

Title: Constraining a Thin Dark Matter Disk with Gaia

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Abstract: <p>If a component of the dark matter has dissipative interactions, it could collapse to form a thin dark disk in our Galaxy coincident with the baryonic disk. It has been suggested that dark disks could explain a variety of observed phenomena, including periodic comet impacts. Using the first data release from the Gaia mission, we search for a dark disk via its effect on stellar kinematics in the Milky Way. I will present new limits on the presence of a thin dark matter disk, as well as measurements on the matter density in the solar neighborhood.</p>

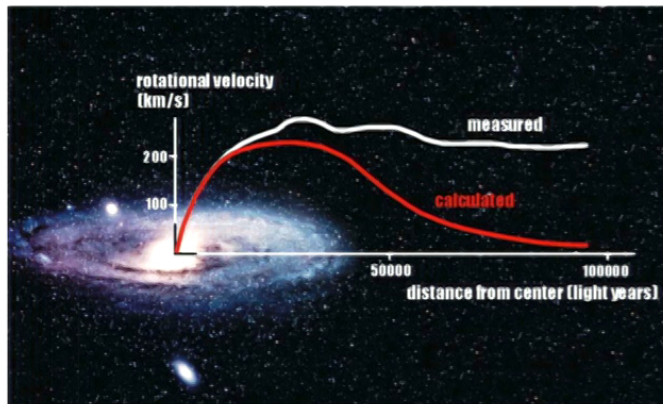
# Constraining a thin dark matter disk with Gaia

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December 12, 2017  
PI Seminar

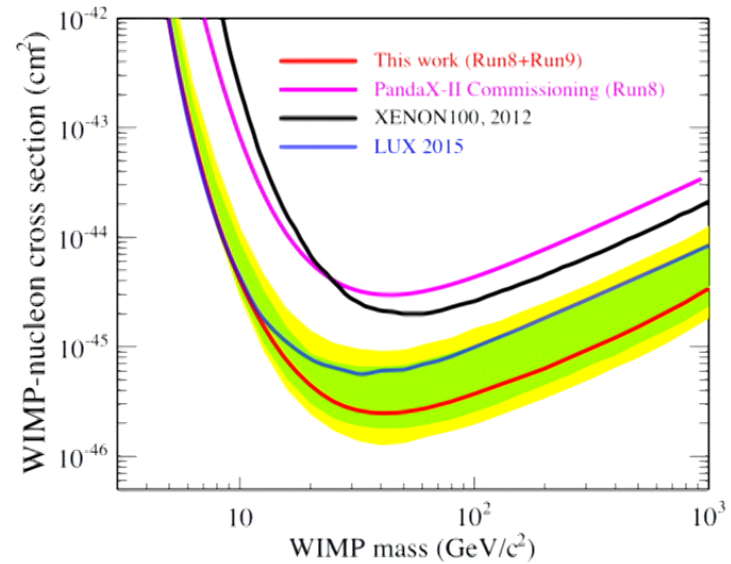
1711.03103 with Katelin Schutz, Ben Safdi, and Chih-Liang Wu

# Particle models for dark matter

Evidence for gravitation interactions



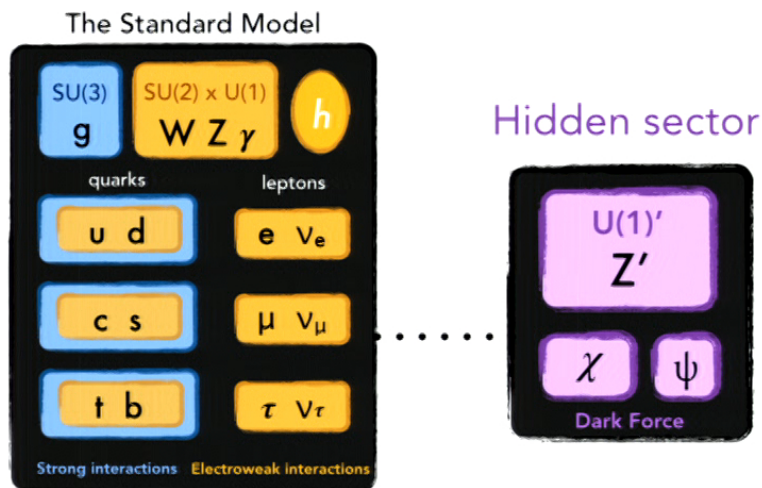
Non-evidence for SM gauge interactions



We can also use gravitational interactions  
to test gauge interactions of DM

# Hidden Sectors

By analogy with the SM, a compelling possibility that dark matter has its own gauge interactions

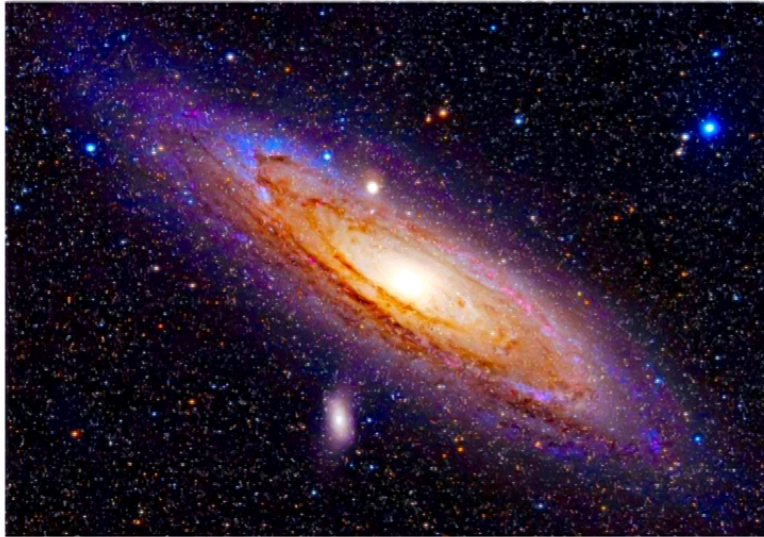


Models containing new light mediators (Z') could explain:

- small scale structure and SIDM
- large scale structure and H0
- Galactic center GeV excess
- ...

where the interacting hidden sector may be a fraction of the total DM

# Why dark disks?



Disk shapes form due to cooling...

By analogy with the Standard Model, if some of the dark matter interacts with "dark forces", this could also lead to cooling and flattening

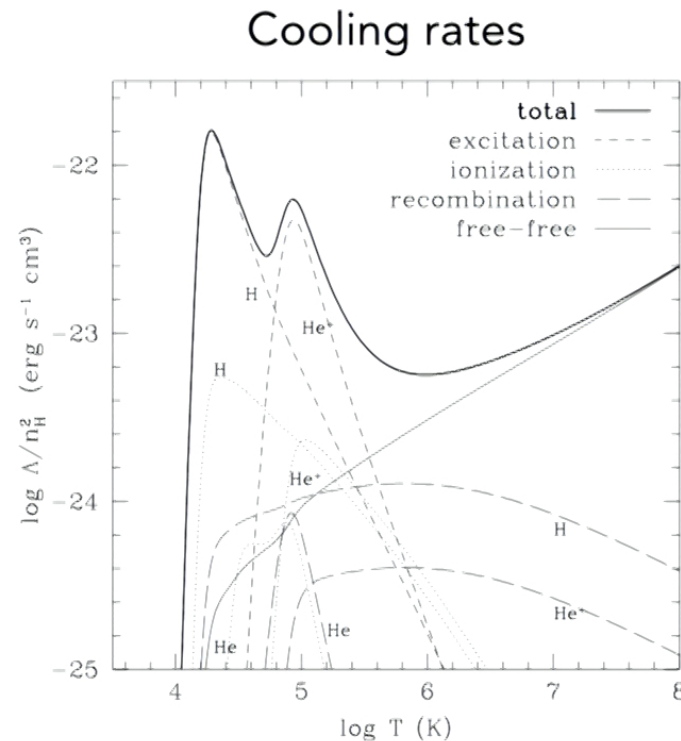
Very rough estimate:

$$z_d \approx 2.5 \text{ pc} \left( \frac{\alpha_D}{0.02} \right)^2 \frac{m_C}{1 \text{ MeV}} \frac{100 \text{ GeV}}{m_X}$$

Thin dark matter disks?

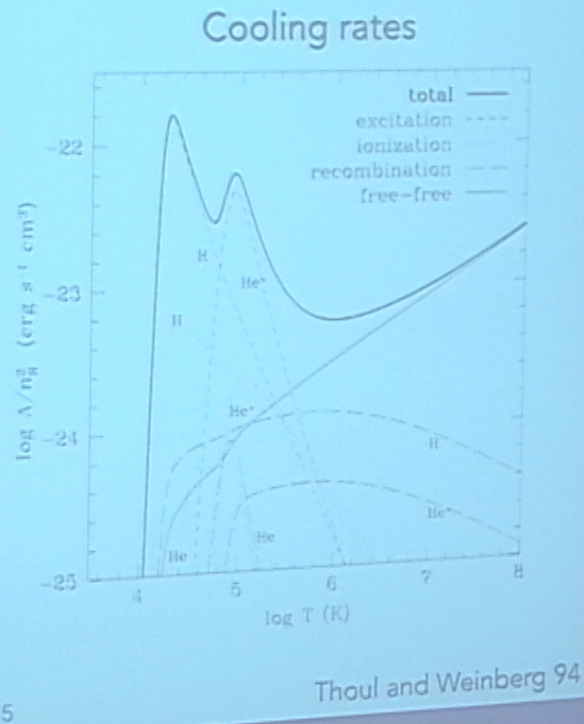
# Conditions for disk formation

- Disk formation: gas collisions produce radiation, which can escape, leading to cooling and collapse
- Fan et al. estimate dark disk scale height using the temperature at which cooling processes complete
- Note: Toomre instability was unaccounted for in estimates of thin dark disk formation



# Conditions for disk formation

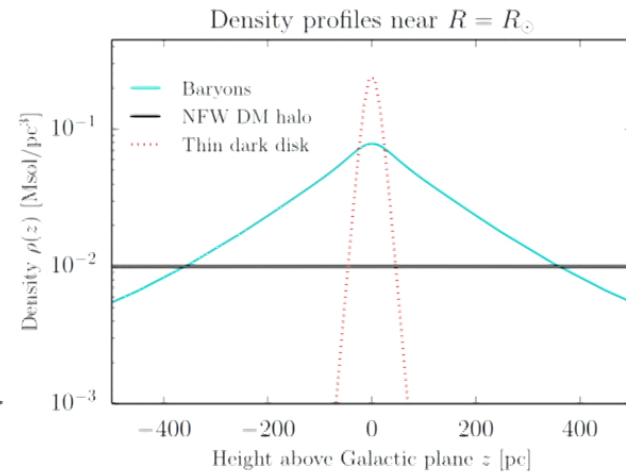
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# Implications of dark disks

- Implications for (in)direct detection (McCullough & Randall, Fan, Katz & Shelton)
- Co-rotation of Andromeda satellites (Randall & Scholtz)
- Periodic disruption of comet trajectories causing mass extinction events (Randall & Reece, Shaviv, Kramer & Ramon)
- Collapsed dark matter objects can account for the point-like nature of the inner galaxy GeV excess (Agrawal & Randall)
- Massive black hole formation (D'Amico et al.)

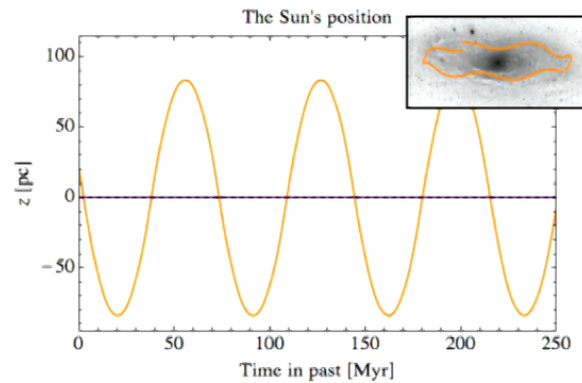


High density dark disk  
considered for these!



# Periodic Comet Impacts

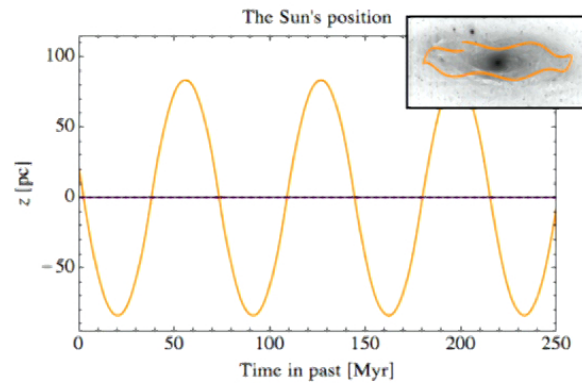
Oscillation of sun about thin dark disk  
can perturb Oort cloud...



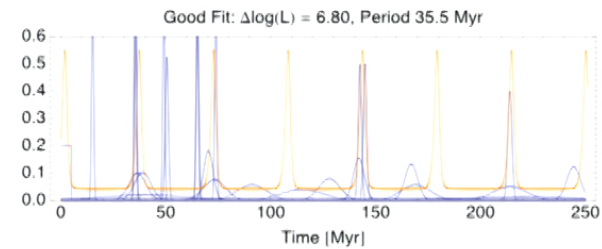
Randall and Reece 2014  
Kramer and Rowan 2016  
Shaviv 2016

# Periodic Comet Impacts

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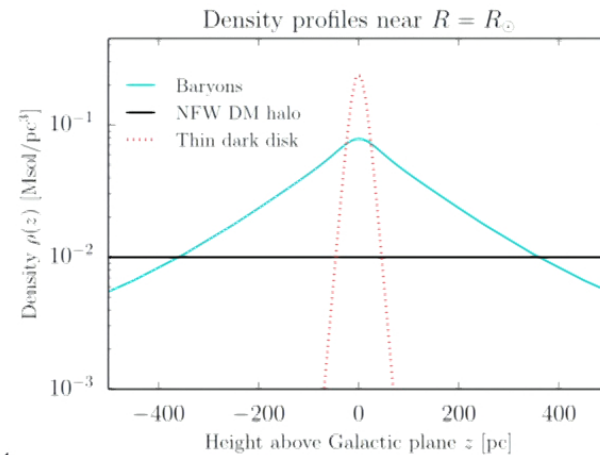
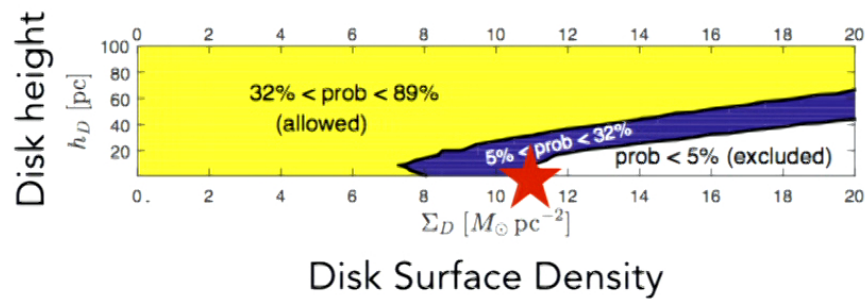


Leading to periodic comet impacts



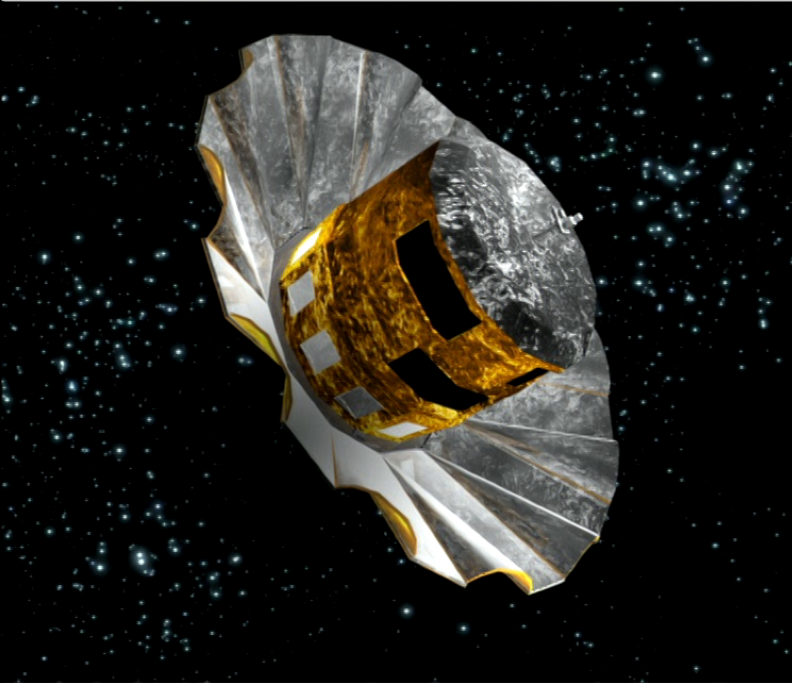
Randall and Reece 2014  
Kramer and Rowan 2016  
Shaviv 2016

# Bounds on a thin dark disk



1604.01407 Eric Kramer and Lisa Randall  
Using kinematics of  $\sim 2000$  stars from Hipparcos

# Gaia satellite



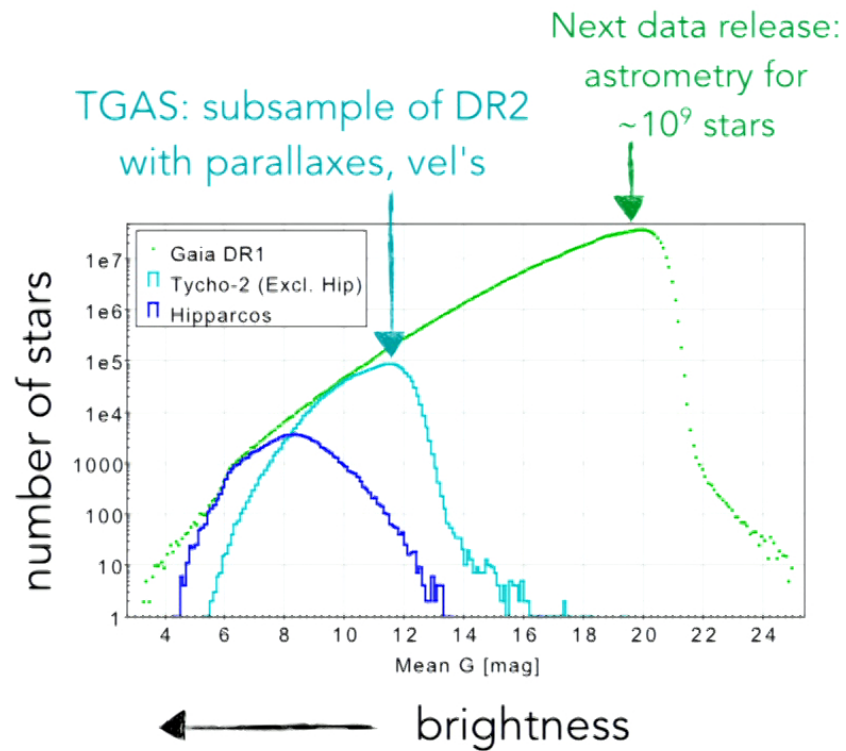
- Full sky view of stellar positions, distances, and velocities
- Launched 2013
- First data release (DR1):  
September 2016 ~  $10^6$  stars
- Next release (DR2):  
April 2018 ~  $10^9$  stars

# Gaia astrometry



With 15  $\mu$ as  
resolution you can  
see a quarter on the  
surface of the moon

# Gaia astrometry

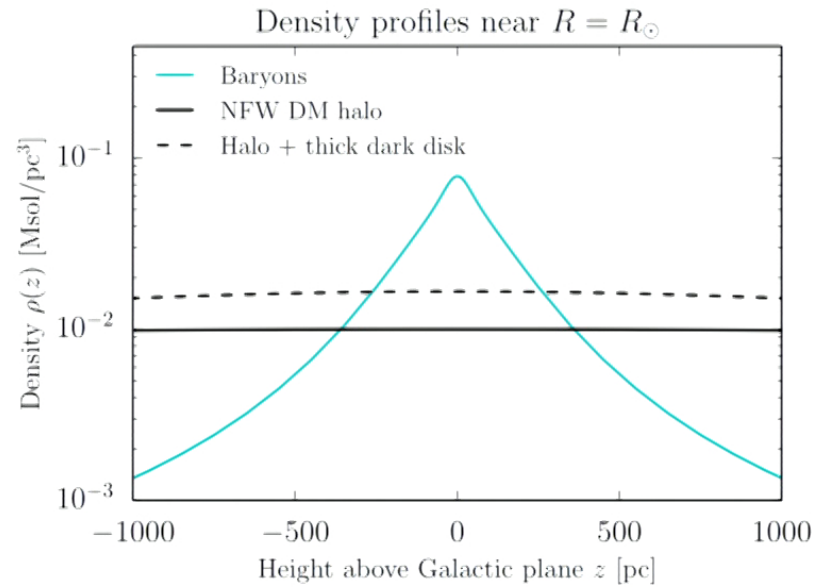
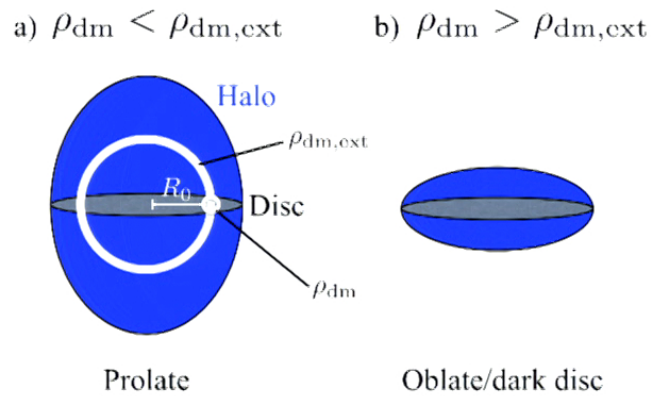


# Dark matter applications

- Local dark matter density *Read 2014, Silverwood et al. 2015*
- Local velocity distribution *Herzog-Arbeitman, Lisanti, Necib 2017*
- Dark matter halo profile *Bovy & Rix 2013 + followups*
- Substructure — dark disks, thick dark matter disks, dark matter subhalos

All important inputs for dark matter  
direct detection, indirect detection

# Thick dark disk



Read et al. 2009  
Silverwood et al. 2015

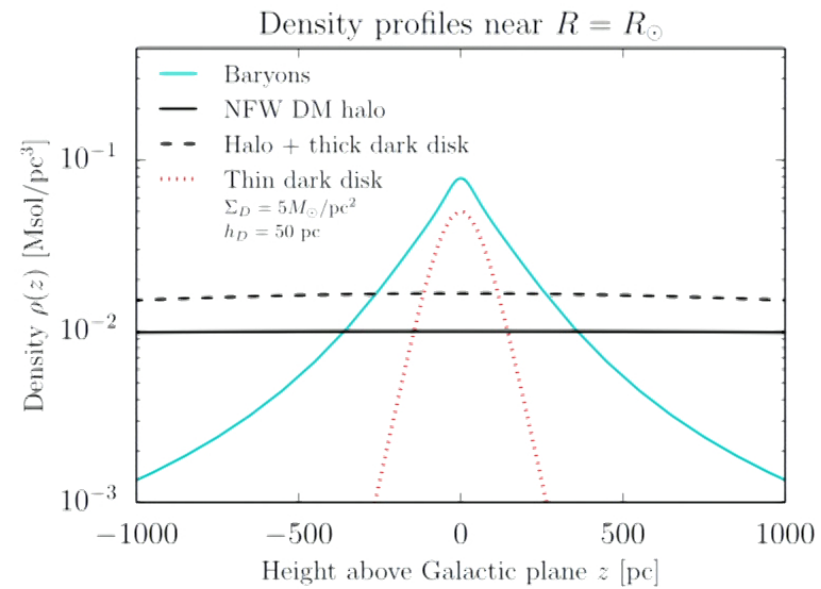


# Thin dark disk

Surface density

$$\rho_{DD}(R_{\odot}, z) = \frac{\Sigma_D}{4h_D} \operatorname{sech}^2\left(\frac{z}{2h_D}\right)$$

Scale height

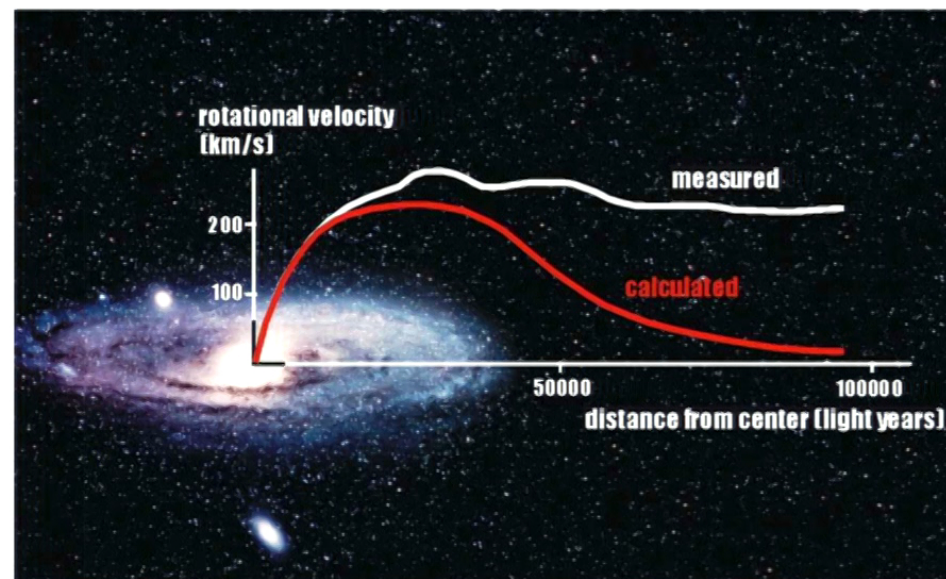


# I. Constraining a dark disk and the local DM density

# Basic idea

Tracer populations (of stars) are sensitive to the total gravitating mass.

Familiar example:



# Basic idea

Tracer populations (of stars) are sensitive to the total gravitating mass.

$$\frac{df_A}{dt} = \partial_t f_A + \partial_{\mathbf{x}} f_A \cdot \mathbf{v} - \partial_{\mathbf{v}} f_A \cdot \partial_{\mathbf{x}} \Phi = 0$$

Under assumptions of separability, equilibrium:

$$v_z \partial_z f_A - \partial_z \Phi \partial_{v_z} f_A = 0$$

For tests of these assumptions,  
see Garbari et al. 2011

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The stellar density is determined by gravitational potential and distribution of mid-plane z-velocities

$$\nu_A(z) = \nu_{A,0} \int dv_z f_{A,0} \left( \sqrt{v_z^2 + 2\Phi(z)} \right)$$

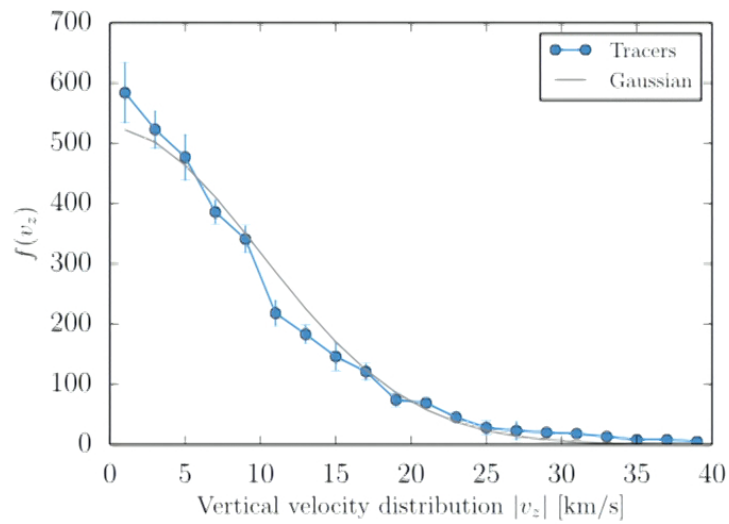
Number density of tracer population

$f(v_z)$  = Distribution of vertical velocities in the plane of the MW.  
Used as input data.

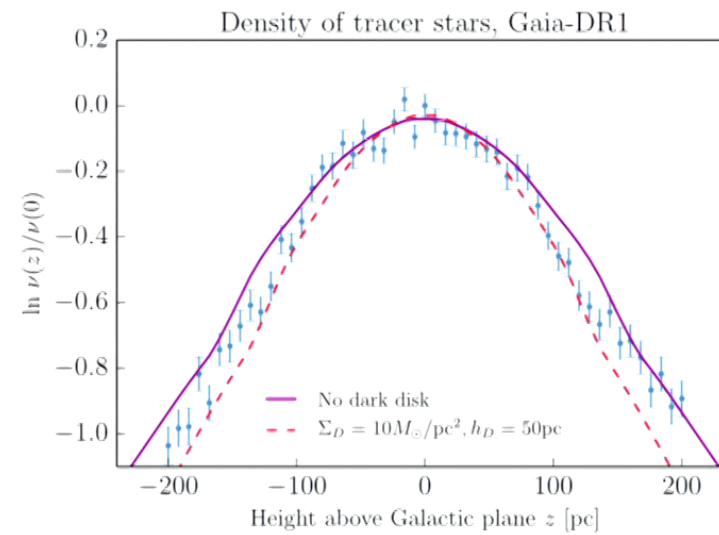
For tests of these assumptions, see Garbari et al. 2011

# Tracer stars

Midplane velocity distribution  $f(v)$



Density fall-off  $v(z)$



# Gravitational Potential

Determining the gravitational potential requires a mass model

Model baryonic mass as a number of isothermal\* components  $\nu_A$

Vertical Jeans equation:  $\frac{\sigma_A^2}{\nu_A} \partial_z \nu_A + \partial_z \Phi = 0 \rightarrow \nu_A(z) = \nu_A(0) e^{-\Phi(z)/\sigma_A^2}$

\*Velocity dispersion doesn't depend on z

Connect this to the mass density profile with Poisson equation

$$\partial_z^2 \Phi = 4\pi G \rho$$

Inject: baryons, smooth dark matter halo, and dark disk

For tests of these assumptions,  
see Garbari et al. 2011

# Baryons

Biggest uncertainties  
are in the gas density

Stellar midplane densities  
are improving with Gaia  
(Bovy 2017)

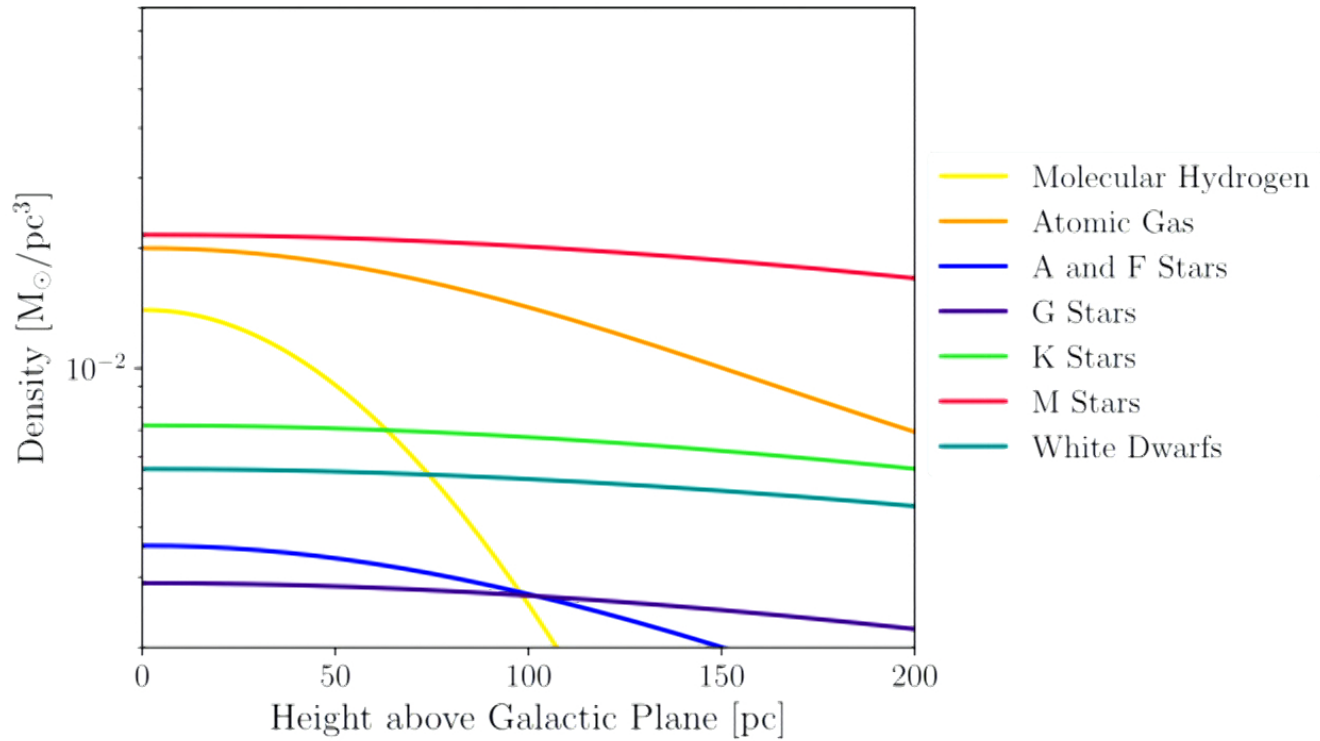
## Isothermal components of Baryonic mass

Baryonic Component	$\rho(0)$ [ $M_{\odot}/\text{pc}^3$ ]	$\sigma$ [km/s]
Molecular Gas ( $\text{H}_2$ )	$0.0104 \pm 0.00312$	$3.7 \pm 0.2$
Cold Atomic Gas ( $\text{H}_I(1)$ )	$0.0277 \pm 0.00554$	$7.1 \pm 0.5$
Warm Atomic Gas ( $\text{H}_I(2)$ )	$0.0073 \pm 0.0007$	$22.1 \pm 2.4$
Hot Ionized Gas ( $\text{H}_{II}$ )	$0.0005 \pm 0.00003$	$39.0 \pm 4.0$
Giant Stars	$0.0006 \pm 0.00006$	$15.5 \pm 1.6$
$M_V < 3$	$0.0018 \pm 0.00018$	$7.5 \pm 2.0$
$3 < M_V < 4$	$0.0018 \pm 0.00018$	$12.0 \pm 2.4$
$4 < M_V < 5$	$0.0029 \pm 0.00029$	$18.0 \pm 1.8$
$5 < M_V < 8$	$0.0072 \pm 0.00072$	$18.5 \pm 1.9$
$M_V > 8$ (M Dwarfs)	$0.0216 \pm 0.0028$	$18.5 \pm 4.0$
White Dwarfs	$0.0056 \pm 0.001$	$20.0 \pm 5.0$
Brown Dwarfs	$0.0015 \pm 0.0005$	$20.0 \pm 5.0$
Total	$0.0889 \pm 0.0071$	—

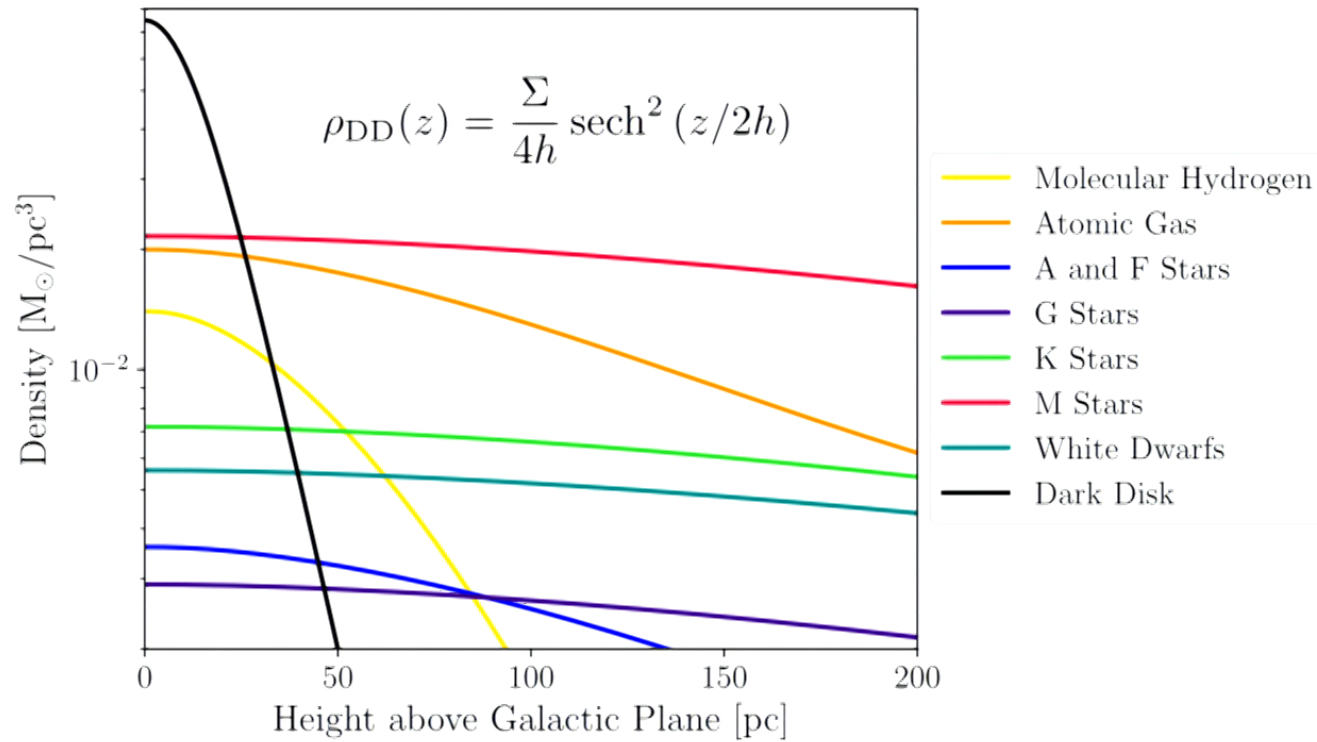
collected from: Flynn et al. 2006, Read 2014,  
McKee et al. 2015, Kramer+Randall 2015, Bovy 2017



# Density profiles

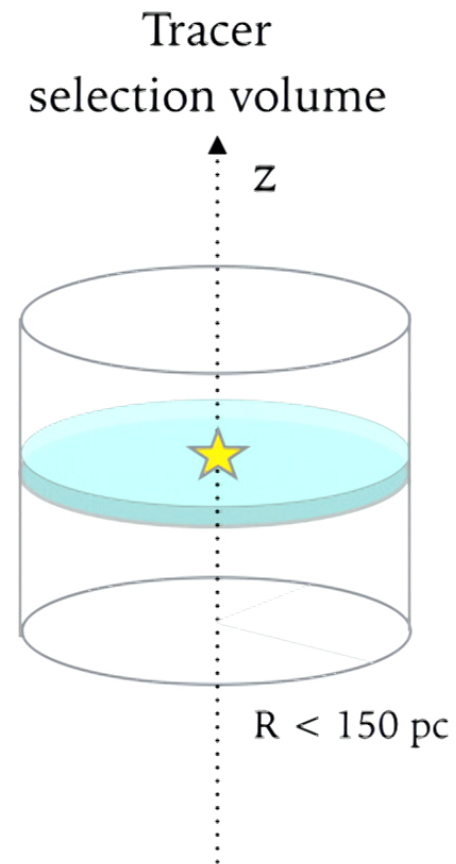


# Density profiles



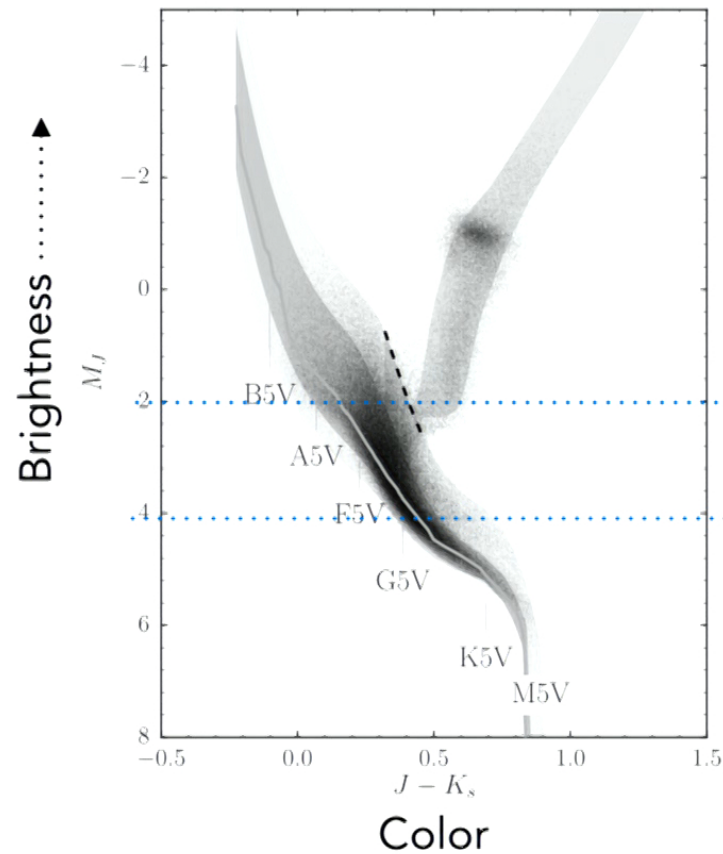
# Analysis

1. Model the baryons at the plane
2. Solve Poisson-Jeans equations for the gravitational potential  $\Phi$ , including dark matter and dark disks
3. Selecting a tracer population of stars, determine the vertical fall-off of stars (using  $\Phi$  and  $f(v)$ )
4. Fit to data, marginalizing over many nuisance parameters



see also: Kramer and Randall, Holmberg and Flynn 2000

# Tracer stars



Gaia measures fainter stars  
out to larger distances

↑ Hipparcos survey (~1990)

↑ Our cuts for Gaia DR1

We divide up tracers into  
three populations:  
A, F, and Early G stars

Bovy 2017

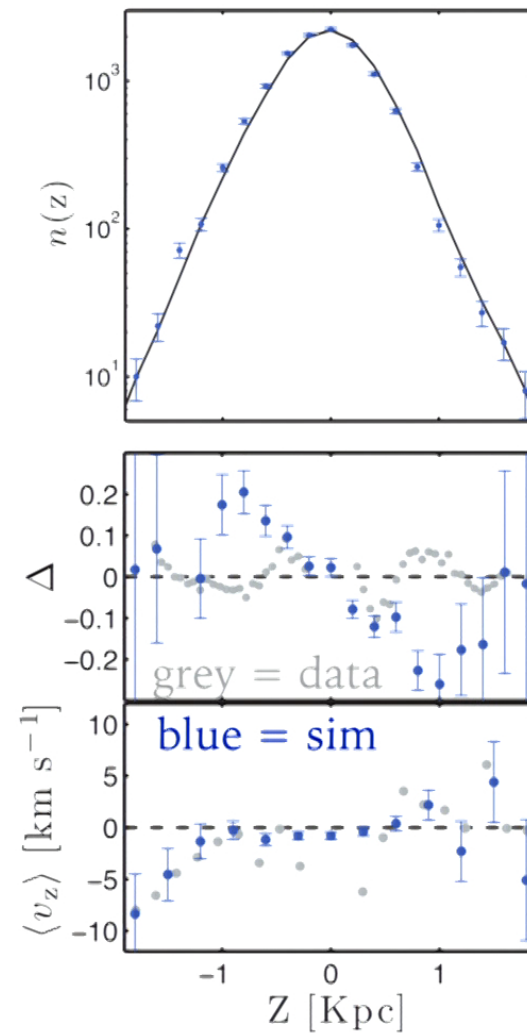
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# Equilibrium?

- The basis of our analysis is that the stellar populations and mass components are close to equilibrium near the mid-plane
- Stellar populations have different lifetimes, and display different departures from equilibrium
- To be conservative, we also include a systematic from asymmetries in  $f(v)$

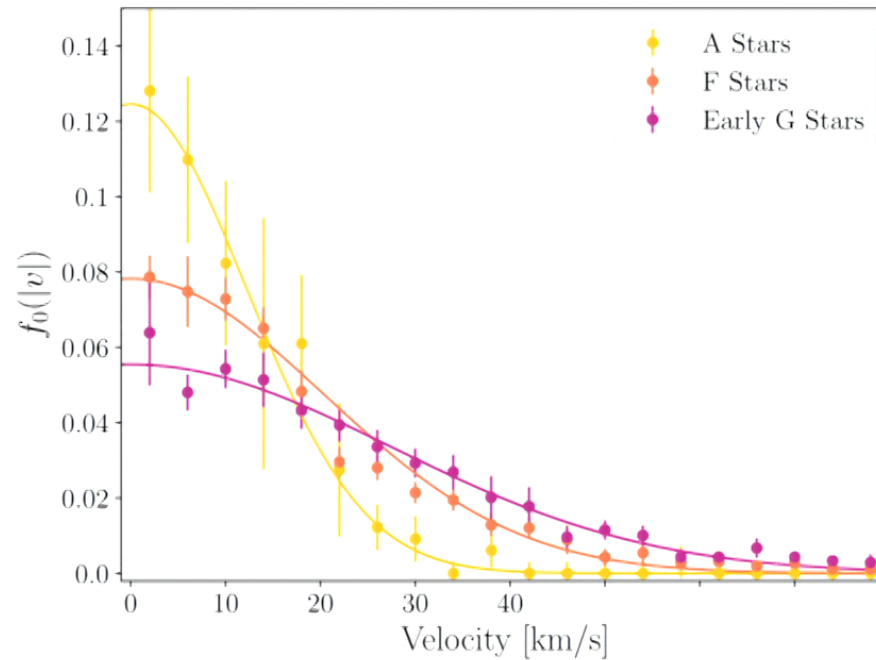
data from Widrow et al. 2012  
Figure from Gomez et al. 2013

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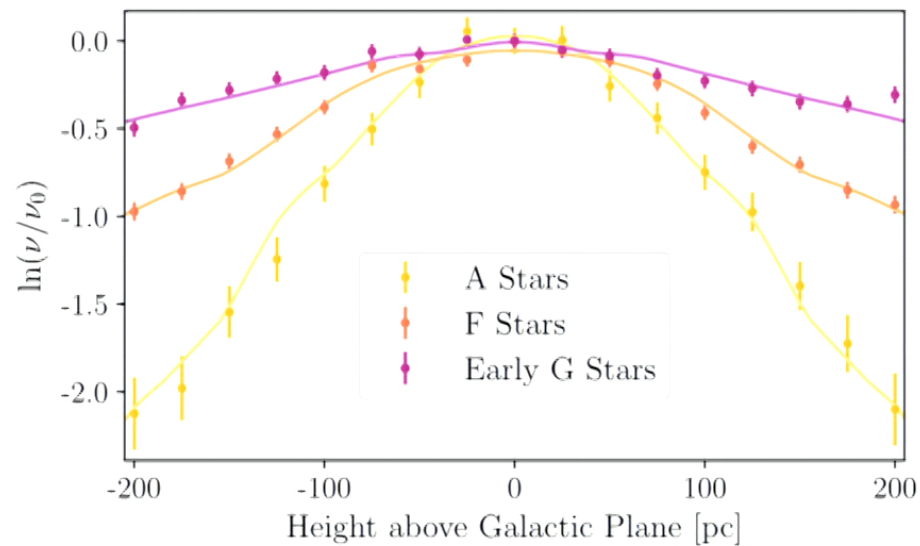


# Velocity distributions

- Best fit  $v_{z,\text{sun}} = 6.8 \pm 0.2$  km/s consistent with independent determinations
- Good fit to Gaussians
- Final systematics are:  
~15%, 5%, 7% for AFG stars



# Number density of tracers



Best fit assuming no dark disk.  
Here  $\rho_{\text{DM}} = 0.014 M_{\odot}/\text{pc}^3 = 0.5 \text{ GeV}/\text{cm}^3$

# Setting limits

In  $L(\chi^2)$  function includes measurements of vertical density profiles;  $f(v)$  for 3 tracer pops; baryons

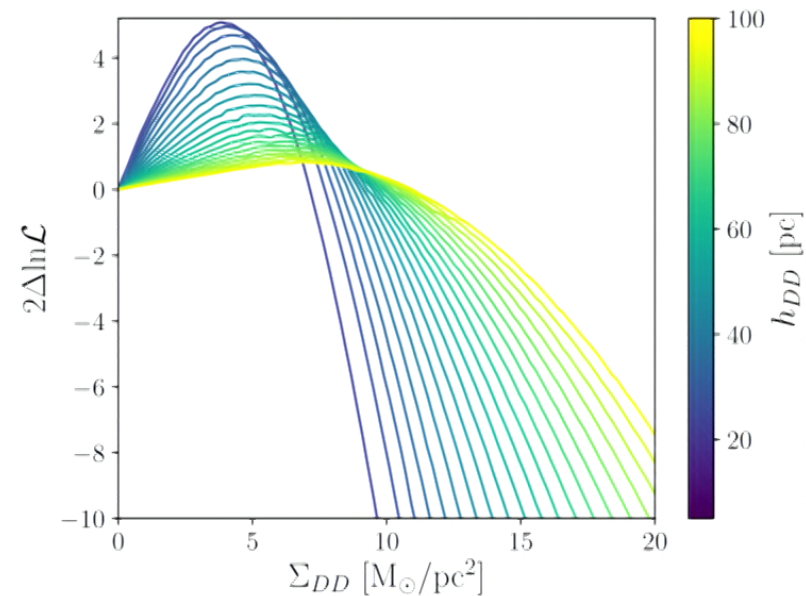
Sample over:

- $\rho_{\text{DM}}$
- $\rho_{\text{b}}$  baryon densities + velocity dispersions
- $f(v)$  in 20 bins for each tracer
- height of sun  $z_{\text{sun}}$
- overall normalization for density profiles

= 89 parameters

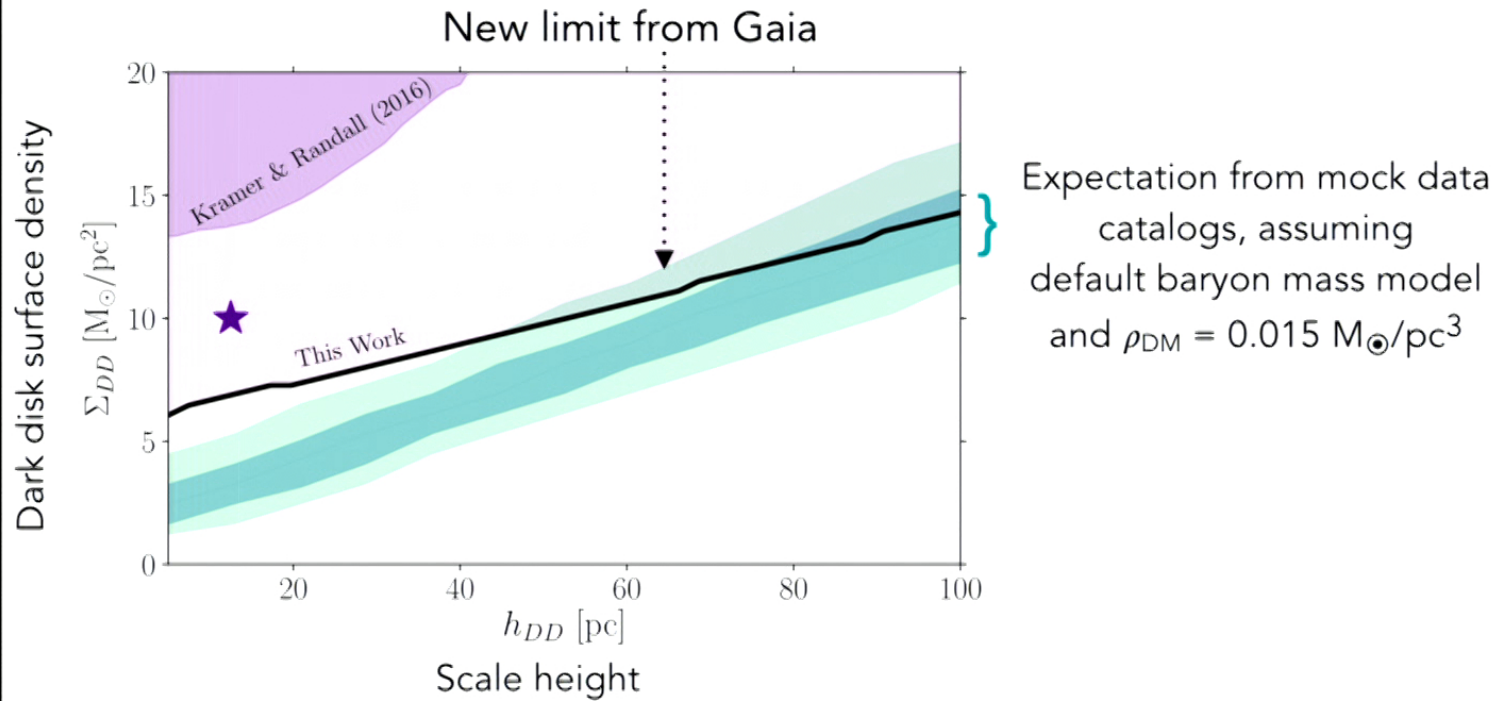
Condition for 95% one-sided limit:

$$\lambda(\Sigma_{DD}) = -2.71$$

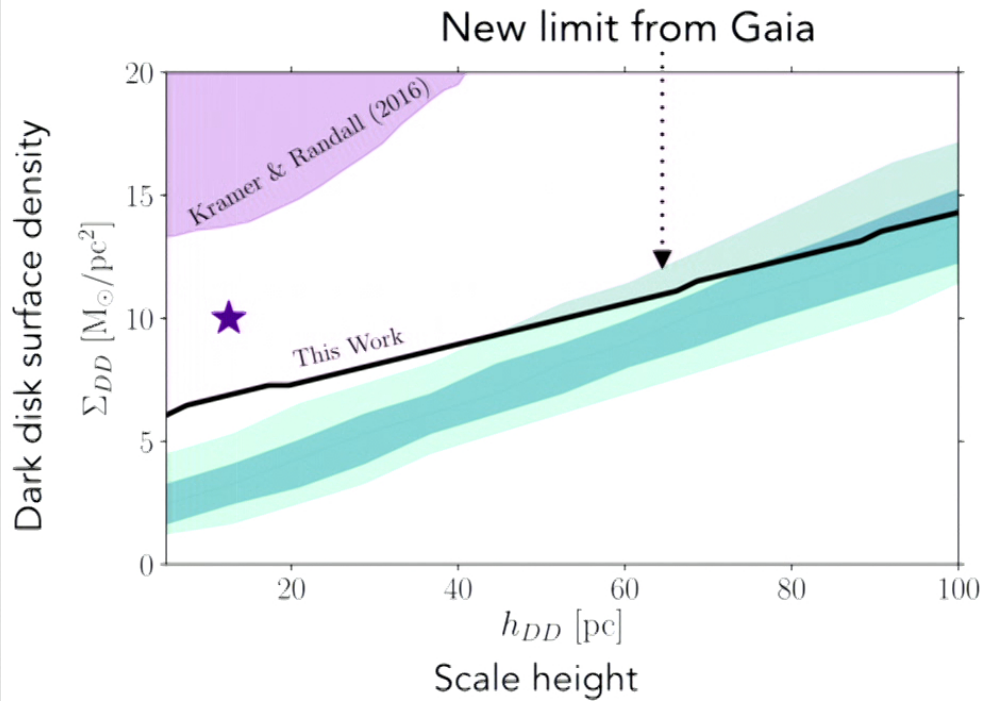




# Dark disk bounds



# Dark disk bounds



## Improvements

Statistics:

~38,000 vs. 2000 stars

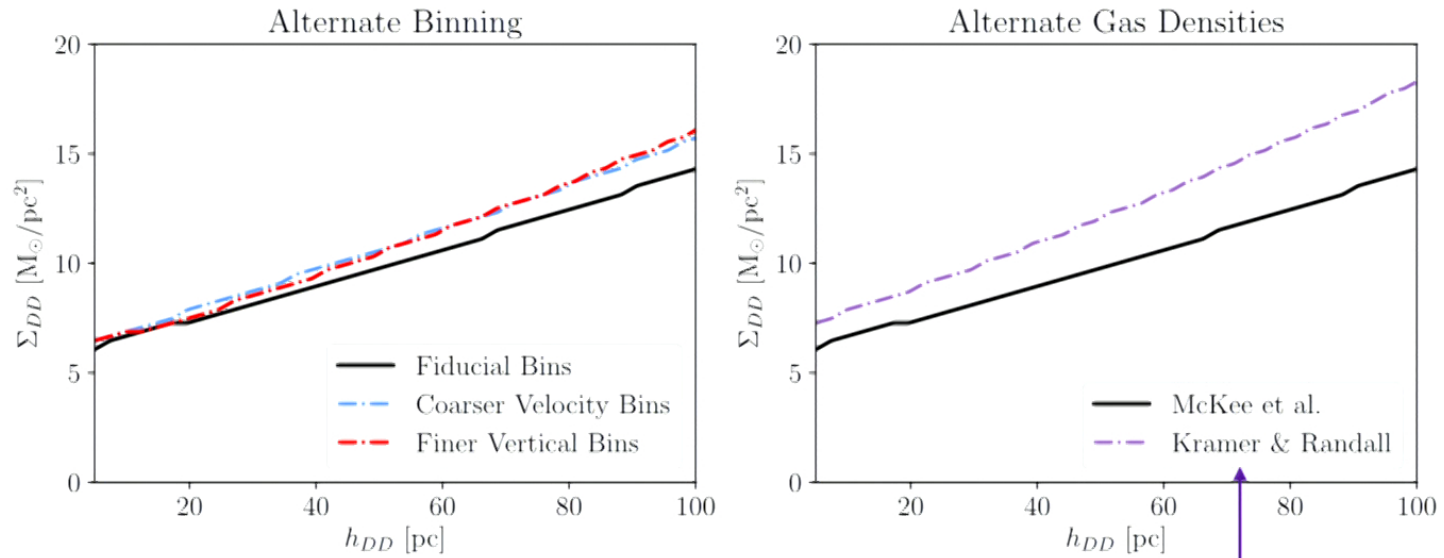
Longer lever arm:

max  $z = 200$  pc vs. 150 pc

Nuisance parameters:

marginalize over many  
nuisance parameters vs.  
"non-equilibrium" method

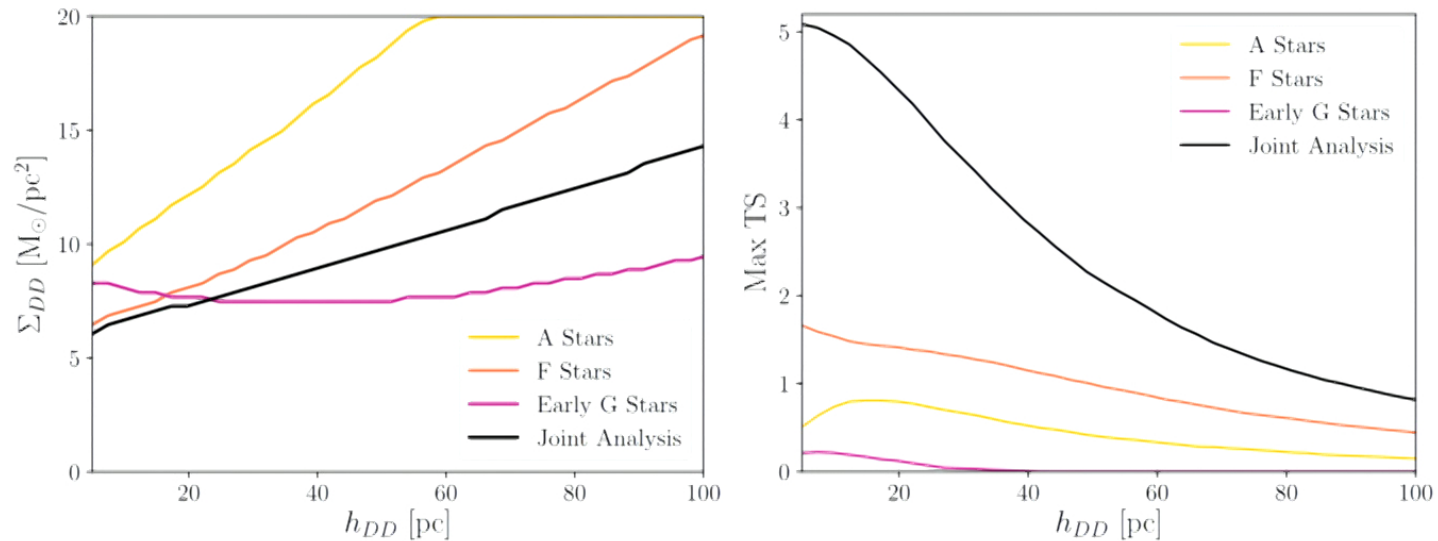
# Robustness of bounds



Gas Component	Kramer & Randall $\rho(0)$ [ $M_{\odot}/pc^3$ ]	McKee et al. $\rho(0)$ [ $M_{\odot}/pc^3$ ]
Molecular Gas ( $H_2$ )	$0.014 \pm 0.005$	$0.01 \pm 0.003$
Cold Atomic Gas ( $H_I(1)$ )	$0.015 \pm 0.003$	$0.028 \pm 0.006$
Warm Atomic Gas ( $H_I(2)$ )	$0.005 \pm 0.001$	$0.007 \pm 0.001$
Hot Ionized Gas ( $H_{II}$ )	$0.0011 \pm 0.0003$	$0.0005 \pm 0.00002$

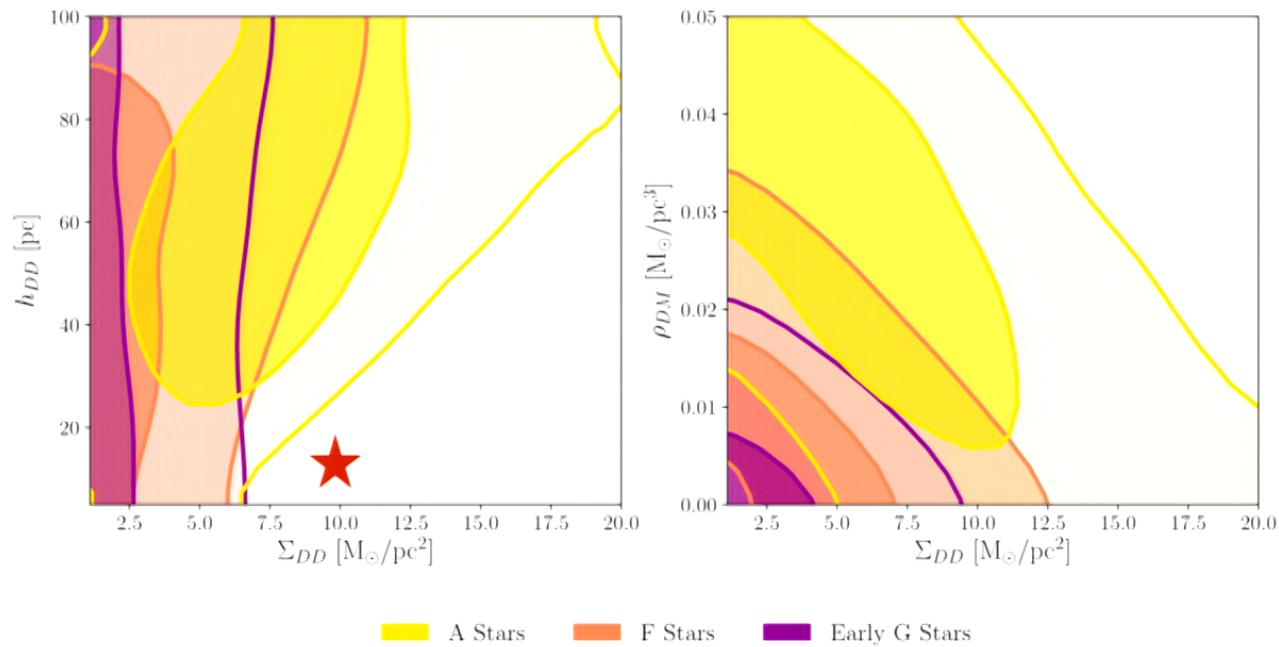
Kramer & Randall gas densities lower by ~25%

# Bounds by population



F and Early G stars have similar statistics: differences may be due to statistical fluctuations, measurement error, different dynamical densities seen by older stars, effect of non-equilibrium perturbations?

# Dark disk parameters



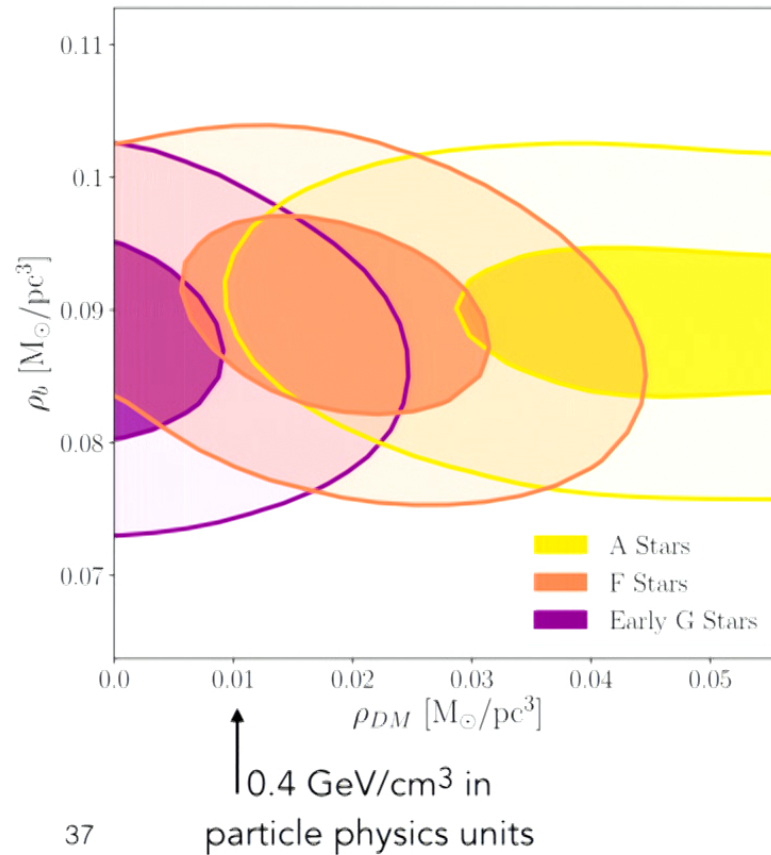
Bayesian analysis gives similar results on the dark disk  
posterior distributions for 3 tracer populations

# Local dark matter density

Measurements of midplane  
density in  $M_{\odot}/\text{pc}^3$

Component	A Stars	F Stars	Early G Stars
baryons	$0.088^{+0.006}_{-0.006}$	$0.088^{+0.007}_{-0.007}$	$0.085^{+0.007}_{-0.006}$
DM	$0.038^{+0.012}_{-0.015}$	$0.019^{+0.012}_{-0.011}$	$0.004^{+0.01}_{-0.004}$

First measurements of  
\*local\* DM density (usually  
obtained at 1.1 kpc).



# Dark disk implications, revisited

- Dynamical influence on local stars in the Milky Way (Kramer/Randall non-equilibrium method):  
 $\Sigma \sim 12 M_{\odot}/\text{pc}^2$ ,  $h \sim 10 \text{ pc}$
- Enhanced direct detection signal:  $O(1)$  level
- Periodic disruption of comet trajectories causing mass extinction events:  
 $\Sigma \sim 10\text{-}20 M_{\odot}/\text{pc}^2$ ,  $h \sim 10\text{-}30 \text{ pc}$
- Collapsed dark matter objects can account for the point-like nature of the inner galaxy GeV excess:  
 $\Sigma \sim 10 M_{\odot}/\text{pc}^2$ ,  $h \sim 10 \text{ pc}$
- Co-rotation of Andromeda satellites  
 $\Sigma \sim 10 M_{\odot}/\text{pc}^2$ ,  $h \sim 50 \text{ pc}$

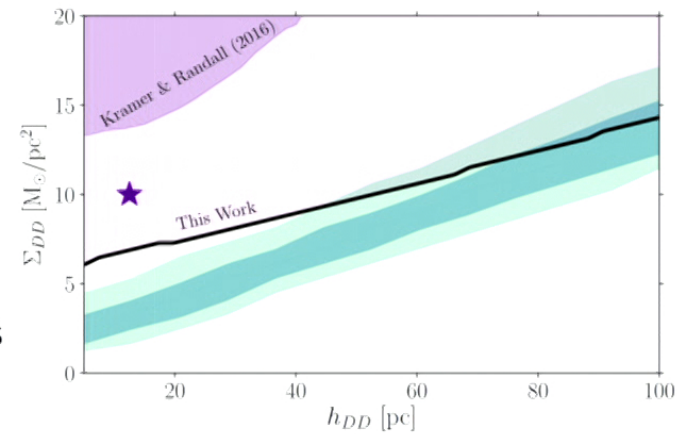


New  
constraints  
on these  
scenarios

Important to note that dissipative DM is not ruled out!

# Conclusions

- Dark disks can arise in models of dissipative dark matter
- Using Gaia TGAS data, we obtain
  - + new bounds on the presence of thin dark matter disks
  - + measurements of midplane densities and first local dark matter density
- Exciting time to better understand dark matter in the Milky Way with Gaia data!



Thanks!