

Title: PSI 17/18 - Statistical Mechanics - Lecture 7

Date: Oct 18, 2017 10:45 AM

URL: <http://pirsa.org/17100040>

Abstract:

FROM ISING TO FIELD THEORY

$$Z = \sum_{\{\sigma_i\}} e^{\frac{1}{2T} \sum_{ij} J_{ij} \sigma_i \sigma_j} = \int D\varphi e^{-S[\varphi]}$$

$$S[\varphi] = \frac{1}{2T} \sum_{ij} J_{ij} \varphi_i \varphi_j - \sum_i \log(2 \cosh(\varphi_i/T)) \approx \text{EXPANSION NEAR CRIT. POINT}$$

$$\int d^d x \left[\frac{1}{2} (\nabla\varphi)^2 + \frac{r}{2} \varphi^2 + \frac{u}{4!} \varphi^4 \right] = S[\varphi]$$

LANDAU-GINZBURG FUNCTIONAL

FROM ISING TO FIELD THEORY

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$$S[\varphi] = \frac{1}{2T} \varphi_i J_{ij}^{-1} \varphi_j - \sum_i \log(2 \cosh(\varphi_i/T))$$

EXPANSION
NEAR CRIT. POINT

TRANSITION
AT CRIT. POINT

$$\int d^d x \left[\frac{1}{2} |\nabla \varphi|^2 + \frac{\mu}{2} \varphi^2 + \frac{\lambda}{4!} \varphi^4 \right] = S[\varphi]$$

LANDAU-GINSBURG FUNCTIONAL

$$\mu > 0, \lambda > 0$$

FROM ISING TO FIELD THEORY

$$Z = \sum_{\{\sigma_i\}} e^{\frac{1}{2T} \sum_{ij} J_{ij} \sigma_i \sigma_j} \approx \int D\varphi e^{-S[\varphi]}$$

$$S[\varphi] = \frac{1}{2T} \varphi_i J_{ij}^{-1} \varphi_j - \sum_i \log(2 \cosh(\varphi_i/T)) \approx$$

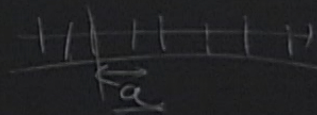
EXPANSION
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$$\int d^d x$$



$\beta \rightarrow \infty$, g .. SMALL

$u > 0$



$S[\varphi]$

(φ_i/T)

EXPANSION
NEAR CRIT. POINT

D-SPACE DIMENSIONS

$$\int d^D x \left[\frac{1}{2} |\nabla \varphi|^2 + \frac{r}{2} \varphi^2 + \frac{u}{4!} \varphi^4 \right] = S[\varphi]$$

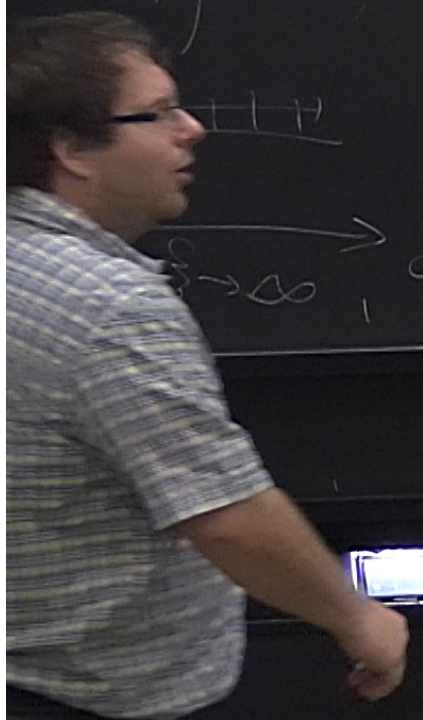
LANDAU-GINSBURG FUNCTIONAL

$$u > 0, \quad r \propto t = \frac{T - T_c}{T_c}$$

|||||

$\xi \rightarrow \infty$

g - SMALL



$S[\varphi]$

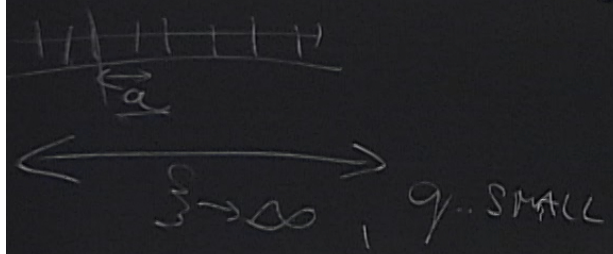
(φ_i/T)

EXPANSION
NEAR CRIT. POINT

D-SPACE DIMENSIONS

$$\int d^D x \left[\frac{1}{2} |\nabla \varphi|^2 + \frac{r}{2} \varphi^2 + \frac{u}{4!} \varphi^4 \right] = S[\varphi]$$

LANDAU-GINSBURG FUNCTIONAL



$$u > 0, r \propto t = \frac{T - T_c}{T_c}$$

HIGHER POWERS OF φ OR ∇ ARE IRRELEVANT
AND SO THIS IS EVERYTHING YOU CAN
WRITE DOWN,

DIMENSIONAL ANALYSIS "UNITS OF ENERGY"

$$[k] = \dim(k) = +1$$

$$[x] = -1$$

∞ , g -SMALL

$$[S] = 0 = \underbrace{[dx^d]}_{-d} + \underbrace{[\nabla^2]}_{+2} + 2 \times [\varphi]$$

$$\Rightarrow [\varphi] = \frac{d-2}{2}$$

$$[n] = +2$$

$$[dx^d] + [n] + 4 \times [\varphi] = 0$$
$$-d + [n] + 4 \times \frac{d-2}{2} = 0$$

$$[n] = 4 - d = \mathcal{E}$$

DEF: COUPLING g IS

RELEVANT	$[g] > 0$
MARGINAL	$[g] = 0$
IRRELEVANT	$[g] < 0$

$d=4$

$[n] = 2$
 $[m] = 0$

$$[m] = 4 - d = \varepsilon$$

$$d=4$$

$$[n] = 2$$

$$[m] = 0$$

$$+ 15 \varphi^6$$

$$[N] = 6 - 2d = \underline{-2}$$

L-G AS TRUNCATED HERE IS ONLY GOOD IN d=4

$$[m] = 4 - d = \varepsilon$$

$$d=4$$

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L-G AS TRUNCATED HERE IS ONLY GOOD IN d=4

• HOW TO CALCULATE ε ?

$$[m] = 4 - d = \epsilon$$

$$d=4$$

$$[n] = 2$$

$$[m] = 0$$

$$+15 \varphi^6$$

$$[N] = 6 - 2d = -2$$

L-G AS TRUNCATED HERE IS ONLY GOOD IN $d=4$

- HOW TO CALCULATE Z ?
 - PERT. THEORY IN μ
 - P.T. FLUCTUATIONS $\delta\phi$
- RIGHT WAY: PERT. TH. IN SMALL ϵ .

d) WILSONIAN (MOMENTUM SPACE) RENORMALIZATION

PART I. MAIN IDEA

"FORMAL WAY OF COARSE GRAINING"

WE WANT:

$$Z = \int D\psi e^{-S[\psi]}$$

IN SMALL STEPS

d) WILSONIAN (MOMENTUM SPACE) RENORMALIZATION

PART I. MAIN IDEA

"FORMAL WAY OF COARSE GRAINING"

WE WANT:

$$Z = \int D\psi e^{-S[\psi]}$$

IN SMALL STEPS AVOIDING DIVERGENCIES
IN SMALL l

$$\psi = \psi_{<} + \psi_{>}$$

$$\psi_{<}(k) = \begin{cases} \psi(k) \\ 0 \end{cases}$$

$$0 < k < N/2$$

$$k > N/2$$

$$\psi_{>}(k) = \begin{cases} 0 \\ \psi(k) \end{cases}$$

$$0 < k < N/2$$

$$k > N/2$$

TRICK: INTEGRATE FAST ONLY,

• $S = S_0 + S_{int}$

$$S_0 = \frac{1}{2} \int d^d x (|\partial\psi|^2 + n\psi^2) = \frac{1}{2} \int \frac{d^d k}{(2\pi)^d} (k^2 + n) |\psi(k)|^2$$

b $= S_0[\psi_<] + S_0[\psi_>]$

NO MIXING OF FAST & SLOW

TRICK: INTEGRATE FAST ONLY,

$$S = S_0 + S_{int}$$

$$S_0 = \frac{1}{2} \int d^d x (|\nabla \psi|^2 + \mu \psi^2) = \frac{1}{2} \int \frac{d^d k}{(2\pi)^d} (k^2 + \mu) |\psi(k)|^2$$

$$= S_0[\psi_<] + S_0[\psi_>]$$

NO REFAST or K-SLOW MIXING.

TRICK: INTEGRATE FAST ONLY,

$$S = S_0 + S_{\text{INT}}$$

$$S_0 = \frac{1}{2} \int d^d x \left(|\nabla \varphi|^2 + \mu \varphi^2 \right) = \frac{1}{2} \int \frac{d^d k}{(2\pi)^d} \left(k^2 + \mu \right) |\varphi(k)|^2$$

$$= S_0[\varphi_<] + S_0[\varphi_>]$$

NO FAST or SLOW MIXING.

$$S_{\text{INT}} = \frac{\mu}{4!} \int d^d x \varphi^4 = \frac{\mu}{4!} \int \frac{d^d k_1}{(2\pi)^d} \dots \frac{d^d k_4}{(2\pi)^d} \varphi(k_1) \dots \varphi(k_4) (2\pi)^d \delta(k_{1+} + k_4)$$

$$= S_{\text{INT}}[\varphi_<, \varphi_>]$$

WANT TO INTEGRATE OVER FAST MODES

$$Z = \int D\psi \int D\psi_L e^{-S[\psi, \psi_L]} = \int D\psi_L e^{-\tilde{S}[\psi_L]}$$

WANT TO INTEGRATE OVER FAST MODES

$$Z = \int D\psi_D D\psi_K e^{-S[\psi_D, \psi_K]} = \int D\psi_K e^{-\tilde{S}[\psi_K]}$$

$\tilde{S}[\psi_K]$... HAS TO HAVE THE SAME FORM WE STARTED WITH.

$$\tilde{S}[\psi_K] = \frac{1}{2} \int_0^{N_b} \frac{d^d k}{(2\pi)^d} |\psi_K(k)|^2 ($$

WANT TO INTEGRATE OVER FAST MODES

$$Z = \int D\psi_D D\psi_K e^{-S[\psi_D, \psi_K]} = \int D\psi_K e^{-\tilde{S}[\psi_K]}$$

$\tilde{S}[\psi_K]$... HAS TO HAVE THE SAME FORM WE STARTED WITH.

$$\tilde{S}[\psi_K] = \frac{1}{Z} \int_0^{N_b} \frac{d^d k}{(2\pi)^d} |\psi_K(k)|^2 (\tilde{\mu}_2 k^2 + \tilde{\mu}_0)$$

2
- 3 [K]

WE STARTED WITH, (CAN GENERATE MORE TERM BUT THESE ARE IRRELEVANT)

$$\approx \frac{\Lambda^4}{4!} \int_0^{\Lambda/b} \frac{d^d k_1 \dots d^d k_4}{(2\pi)^{4d}} \varphi_k(k_1) \dots \varphi_k(k_4) (2\pi)^d \delta(k_1 + k_2 + k_3 + k_4)$$

• TO BRING THIS TO EXACTLY THE SAME FORM WE
STARTED WITH WE SCALE:

$$\tilde{r} = kb, \quad \tilde{\psi}(\tilde{r}) = \frac{\psi(r)}{r}$$

WAVE FUNCTION

$[X] = -1$

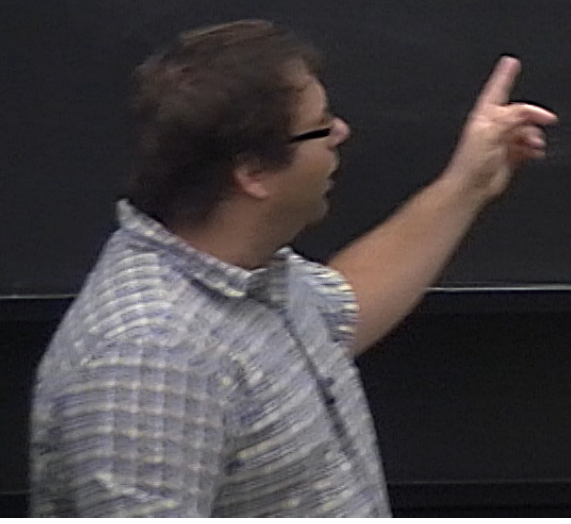
$L(\lambda) = \dots$

• TO BRING THIS TO EXACTLY THE SAME FORM WE STARTED WITH WE SCALE:

$\tilde{l} = lb, \quad \tilde{\psi}(l) = \frac{\psi(l)}{Z}$

LEAVE FUNCTION RENORMALIZATION

DIMENSIONS



$[X] = -1$

$[L] = -1$

• TO BRING THIS TO EXACTLY THE SAME FORM WE STARTED WITH WE SCALE:

$\tilde{k} = kb, \quad \tilde{\varphi}(\tilde{k}) = \frac{\varphi(k)}{Z}$ ← WAVE FUNCTION RENORMALIZATION

$$S[\varphi] = \frac{1}{2} \int_0^\Lambda \frac{d^d k}{(2\pi)^d} b^{-d} \left(\tilde{\Lambda}^2 b^{-2} k^2 + \tilde{\Lambda}_0 \right) |\varphi(k)|^2 Z^2$$

DIMENSIONS

$$L(x) = \dots$$

$$[x] = -1$$

$$L(x) = \dots$$

• TO BRING THIS TO EXACTLY THE SAME FORM WE STARTED WITH WE SCALE:

$$\tilde{k} = kb, \quad \tilde{\varphi}(\tilde{k}) = \frac{\varphi(k)}{Z}$$

← WAVE FUNCTION RENORMALIZATION

$$S[\varphi] = \frac{1}{2} \int_0^\Lambda \frac{d^d k}{(2\pi)^d} \left(b^{-d} Z^{-2} \Lambda^2 b^{-2} k^2 + Z \Lambda_0 b^d \right) |\varphi(k)|^2$$

DIMENSIONS



$$[X] = -1$$

$$L(\psi) = \dots$$

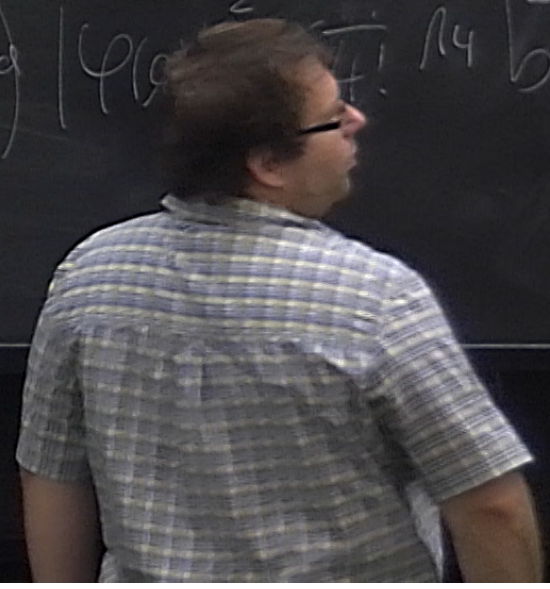
• TO BRING THIS TO EXACTLY THE SAME FORM WE STARTED WITH WE SCALE:

$$\tilde{k} = kb, \quad \tilde{\psi}(\tilde{k}) = \frac{\psi(k)}{Z}$$

WAVE FUNCTION RENORMALIZATION

$$S[\psi] = \frac{1}{2} \int_0^\Lambda \frac{d^d k}{(2\pi)^d} \left(b^{-d} Z^{-2} \tilde{\Lambda}^2 b^{-2} k^2 + Z \tilde{\Lambda}^d \right) |\psi(k)|^2$$

DIMENSIONS



STARTED WITH, (CAN GENERATE MORE TERM BUT THESE ARE IRRELEVANT)

$$+ \frac{\Lambda^4}{4!} \int_0^{\Lambda/b} \frac{d^d k_1 \dots d^d k_4}{(2\pi)^{4d}} \varphi(k_1) \dots \varphi(k_4) (2\pi)^d \delta(k_1 + k_2 + k_3 + k_4)$$
$$\delta(f(x)) = \sum \frac{\delta(x-x_i)}{|f'(x_i)|}$$

$$1 = b^{-d-2} z^2 \tilde{\Lambda}_2$$

$$\tilde{\Lambda}_2 = b^{\gamma}$$

$$1 = b^{-d-2} z^2 \tilde{\Lambda}_2$$

$$= b^{-d-2+\eta} z^2$$

$$\tilde{\Lambda}_2 = b^\eta$$

→ ANOMALOUS
DIMENSION

$$\Rightarrow z = b^{\frac{d+2-\eta}{2}}$$

$$\varphi(k) = b^{\frac{-d+2-\eta}{2}} \varphi_c(k)$$

$$\tilde{\Lambda}_2 = z^2$$

b^η → ANOMALOUS
DIMENSION

$$\tilde{n} = z^2 n_0 b^{-d} = b^{d+2-\eta-d} n_0 = b^{2-\eta} n_0$$

$$b^{\frac{d+2-\eta}{2}}$$

$$\tilde{n} = z^{2\tilde{n}} \Lambda_0 b^{-d} = b^{d+2\tilde{n}-d}$$

$$\Lambda_0 = \left[b^{2-\gamma} \tilde{n} \Lambda_0 = \tilde{n} \right]$$

$$\tilde{u} = b^{4-2\gamma-d} \tilde{n} \Lambda_4$$

$$\tilde{\mu} = z^2 \tilde{\Lambda}_0 b^{-d} = b^{d+2-\eta-d} \tilde{\Lambda}_0 = \boxed{b^{2-\eta} \tilde{\Lambda}_0 = \tilde{\mu}}$$

$$\tilde{\mu} = b^{4-2\eta-d} \tilde{\Lambda}_4$$

HAVE NOW

$$S[\varphi] = \int d^d x \left[\frac{1}{2} (\nabla \varphi)^2 + \frac{\tilde{\mu}}{2} \varphi^2 + \frac{\tilde{\mu}}{4!} \varphi^4 \right]$$

• EFFECTIVELY,

$$(n, m) \rightarrow (\tilde{n}, \tilde{m})$$

WE DO THIS REPEATEDLY

• EFFECTIVELY,

$$(n, m) \rightarrow (\tilde{n}, \tilde{m})$$

WE DO THIS REPEATEDLY

WE MAY FIND A FIXED POINT (n^*, m^*)

AT EACH STEP THE CORRELATION LENGTH

$$\xi = \xi(n, m) \rightarrow \xi(\tilde{n}, \tilde{m})$$

PART 2 EXPLICIT CALCULATION

$$Z = \int D\psi D\bar{\psi} e^{-S} = \int D\psi e^{-\tilde{S}[\psi]}$$

$$e^{-S[\psi]} =$$

$$S = S$$

$$\begin{aligned}
 D\psi_k e^{-S} &= \int D\psi_k e^{-\tilde{S}[\psi_k]} \\
 e^{-S_0[\psi_k]} \int D\psi_b e^{-S_0[\psi_b] - S_{int}[\psi_k, \psi_b]} &\times \int D\psi_b e^{-S_0[\psi_b]}
 \end{aligned}$$

$$S = S_0 + S_{int} = S_0(\psi_k) + S_0(\psi_b)$$

$$\tilde{S}(k) = S_0(k) - \log \left\langle e^{-Sint} \right\rangle_{\text{FAST}} - \log Z_0$$

↓
↓
 CONSTANT

CUMULANT EXPANSION

Ω RANDOM VARIABLE

$$\left\langle e^{\Omega} \right\rangle = \left\langle \Omega \right\rangle + \frac{1}{2} \left[\left\langle \Omega^2 \right\rangle - \left\langle \Omega \right\rangle^2 \right] + \dots$$

$$\tilde{S}(k) = S_0(k) - \log \left\langle e^{-S_{int}} \right\rangle_{\text{FAST}} - \log Z_0 = S_0(k) + \langle S_{int} \rangle$$

CONSTANT

CUMULANT EXPANSION

R RANDOM VARIABLE

$$\left\langle e^{R} \right\rangle = e^{\langle R \rangle + \frac{1}{2} [\langle R^2 \rangle - \langle R \rangle^2]} + \dots$$

