Title: Long-Lived Inverse Chirp Signals from Core-Collapse in Massive Scalar-Tensor Gravity

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Abstract: We model stellar core collapse in massive scalar-tensor theories of gravity.

Pirsa: 17110100 Page 1/16

Long-Lived Inverse Chirp Signals from Core-Collapse in Massive Scalar-tensor Gravity

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Quantum Black Holes in the Sky? Perimeter Institute, 10 Nov 2017









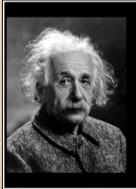


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Pirsa: 17110100 Page 2/16

Do we need a theory beyond GR?

When asked what he would do if Eddington's mission failed...



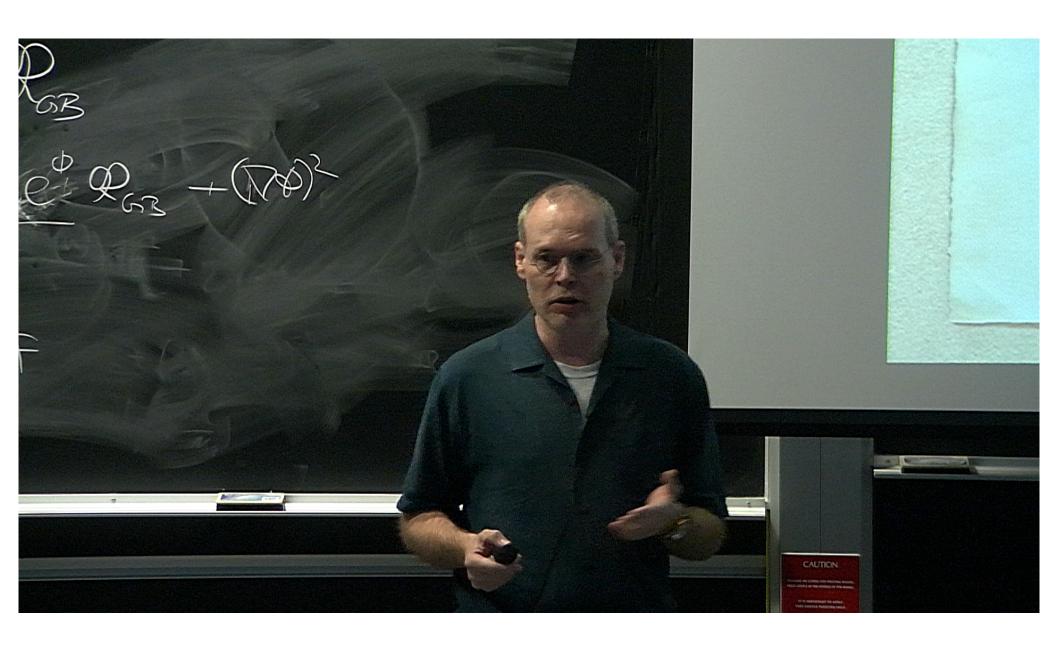
Then I would feel sorry for the good Lord. The theory is correct anyway.

(Albert Einstein)

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- But we have reasons to search for "beyond GR"
 - Renormalization: Requires, e.g., higher curvature terms.
 - → GR is low-energy limit of more fundamental theory
 - \bigcirc Dark energy: Why is Λ so small and why $ho_{
 m dark} \sim
 ho_{
 m mat}$
 - Dark matter: "Neptune" or "Vulcan" ?

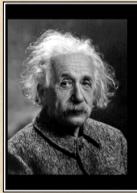
Pirsa: 17110100 Page 3/16



Pirsa: 17110100 Page 4/16

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Pirsa: 17110100 Page 5/16

Scalar tensor theory of gravity

- Scalars appear naturally in extra-dimensional theories
- Scalars prominent in cosmology
- ST theory well-posed; fairly well understood mathematically
- No-hair theorems limit potential of black-hole spacetimes
 ⇒ Matter: Neutron stars, core-collapse
- Best example of smoking gun to date:
 Spontaneous scalarization Damour & Esposito-Farese PRL 1993
- Collapse studies in massless case

Novak PRD 1998/1999 Novak & Ibanez ApJ 2000, Gerosa+ CQG 2016

Pirsa: 17110100 Page 6/16

Core-collapse scenario to 0th order

- ullet Massive stars: $M_{\mathrm{ZAMS}} = 8 \dots 100 \ M_{\odot}$
- ullet Core compressed from $\sim 1\,500~{
 m km}$ to $\sim 15~{
 m km}$ $\sim 10^{10}~{
 m g/cm^3}$ to $\gtrsim 10^{15}~{
 m g/cm^3}$
- ullet Released gravitational energy: $\mathcal{O}(10^{53})~\mathrm{erg}$ $\sim 99~\%$ in neutrinos, $\sim 10^{51}~\mathrm{erg}$ in outgoing shock, explosion
- Explosion mechanism: still uncertainties...
- Failed explosions lead to BH formation
 "Collapsar": possible engine for long-soft GRBs
- All of this handled for us by Woosley & Heger Phys.Rept. 2007
 - → Initial pre-collapse profile

Pirsa: 17110100 Page 7/16

Theoretical framework

Einstein frame: conformal metric $\bar{g}_{\mu\nu} = F(\varphi) g_{\mu\nu}$

Action

$$S = \frac{1}{16\pi} \int dx^4 \sqrt{-\bar{g}} \left[\bar{R} - 2\bar{g}^{\mu\nu} \partial_{\mu} \varphi \, \partial_{\nu} \varphi - 4V(\varphi) \right] + S_m [\psi_m, \bar{g}_{\mu\nu}/F(\varphi)]$$

- ullet Energy momentum tensor: $T_{\alpha\beta}=
 ho h u_{\alpha}u_{\beta}+Pg_{\alpha\beta}$
- Spherical symmetry: $d\bar{s}^2 = \bar{g}_{\mu\nu} dx^{\mu} dx^{\nu} = -F\alpha^2 dt^2 + FX^2 dr^2 + r^2 d\Omega^2$ $u^{\alpha} = \frac{1}{\sqrt{1-v^2}} [\alpha^{-1}, \ vX^{-1}, \ 0, \ 0]$
- Equations (gravity): $\partial_r \alpha = \dots$, $\partial_r X = \dots$ $\partial_t \partial_t \varphi = \dots$
- ullet Equations (matter): $(\rho,\,h,\,v) \leftrightarrow (D,\,S^r,\, au) \Rightarrow {\sf HRSC}$ GR1D code O'Connor & Ott CQG 2009

Equation of state

- ullet Pressure: "cold" + "thermal" contribution: $P=P_{
 m c}+P_{
 m th}$
- Internal energy from 1st law: $\epsilon_c = \begin{cases} \frac{K_1}{\Gamma_1 1} \rho^{\Gamma_1 1} & \text{if} \quad \rho \leq \rho_{\text{nuc}} \\ \frac{K_2}{\Gamma_2 1} \rho^{\Gamma_2 1} + E_3 & \text{if} \quad \rho > \rho_{\text{nuc}} \end{cases}$
- Thermal pressure: $P_{\rm th} = (\Gamma_{\rm th} 1)\rho(\epsilon \epsilon_{\rm th})$
- ullet Parameters: $\Gamma_1=1.3\,,\quad \Gamma_2=2.5\,,\quad \Gamma_{\mathrm{th}}=1.35\,$

$$K_1 = 4.9345 \times 10^{14} \text{ [cgs]}, \quad \rho_{\text{nuc}} = 2 \times 10^{14} \text{ g cm}^{-3}$$

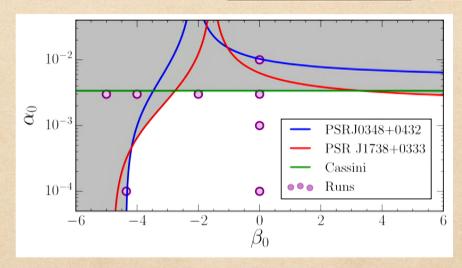
 K_2 , E_3 from continuity at $\rho = \rho_{\rm nuc}$

The coupling function and potential

Coupling function, potential: $F(\varphi) = e^{-2\alpha_0 \varphi - \beta_0 \varphi^2} \ | V(\varphi) = \frac{1}{2} \mu^2 \varphi^2$

$$F(\varphi) = e^{-2\alpha_0 \varphi - \beta_0 \varphi^2}$$

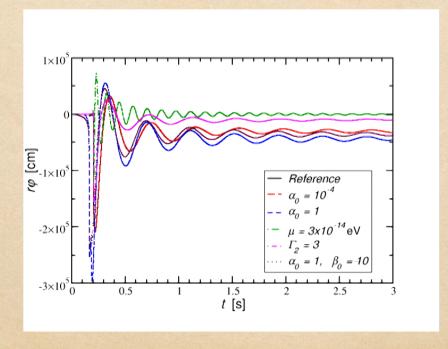
$$V(\varphi) = \frac{1}{2}\mu^2\varphi^2$$



- Only for $\mu \lesssim 10^{-19}\,\mathrm{eV}$!! Here: $\mu[\mathrm{eV}] \in [10^{-15}, 10^{-12}]$ Ramazanoglu & Pretorius PRD 2016
- Free parameters: μ , α_0 , β_0 , Γ_1 , Γ_2 , Γ_{th}

Waveforms "close to" the source

For $\mu = 10^{-14} \, \text{eV}$, $\alpha_0 = 10^{-2}$, $\beta_0 = -20$ $\Gamma_1 = 1.3$, $\Gamma_2 = 2.5$, $\Gamma_{\text{th}} = 1.35$



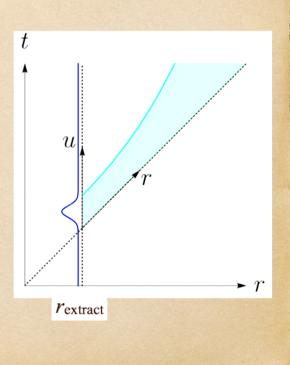
 \bullet $r\varphi \gg$ massless case; fairly insensitive to parameters; dispersion!

Waveforms "far from" the source

- LIGO will observe the above scalar profiles after they propagate to large distances
- In the massless case this is almost trivial $\varphi(t;r) = \frac{1}{r} \varphi(t-r;r_{\rm extract})$
- In the massive case things are more complicated: signals propagate with

dispersion

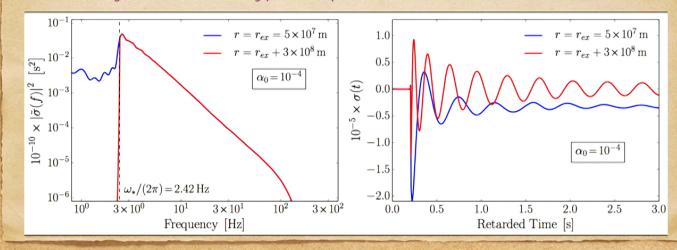
rextract



Waveforms "far from" the source

- Far from the source, scalar dynamics are governed by the flat-space Klein-Gordon wave equation $\ \partial_t^2 \varphi \nabla^2 \varphi + \omega_*^2 \varphi = 0$
- ullet Easier to work with the radially rescaled field $\sigma \equiv r arphi$
- As the signal propagates outwards:
 - Low frequencies are suppressed
 - High frequency power spectrum is unaffected
 - Signal spreads out in time
 - High frequencies arrive earlier than low frequencies
 - Signal becomes increasingly oscillatory

The scalar field mass has a natural frequency $\omega_* = c^2 \mu/\hbar$



Pirsa: 17110100 Page 13/16

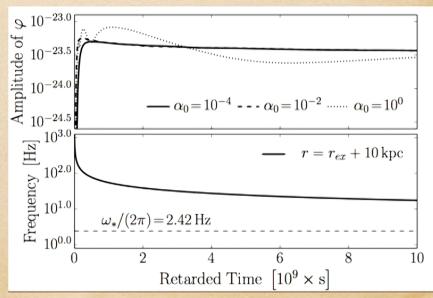
Waveforms "far from" the source

- Signals become more oscillatory as they propagate outwards
- In the large-distance limit the stationary phase approximation applies → analytic expression for the time domain signal
- Signals have a characteristic "inverse chirp" lasting many years

SPA frequency as function of time (Inverse Chirp)

$$F(t) = \frac{\omega_*}{2\pi} \frac{1}{\sqrt{1 - (d/t)^2}}$$

Distance to source $d = 10 \,\mathrm{kpc}$



Detection with LIGO-Virgo

GWs from core-collapse in ST gravity may fall into 3 classes:

- Burst signals: For light scalars $(\mu < 10^{-20}\,\mathrm{eV})$ and short distances $(10\,\mathrm{kpc})$, the pulse does not disperse significantly; will look like a $< 1\,\mathrm{s}$ burst
- Continuous wave signal: for heavier scalars, long dispersion turns pulse into a quasi-monochromatic signal
 - → capture using standard directed CW searches, assuming EM counterpart; e.g. SN1987A, Kepler1604
- Stochastic background:
 - Many quiet sources + very long duration (superposed)
 - Cosmological redshift + mass variation → smeared low-f cutoff
 - Characteristic "bump" in background, peaking at $\sim \omega_*$
 - Well in reach for aLIGO/AdVirgo stochastic searches

Pirsa: 17110100 Page 15/16

Conclusions

- We have simulated stellar core collapse in massive ST theory
- Explored combined parameter space of EOS and ST theory parameters
- Spontaneous scalarization occurs as in massless case, but effect can be more dramatic because the scalar mass "screens" the effect of the scalar, allowing larger values of α_0 , β_0 to be compatible with binary pulsar observations
- Signals propagate with dispersion, signals can last for years to centuries at kpc distances
- Signals can show up in LIGO/Virgo burst, CW or stochastic searches

Pirsa: 17110100 Page 16/16