

Title: An asymptotically safe point of view on the gravitational path integral

Speakers: Astrid Eichhorn

Series: Cosmology, Quantum Gravity

Date: November 13, 2017 - 2:00 PM

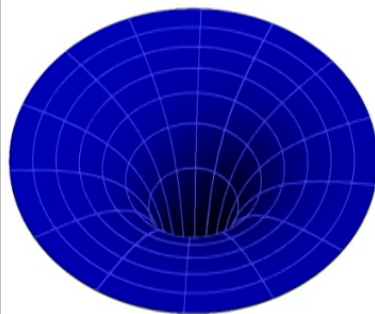
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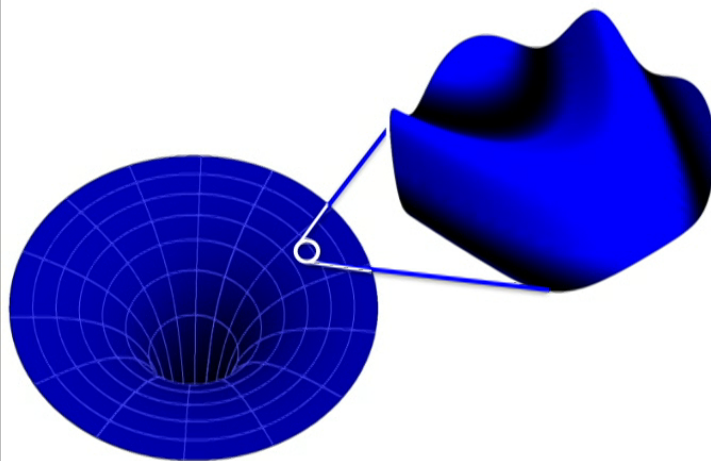
Renormalization Group flows of spacetime

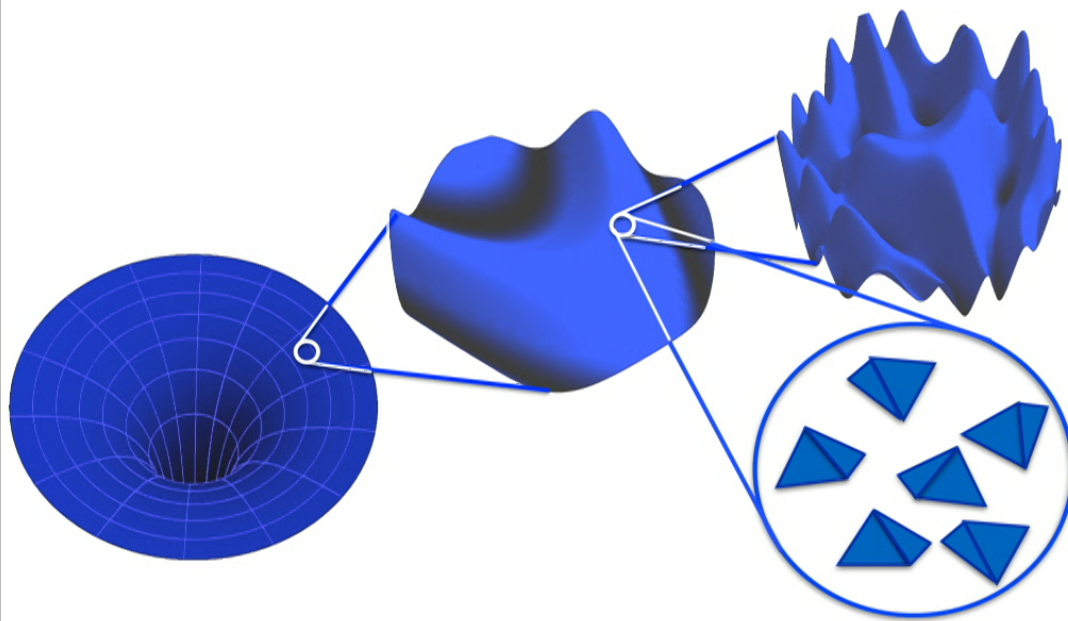
Astrid Eichhorn
University of Heidelberg

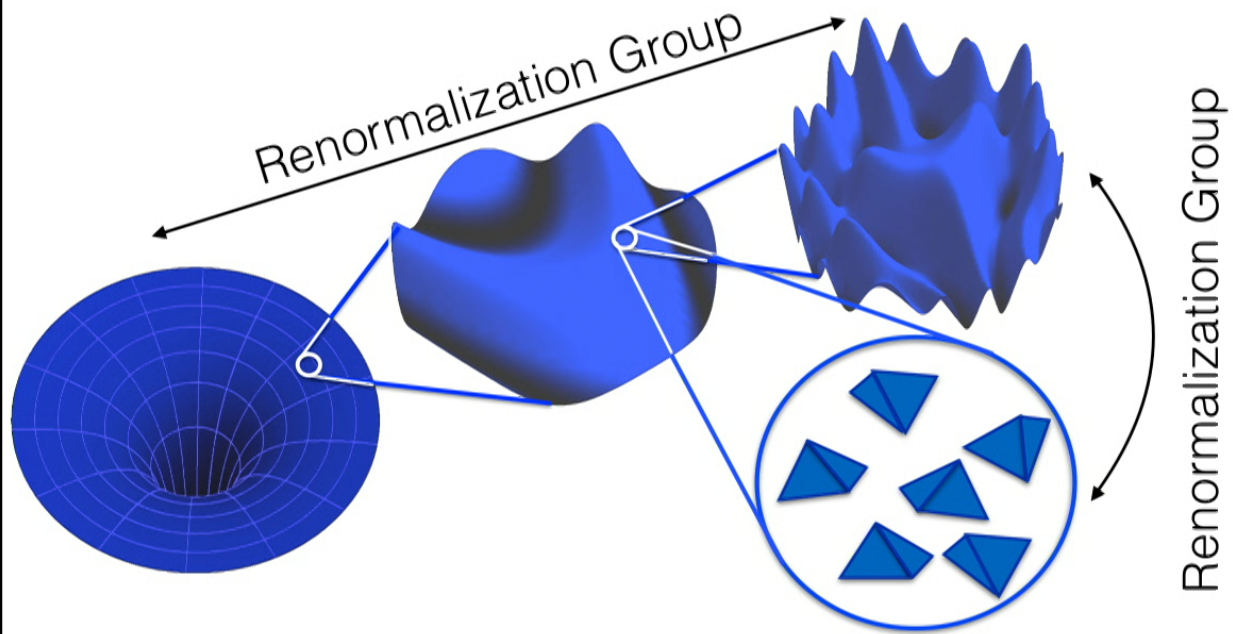


May 16th,
Perimeter Institute for Theoretical Physics
Workshop on Shape Dynamics



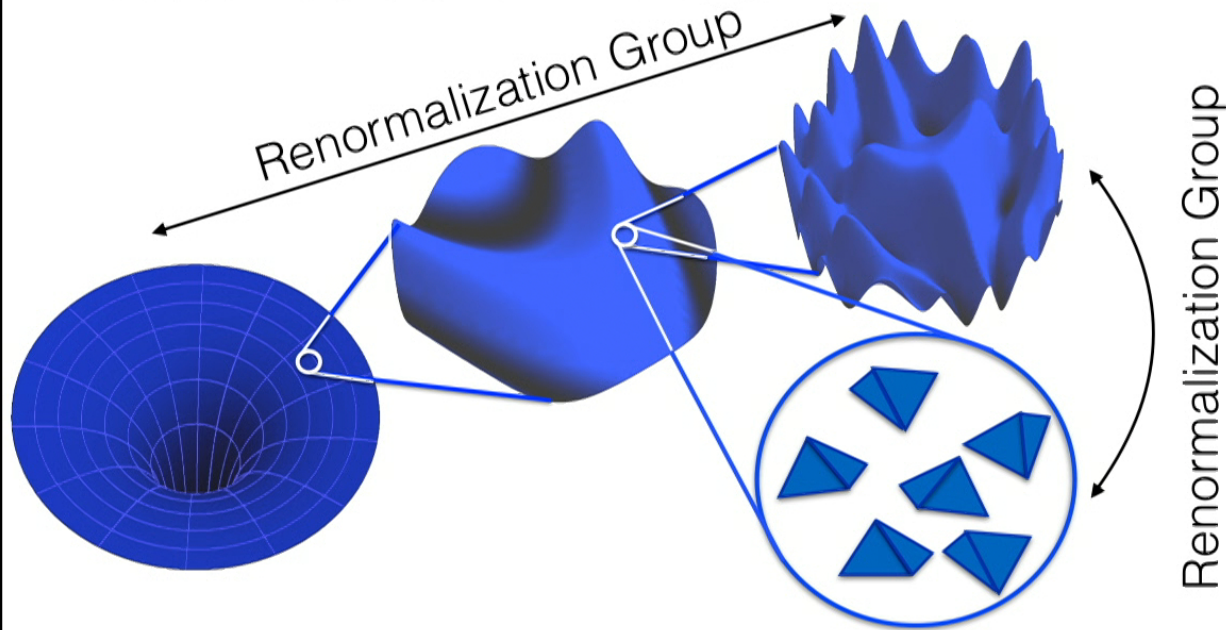






Outline

- 1) Functional RG: Probing the scale dependence of QFTs
- 2) RG for spacetime in the continuum:
Setting a scale in Quantum Gravity
- 3) RG for discrete spacetimes:
Scales in a pre-geometric setting?



Functional Renormalization Group

$$e^{-\Gamma[\phi]} = \int \mathcal{D}\varphi e^{-S[\varphi]}$$

$\phi = \langle \phi \rangle$

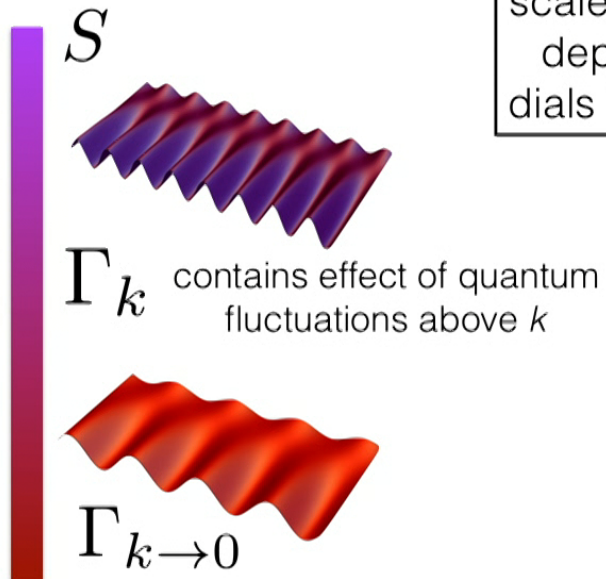
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Wetterich '93; Morris, '94

scale- and momentum-
dependent "mass":
dials "resolution scale"

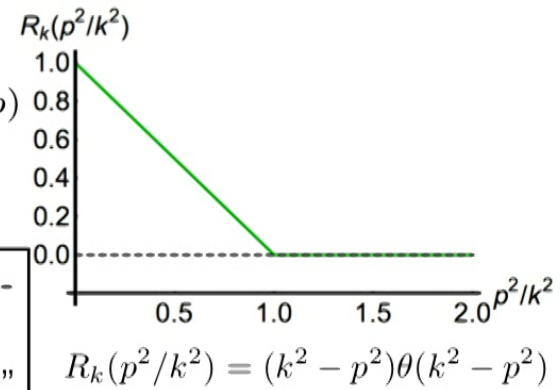


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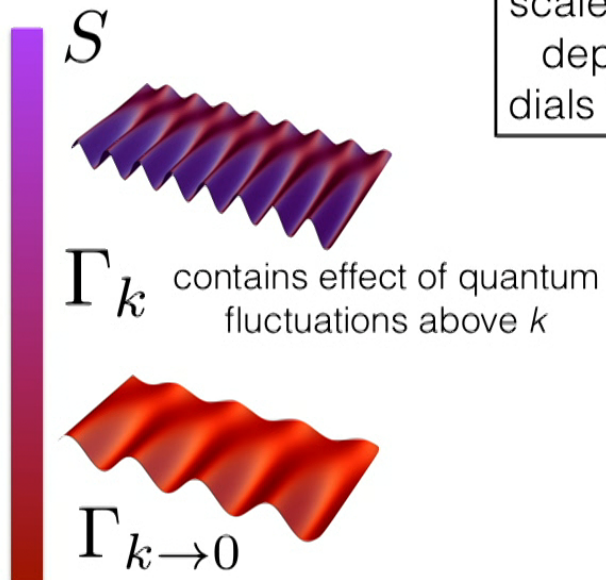
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Litim '01

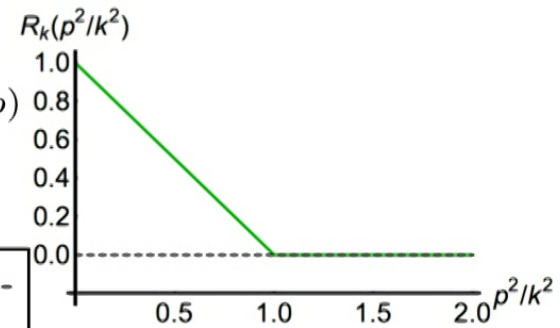


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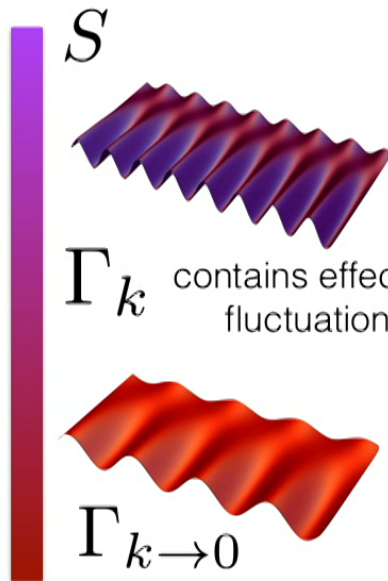
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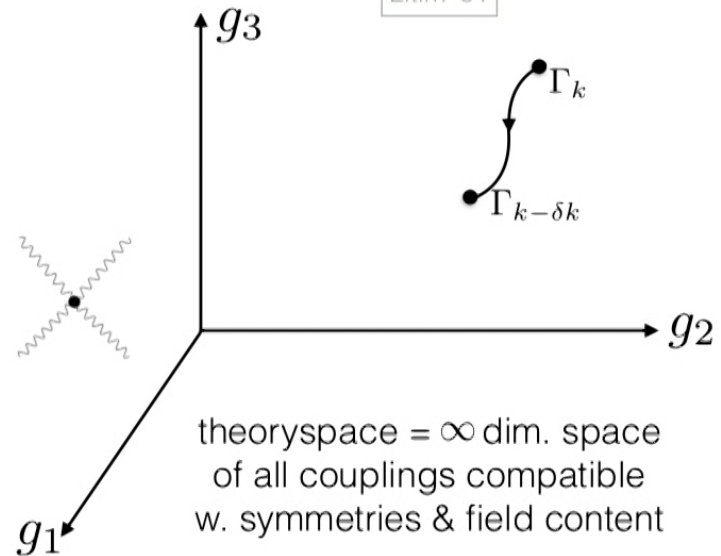
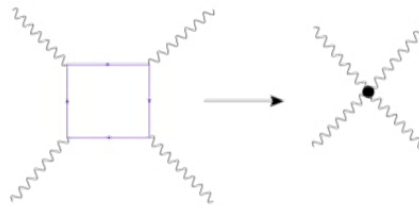
$$R_k(p^2/k^2) = (k^2 - p^2)\theta(k^2 - p^2)$$

Litim '01



Γ_k contains effect of quantum
fluctuations above k

$\Gamma_{k \rightarrow 0}$



theoryspace = ∞ dim. space
of all couplings compatible
w. symmetries & field content

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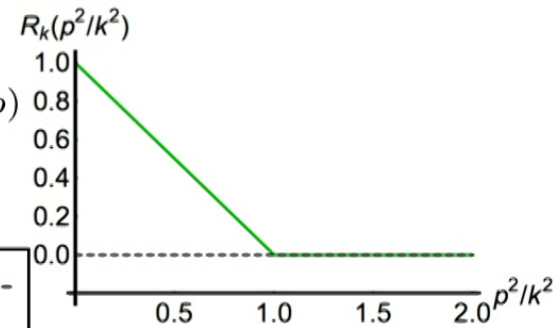
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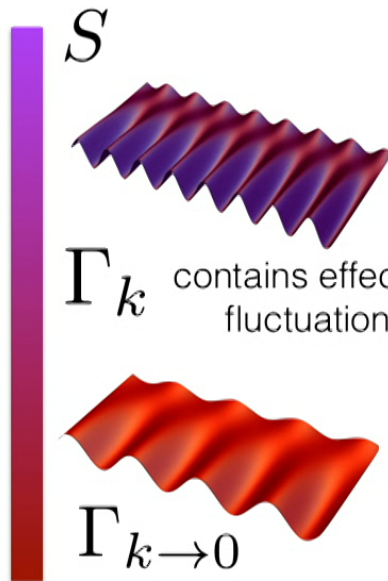
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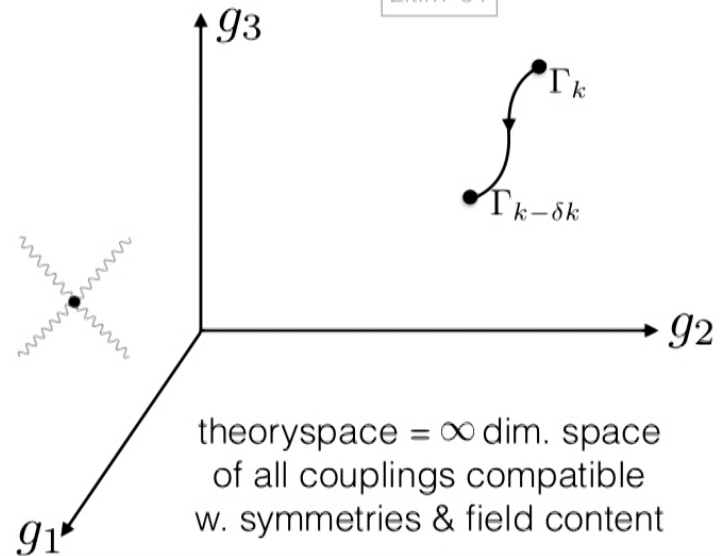
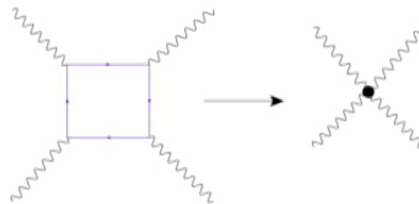
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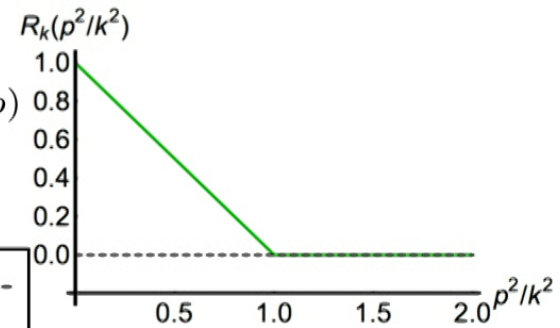
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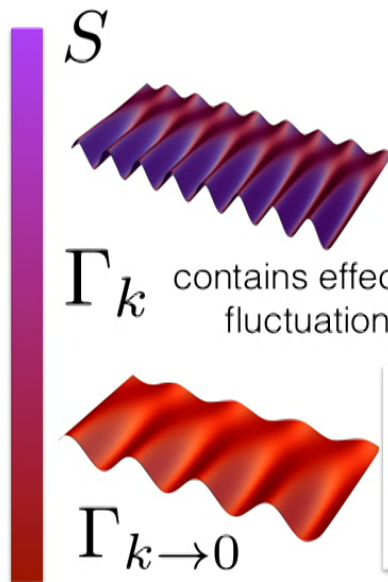
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$$\partial_k \Gamma_k = \frac{1}{2} \text{STr} \left(\Gamma_k^{(2)} + R_k \right)^{-1} \partial_k R_k =$$



g_1

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Γ_k
 $\Gamma_{k-\delta k}$

g_2

$$\phi = \langle \psi \rangle$$

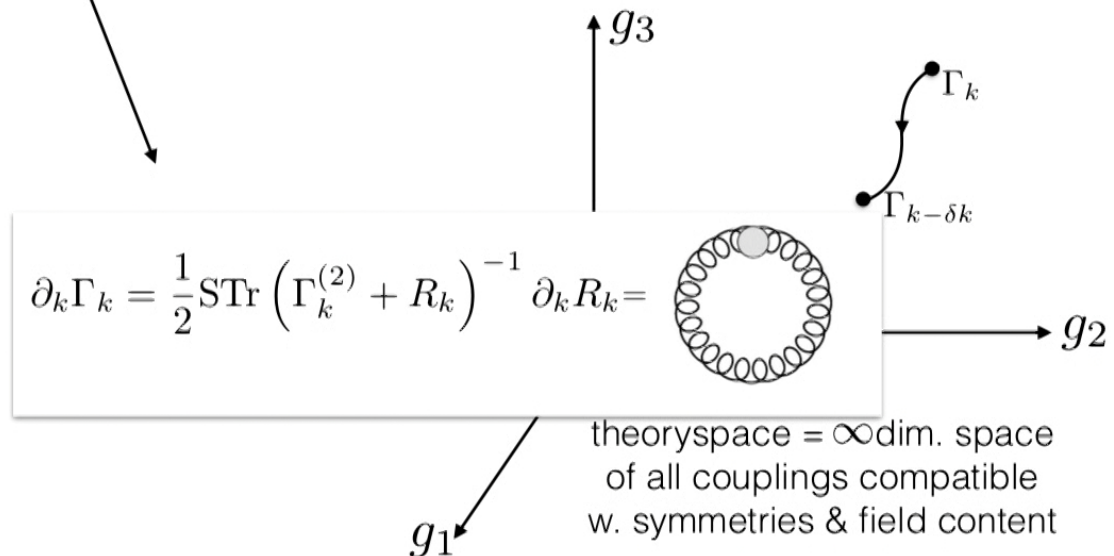


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Diagram illustrating the flow of the effective action Γ_k with respect to the scale k . The diagram shows a loop structure with vertices labeled Γ_k and $\Gamma_{k-\delta k}$, and a wavy line representing the regulator R_k . The flow is indicated by arrows labeled g_1 , g_2 , and g_3 .

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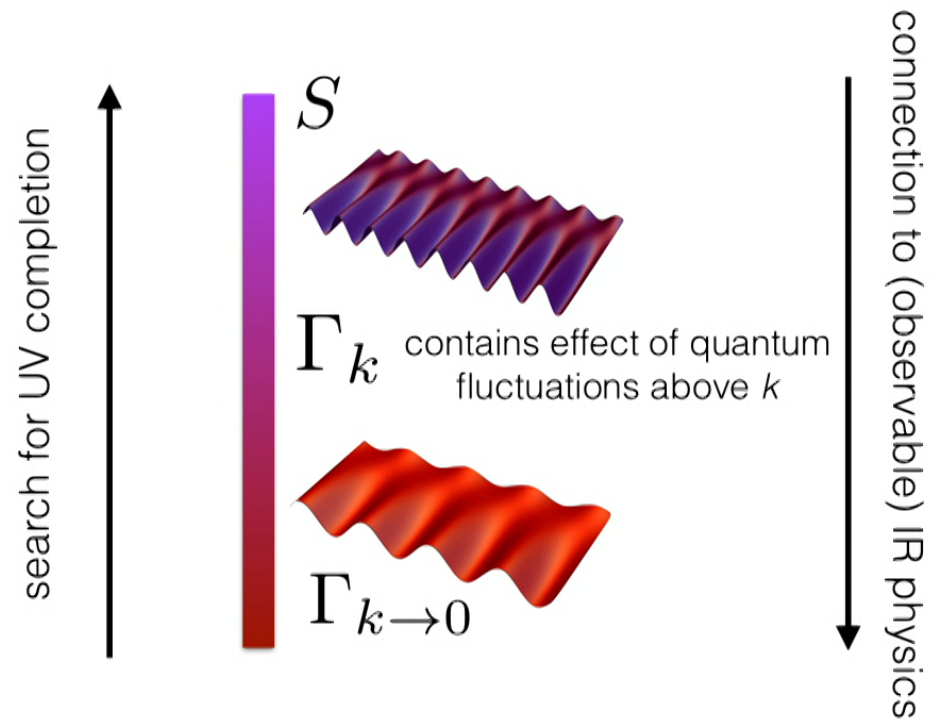
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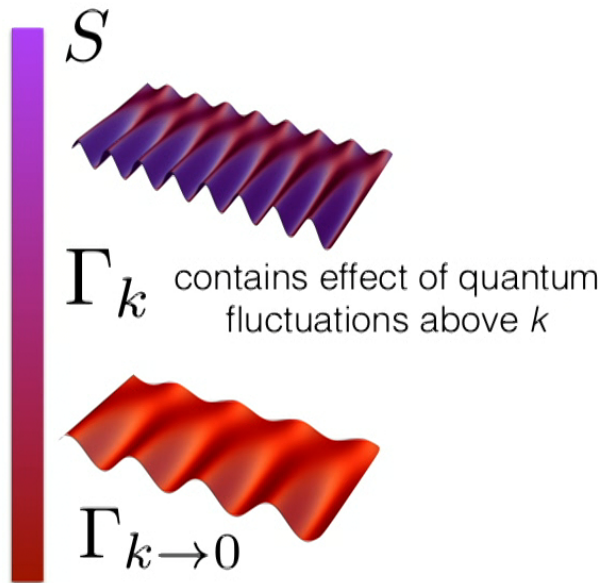

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Renormalization Group: uses in Quantum Gravity

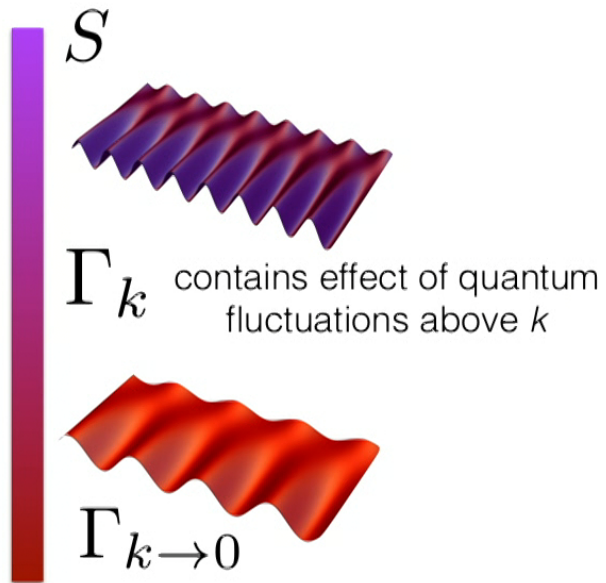


Search for UV completion in quantum gravity



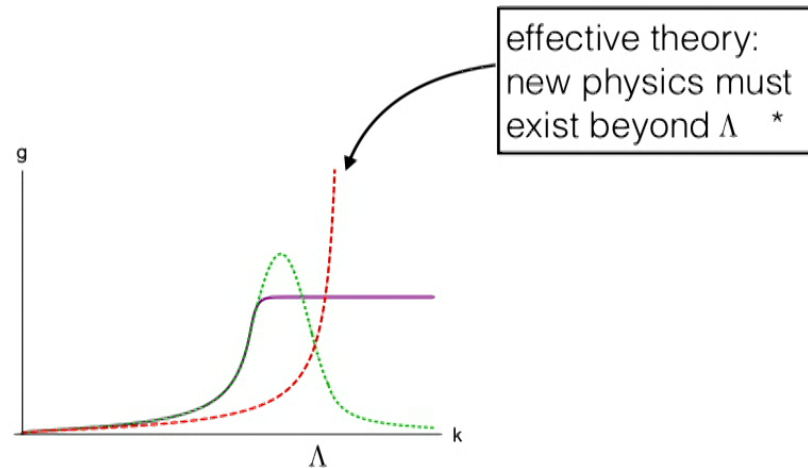
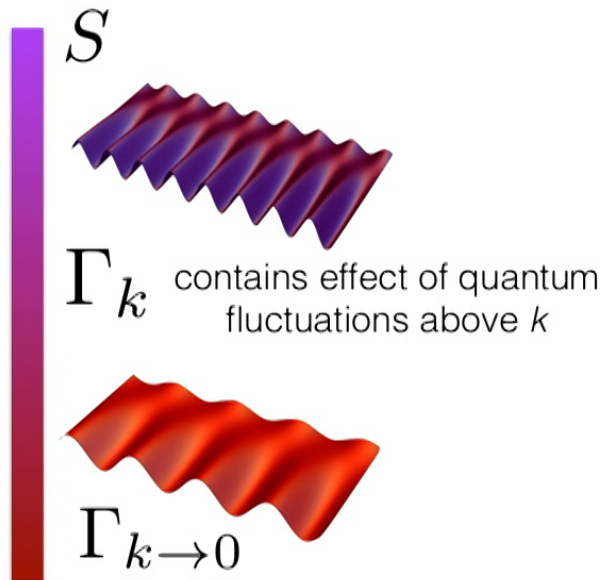
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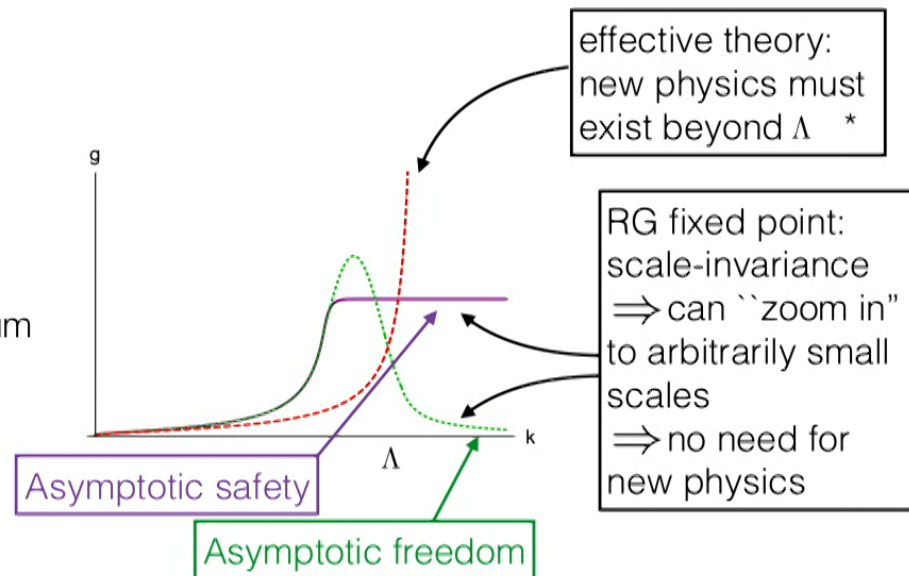
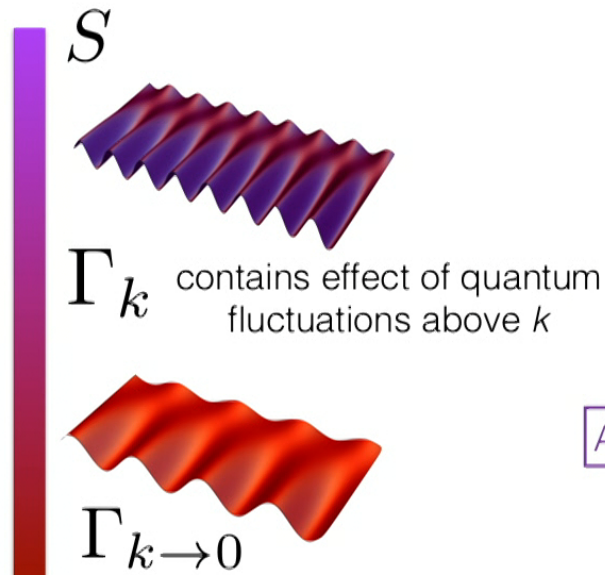


effective theory:
new physics must
exist beyond Λ *

* note: perturbative renormalizability not sufficient for fundamental theory

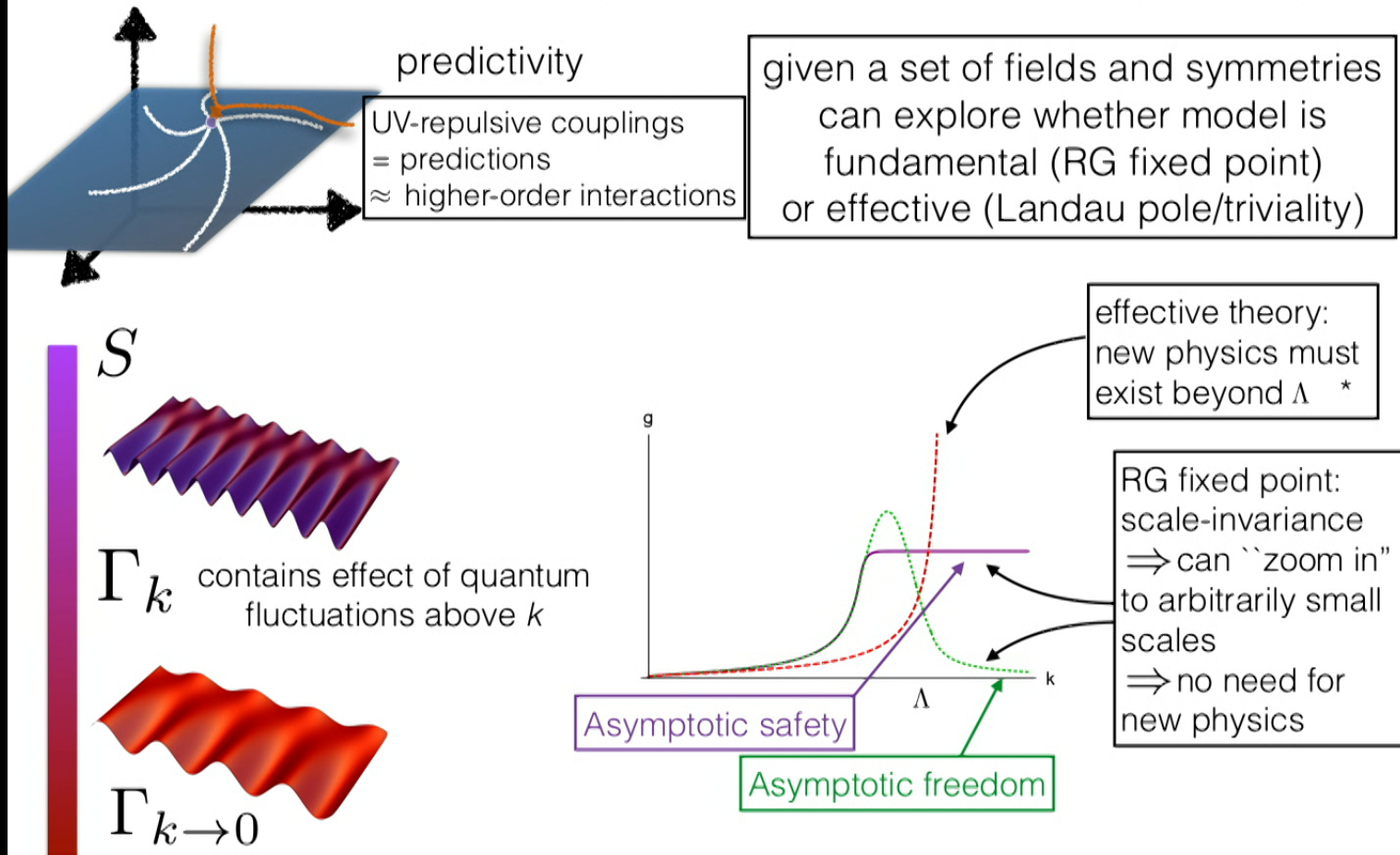
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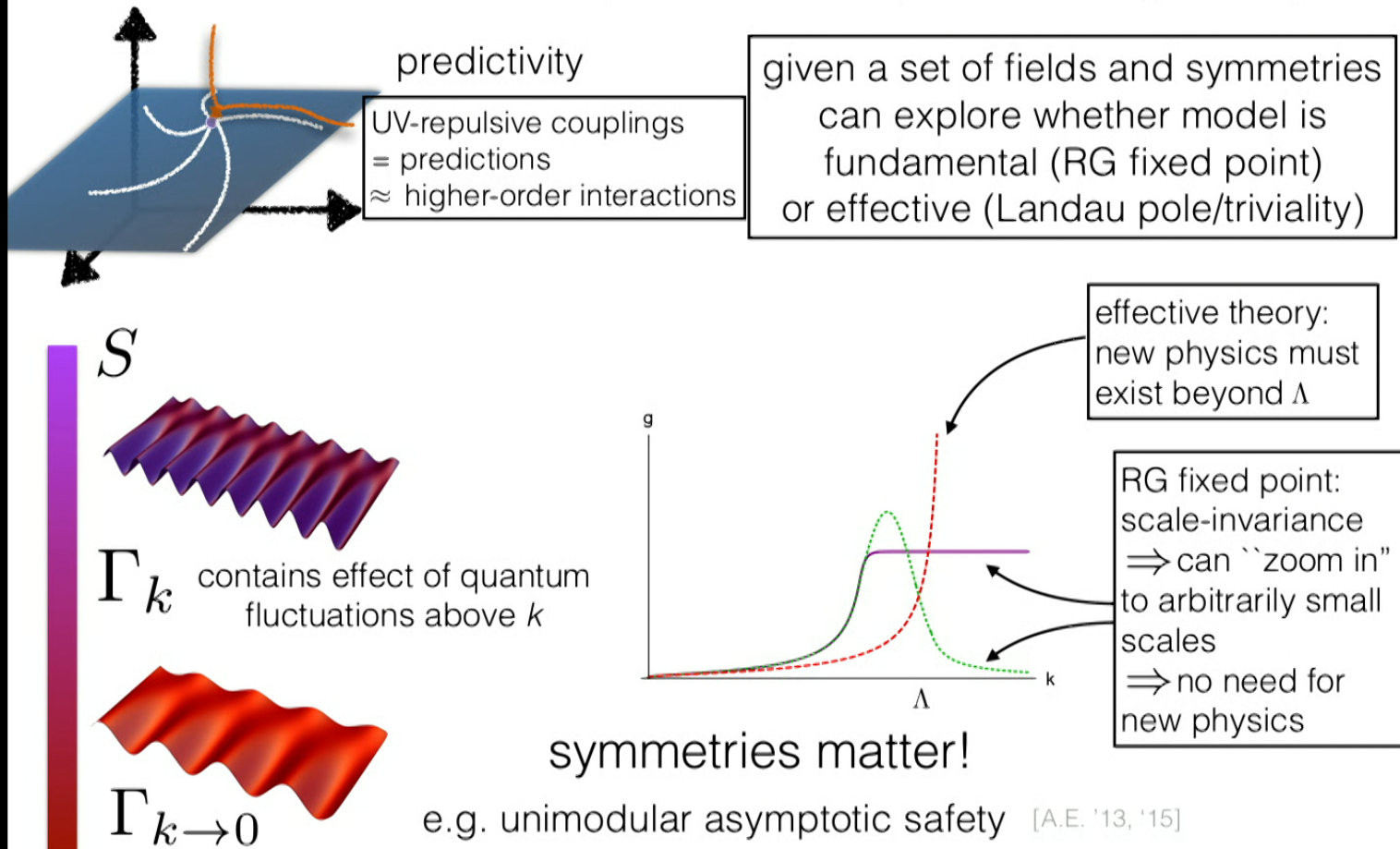
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Setting a scale in a geometric setting $\int \mathcal{D}g_{\mu\nu}$

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covariant generalization of momentum: (eigenvalues of) Laplacian

$g_{\mu\nu} R_k^{\mu\nu\kappa\lambda} (-D^2) g_{\kappa\lambda} \rightarrow$ one-loop structure of flow equation 

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$\rightarrow$ background field method

linear split: $g_{\mu\nu} = \bar{g}_{\mu\nu} + h_{\mu\nu}$ Reuter '96, ...

exponential split: $g_{\mu\nu} = \bar{g}_{\mu\kappa} (\exp(h_{\kappa\nu}))^\kappa_\nu$

A.E. '13, Percacci, Vacca '15, ...

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$$h_{\mu\nu} R_k^{\mu\nu\kappa\lambda} (-\bar{D}^2) h_{\kappa\lambda} \rightarrow \text{one-loop structure of flow equation intact}$$

No free lunch

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shift symmetry: (“there is only one metric”)

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- if the regulator breaks a symmetry, the flow no longer preserves the symmetry, even if the initial condition does
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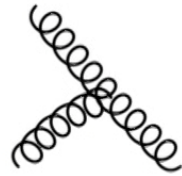
$\rightarrow$ enlarged theory space

$$\begin{aligned} S[g_{\mu\nu}] &= -\frac{1}{16\pi G_N} \int d^4x \sqrt{g} R \\ &= -\frac{1}{16\pi \bar{G}_N} \int d^4x \sqrt{\bar{g}} \bar{R} \\ &\quad - \frac{1}{16\pi G_1} \int d^4x \sqrt{\bar{g}} h_{\mu\nu} K^{\mu\nu}[\bar{g}] + \frac{1}{16\pi G_2} \int d^4x \sqrt{\bar{g}} h_{\mu\nu} h_{\kappa\lambda} K^{\mu\nu\kappa\lambda}[\bar{g}] \\ &\quad + \frac{1}{16\pi G_3} \int d^4x \sqrt{\bar{g}} h_{\mu\nu} h_{\kappa\lambda} h_{\rho\sigma} K^{\mu\nu\kappa\lambda\rho\sigma}[\bar{g}] + \dots \end{aligned}$$

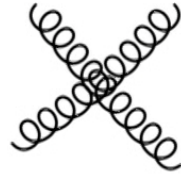
$\beta_{\bar{G}_N} \neq \beta_{G_1} \neq \beta_{G_2} \dots$

How expensive is it?

“Newton couplings” in gravity-matter systems



$$\sqrt{G_{(3,0)}}$$



$$G_{(4,0)}$$



$$\sqrt{G_{(1,2)}}$$



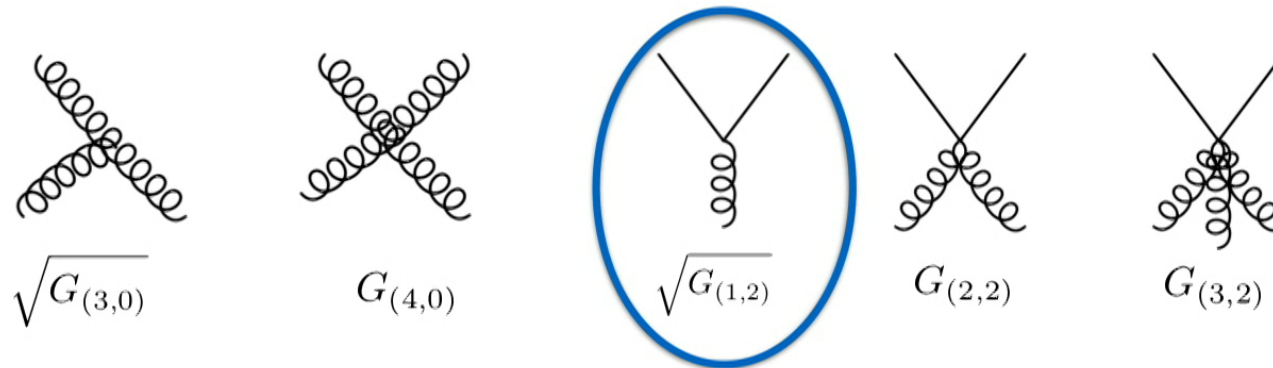
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background Newton coupling:

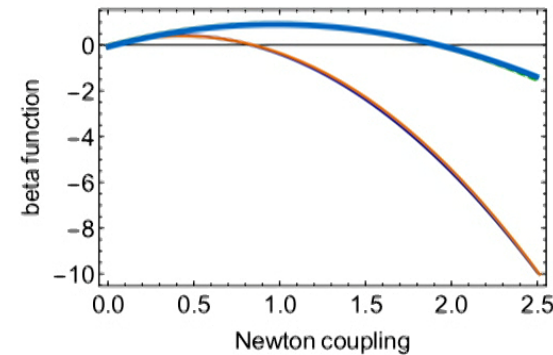
$$\beta_G = 2G + \frac{G^2}{6\pi} (-46 + N_S) + \dots$$

[Dona, A.E., Percacci, '13]

“graviton”-scalar coupling:

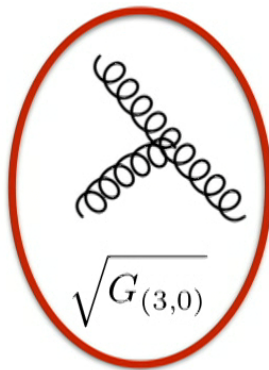
$$\beta_{G_{(1,2)}} = 2G_{(1,2)} + \frac{G_{(1,2)}^2}{3\pi} \left(-10 + \frac{N_S}{8} \right) + \dots$$

[Dona, A.E., Labus, Percacci, '15]

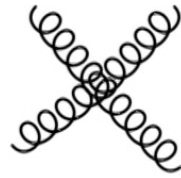


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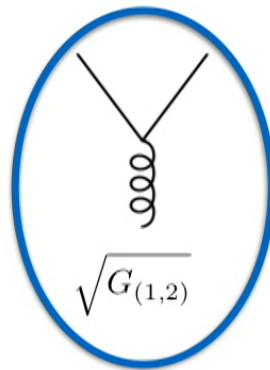
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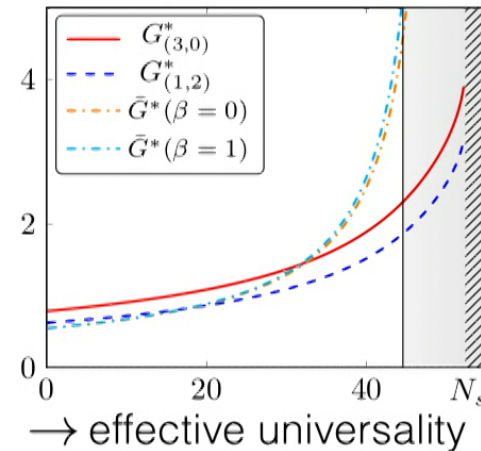
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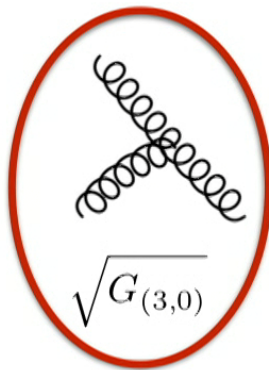
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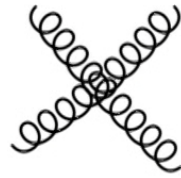
[A.E., Labus, Pawłowski, Reichert, to appear]

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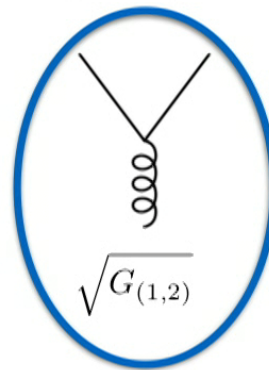
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quantitative difference but qualitative agreement

scalars drive gravity into
the “strong-gravity” regime

⇒ violation of weak-gravity bound

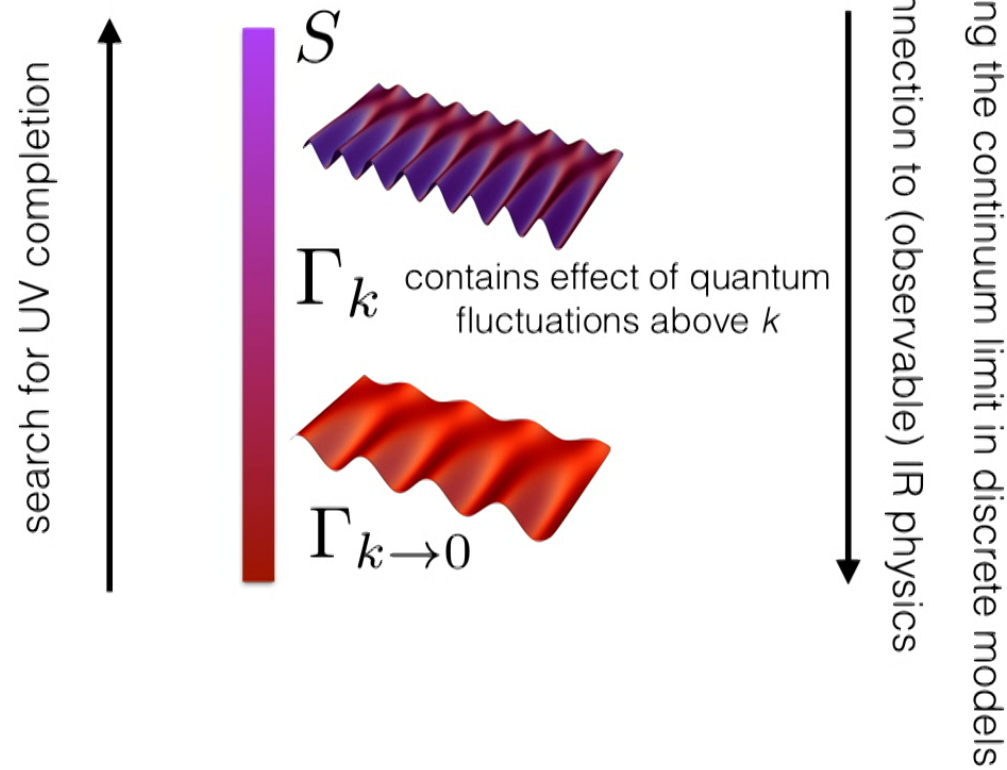
[A.E., Held, Pawłowski '16; Christiansen, A.E. '17, A.E., Held '17]

⇒ asymptotic safety only for $N_S < N_{S \text{ crit}}$

→ contact of QG to BSM pheno.

→ what about qm shape dynamics?

Renormalization Group: uses in Quantum Gravity



Setting a scale in a pre-geometric setting

RG flow: from many degrees of freedom to
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example: $N \times \dots \times N$ matrix/tensor models

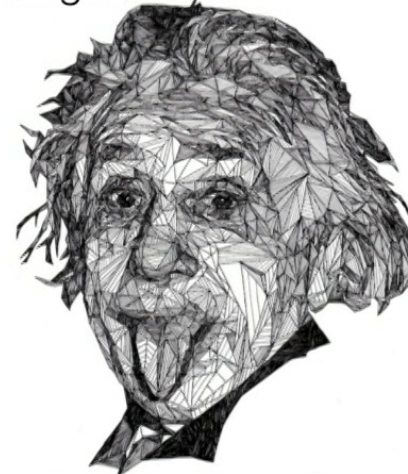
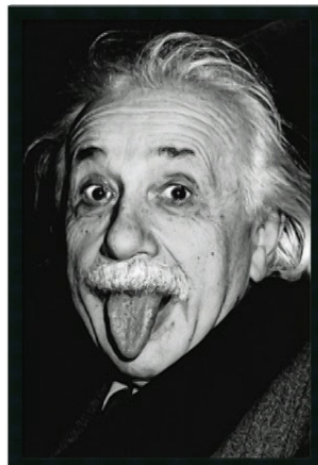
$$\int dT_{i_1 \dots i_d} d\bar{T}_{i_1 \dots i_d} e^{-S[\bar{T}, T]}$$

[Weingarten, Ambjorn, Durhuus, Frohlich,
Kazakov, Migdal, Boulatov, 1980's]

[Sasakura '91, Gross '91, Ambjorn, Durhuus, Jonsson '91,
Gurau '11, Gurau, Rivasseau '11, Gurau '12 ...]

trace invariants: $T_{i_1 \dots i_d} \bar{T}_{i_1 \dots i_d} + T_{i_1 \dots i_d} \bar{T}_{i_1 j_2 \dots j_d} T_{j_1 \dots j_d} \bar{T}_{j_1 i_2 \dots i_d} + \dots$

physics: discrete path integral over spacetime configurations



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[Sasakura '91, Gross '91, Ambjorn, Durhuus, Jonsson '91,
Gurau '11, Gurau, Rivasseau '11, Gurau '12 ...]

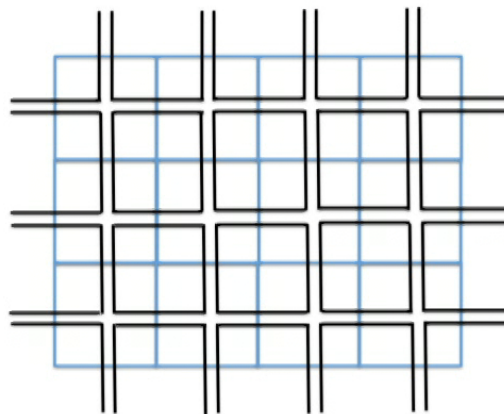
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physics: discrete path integral over spacetime configurations

matrix model:

$$S = \text{Tr} T^2 + g \text{Tr} T^4$$

Feynman-diag. expansion
= all possible tessellations



continuum limit =
double-scaling limit:

$$N \rightarrow \infty \quad g \rightarrow g_c$$

$$N (g - g_c)^{-\frac{1}{\theta}} = \text{const}$$

Setting a scale in a pre-geometric setting

RG flow: from many degrees of freedom to
fewer (coarser) degrees of freedom

example: $N \times \dots \times N$ matrix/tensor models

$$\int dT_{i_1 \dots i_d} d\bar{T}_{i_1 \dots i_d} e^{-S[\bar{T}, T]}$$

[Weingarten, Ambjorn, Durhuus, Frohlich,
Kazakov, Migdal, Boulatov, 1980's]

[Sasakura '91, Gross '91, Ambjorn, Durhuus, Jonsson '91,
Gurau '11, Gurau, Rivasseau '11, Gurau '12 ...]

trace invariants: $T_{i_1 \dots i_d} \bar{T}_{i_1 \dots i_d} + T_{i_1 \dots i_d} \bar{T}_{i_1 j_2 \dots j_d} T_{j_1 \dots j_d} \bar{T}_{j_1 i_2 \dots i_d} + \dots$

double-scaling limit:

$$N \rightarrow \infty, g \rightarrow g_c$$

$$N (g - g_c)^{-\frac{1}{\theta}} = \text{const}$$

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double-scaling limit:

$$N \rightarrow \infty, g \rightarrow g_c$$

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universal scaling near RG fixed points:

$$\beta_g = \underbrace{\frac{\partial \beta_g}{\partial g} \Big|_{g=g_*}}_{=-\theta} (g - g_*) + \dots \rightarrow g(k) - g_* = \left(\frac{k}{k_0} \right)^{-\theta}$$

Setting a scale in a pre-geometric setting

RG flow: from many degrees of freedom to
fewer (coarser) degrees of freedom

example: $N \times \dots \times N$ matrix/tensor models

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double-scaling limit:

$$N \rightarrow \infty, g \rightarrow g_c$$

$$N (g - g_c)^{-\frac{1}{\theta}} = \text{const}$$

universal scaling near RG fixed points:

$$\beta_g = \underbrace{\frac{\partial \beta_g}{\partial g} \Big|_{g=g_*}}_{=-\theta} (g - g_*) + \dots \rightarrow g(k) - g_* = \left(\frac{k}{k_0} \right)^{-\theta}$$

double-scaling limit = universal scaling near RG fixed point in N
[Brezin, Zinn-Justin, '92]

No free lunch

example: $N' \times \dots \times N'$ matrix/tensor models

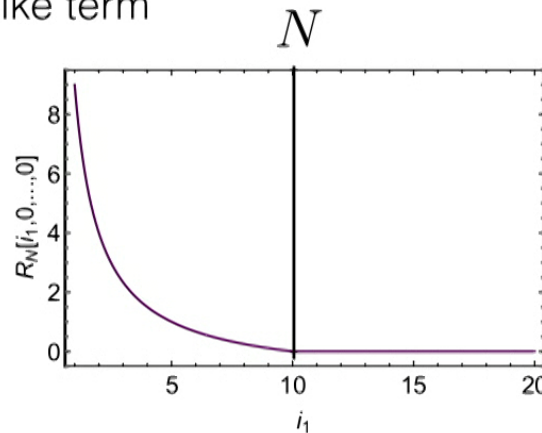
$$\int dT_{i_1 \dots i_d} d\bar{T}_{i_1 \dots i_d} e^{-S[\bar{T}, T]}$$

trace invariants: $T_{i_1 \dots i_d} \bar{T}_{i_1 \dots i_d} + T_{i_1 \dots i_d} \bar{T}_{i_1 j_2 \dots j_d} T_{j_1 \dots j_d} \bar{T}_{j_1 i_2 \dots i_d} + \dots$

has $U(N') \otimes \dots \otimes U(N')$ invariance

$$\text{Regulator: } R_N[i_1, \dots, i_d] = \left(\frac{N}{i_1 + \dots + i_d} - 1 \right) \theta(N - (i_1 + \dots + i_d))$$

index- dependent mass-like term



A.E., Kosłowski '13, '14, '17

No free lunch

example: $N' \times \dots \times N'$ matrix/tensor models

$$\int dT_{i_1 \dots i_d} d\bar{T}_{i_1 \dots i_d} e^{-S[\bar{T}, T]}$$

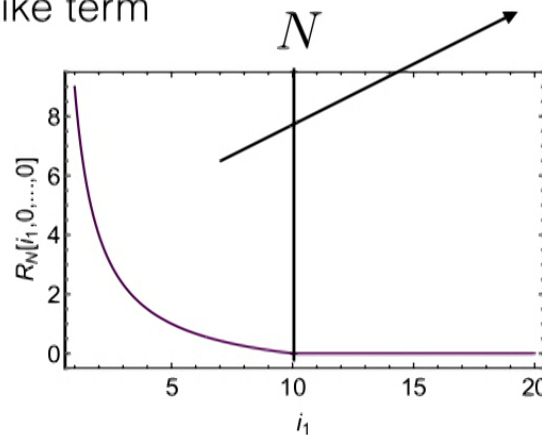
trace invariants: $T_{i_1 \dots i_d} \bar{T}_{i_1 \dots i_d} + T_{i_1 \dots i_d} \bar{T}_{i_1 j_2 \dots j_d} T_{j_1 \dots j_d} \bar{T}_{j_1 i_2 \dots i_d} + \dots$

has $U(N') \otimes \dots \otimes U(N')$ invariance

$$\text{Regulator: } R_N[i_1, \dots, i_d] = \left(\frac{N}{i_1 + \dots + i_d} - 1 \right) \theta(N - (i_1 + \dots + i_d))$$

index- dependent mass-like term

$U(N') \otimes \dots \otimes U(N')$
broken



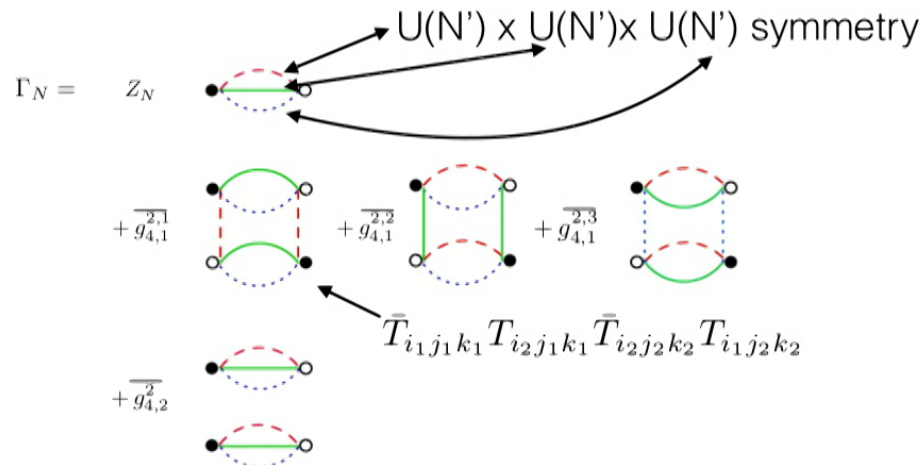
A.E., Koslowski '13, '14, '17

RG fixed points in rank 3 hermitian tensor models

the model

$T : \circ$

$\bar{T} : \bullet$



scaling dimensionality: $\bar{g}_i = N^{d_{\bar{g}_i}} g_i$

consistently determined by RG equation (not just in tensor models)

$$\beta_{g_i} = -d_{\bar{g}_i} g_i + \#_1(N) g_i g_4 + \#_2(N) g_{i+2} + \dots$$

$$\beta_{g_{i+2}} = -d_{\bar{g}_{i+2}} g_{i+2} + \#_3(N) g_{i+2} g_4 + \#_4(N) g_{i+4} + \dots$$

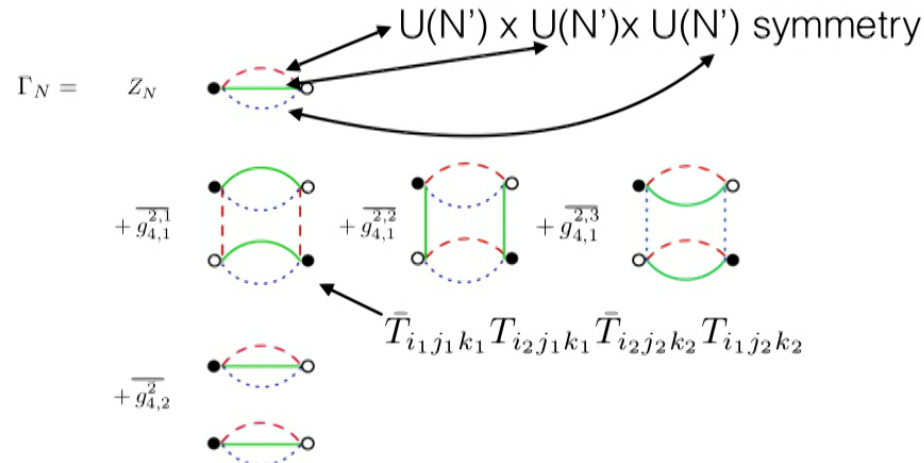
[A.E., Koslowski '17, accepted in Annales de l'Institut Henri Poincaré D]

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$$\beta_{g_i} = -d_{\bar{g}_i} g_i + \#_1(N) g_i g_4 + \#_2(N) g_{i+2} + \dots$$

$$\beta_{g_{i+2}} = -d_{\bar{g}_{i+2}} g_{i+2} + \#_3(N) g_{i+2} g_4 + \#_4(N) g_{i+4} + \dots$$

[A.E., Koslowski '17, accepted in Annales de l'Institut Henri Poincaré D]

RG fixed points in rank 3 hermitian tensor models

the model

$T : \circ$

$\bar{T} : \bullet$

$$\Gamma_N = Z_N \quad \begin{array}{c} \nearrow U(N') \times U(N') \times U(N') \text{ symmetry} \\ \searrow \end{array}$$

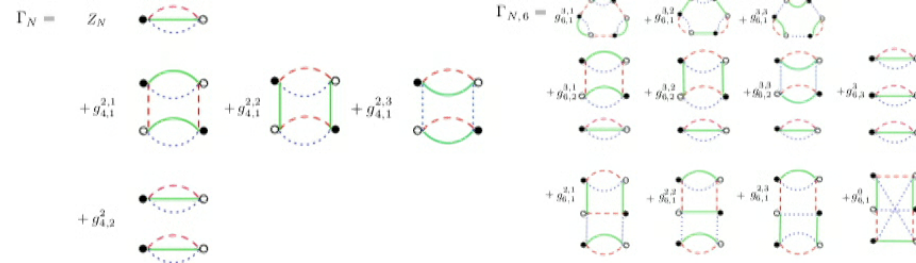
$$+ \overline{g_{4,1}^{2,1}} + \overline{g_{4,1}^{2,2}} + \overline{g_{4,1}^{2,3}} + \overline{g_{4,2}^2} \rightarrow \text{higher-order terms generated by the flow}$$

$\bar{T}_{i_1 j_1 k_1} T_{i_2 j_1 k_1} \bar{T}_{i_2 j_2 k_2} T_{i_1 j_2 k_2}$

$$\Gamma_{N,6} = g_{6,1}^{3,1} + g_{6,1}^{3,2} + g_{6,1}^{3,3} + g_{6,2}^{3,1} + g_{6,2}^{3,2} + g_{6,2}^{3,3} + g_{6,3}^3 + g_{6,1}^{2,1} + g_{6,1}^{2,2} + g_{6,1}^{2,3} + g_{6,1}^0$$

[A.E., Koslowski '17, accepted in Annales de l'Institut Henri Poincaré D]

Fixed-point results



scheme	$g_{4,1}^{2,1}$	$g_{4,2}^{2,1}$	$g_{6,1}^{3,1}$	$g_{6,1}^{2,1}$	$g_{6,2}^{3,1}$	$g_{6,3}^{3,1}$
full	-1.94	0	—	—	—	—
semi-pert	-2.14	0	—	—	—	—
pert	-4.62	0	—	—	—	—
full	-1.37	0	-2.14	0	0	0
semi-pert	-1.47	0	-2.46	0	0	0
pert	-2.14	0	-6.12	0	0	0

$$\theta = 1$$

[Bonzom, Gurau, Ryan, Tanasa '14
Dartois, Gurau, Rivasseau '13]

anomalous dimension

$$\eta = -\partial_t \ln Z_N$$

$$\theta : \beta_g = \beta_g(g, \eta(g))$$

$$\theta' : \beta_g = \beta_g(g, \eta(g)_* = \text{const})$$

scheme	θ_1	θ_2	θ_3	θ_4	θ_5	θ_6	θ_7	θ_8	θ_9	θ_{10}	θ_{11}	θ_{12}	θ_{13}	θ_{14}	η
full	2.21	-0.24	-0.93	-0.93	—	—	—	—	—	—	—	—	—	—	-0.54
semi-pert	2	-0.21	-0.68	-0.68	—	—	—	—	—	—	—	—	—	—	-0.54
pert	2	0.69	-2	-2	—	—	—	—	—	—	—	—	—	—	-0.54
full	2.14	-0.59	-1.26	-1.26	-2.24	-2.54	-2.89	-2.89	-3.15	-3.26	-3.26	-3.31	-3.31	-3.89	-0.37
semi-pert	2	-0.58	-1.26	-1.26	-2.24	-2.54	-2.90	-2.90	-3.13	-3.26	-3.26	-3.31	-3.31	-3.90	-0.37
pert	2	-0.35	-2	-2	-3.23	-3.37	-3.43	-4	-4	-4.07	-4.07	-4.14	-4.14	-5	0

scheme	θ'_1	θ'_2	θ'_3	θ'_4	θ'_5	θ'_6	θ'_7	θ'_8	θ'_9	θ'_{10}	θ'_{11}	θ'_{12}	θ'_{13}	θ'_{14}	η
full	0.93	-0.24	-0.93	-0.93	—	—	—	—	—	—	—	—	—	—	-0.54
semi-pert	0.93	-0.21	-0.93	-0.93	—	—	—	—	—	—	—	—	—	—	-0.54
full	1.33	-0.59	-1.26	-1.26	-2.30	-2.54	-2.89	-2.89	-3.15	-3.26	-3.26	-3.31	-3.31	-3.89	-0.37
semi-pert	1.33	-0.58	-1.26	-1.26	-2.31	-2.54	-2.90	-2.90	-3.13	-3.26	-3.26	-3.31	-3.31	-3.90	-0.37

[A.E., Koslowski '17, accepted in Annales de l'Institut Henri Poincaré D]

Summary

1) RG: Probing the scale dependence of QFTs

2) RG for spacetime in the continuum:

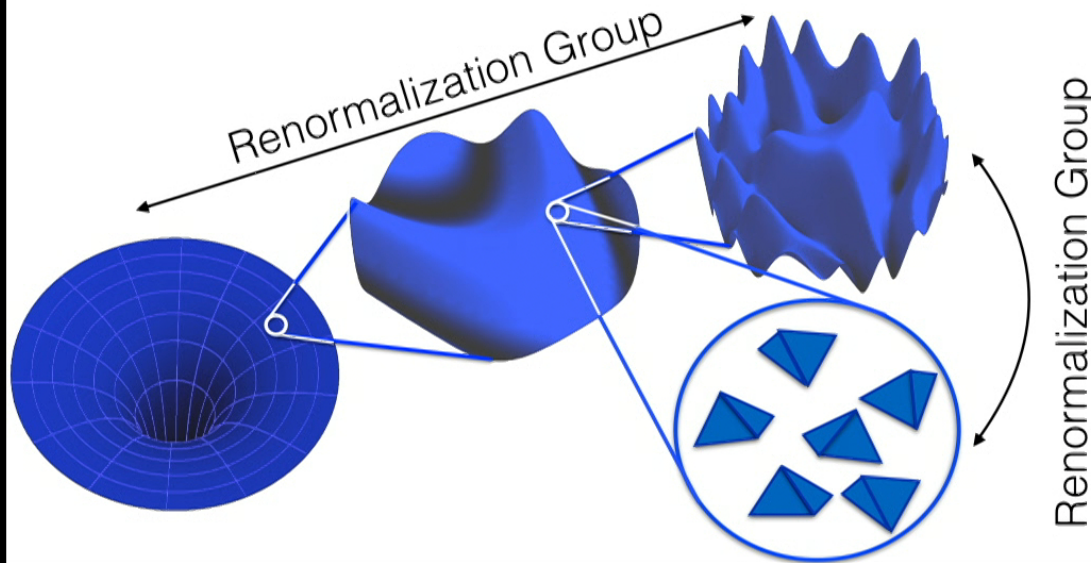
The price for setting a scale in Quantum Gravity

→ effective universality
in asymptotically
safe qg

3) RG for discrete spacetimes:

The price for setting a scale in Quantum Gravity

→ reason behind
large error in
estimate of
scaling exponent
in tensor models?



Fixed-point results

scheme	$g_{4,1}^{2,1}$	$g_{4,2}^{2,1}$	$g_{6,1}^{3,1}$	$g_{6,1}^{2,1}$	$g_{6,2}^{3,1}$	$g_{6,3}^{3,1}$
full	-1.94	0	—	—	—	—
semi-pert	-2.14	0	—	—	—	—
pert	-4.62	0	—	—	—	—
full	-1.37	0	-2.14	0	0	0
semi-pert	-1.47	0	-2.46	0	0	0
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$$\theta = 1$$

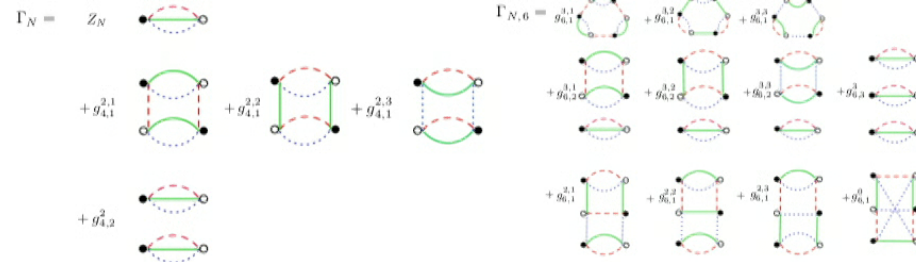
[Bonzom, Gurau, Ryan, Tanasa '14
Dartois, Gurau, Rivasseau '13]

anomalous dimension

$$\eta = -\partial_t \ln Z_N$$

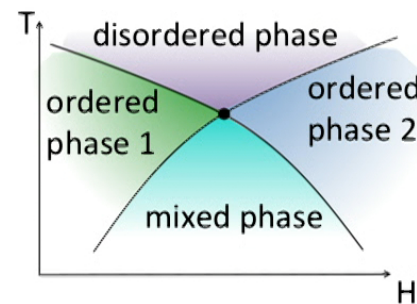
$$\theta : \beta_g = \beta_g(g, \eta(g))$$

$$\theta' : \beta_g = \beta_g(g, \eta(g)_* = \text{const})$$



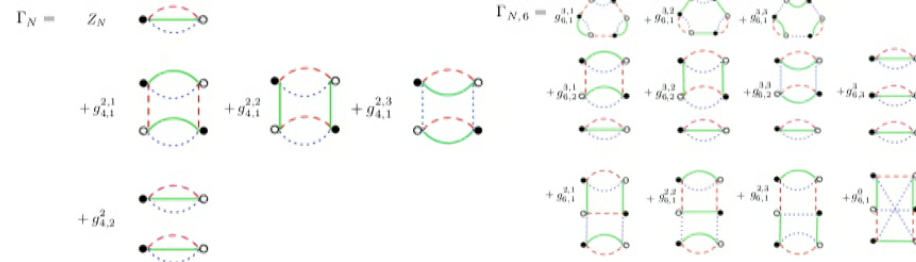
Observation:

Second scheme in accordance with scaling relation from ϵ expansion at multicritical points in systems with competing orders



[A.E., Mesterhazy, Scherer '13;
Boettcher '14]

Fixed-point results



scheme	$g_{4,1}^{2,1}$	$g_{4,2}^{2,1}$	$g_{6,1}^{3,1}$	$g_{6,2}^{2,1}$	$g_{6,3}^{3,1}$	$g_{6,4}^{3,1}$
full	-1.94	0	—	—	—	—
semi-pert	-2.14	0	—	—	—	—
pert	-4.62	0	—	—	—	—
full	-1.37	0	-2.14	0	0	0
semi-pert	-1.47	0	-2.46	0	0	0
pert	-2.14	0	-6.12	0	0	0

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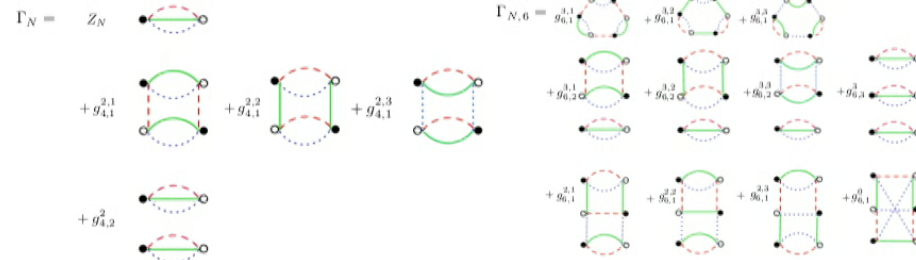
$$\theta' : \beta_g = \beta_g(g, \eta(g)_* = \text{const})$$

scheme	θ_1	θ_2	θ_3	θ_4	θ_5	θ_6	θ_7	θ_8	θ_9	θ_{10}	θ_{11}	θ_{12}	θ_{13}	θ_{14}	η
full	2.21	-0.24	-0.93	-0.93	—	—	—	—	—	—	—	—	—	—	-0.54
semi-pert	2	-0.21	-0.68	-0.68	—	—	—	—	—	—	—	—	—	—	-0.54
pert	2	0.69	-2	-2	—	—	—	—	—	—	—	—	—	—	-0.54
full	2.14	-0.59	-1.26	-1.26	-2.24	-2.54	-2.89	-2.89	-3.15	-3.26	-3.26	-3.31	-3.31	-3.89	-0.37
semi-pert	2	-0.58	-1.26	-1.26	-2.24	-2.54	-2.90	-2.90	-3.13	-3.26	-3.26	-3.31	-3.31	-3.90	-0.37
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full	0.93	-0.24	-0.93	-0.93	—	—	—	—	—	—	—	—	—	—	-0.54
semi-pert	0.93	-0.21	-0.93	-0.93	—	—	—	—	—	—	—	—	—	—	-0.54
full	1.33	-0.59	-1.26	-1.26	-2.30	-2.54	-2.89	-2.89	-3.15	-3.26	-3.26	-3.31	-3.31	-3.89	-0.37
semi-pert	1.33	-0.58	-1.26	-1.26	-2.31	-2.54	-2.90	-2.90	-3.13	-3.26	-3.26	-3.31	-3.31	-3.90	-0.37

[A.E., Koslowski '17, accepted in Annales de l'Institut Henri Poincaré D]

Fixed-point results



scheme	$g_{4,1*}^{2,1}$	$g_{4,2*}^2$	$g_{6,1*}^{3,1}$	$g_{6,1*}^{2,1}$	$g_{6,2*}^{3,1}$	$g_{6,3*}^3$
full	-1.94	0	—	—	—	—
semi-pert	-2.14	0	—	—	—	—
pert	-4.62	0	—	—	—	—
full	-1.37	0	-2.14	0	0	0
semi-pert	-1.47	0	-2.46	0	0	0
pert	-2.14	0	-6.12	0	0	0

$$\theta = 1$$

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towards convergence in extended truncations?

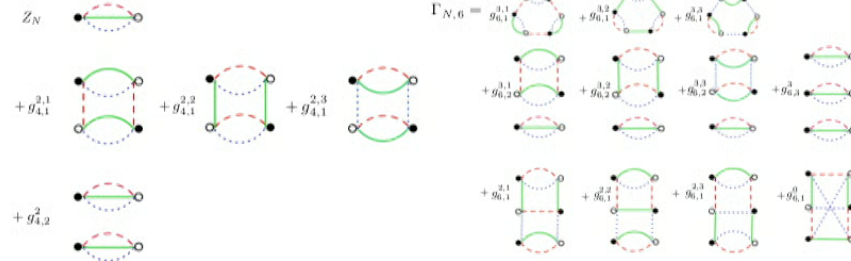
scheme	θ_1	θ_2	θ_3	θ_4	θ_5	θ_6	θ_7	θ_8	θ_9	θ_{10}	θ_{11}	θ_{12}	θ_{13}	θ_{14}	η
full	2.21	-0.24	-0.93	-0.93	—	—	—	—	—	—	—	—	—	—	-0.54
semi-pert	2	-0.21	-0.68	-0.68	—	—	—	—	—	—	—	—	—	—	-0.54
pert	2	0.69	-2	-2	—	—	—	—	—	—	—	—	—	—	-0.54
full	2.14	-0.59	-1.26	-1.26	-2.24	-2.54	-2.89	-2.89	-3.15	-3.26	-3.26	-3.31	-3.31	-3.89	-0.37
semi-pert	2	-0.58	-1.26	-1.26	-2.24	-2.54	-2.90	-2.90	-3.13	-3.26	-3.26	-3.31	-3.31	-3.90	-0.37
pert	2	-0.35	-2	-2	-3.23	-3.37	-3.43	-4	-4	-4.07	-4.07	-4.14	-4.14	-5	0

scheme	θ'_1	θ'_2	θ'_3	θ'_4	θ'_5	θ'_6	θ'_7	θ'_8	θ'_9	θ'_{10}	θ'_{11}	θ'_{12}	θ'_{13}	θ'_{14}	η
full	0.93	-0.24	-0.93	-0.93	—	—	—	—	—	—	—	—	—	—	-0.54
semi-pert	0.93	-0.21	-0.93	-0.93	—	—	—	—	—	—	—	—	—	—	-0.54
full	1.33	-0.59	-1.26	-1.26	-2.30	-2.54	-2.89	-2.89	-3.15	-3.26	-3.26	-3.31	-3.31	-3.89	-0.37
semi-pert	1.33	-0.58	-1.26	-1.26	-2.31	-2.54	-2.90	-2.90	-3.13	-3.26	-3.26	-3.31	-3.31	-3.90	-0.37

[A.E., Koslowski '17, accepted in Annales de l'Institut Henri Poincaré D]

Beyond double scaling?

scheme	$g_{4,1}^{2,1}$	$g_{4,2}^2$	$g_{6,1}^{3,1}$	$g_{6,1}^{2,1}$	$g_{6,2}^{3,1}$	$g_{6,3}^3$
full	-1.05	-1.33	—	—	—	—
semi-pert	-1.58	-1.05	—	—	—	—
pert	-4.62	1.73	—	—	—	—
full	-0.53	-2.11	-0.14	0	-0.39	—
semi-pert	-0.63	-2.35	-0.20	0	-0.57	—
pert	-2.01	-1.64	-4.43	0	-4.84	—
full	-1.04	-0.84	-0.97	0	-0.64	-0.99
semi-pert	-1.15	-0.89	-1.21	0	-0.78	-1.14
pert	-2.10	-0.54	-5.54	0	-1.68	-0.47



scheme	θ_1	θ_2	θ_3	θ_4	θ_5	θ_6	θ_7	θ_8	θ_9	θ_{10}	θ_{11}	θ_{12}	θ_{13}	θ_{14}	η
full	2.64	0.14	-0.65	-0.65	—	—	—	—	—	—	—	—	—	—	-0.67
semi-pert	2.28	0.14	-0.68	-0.68	—	—	—	—	—	—	—	—	—	—	-0.66
pert	2	-0.69	-2	-2	—	—	—	—	—	—	—	—	—	—	0
full	3.10	0.25	-0.48	-0.48	$-1.22 \pm i 0.44$	-1.53	-1.53	-1.72	-1.72	-2.46	-2.46	-2.72	—	—	-0.76
semi-pert	2.68	0.26	-0.51	-0.51	$-1.24 \pm i 0.45$	-1.57	-1.57	-1.76	-1.76	-2.49	-2.49	-2.76	—	—	-0.75
pert	2.11	0.34	-2	-2	-2.99	-2.99	-3.54	-3.54	-4	-4	-4.13	-4.13	-5	—	0
full	2.56	0.44	-0.97	-0.97	$-1.80 \pm i 0.42$	-2.45	-2.45	-2.63	-2.63	-2.95	-2.95	-3.06	-3.45	—	-0.52
semi-pert	2.30	0.44	-0.98	-0.98	$-1.81 \pm i 0.42$	-2.47	-2.47	-2.65	-2.65	-2.97	-2.97	-3.06	-3.47	—	-0.51
pert	2.04	0.37	-2	-2	-2.68	-3.16	-3.87	-3.95	-3.95	-4	-4	-4.09	-4.09	-5	0

scheme	θ'_1	θ'_2	θ'_3	θ'_4	θ'_5	θ'_6	θ'_7	θ'_8	θ'_9	θ'_{10}	θ'_{11}	θ'_{12}	θ'_{13}	θ'_{14}	η
full	0.65	0.47	-0.65	-0.65	—	—	—	—	—	—	—	—	—	—	-0.67
semi-pert	0.68	0.42	-0.68	-0.68	—	—	—	—	—	—	—	—	—	—	-0.66
full	1.03	0.43	-0.48	-0.48	$-1.23 \pm i 0.46$	-1.53	-1.53	-1.72	-1.72	-2.46	-2.46	-2.72	—	—	-0.76
semi-pert	1.04	0.44	-0.51	-0.51	$-1.25 \pm i 0.47$	-1.57	-1.57	-1.76	-1.76	-2.49	-2.49	-2.76	—	—	-0.75
full	1.17	0.64	-0.97	-0.97	$-1.82 \pm i 0.42$	-2.45	-2.45	-2.63	-2.63	-2.95	-2.95	-3.10	-3.45	—	-0.52
semi-pert	1.17	0.63	-0.98	-0.98	$-1.83 \pm i 0.42$	-2.47	-2.47	-2.65	-2.65	-2.97	-2.97	-3.10	-3.47	—	-0.51

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Summary

1) RG: Probing the scale dependence of QFTs

2) RG for spacetime in the continuum:

The price for setting a scale in Quantum Gravity

→ effective universality
in asymptotically
safe qg

3) RG for discrete spacetimes:

The price for setting a scale in Quantum Gravity

→ reason behind
large error in
estimate of
scaling exponent
in tensor models?

