

Title: Bounds on new physics from electric dipole moments

Date: Oct 31, 2017 01:00 PM

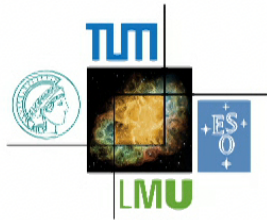
URL: <http://pirsa.org/17100068>

Abstract: 

Electric dipole moments are extremely sensitive probes for additional sources of CP violation in new physics models. The multi-scale problem of relating the high precision measurements with neutrons, atoms and molecules to fundamental parameters can be approached model-independently to a large extent; however, care must be taken to include the uncertainties from especially nuclear and QCD calculations properly. The resulting bounds on fundamental parameters are illustrated in the context of Two-Higgs-Doublet models.

# Bounds on new physics from EDMs

Martin Jung

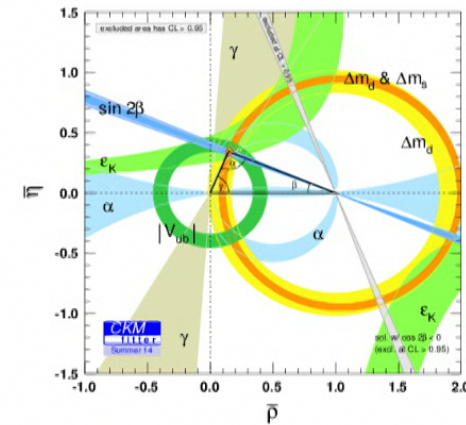


Seminar  
31st of October 2017  
PI, Waterloo, ON, Canada

## Motivation

Quark-flavour and CP violation in the SM:

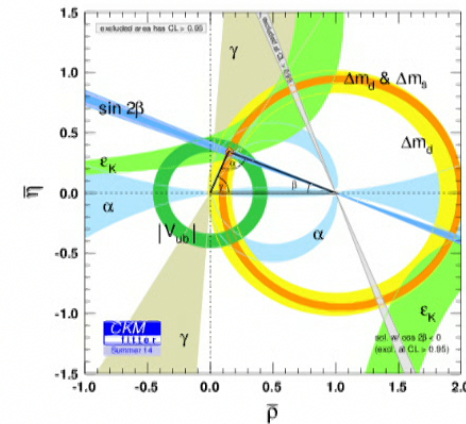
- CKM describes flavour **and** CP violation
- Extremely constraining, one phase
- Especially,  $K$  and  $B$  physics agree
- Only tensions so far  
( $R_{K,K^*}, P'_5, B \rightarrow D^{(*)} \tau \nu, g_{\mu} - 2, \dots$ )
- ➡ Works well!



## Motivation

Quark-flavour and CP violation in the SM:

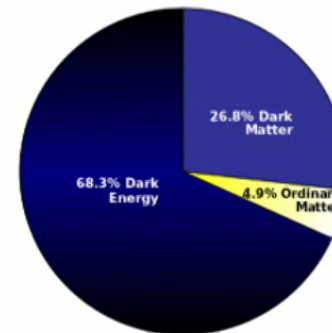
- CKM describes flavour **and** CP violation
  - Extremely constraining, one phase
  - Especially,  $K$  and  $B$  physics agree
  - Only tensions so far  
( $R_{K,K^*}, P'_5, B \rightarrow D^{(*)}TV, g_{\mu} - 2, \dots$ )
- ➔ Works **too** well!



We expect new physics (ideally at the (few-)TeV scale):

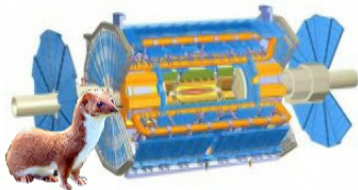
- Baryon asymmetry of the universe
- Hierarchy problem
- Dark matter and energy
- ...

➔ So where is it?



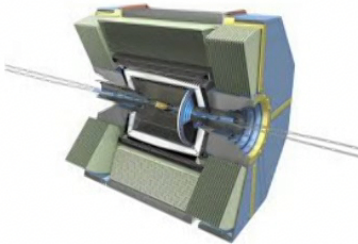
## The Quest for New Physics

Three of the main strategies (missing are e.g.  $\nu$ , DM, astro, ...):



### Direct search:

- Tevatron, LHC
- Maximal energy fixed



### Indirect search, flavour violating:

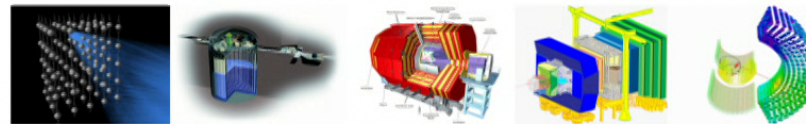
- LHCb, Belle II, BES III, NA62, MEG, ...
- Maximal reach flexible



### Indirect search, flavour diagonal:

- **EDM experiments**, g-2, ...
- Maximal reach flexible, complementary to flavour-violating searches

**A new era in  
particle physics!**

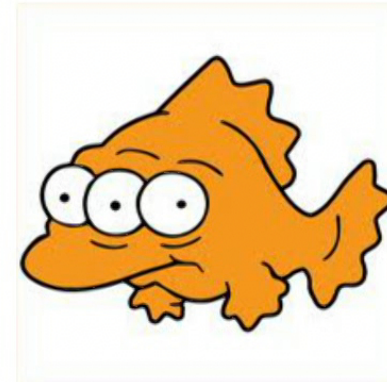


## The curious case of the One-Higgs-Doublet Model

EDMs are finite in the SM...

...but flavour-sector of the SM is special ( $\rightarrow$ ):

- Unique connection between Flavour- and CP-violation
- FCNCs highly suppressed,  $\sim \Delta m^2 / M_W^2$   
 $\rightarrow \Delta m^2 / M_W^2 \sim 10^{-25}$  for  $\nu$  in the loop!
- FC**onserving**NCs with CPV as well:  
 $\rightarrow d_e^{SM} \lesssim 10^{-38} e \text{ cm}$  [Khriplovich/Pospelov '91]



EDMs are quasi-nulltests of the SM!

NP models typically do **not** exhibit such strong cancellations

- $\rightarrow$  Background-free precision-laboratories for NP  
(assuming dynamical solution for strong CP)
- $\rightarrow$  EDMs  $\sim CPV / \Lambda^2$  (interference with SM, e.g. LFV  $\sim 1 / \Lambda^4$ )

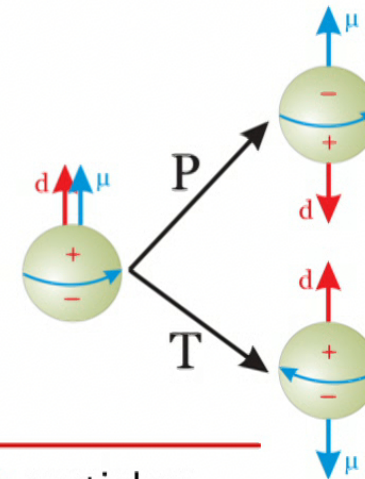
Here: focus as much as possible on model-independent statements

## Back to basics: EDMs

Classically:  $\mathbf{d} = \int d^3r \rho(\mathbf{r})\mathbf{r}$ ,  $U = \mathbf{d} \cdot \mathbf{E}$

QM: non-degenerate ground state implies  $\mathbf{d} \sim \mathbf{j}$

- ➔  $\mathbf{d} \neq \mathbf{0}$  implies T- and P-violation!
- ➔ CP-violation for conserved CPT
- ➔ Search for linear shift  $U = d\mathbf{j} \cdot \mathbf{E}$



**Non-relativistic** neutral system of **point-like** particles:  
Potential EDMs of constituents are shielded! [Schiff'63]

- ➔ Sensitivity stems from violations of the assumptions
  - Paramagnetic systems: relativistic enhancement
  - Diamagnetic systems: finite-size effects

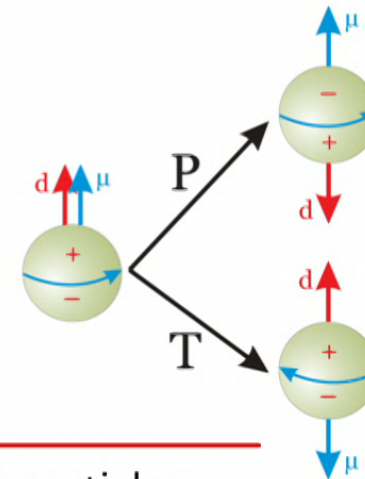
Shielding can be reversed, e.g.  $d_A^{\text{para}} \sim \mathcal{O}(100) \times d_e!$   
[Sandars'65, '66]

## Back to basics: EDMs

Classically:  $\mathbf{d} = \int d^3r \rho(\mathbf{r}) \mathbf{r}$ ,  $U = \mathbf{d} \cdot \mathbf{E}$

QM: non-degenerate ground state implies  $\mathbf{d} \sim \mathbf{j}$

- ➔  $\mathbf{d} \neq \mathbf{0}$  implies T- and P-violation!
- ➔ CP-violation for conserved CPT
- ➔ Search for linear shift  $U = d \mathbf{j} \cdot \mathbf{E}$



Non-relativistic neutral system of point-like particles:  
Potential EDMs of constituents are shielded! [Schiff'63]

- ➔ Sensitivity stems from violations of the assumptions
  - Paramagnetic systems: relativistic enhancement
  - Diamagnetic systems: finite-size effects

Shielding can be reversed, e.g.  $d_A^{\text{para}} \sim \mathcal{O}(100) \times d_e!$   
[Sandars'65, '66]

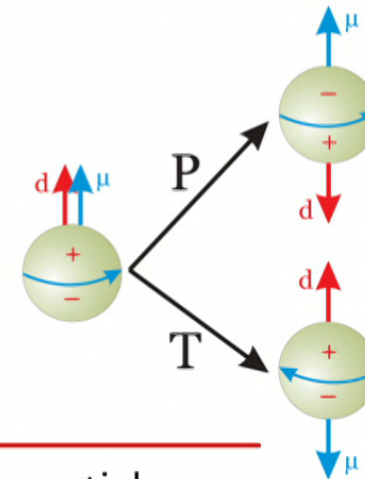


## Back to basics: EDMs

Classically:  $\mathbf{d} = \int d^3r \rho(\mathbf{r}) \mathbf{r}$ ,  $U = \mathbf{d} \cdot \mathbf{E}$

QM: non-degenerate ground state implies  $\mathbf{d} \sim \mathbf{j}$

- ➔  $\mathbf{d} \neq \mathbf{0}$  implies T- and P-violation!
- ➔ CP-violation for conserved CPT
- ➔ Search for linear shift  $U = d \mathbf{j} \cdot \mathbf{E}$



Non-relativistic neutral system of point-like particles:  
Potential EDMs of constituents are shielded! [Schiff'63]

- ➔ Sensitivity stems from violations of the assumptions
  - Paramagnetic systems: relativistic enhancement
  - Diamagnetic systems: finite-size effects

Shielding can be reversed, e.g.  $d_A^{\text{para}} \sim \mathcal{O}(100) \times d_e!$   
[Sandars'65, '66]

## EDMs and New Physics: Generalities

Sakharov's conditions ('67):  
NP models necessarily involve new sources of CPV!

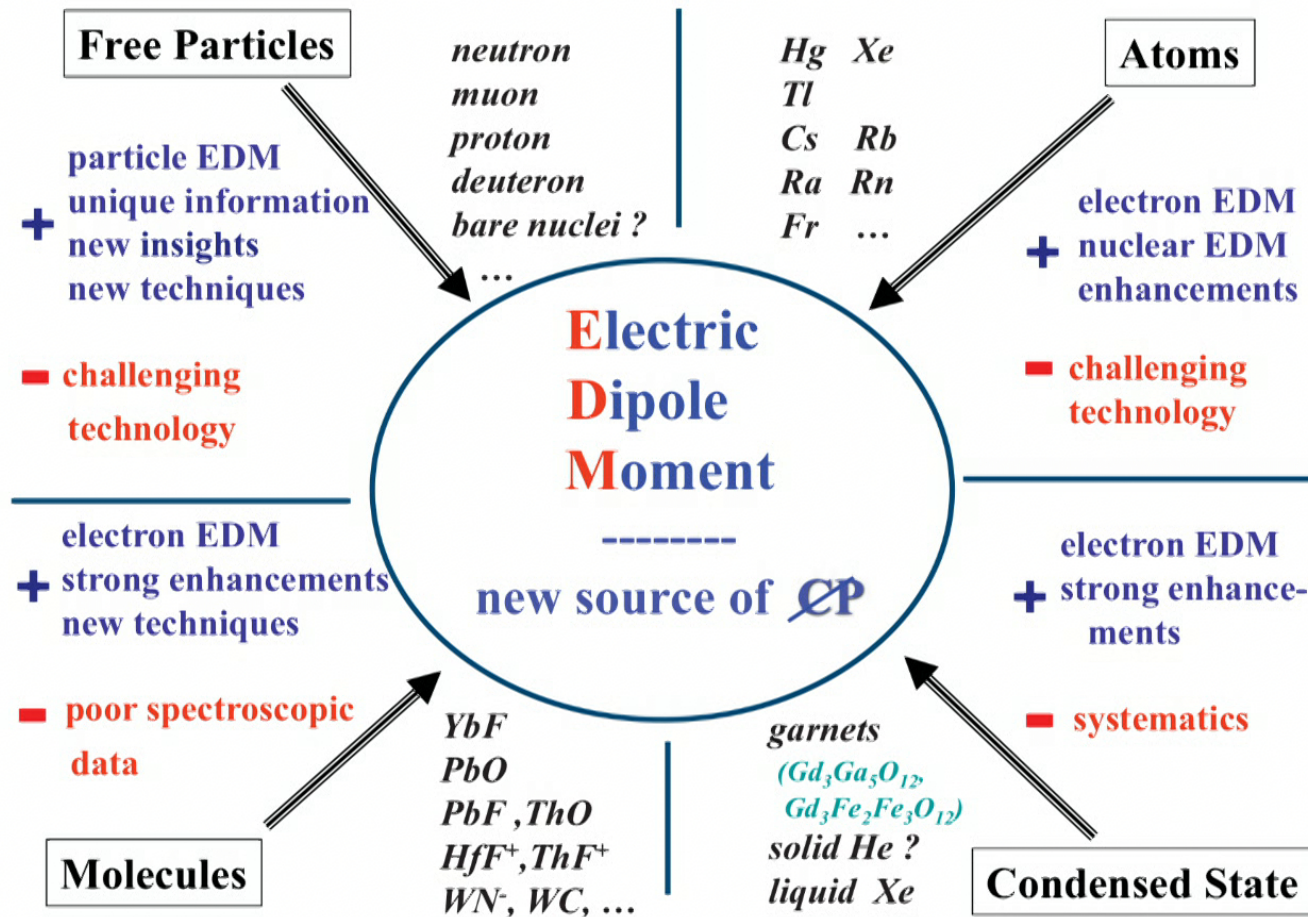
- This does not *imply* sizable EDMs
- However, typically (too) large EDMs in NP models
- Generic one-loop contributions excluded  
(→ SUSY CP-problem)
- EDMs test combination of flavour- and CPV-structure

EDMs important on two levels:

- “Smoking-gun-level”: Visible EDMs proof for NP
- Quantitative level:  
Setting limits/determining parameters
- Theory uncertainties are important!

# Experimental approaches [K. Jungmann'13 in Annalen der Physik]

## Lines of attack towards an EDM



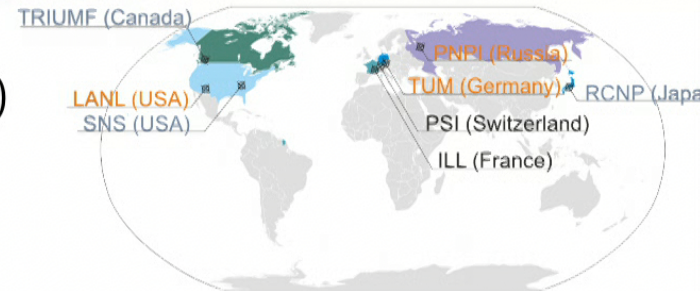
## Experimental status

### Neutron EDM:

- $|d_n| \leq 3.6 \times 10^{-26} e \text{ cm}$  (95%CL)

[Pendlebury+'15, Baker+'06]

- Worldwide effort aiming at  $(10 \rightarrow 0.1) \times 10^{-27} e \text{ cm}$
- UCN sources critical problem



[P.Schmidt-Wellenburg'16]

### Paramagnetic systems:

- Atomic:  $|d_{\text{Tl}}| \leq 9.6 \times 10^{-25} e \text{ cm}$  (95% CL) [Regan+'02]
- Molecular:  $|\omega_{\text{ThO}}| \leq 11.1 \text{ mrad/s}$  (90% CL) [Baron+'13]
- Naive interpretation:  $|d_e| \leq 8.7 \times 10^{-29} e \text{ cm}$
- Ongoing: ThO, YbF, Cs, Fr, Rb, HfF<sup>+</sup>...

### Diamagnetic systems:

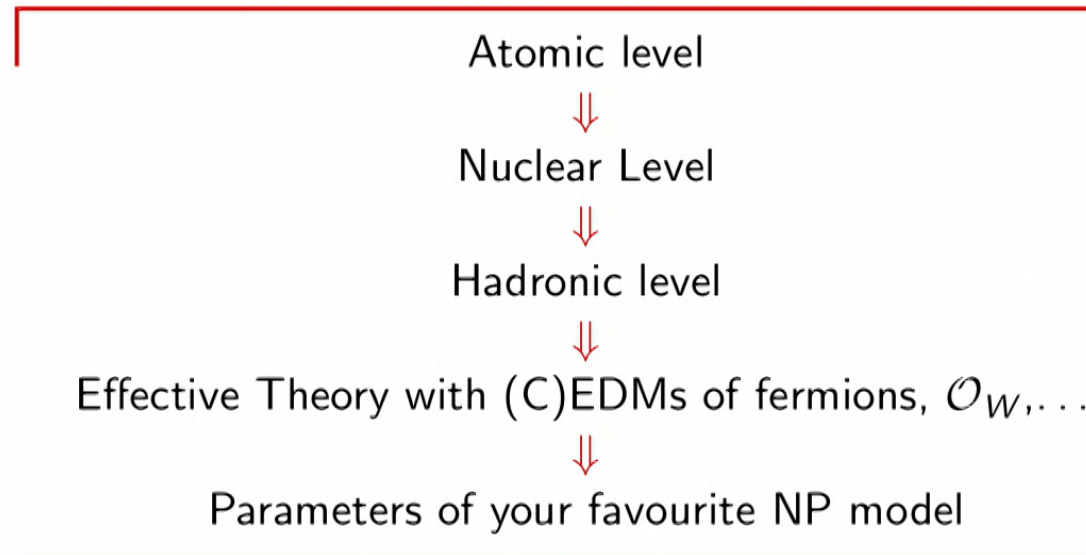
- $|d_{\text{Hg}}| \leq 7.4 \times 10^{-30} e \text{ cm}$  (95% CL) [Graner+'16]
- Ongoing: exploit **octupole deformation**, e.g. Ra, Rn,...

**Solid state systems:**  $|d_e| \leq 6.1 \times 10^{-24-25} e \text{ cm}$  [Eckel+'12, Kim+'15]

**Storage rings:**  $|d_\mu| \leq 1.9 \times 10^{-19} e \text{ cm}$  [Bennett+'08]

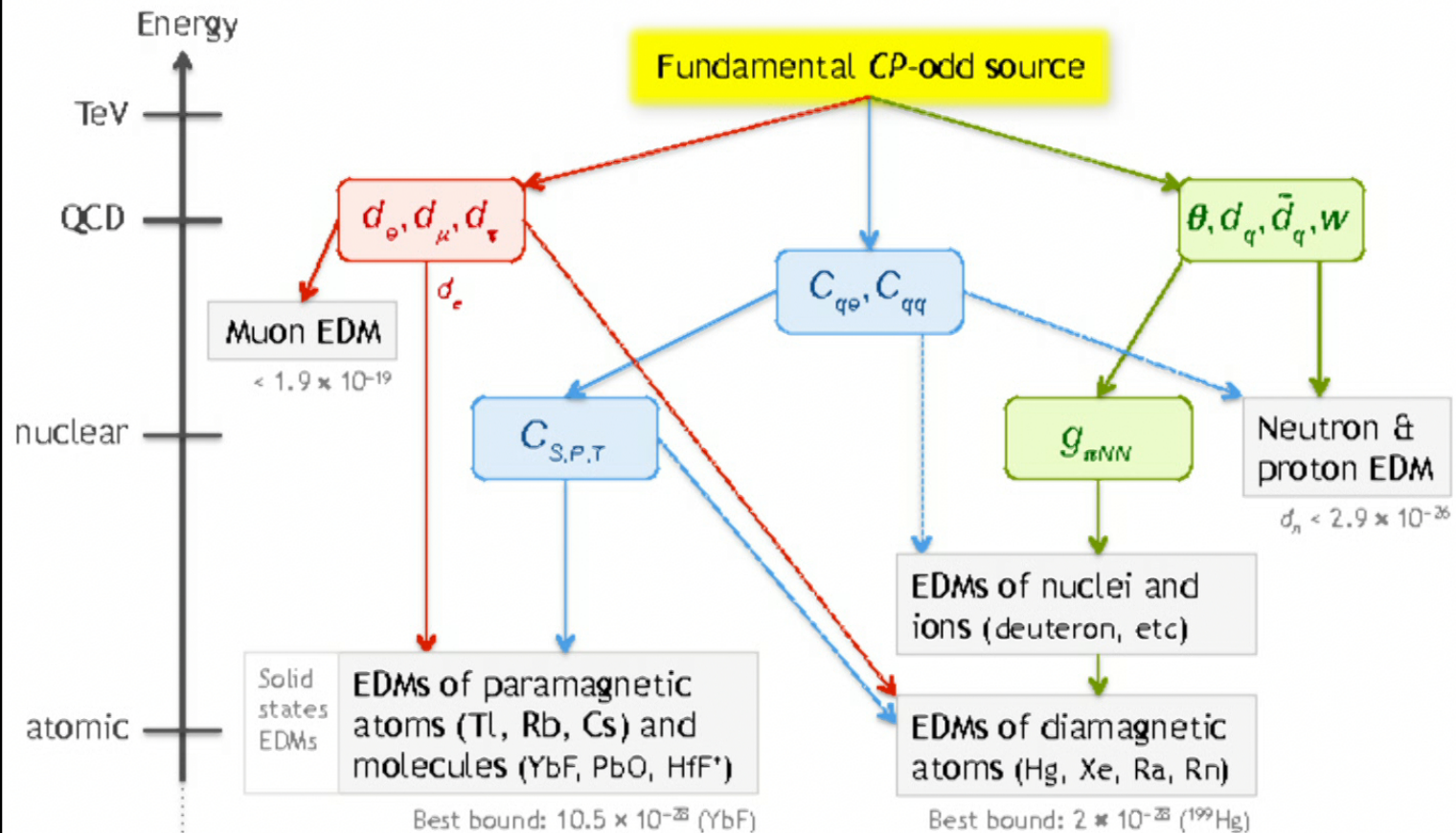
## Relating NP parameters and experiment

- Most stringent constraints from neutron, atoms and molecules
  - ➔ Shielding typically applies



- Each step potentially involves large uncertainties!
- 4/5 model-independent  $\Rightarrow$  series of EFTs [e.g. deVries+'11]
- Limits usually displayed as allowed regions
  - ➔ Conservative uncertainty estimates important

# Schematic EFT framework [Pospelov/Ritz'05,Hoecker'12]



## The EDM in heavy paramagnetic systems

Two main contributions, enhanced by  $Z^3$ : [Sandars'65, Flambaum'76]

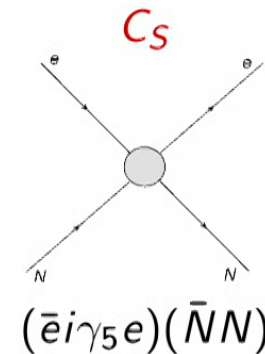
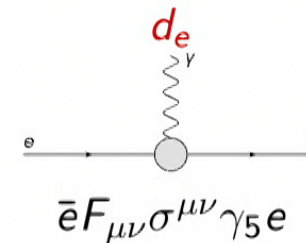
➡ A single measurement does **not** restrict  $d_e$  directly

- $C_S$ : CP-odd electron-nucleon interaction
- Atoms: typically polarized in external field
- Molecules: aligned in external field
  - ➡ Exploit huge internal field

For molecules: energy shift  $\Delta E = \hbar\omega$  with

$$\omega = 2\pi \left( \frac{W_d^M}{2} d_e + \frac{W_c^M}{2} C_S \right) .$$

Molecule	$W_d^M/10^{25}\text{Hz}/e\text{ cm}$	$W_c^M/\text{kHz}$
YbF	$-1.3 \pm 0.1$	$-92 \pm 9$
ThO	$-3.67 \pm 0.18$	$-598 \pm 90$

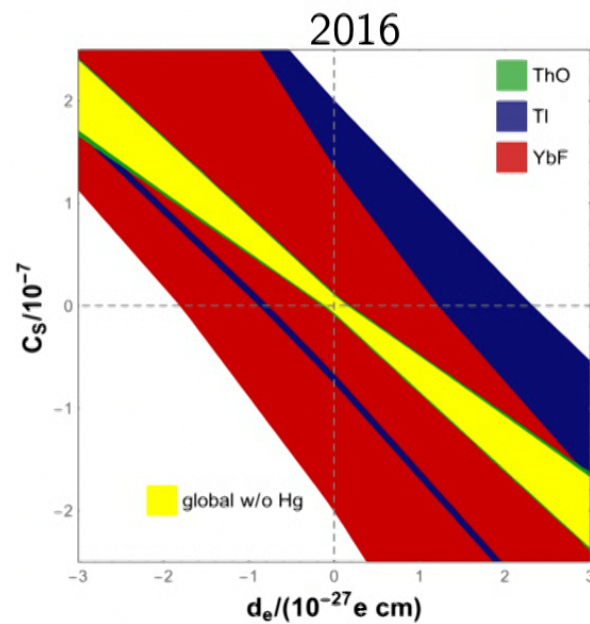


[Results entering: Nayak/Chaudhuri'07,'08,'09; Dzuba et al.'11, Meyer/Bohn'08, Skripnikov et al.'13, Fleig/Nayak'14; Averages: MJ'13, MJ/Pich'14]

## Model-independent extraction of $d_e$ and $C_S$

In principle: two unknowns, three measurements (TI, YbF, ThO)

➔ Extract  $d_e$ ,  $C_S$  model-independently [Dzuba et al.'11, MJ'13]



**Problem:** Aligned constraints

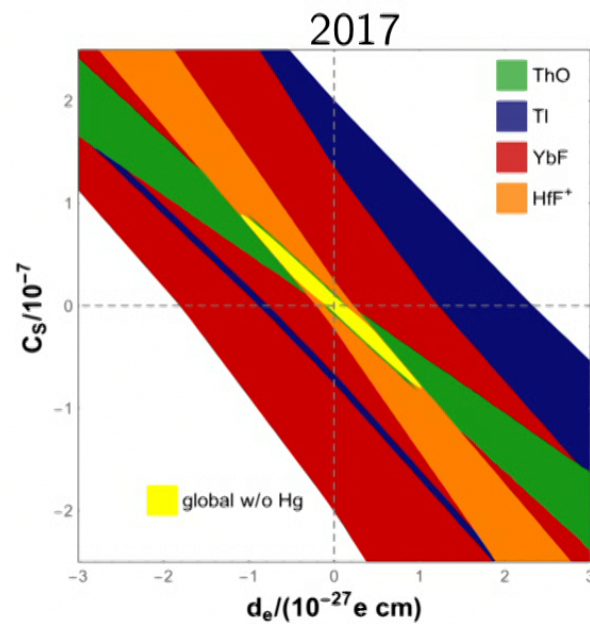
➔ weak limits



## Model-independent extraction of $d_e$ and $C_S$

In principle: two unknowns, three measurements (TI, YbF, ThO)

➔ Extract  $d_e$ ,  $C_S$  model-independently [Dzuba et al.'11, MJ'13]



**Problem:** Aligned constraints

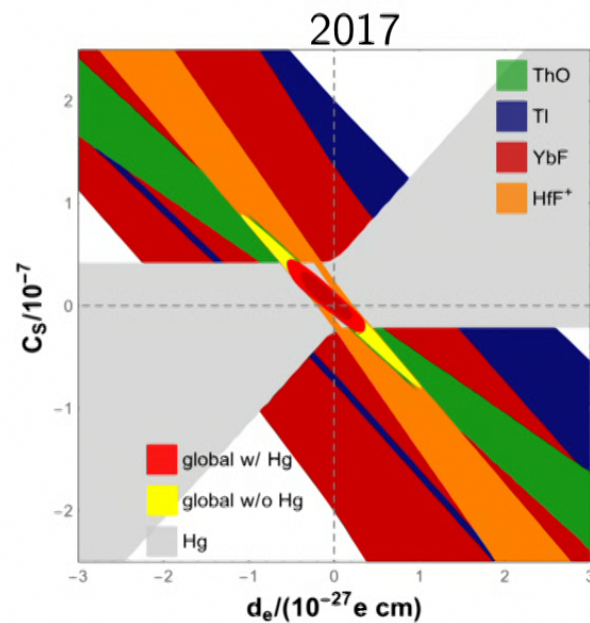
➔ weak limits

**Partial resolution:** HfF<sup>+</sup> result

## Model-independent extraction of $d_e$ and $C_S$

In principle: two unknowns, three measurements (Tl, YbF, ThO)

➔ Extract  $d_e$ ,  $C_S$  model-independently [Dzuba et al.'11, MJ'13]



**Problem:** Aligned constraints

➔ weak limits

**Partial resolution:** HfF<sup>+</sup> result

**Mercury bound  $\sim$  orthogonal!**

Assumption:  $C_S$ ,  $d_e$  saturate  $d_{\text{Hg}}$

➔ Conservative

[Fleig, MJ('17) preliminary]

$$d_e \leq 3.9 \times 10^{-28} e \text{ cm}$$

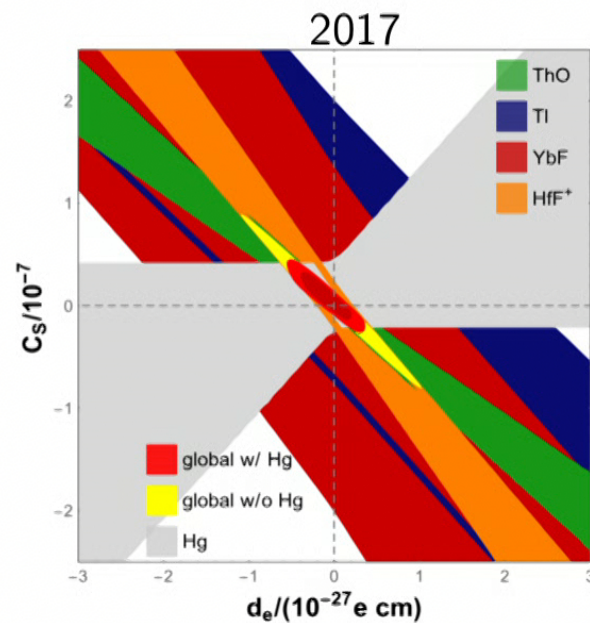
$$C_S \leq 3.2 \times 10^{-8}$$

Yields model-independent limit on **every** paramagnetic system!

## Model-independent extraction of $d_e$ and $C_S$

In principle: two unknowns, three measurements (Tl, YbF, ThO)

➔ Extract  $d_e$ ,  $C_S$  model-independently [Dzuba et al.'11, MJ'13]



**Problem:** Aligned constraints

➔ weak limits

**Partial resolution:** HfF<sup>+</sup> result

**Mercury bound ~ orthogonal!**

Assumption:  $C_S$ ,  $d_e$  saturate  $d_{\text{Hg}}$

➔ Conservative

[Fleig, MJ('17) preliminary]

$$d_e \leq 3.9 \times 10^{-28} e \text{ cm}$$

$$C_S \leq 3.2 \times 10^{-8}$$

Yields model-independent limit on **every** paramagnetic system!

Future measurements aim at precision beyond present constraints!

➔ Help to resolve the alignment problem

➔ Requires precision measurements of low-Z and high-Z elements

## EDMs of diamagnetic systems and nucleons

Situation more complicated than for paramagnetic systems:

- Potential SM contribution:  $\bar{\theta}$  ( $\rightarrow$  strong CP puzzle)
- Contributions from  $\bar{\theta}$ ,  $d_q$ ,  $\tilde{d}_q$ ,  $w$ ,  $C_{S,P,T}$ ,  $C_{qq}$ 
  - $\rightarrow$  Interpretation usually model-dependent  
(for model-independent prospects: [Chupp/Ramsey-Musolf'14] )

Complementary measurements, different sources possible/likely

- $|d_{Hg}| \leq 7.4 \times 10^{-30} e \text{ cm}$  [Graner et al. '16] , very constraining  
Problem: QCD and nuclear theory uncertainties ( $\times 100\%$ )  
 $\rightarrow$  No conservative constraint on CEDMs left! [MJ/Pich'13]
- $|d_n| \leq 3.6 \times 10^{-26} e \text{ cm}$  [Pendlebury'15]  
Theory in better shape, still  $\mathcal{O}(100\%)$  uncertainties  
[Pospelov/Ritz'01, Hisano et al'12, Demir et al'03,'04, de Vries et al'11]

Progress in theory necessary to fully exploit these measurements  
Unique: orders-of-magnitude improvement **w/o new measurement!**

## The role of Mercury in determining the electron EDM

Mercury is a diamagnetic system, many contributions

- Why is it shown in the paramagnetic global fit? [MJ'13]
  - Shielding of  $C_S$  and  $d_e$  effective (even vanishing at LO)
    - Schiff moment contribution expected to be dominant
    - $d_e, C_S$  only a fraction of the total EDM
  - Assuming  $d_e, C_S$  to saturate the exp. limit is **conservative**

**New calculation of the  $C_S$  coefficient** [Fleig/MJ('17),preliminary]

LO contribution vanishes

- Triple perturbative expansion necessary:
  1. External electric field (here: included in basis set)
  2. Hyperfine splitting
  3.  $d_e/C_S$

$$\alpha_{C_S} = -2.8(6) \times 10^{-22} e \text{ cm} \quad (\text{preliminary!})$$

$\alpha_{d_e}$  w.i.p., so far old calculation [Martensson-Pendrill/Oster'85] + conservative error estimate

## The importance of multiple measurements

Only **pattern** of CPV observables allows for model-differentiation!

➔ There is no single “best” measurement!

Paramagnetic systems:

- 1 significant measurement NP
- 2 determine ideally  $d_e$  and  $C_S$
- More for consistency (unless MQM is relevant)

Diamagnetic systems, nucleons, light nuclei:

- 1 significant measurement:  $\bar{\theta}$  possible explanation
- 2 should tell  $\bar{\theta}$  from other sources
- Many more to identify model-independently CPV structure

➔ We need as many measurement as possible!

➔ Ideally very different systems

➔ Try to find P-,T-odd measurements besides EDMs

## EDMs in NP Models

EDM constraints forbid generic CPV contributions up to two loops

↳ huge scales or highly specific structure!

- hardly testable elsewhere
- simple power-counting insufficient (UV sensitivity)
- ↳ Model-independent analyses difficult

EDMs unique, both blessing and curse

- some model-independent relations exist, e.g. to  $\beta$  decay [Khriplovich'91, see also e.g. Dekens/Vos'15]
- strong (model-dependent) constraints of related observables
- ↳ Consider models or subsets of model-independent framework

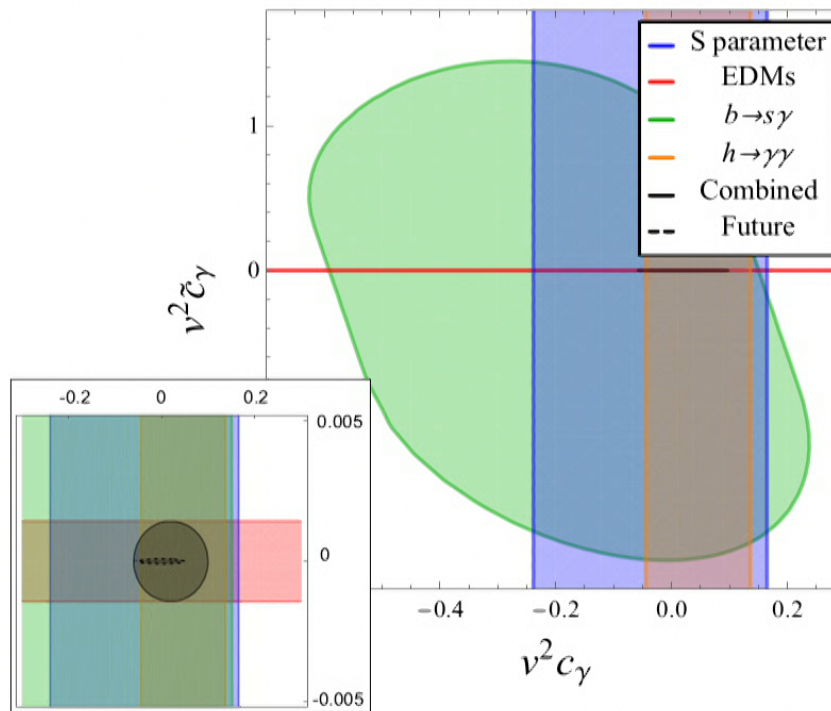


## Connecting high- and low-energy observables with EFTs

Example from [Cirigliano et al.'16] :

Consider chirality-flipping SMEFT operators with top and Higgs

➔ Affect EDMs, Higgs observables, flavour, ...



[Cirigliano+'16]

- CPV dominated by EDMs
- collider observables complementary
- ➔ significant progress expected for both



## Why 2HDM?

Model-independent NP analysis: Too many parameters in general

EW symmetry breaking mechanism still not completely fixed:

- 1HDM minimal and elegant, but “unlikely” (SUSY, GUTs, ...)
- 2HDM “next-to-minimal”:
  - $\rho$ -parameter “implies” doublets
  - low-energy limit of more complete NP models
    - ➔ Model-independent element
  - simple structure, but interesting phenomenology
  - important effects in flavour observables
- Plethora of 2HDMs:
  - ➔ differ in their suppression mechanism for FCNCs

Could explain tensions in the flavour sector (e.g.  $B \rightarrow D^{(*)} \tau \nu$ )



Not an attempt at a complete theory!

## Framework for 2HDM contributions

The CPV interactions of the 2nd doublet can generate EDMs

**General** parametrization for  $H^\pm$  Yukawas,  $\varsigma_i$  **complex matrices**:

$$\mathcal{L}_Y^{H^\pm} = -\frac{\sqrt{2}}{v} H^\pm \left\{ \bar{u} \left[ V \varsigma_d M_d \mathcal{P}_R - \varsigma_u M_u^\dagger V \mathcal{P}_L \right] d + \bar{\nu} \varsigma_l M_l \mathcal{P}_R \right\} + \text{h.c.}$$

- Induce couplings like  $W$ -exchange, just with a charged Higgs ( $M_{H^\pm} \gtrsim m_t$ )
- Easily matched on your favourite model
  - ➡  $M_i$  only choice of normalization
- $\varsigma_i \rightarrow$  **numbers**: Aligned 2HDM [Pich/Tuzon'09, MJ/Pich/Tuzon'10]
  - ➡ Comparisons with flavour data in this model

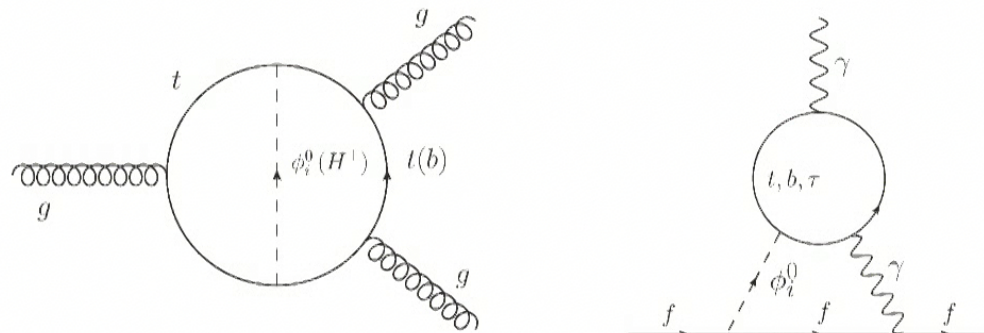
Neutral Higgs exchanges: couplings  $y_i^0(\varsigma_i, V)$

- ➡ Additional CPV contributions from the potential
- ➡ Analysis depends on many unknown parameters

## EDMs in 2HDMs

From necessary **flavour suppression** for a viable model:

- One-loop (C)EDMs: controlled (not tiny) [e.g. Buras et al. '10]
  - 4-quark operators small (no  $\tan^3\beta$ -enhancement)
  - Two-loop graphs dominant
    - Weinberg diagram important for neutron EDM
    - Barr-Zee(-like) diagrams dominate other EDMs
- [Weinberg '89, Dicus '90, Barr/Zee '90, Gunion/Wyler '90, ...]



Paramagnetic systems: tree-level can be relevant ( $C_S \times Z^3$ )  
(light-quark mass  $\times$  tree) vs. (top mass  $\times$  two-loop)

## Neutral Higgs contributions in general 2HDMs [MJ/Pich'13]

Contributions typically involve the following sum:  
( $f, f'$ : fermions,  $F(f)$ : family of the fermion)

$$\sum_i \operatorname{Re} \left( y_f^{\varphi_i^0} \right) \operatorname{Im} \left( y_{f'}^{\varphi_i^0} \right) = \pm \operatorname{Im} \left[ (\zeta_{F(f)}^*)_{ff} (\zeta_{F(f')})_{f'f'} \right]$$

- R.h.s. independent of the Higgs potential
- Vanishes for equal fermions (universality: equal family)
- Modified by mass-dependent weight factors. . .
- ➡ but holds for degenerate masses **and** decoupling limit

CPV in the potential tends to have smaller impact

➡ Approximation for phenomenological analysis:

$$\sum_i f(M_{\varphi_i^0}) \operatorname{Re} \left( y_f^{\varphi_i^0} \right) \operatorname{Im} \left( y_{f'}^{\varphi_i^0} \right) \rightarrow \pm f(\overline{M}_\varphi) \operatorname{Im} \left[ (\zeta_{F(f)}^*)_{ff} (\zeta_{F(f')})_{f'f'} \right] .$$

## Conclusions

- EDMs unique tests of NP models
- Model-independent constraints on NP parameters difficult
  - ➔ Need (at least) as many experiments as (eff.) parameters
- Quantitative results require close look at theory uncertainties
  - ➔ Use conservative limits, allowing for cancellations
  - ➔ For e.g.  $d_n, d_{\text{Hg}}$  bottleneck! **Chance for nuclear theory**
- Robust, model-independent limit on electron EDM ( $C_5$  not model-independently negligible):

$$|d_e| \leq 3.9 \times 10^{-28} e \text{ cm} \quad (95\% \text{ CL})$$

- General discussion of 2HDM constraints possible
  - ➔  $\varsigma_i$  key parameters, CPV from potential suppressed
- Interplay of EDMs with flavour physics
  - ➔ Flavour suppression just sufficient
  - ➔ CPV in other observables strongly restricted
- Plethora of new results to come
  - ➔ Might turn limits into determinations!