

Title: PSI 17/18 - Condensed Matter - Lecture 8

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Abstract:

$$\langle E^{\text{ions}} \rangle_{\text{thermal}} = \sum_{|\mathbf{k}| < k_D} \langle \hbar \omega_{\mathbf{k}} (b_{\mathbf{k}}^{\dagger} b_{\mathbf{k}} + \frac{1}{2}) \rangle ; C_V = \frac{d}{dT} \langle E^{\text{ions}} \rangle$$

$$C_V = \frac{d}{dT} \int \frac{d^3 \mathbf{k}}{(2\pi)^3} \Omega \hbar \omega_{\mathbf{k}} \langle \hat{n}_{\mathbf{k}} \rangle ; \langle \hat{n}_{\mathbf{k}} \rangle = \frac{1}{e^{\beta \epsilon_{\mathbf{k}}} - 1} \quad \epsilon_{\mathbf{k}} = \hbar \omega_{\mathbf{k}}$$

$$= \frac{d}{dT} \int \frac{d^3 \mathbf{k}}{(2\pi)^3} \Omega \int d\epsilon \delta(\epsilon - \epsilon_{\mathbf{k}}) \frac{\epsilon}{e^{\beta \epsilon} - 1}$$

$$= \frac{d}{dT} \int d\epsilon \Omega \underbrace{\int \frac{d^3 \mathbf{k}}{(2\pi)^3} \delta(\epsilon - \epsilon_{\mathbf{k}})}_{\rho(\epsilon)} \frac{\epsilon}{e^{\beta \epsilon} - 1} = \frac{d}{dT} \int d\epsilon \Omega \rho(\epsilon) \frac{\epsilon}{e^{\beta \epsilon} - 1}$$

Electron phonon interaction

$$\hat{V}_{\text{ion-el}} = \sum_{i,n} V(|r_i - R_n|)$$

$$= \sum_{i,n} \frac{ze^2}{4\pi\epsilon_0} \frac{1}{|r_i - R_n|}$$

$r_i \equiv$ electronic position operator
 $R_n \equiv$ ionic " "

$$V(r_i, R_n) |\psi_m, \psi_n\rangle = \sum_{j,k} V_{jk}^{mn} |\psi_j, \psi_k\rangle$$

$$\langle \psi_j, \psi_k | V(r_i, R_n) | \psi_m, \psi_n \rangle$$

$$\hat{V}_{\text{tot}} = \sum_{m,n,j,k} V_{jk}^{mn} b_j^\dagger b_k^\dagger b_m b_n$$

$$\bar{R}_n = \bar{R}_n^0 + \bar{u}_n$$

$$V_{\text{ion-el}} = \underbrace{\sum_{i,n} V(\bar{r}_i - \bar{R}_n^0)}_{\text{Bloch bands}} - \underbrace{\sum_{i,n} \bar{u}_n \cdot \nabla_{\bar{R}_n} V(\bar{r}_i - \bar{R}_n)}_{\text{electron phonon interactions}} \Big|_{\bar{R}_n = \bar{R}_n^0}$$

$$\hat{T}_{\text{tot}} = \sum_{j,k} T_{jk} b_j^\dagger b_k = \sum_{\bar{k}} \frac{\hbar^2 |\bar{k}|^2}{2m} C_{\bar{k}}^\dagger C_{\bar{k}}$$

$$H_{el-ph} = - \sum_n \sum_{\substack{k' \sigma' \\ k \sigma}} \bar{U}_n \langle k' \sigma' | \nabla_{\bar{R}_n} V(\bar{r} - \bar{R}_n) |_{\bar{R}_n = \bar{R}_n^0} k \sigma \rangle C_{k' \sigma'}^{\dagger} C_{k \sigma}$$

$$= - \sum_n \bar{U}_n \sum_{\substack{k k' \\ \sigma}} \int d^3 \bar{r} \varphi_{k'}^*(\bar{r}) \nabla_{\bar{R}_n} V(\bar{r} - \bar{R}_n) |_{\bar{R}_n = \bar{R}_n^0} \varphi_k(\bar{r}) C_{k' \sigma}^{\dagger} C_{k \sigma}$$

$$H_{\text{el-ph}} = - \sum_n \bar{U}_n \langle k' \sigma' | \nabla_{\bar{R}_n} V(\bar{r} - \bar{R}_n) |_{\bar{R}_n = \bar{R}_n^0} | k \sigma \rangle C_{k' \sigma'}^+ C_{k \sigma}$$

$$= - \sum_n \bar{U}_n \sum_{k k' \sigma} \int d^3 r \varphi_{k'}^*(\bar{r}) \nabla_{\bar{R}_n} V(\bar{r} - \bar{R}_n) |_{\bar{R}_n = \bar{R}_n^0} \varphi_k(\bar{r}) C_{k' \sigma}^+ C_{k \sigma}$$

$$H_{\text{el-ph}} = - \sum_n \sum_{\substack{k' \sigma' \\ k \sigma}} \bar{U}_n \langle k' \sigma' | \nabla_{\bar{R}_n} V(\bar{r} - \bar{R}_n) |_{\bar{R}_n = \bar{R}_n^0} k \sigma \rangle C_{k' \sigma'}^{\dagger} C_{k \sigma}$$

$$= - \sum_n \bar{U}_n \sum_{\substack{k k' \\ \sigma}} \int d^3 \bar{r} \varphi_{k'}^*(\bar{r}) \nabla_{\bar{R}_n} V(\bar{r} - \bar{R}_n) |_{\bar{R}_n = \bar{R}_n^0} \varphi_k(\bar{r}) C_{k' \sigma}^{\dagger} C_{k \sigma}$$

$$\varphi_{\bar{k}}(\bar{r} + \bar{R}_n^0) = e^{i \bar{k} \cdot \bar{R}_n^0} \varphi_{\bar{k}}(\bar{r})$$

$$\langle k' | \nabla_{\vec{R}_n} V(|\vec{r} - \vec{R}_n|) |_{\vec{R}_n = \vec{R}_n^0} | k \rangle$$

↑ Replace $\vec{r} = \vec{r} - \vec{R}_n^0$

$$= e^{i(\vec{k} - \vec{k}') \cdot \vec{R}_n^0} \int d^3(\vec{r} - \vec{R}_n^0) \varphi_{\vec{k}'}^*(\vec{r} - \vec{R}_n^0) \nabla_{\vec{R}_n} V(\vec{r} - \vec{R}_n) |_{\vec{R}_n = \vec{R}_n^0} \varphi_{\vec{k}}(\vec{r} - \vec{R}_n^0)$$

$$= e^{i(\vec{k} - \vec{k}') \cdot \vec{R}_n^0} \sum_{\vec{k}, \vec{k}'}$$

$$H_{el-ph} = - \sum_{\vec{k}, \vec{k}'} S_{\vec{k}, \vec{k}'} \sum_n e^{i(\vec{k} - \vec{k}') \cdot \vec{R}_n^0} \bar{u}_n c_{\vec{k}'}^+ c_{\vec{k}0}$$

$$U_n = \frac{1}{\sqrt{N}} \sum_{\mathbf{k}} e^{i\mathbf{k}\cdot\mathbf{R}_n} Q_{\mathbf{k}} \quad \text{and} \quad Q_{\mathbf{k}} = \sqrt{\frac{\hbar}{2M\omega_{\mathbf{k}}}} (b_{-\mathbf{k}}^{\dagger} + b_{\mathbf{k}})$$

$$\bar{U}_n = \frac{1}{\sqrt{N}} \sum_{\substack{\mathbf{q}, \lambda \\ \mathbf{q} \in \text{FBZ}}} \sqrt{\frac{\hbar}{2M\omega_{\lambda}(\mathbf{q})}} e^{i\mathbf{q}\cdot\mathbf{R}_n} (b_{-\mathbf{q}, \lambda}^{\dagger} + b_{\mathbf{q}, \lambda}) \bar{A}_{\lambda}$$

$\lambda \in$ All phonon branches

$$\bar{R}_n = \bar{R}_n^0 \quad \varphi_{\mathbf{k}}(\bar{\mathbf{r}} - \bar{R}_n^0)$$

$$+ \quad k_0^{\dagger} c_{k_0}$$

$$U_n = \frac{1}{\sqrt{N}} \sum_{\mathbf{k}} e^{i\mathbf{k}\cdot\mathbf{r}_n} Q_{\mathbf{k}} \quad \text{and} \quad Q_{\mathbf{k}} = \sqrt{\frac{\hbar}{2M\omega_{\mathbf{k}}}} (b_{-\mathbf{k}}^{\dagger} + b_{\mathbf{k}})$$

$$\bar{U}_n = \frac{1}{\sqrt{N}} \sum_{\substack{q\lambda \\ q \in \text{FBZ}}} \sqrt{\frac{\hbar}{2M\omega_{\lambda}(q)}} e^{i\bar{q}\cdot\bar{\mathbf{r}}_n} (b_{-q,\lambda}^{\dagger} + b_{q,\lambda}) \bar{A}_{\lambda}$$

$\lambda \in$ All phonon branches

$\bar{A}_{\lambda} \equiv$ polarization for the λ^{th} branch

$$\begin{aligned} & \sum_{\mathbf{G}} \sum_{\mathbf{k}, \mathbf{k}', \mathbf{G}} \\ &= \sum_{\mathbf{G}} \frac{1}{N} \sum_n e^{i(\mathbf{k}-\mathbf{k}'+\mathbf{G})\cdot\bar{\mathbf{r}}_n} \\ &= \frac{1}{N} \sum_{\mathbf{G}} \sum_n e^{i(\mathbf{k}-\mathbf{k}')\cdot\bar{\mathbf{r}}_n} \end{aligned}$$

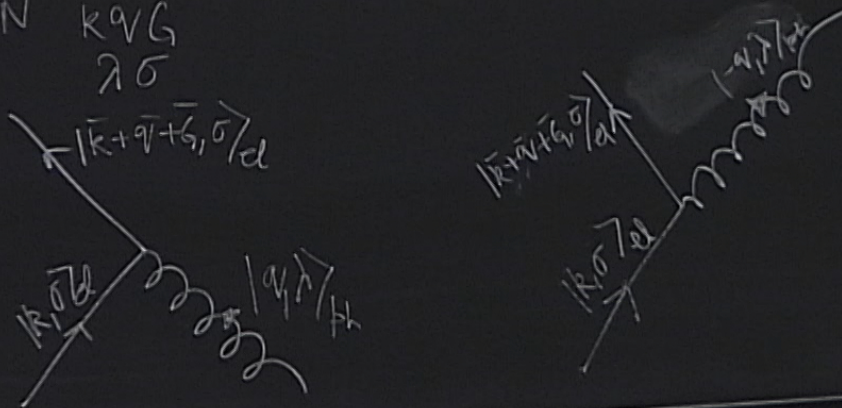
$$H_{el-ph} = \frac{1}{\sqrt{N}} \sum_{\substack{k, k', q \\ n, \lambda, \sigma}} S_{kk'} \cdot \bar{A}_{\lambda} \sqrt{\frac{\hbar}{2M\omega_{\lambda}(q)}} e^{i(\bar{k}-\bar{k}'+\bar{q}) \cdot \bar{R}_n^0} (b_{-q, \lambda}^{\dagger} + b_{q, \lambda}) c_{k\sigma}^{\dagger} c_{k\sigma}$$

$$= \frac{1}{\sqrt{N}} \sum_{\substack{\mathbf{k}, \mathbf{k}' \\ n, \lambda, \sigma}} S_{\mathbf{k}\mathbf{k}'} \cdot \bar{A}_{\lambda} \sqrt{\frac{\hbar}{2M\omega_{\lambda}(q)}} e^{i(\bar{\mathbf{k}} - \bar{\mathbf{k}}' + \bar{\mathbf{q}}) \cdot \bar{\mathbf{R}}_n^0} (b_{-\mathbf{q}, \lambda}^{\dagger} + b_{\mathbf{q}, \lambda}) c_{\mathbf{k}'\sigma}^{\dagger} c_{\mathbf{k}\sigma}$$

$$= \frac{1}{\sqrt{N}} \sum_{\substack{\mathbf{k}, \mathbf{q}, \mathbf{G} \\ \lambda, \sigma}} S_{\bar{\mathbf{k}}, \bar{\mathbf{k}} + \bar{\mathbf{q}} + \bar{\mathbf{G}}} \cdot \bar{A}_{\lambda} \sqrt{\frac{\hbar}{2M\omega_{\lambda}(q)}} (b_{-\mathbf{q}, \lambda}^{\dagger} + b_{\mathbf{q}, \lambda}) c_{\bar{\mathbf{k}} + \bar{\mathbf{q}} + \bar{\mathbf{G}}, \sigma}^{\dagger} c_{\mathbf{k}\sigma}$$

$$H_{el-ph} = \frac{1}{\sqrt{N}} \sum_{\substack{k, k', q \\ n, \lambda, \sigma}} S_{kk'} \cdot \bar{A}_\lambda \sqrt{\frac{\hbar}{2M\omega_\lambda(q)}} e^{i(\bar{k}-\bar{k}'+\bar{q}) \cdot \bar{R}_n^0} (b_{-q, \lambda}^+ + b_{q, \lambda}) C_{k\sigma}^+ C_{k\sigma}$$

$$= \frac{1}{\sqrt{N}} \sum_{\substack{k, q, G \\ \lambda, \sigma}} S_{\bar{k}, \bar{k}+\bar{q}+\bar{G}} \cdot \bar{A}_\lambda \sqrt{\frac{\hbar}{2M\omega_\lambda(q)}} (b_{-q, \lambda}^+ + b_{q, \lambda}) C_{\bar{k}+\bar{q}+\bar{G}, \sigma}^+ C_{k\sigma}$$



$c_{k\sigma}^\dagger c_{k\sigma}$

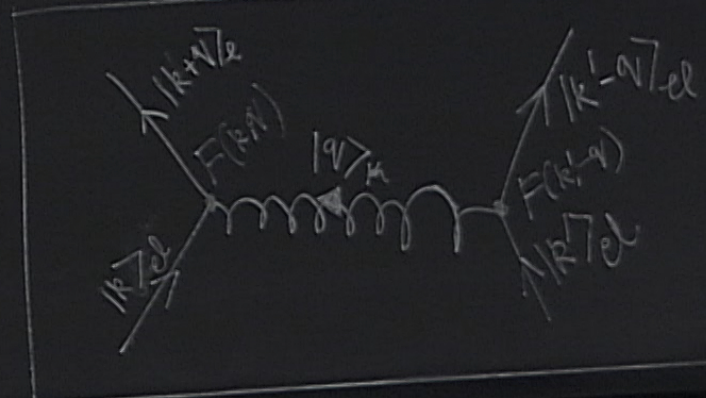
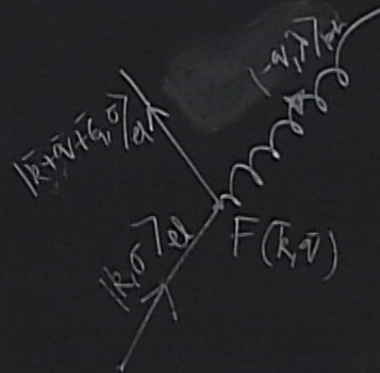
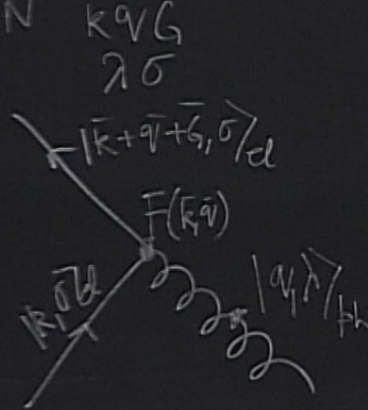
$c_{k\sigma}^\dagger c_{k\sigma}$

$$\sum_{\vec{G}} = (\vec{G}=0 \text{ term}) + \sum_{\vec{G} \neq 0} \text{(Umklapp processes)} \propto \frac{1}{(q+\vec{G})^2}$$

$$\tilde{V}(q) = \frac{1}{\Omega} \int d^3r \frac{V(\vec{r})}{|\vec{r}|} e^{i\vec{q} \cdot \vec{r}} \propto \frac{1}{q^2}$$

$$= \frac{1}{\sqrt{N}} \sum_{\substack{\mathbf{k}, \mathbf{k}' \\ n, \lambda, \sigma}} S_{\mathbf{k}\mathbf{k}'} \cdot \bar{A}_\lambda \sqrt{\frac{\hbar}{2M\omega_\lambda(q)}} e^{i(\bar{\mathbf{k}} - \bar{\mathbf{k}}' + \bar{\mathbf{q}}) \cdot \bar{\mathbf{R}}_n^0} (b_{-\mathbf{q}, \lambda}^\dagger + b_{\mathbf{q}, \lambda}) c_{\mathbf{k}'\sigma}^\dagger c_{\mathbf{k}\sigma}$$

$$= \frac{1}{\sqrt{N}} \sum_{\substack{\mathbf{k}, \mathbf{q}, \mathbf{G} \\ \lambda, \sigma}} S_{\bar{\mathbf{k}}, \bar{\mathbf{k}} + \bar{\mathbf{q}} + \bar{\mathbf{G}}} \cdot \bar{A}_\lambda \sqrt{\frac{\hbar}{2M\omega_\lambda(q)}} (b_{-\mathbf{q}, \lambda}^\dagger + b_{\mathbf{q}, \lambda}) c_{\bar{\mathbf{k}} + \bar{\mathbf{q}} + \bar{\mathbf{G}}, \sigma}^\dagger c_{\mathbf{k}\sigma}$$



ko ko

Schrieffer-Wolff transformation

$$U = e^S, \quad U^\dagger = e^{S^\dagger} = e^{-S} \quad \boxed{S^\dagger = -S}$$

$$\tilde{H} = U^\dagger H U = e^{-S} H e^S = H - [S, H] + \frac{(-1)^2}{2} [S, [S, H]] + (\dots)$$

k_0 k_0

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$$H = H_0 + H_{\text{el-ph}} \quad \text{where} \quad H_0 = H_{\text{el}} + H_{\text{ph}}$$

$$[S, H_0] = H_{\text{el-ph}} \quad \left| \quad \tilde{H} = H_0 - [S, H_{\text{el-ph}}] + \frac{1}{2} [S, H_{\text{el-ph}}] + \frac{1}{2} [S, [S, H_{\text{el-ph}}]] + (\dots)$$

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$$[S, H_0] = H_{\text{el-ph}} \quad \left| \quad \tilde{H} = H_0 - [S, H_{\text{el-ph}}] + \frac{1}{2} [S, H_{\text{el-ph}}] + \frac{1}{2} [S, [S, H_{\text{el-ph}}]] + (\dots) \right. \quad \left. \begin{matrix} O(F^3) \\ \end{matrix} \right.$$

$$\hat{H} = H_0 - \frac{1}{2} [S, H_{d-ph}] + O(P^3 \text{ or higher})$$

$$S = \sum_{\vec{k}, \vec{q}} F(\vec{k}, \vec{q}) C_{\vec{k}+\vec{q}}^\dagger C_{\vec{k}} (y b_{-\vec{q}}^\dagger + x b_{\vec{q}})$$

Using this Ansatz & demanding $[S, H_0] = H_{d-ph}$

$$x = \frac{1}{\epsilon_{\vec{k}+\vec{q}} - \epsilon_{\vec{k}} - \hbar\omega(\vec{q})} \quad \text{and} \quad y = \frac{1}{\epsilon_{\vec{k}+\vec{q}} - \epsilon_{\vec{k}} + \hbar\omega(\vec{q})}$$

$$\hat{H} = H_0 - \frac{1}{2} [S, H_{d-ph}] + O(\text{P}^3 \text{ or higher})$$

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$$-\frac{1}{2} [S, H_{cl-ph}]$$

$$[b_{q_1}, b_{-q}^+] = \delta_{q_1, -q}$$

$$\hat{H} = H_{cl} + H_{ph} + \sum_{k, k', q} \frac{F(k, q) F(k' - q) (\hbar \omega(q))}{(\epsilon_{k+q} - \epsilon_k)^2 - (\hbar \omega(q))^2} c_{k+q}^+ c_{k'-q}^+ c_{k'} c_k$$

$$-\frac{1}{2} [S, H_{d-ph}]$$

$$\hat{H} = H_d + H_{ph} + \sum_{\mathbf{k}, \mathbf{k}+\mathbf{q}} \frac{F(\mathbf{k}, \mathbf{q}) F(\mathbf{k}-\mathbf{q}) (\hbar \omega(\mathbf{q}))}{(\epsilon_{\mathbf{k}+\mathbf{q}} - \epsilon_{\mathbf{k}})^2 - (\hbar \omega(\mathbf{q}))^2} c_{\mathbf{k}+\mathbf{q}}^\dagger c_{\mathbf{k}-\mathbf{q}}^\dagger c_{\mathbf{k}} c_{\mathbf{q}}$$

$$[b_{\mathbf{q}}, b_{-\mathbf{q}}^\dagger] = \delta_{\mathbf{q}, -\mathbf{q}}$$

Consider $q \rightarrow 0$, $\omega(q) = \text{constant} \neq 0$ for optical phonons

Consider TWO electrons close to the Fermi Sea

$$(\epsilon_{\mathbf{k}+\mathbf{q}} - \epsilon_{\mathbf{k}})^2 = \frac{\hbar^2 (\mathbf{k}_F + \mathbf{q})^2}{2m} - \frac{\hbar^2 \mathbf{k}_F^2}{2m} \rightarrow 0 \text{ as } q \rightarrow 0$$



$$\frac{1}{2} [S, H_{el-ph}]$$

$$\hat{H} = H_{el} + H_{ph} + \sum_{\mathbf{k}, \mathbf{q}} \frac{F(\mathbf{k}, \mathbf{q}) F(\mathbf{k}-\mathbf{q}) (\hbar \omega(\mathbf{q}))}{(\epsilon_{\mathbf{k}+\mathbf{q}} - \epsilon_{\mathbf{k}})^2 - (\hbar \omega(\mathbf{q}))^2} c_{\mathbf{k}+\mathbf{q}}^\dagger c_{\mathbf{k}-\mathbf{q}}^\dagger c_{\mathbf{k}} c_{\mathbf{q}}$$

$$[b_{\mathbf{q}}, b_{-\mathbf{q}}^\dagger] = \delta_{\mathbf{q}, -\mathbf{q}}$$

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$$(\epsilon_{\mathbf{k}+\mathbf{q}} - \epsilon_{\mathbf{k}})^2 = \frac{\hbar^2 (\mathbf{k}_F + \mathbf{q})^2}{2m} - \frac{\hbar^2 \mathbf{k}_F^2}{2m} \rightarrow 0 \text{ as } q \rightarrow 0$$

$$E_F = \frac{\hbar^2 k_F^2}{2m}$$



$$S^+ = -S$$

$$[S, H] + \frac{(-1)^2}{2} [S, [S, H]] + (\dots)$$

$$+ \frac{1}{2} [S, H_{\text{cl-ph}}] + \frac{1}{2} [S, [S, H_{\text{cl-ph}}]] + (\dots)$$

$O(F^3)$

