_{+}\rangle+b|-\rangle|O_{-}\rangle)|E_{0}\rangle \rightarrow a|+\rangle|O_{+}\rangle|E_{+}\rangle+b|-\rangle|O_{-}\rangle|E_{-}\rangle$$

Pirsa: 18020064

$$\begin{split} \left(a\big|+\big\rangle\!\big|O_{\scriptscriptstyle{+}}\big\rangle\!+b\big|-\big\rangle\!\big|O_{\scriptscriptstyle{-}}\big\rangle\big)\big|E_{\scriptscriptstyle{0}}\big\rangle & \to a\big|+\big\rangle\!\big|O_{\scriptscriptstyle{+}}\big\rangle\!\big|E_{\scriptscriptstyle{+}}\big\rangle\!+b\big|-\big\rangle\!\big|O_{\scriptscriptstyle{-}}\big\rangle\!\big|E_{\scriptscriptstyle{-}}\big\rangle \\ & = \left(\frac{a\big|+\big\rangle\!\big|E_{\scriptscriptstyle{+}}\big\rangle\!+b\big|-\big\rangle\!\big|E_{\scriptscriptstyle{-}}\big\rangle}{\sqrt{2}}\right)\!\left(\frac{\big|O_{\scriptscriptstyle{+}}\big\rangle\!+\big|O_{\scriptscriptstyle{-}}\big\rangle}{\sqrt{2}}\right) \\ & + \left(\frac{a\big|+\big\rangle\!\big|E_{\scriptscriptstyle{+}}\big\rangle\!-b\big|-\big\rangle\!\big|E_{\scriptscriptstyle{-}}\big\rangle}{\sqrt{2}}\right)\left(\frac{\big|O_{\scriptscriptstyle{+}}\big\rangle\!-\big|O_{\scriptscriptstyle{-}}\big\rangle}{\sqrt{2}}\right) \end{split}$$

Pirsa: 18020064 Page 42/62

$$\begin{split} \left(a\big|+\big\rangle\big|O_{+}\big\rangle+b\big|-\big\rangle\big|O_{-}\big\rangle\Big)\big|E_{0}\big\rangle & \to a\big|+\big\rangle\big|O_{+}\big\rangle\big|E_{+}\big\rangle+b\big|-\big\rangle\big|O_{-}\big\rangle\big|E_{-}\big\rangle \\ & = \left(\frac{a\big|+\big\rangle\big|E_{+}\big\rangle+b\big|-\big\rangle\big|E_{-}\big\rangle}{\sqrt{2}}\right)\!\left(\frac{\big|O_{+}\big\rangle+\big|O_{-}\big\rangle}{\sqrt{2}}\right) \\ & + \left(\frac{a\big|+\big\rangle\big|E_{+}\big\rangle-b\big|-\big\rangle\big|E_{-}\big\rangle}{\sqrt{2}}\right)\!\left(\frac{\big|O_{+}\big\rangle-\big|O_{-}\big\rangle}{\sqrt{2}}\right) \end{split}$$

$$|O_{+}\rangle|E_{0}^{(2)}\rangle \rightarrow |O_{+}\rangle|E_{+}^{(2)}\rangle$$

Pirsa: 18020064

$$\begin{split} \left(a\big|+\big\rangle\big|O_{+}\big\rangle+b\big|-\big\rangle\big|O_{-}\big\rangle\right)\big|E_{0}\big\rangle & \longrightarrow a\big|+\big\rangle\big|O_{+}\big\rangle\big|E_{+}\big\rangle+b\big|-\big\rangle\big|O_{-}\big\rangle\big|E_{-}\big\rangle \\ & = \left(\frac{a\big|+\big\rangle\big|E_{+}\big\rangle+b\big|-\big\rangle\big|E_{-}\big\rangle}{\sqrt{2}}\right)\!\!\left(\frac{\big|O_{+}\big\rangle+\big|O_{-}\big\rangle}{\sqrt{2}}\right) \\ & \quad + \left(\frac{a\big|+\big\rangle\big|E_{+}\big\rangle-b\big|-\big\rangle\big|E_{-}\big\rangle}{\sqrt{2}}\right)\!\!\left(\frac{\big|O_{+}\big\rangle-\big|O_{-}\big\rangle}{\sqrt{2}}\right) \\ & \quad |O_{+}\big\rangle\big|E_{0}^{(2)}\big\rangle \longrightarrow \big|O_{+}\big\rangle\big|E_{+}^{(2)}\big\rangle \\ & \quad \left(\frac{\big|O_{+}\big\rangle+\big|O_{-}\big\rangle}{\sqrt{2}}\right)\!\!\Big|E_{0}^{(2)}\big\rangle \longrightarrow \left(\frac{\big|O_{+}\big\rangle\big|E_{-}^{(2)}\big\rangle+\big|O_{-}\big\rangle\big|E_{-}^{(2)}\big\rangle}{\sqrt{2}}\right) \\ & \quad = \frac{1}{\sqrt{2}}\left(\frac{\big|O_{+}\big\rangle+\big|O_{-}\big\rangle}{\sqrt{2}}\right)\!\!\left(\frac{\big|E_{+}^{(2)}\big\rangle+\big|E_{-}^{(2)}\big\rangle}{\sqrt{2}}\right) + \frac{1}{\sqrt{2}}\left(\frac{\big|O_{+}\big\rangle-\big|O_{-}\big\rangle}{\sqrt{2}}\right)\!\!\left(\frac{\big|E_{+}^{(2)}\big\rangle-\big|E_{-}^{(2)}\big\rangle}{\sqrt{2}}\right) \end{split}$$

Pirsa: 18020064 Page 44/62

$$\begin{split} \left(a\big|+\big\rangle\big|O_{+}\big\rangle+b\big|-\big\rangle\big|O_{-}\big\rangle\Big)\big|E_{0}\big\rangle & \longrightarrow a\big|+\big\rangle\big|O_{+}\big\rangle\big|E_{+}\big\rangle+b\big|-\big\rangle\big|O_{-}\big\rangle\big|E_{-}\big\rangle \\ & = \left(\frac{a\big|+\big\rangle\big|E_{+}\big\rangle+b\big|-\big\rangle\big|E_{-}\big\rangle}{\sqrt{2}}\right)\!\left(\frac{\big|O_{+}\big\rangle+\big|O_{-}\big\rangle}{\sqrt{2}}\right) \\ & + \left(\frac{a\big|+\big\rangle\big|E_{+}\big\rangle-b\big|-\big\rangle\big|E_{-}\big\rangle}{\sqrt{2}}\right)\!\left(\frac{\big|O_{+}\big\rangle-\big|O_{-}\big\rangle}{\sqrt{2}}\right) \end{split}$$

 $\left|O_{_{+}}\right\rangle \left|E_{_{0}}^{(2)}\right\rangle
ightarrow \left|O_{_{+}}\right\rangle \left|E_{_{+}}^{(2)}\right\rangle$ Evolution in this basis is predictable

$$\left(\frac{\left|O_{+}\right\rangle + \left|O_{-}\right\rangle}{\sqrt{2}}\right) \left|E_{0}^{(2)}\right\rangle \rightarrow \left(\frac{\left|O_{+}\right\rangle \left|E_{+}^{(2)}\right\rangle + \left|O_{-}\right\rangle \left|E_{-}^{(2)}\right\rangle}{\sqrt{2}}\right) \\
= \frac{1}{\sqrt{2}} \left(\frac{\left|O_{+}\right\rangle + \left|O_{-}\right\rangle}{\sqrt{2}}\right) \left(\frac{\left|E_{+}^{(2)}\right\rangle + \left|E_{-}^{(2)}\right\rangle}{\sqrt{2}}\right) + \frac{1}{\sqrt{2}} \left(\frac{\left|O_{+}\right\rangle - \left|O_{-}\right\rangle}{\sqrt{2}}\right) \left(\frac{\left|E_{+}^{(2)}\right\rangle - \left|E_{-}^{(2)}\right\rangle}{\sqrt{2}}\right)$$

Evolution in this basis is unpredictable

Response (drawn primarily from the work of David Wallace)

No axiom is needed for basis selection because real things (macroscopic objects and worlds) are emergent patterns.



real, but not directly represented in the axioms

The basis picked out by decoherence admits of patterns that have explanatory and predictive power, such as tigers. Such patterns define the ontology.

Patterns are not precisely defined, but this need not detract from their reality (consider a mountain, or a species)

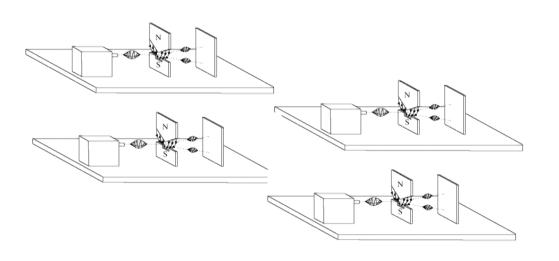
Pirsa: 18020064 Page 46/62

The problem with probabilities

How can anything "be probability" in a deterministic theory where all possible outcomes occur? There are no propensities and there is nothing to be ignorant about.

Pirsa: 18020064 Page 47/62

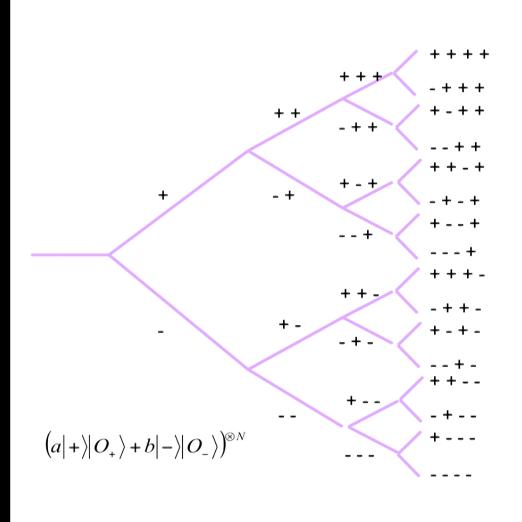
Sequence of measurements:

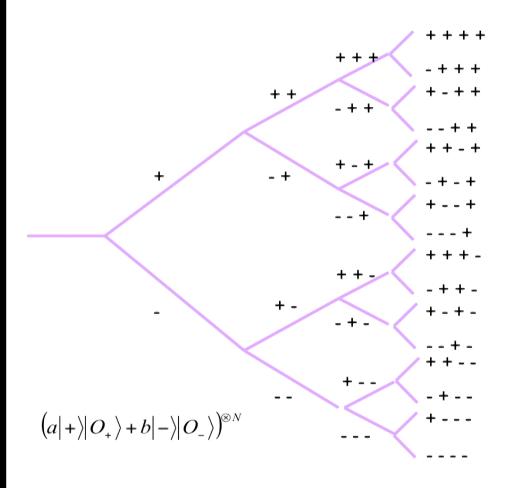




$$(a|+\rangle+b|-\rangle)^{IV}(a|+\rangle+b|-\rangle)^{II}(a|+\rangle+b|-\rangle)^{II}(a|+\rangle+b|-\rangle)^{I}(a|+\rangle+b|-\rangle)^{I}(a|+\rangle+b|-\rangle)^{I}(a|+\rangle+b|-\rangle)^{I}(a|+\rangle+b|-\rangle)^{I}(a|+\rangle+b|-\rangle)^{I}(a|+\rangle+b|-\rangle)^{I}(a|+\rangle^{I}|O_{+}\rangle^{I}+b|-\rangle^{I}|O_{-}\rangle^{I}) |O_{0}\rangle^{II}|O_{0}\rangle^{II}|O_{0}\rangle^{II}|O_{0}\rangle^{II}$$
...
$$(a|+\rangle|O_{+}\rangle+b|-\rangle|O_{-}\rangle)^{\otimes 4}$$

Pirsa: 18020064

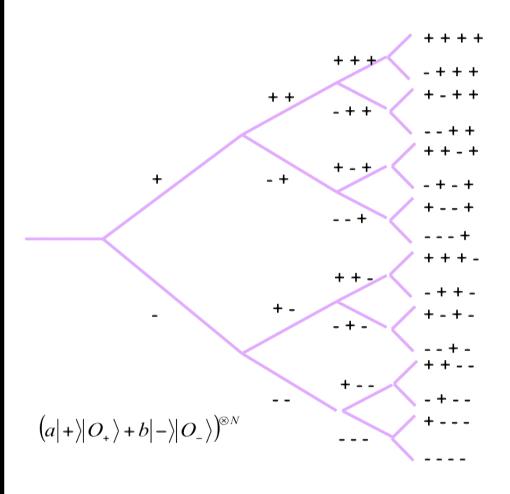




Different branches correspond to different subjective experiences

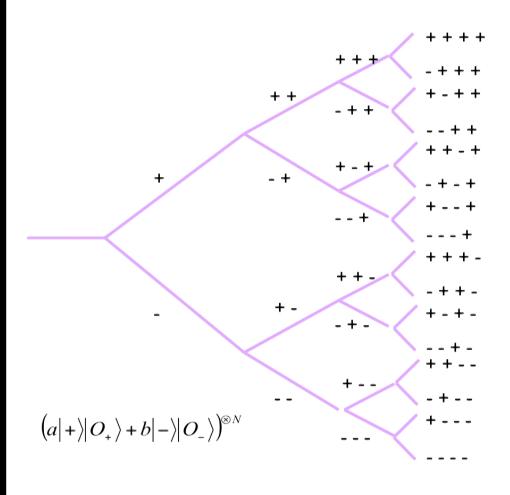
All branches are actual → all experiences occur

Therefore, cannot understand probability in terms of where the "real" me ends up



In the limit $N \rightarrow \infty$, in all branches except a set of measure zero, the frequency of + results is the same.

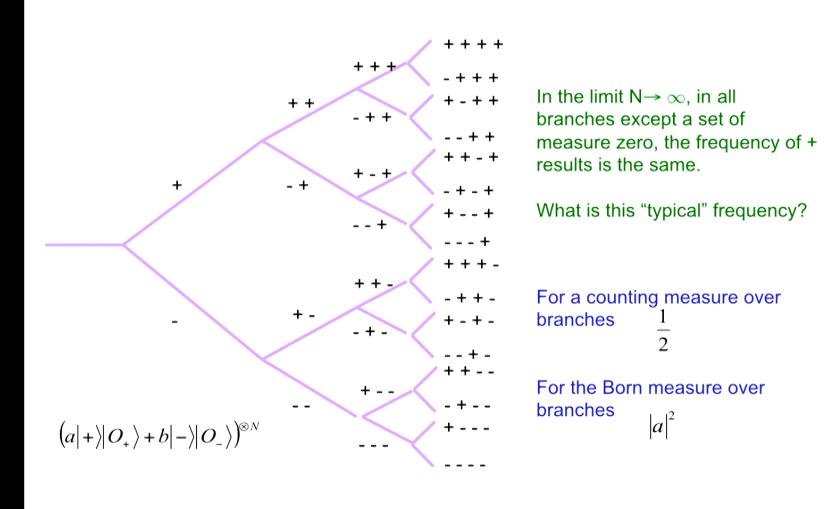
What is this "typical" frequency?



In the limit $N \rightarrow \infty$, in all branches except a set of measure zero, the frequency of + results is the same.

What is this "typical" frequency?

For a counting measure over branches $\frac{1}{2}$



Deutsch's decision-theoretic strategy: Probability gets its meaning through the rational preferences of agents who consider the utility of all future versions of themselves.

Purports to show that a rational agent who knows that the Born-rule weight of an outcome is p is rationally compelled to act as if that outcome had probability p.

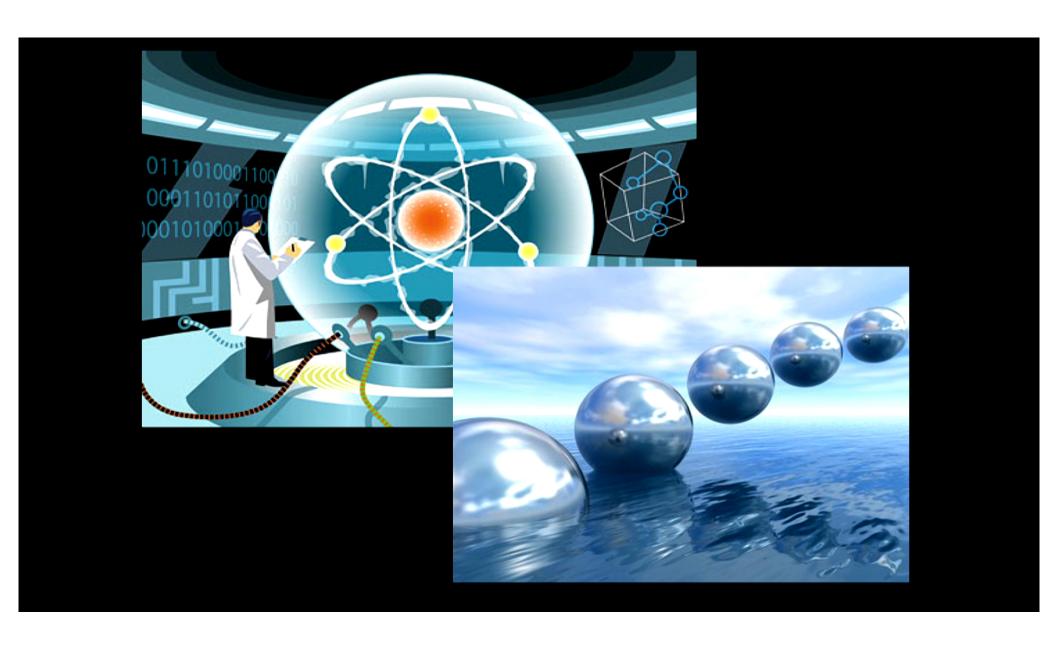
Pirsa: 18020064 Page 54/62

Deutsch's decision-theoretic strategy: Probability gets its meaning through the rational preferences of agents who consider the utility of all future versions of themselves.

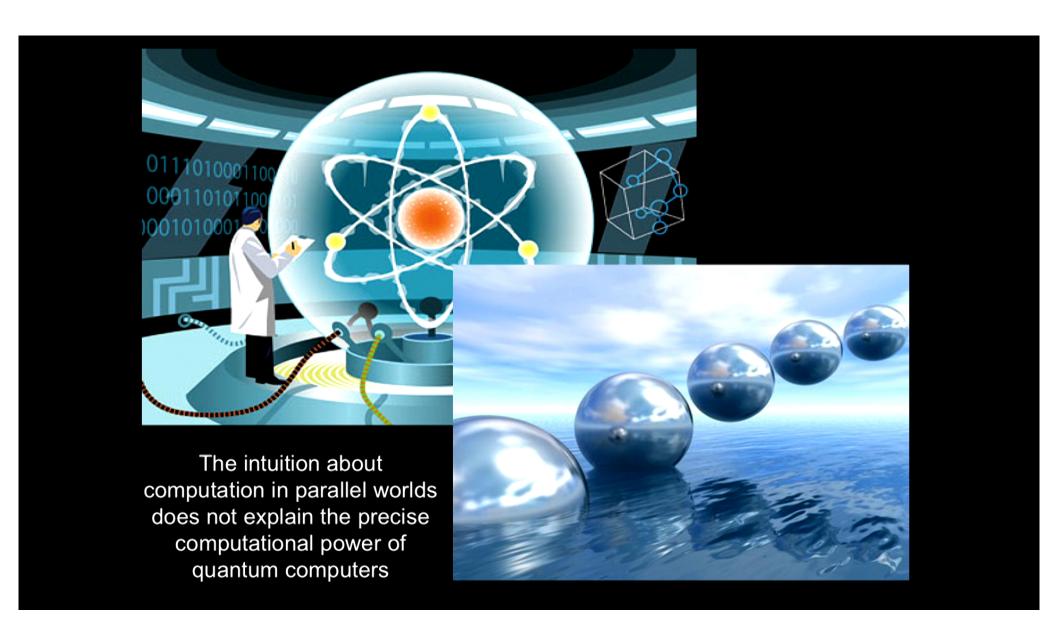
Purports to show that a rational agent who knows that the Born-rule weight of an outcome is p is rationally compelled to act as if that outcome had probability p.

Barnum et al.: Deutsch's proof has a hidden assumption which is akin to applying Laplace's Principle of Insufficient Reason to a set of equal-amplitude alternatives, an application that requires assuming, rather than deriving, how amplitudes are related to probabilities.

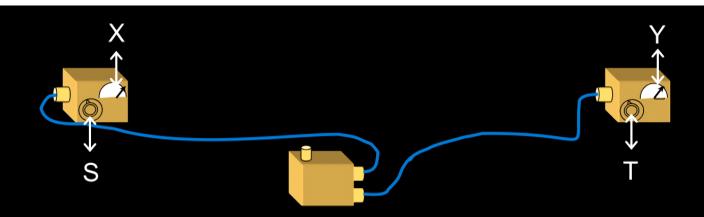
Pirsa: 18020064 Page 55/62



Pirsa: 18020064 Page 56/62



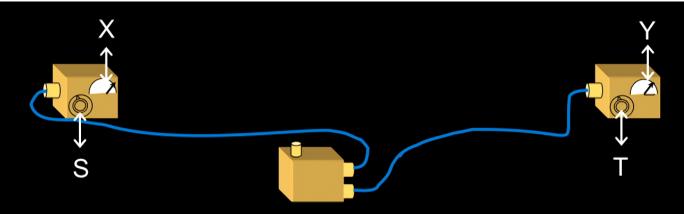
Pirsa: 18020064 Page 57/62



	X=0, Y=0	X=0, Y=1	X=1, Y=0	X=1, Y=1
S=0, T=0	0.427	0.073	0.073	0.427
S=0, T=1	0.427	0.073	0.073	0.427
S=1, T=0	0.427	0.073	0.073	0.427
S=1, T=1	0.073	0.427	0.427	0.073

	X=0, Y=0	X=0, Y=1	X=1, Y=0	X=1, Y=1
S=0, T=0	1/2	0	0	1/2
S=0, T=1	1/2	0	0	1/2
S=1, T=0	1/2	0	0	1/2
S=1, T=1	0	1/2	1/2	0

Pirsa: 18020064 Page 58/62

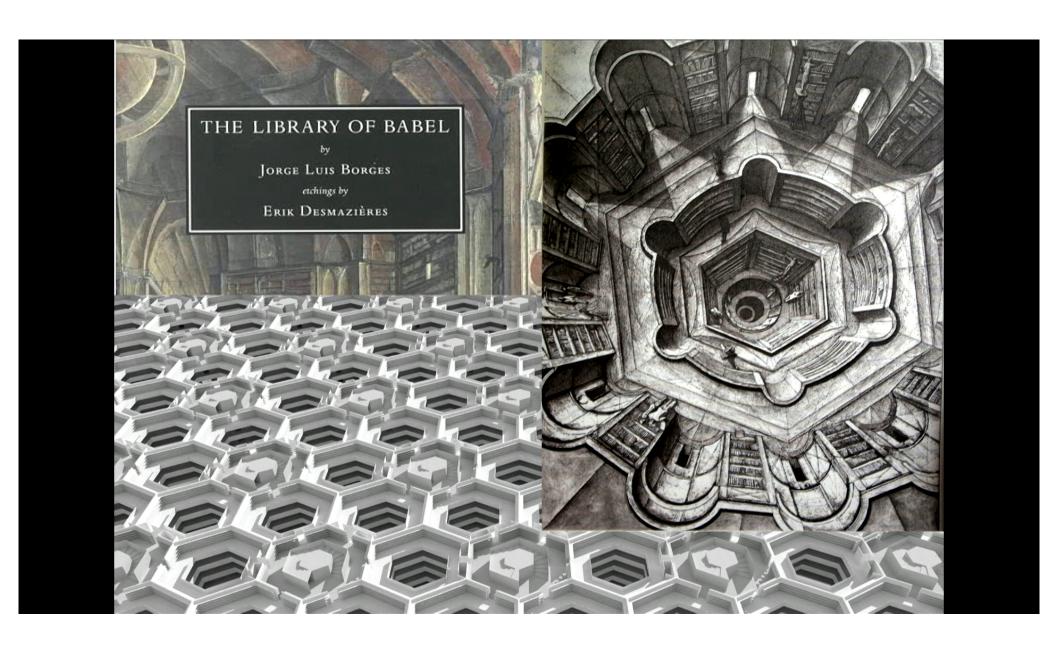


	X=0, Y=0	X=0, Y=1	X=1, Y=0	X=1, Y=1
S=0, T=0	0.427	0.073	0.073	0.427
S=0, T=1	0.427	0.073	0.073	0.427
S=1, T=0	0.427	0.073	0.073	0.427
S=1, T=1	0.073	0.427	0.427	0.073

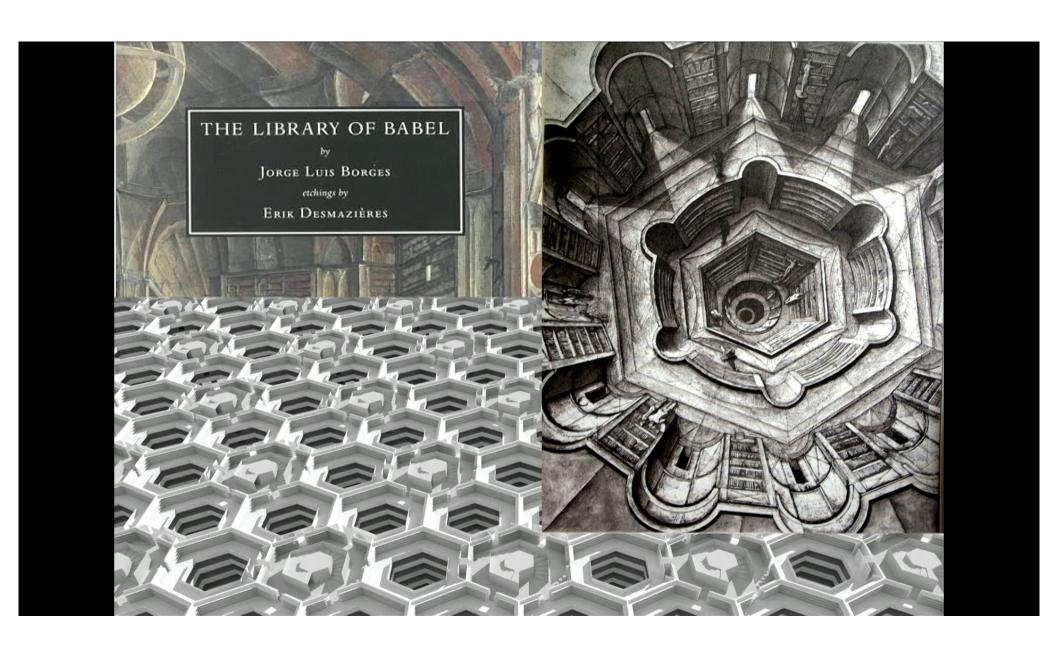
	X=0, Y=0	X=0, Y=1	X=1, Y=0	X=1, Y=1
S=0, T=0	1/2	0	0	1/2
S=0, T=1	1/2	0	0	1/2
S=1, T=0	1/2	0	0	1/2
S=1, T=1	0	1/2	1/2	0

Any generalized probabilistic theory can be given a many worlds interpretation. The interpretation does not explain the precise degree of violation of Bell inequalities seen in quantum theory

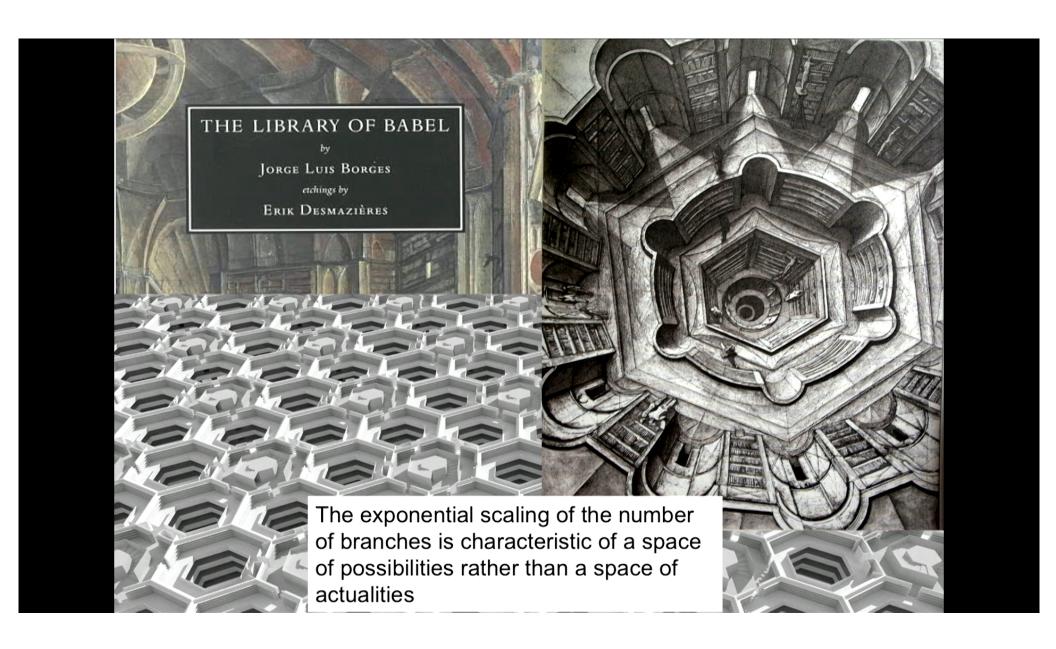
Pirsa: 18020064 Page 59/62



Pirsa: 18020064 Page 60/62



Pirsa: 18020064 Page 61/62



Pirsa: 18020064 Page 62/62