

Title: General Relativity for Cosmology - Lecture 8

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Abstract:

GR for Cosmology, Achim Kempf, Fall 2017, Lecture 9

Recall: So far, we have 2 ways to capture shape:

- Specified  $g$   $\Rightarrow$  specified distances in  $M$   
 $\Rightarrow$  implicitly specified "shape" of  $M$   
(Notice (for essay): See also my new paper 1510.02725)

Then, new:

- Specified  $\nabla$   $\Rightarrow$  specified parallel transport in  $M$   
 $\Rightarrow$  implicitly specified "shape" of  $M$

Question:

How does  $\nabla$  determine shape?



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Note Title

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Then, new:

- Specified  $\nabla \Rightarrow$  specified parallel transport in  $M$   
 $\Rightarrow$  implicitly specified "shape" of  $M$

Question:

How does  $\nabla$  determine "shape"? Through:

Recall:

$$\bar{\Gamma}^r_{ab} = \frac{\partial \bar{x}^r}{\partial x^k} \frac{\partial x^i}{\partial \bar{x}^a} \frac{\partial x^j}{\partial \bar{x}^b} \Gamma^k_{ij} + \frac{\partial \bar{x}^r}{\partial x^k} \frac{\partial^2 x^k}{\partial \bar{x}^a \partial \bar{x}^b}$$

Notice:

The antisymmetric part of  $\Gamma$  transforms tensorially!

$$\left. \begin{aligned} \Gamma^k_{(sym)ij} &:= \frac{1}{2} (\Gamma^k_{ij} + \Gamma^k_{ji}) \\ \Gamma^k_{(asym)ij} &:= \frac{1}{2} (\Gamma^k_{ij} - \Gamma^k_{ji}) \end{aligned} \right\} \Gamma^k_{ij} = \Gamma^k_{(sym)ij} + \Gamma^k_{(asym)ij}$$

$$\Rightarrow \bar{\Gamma}^r_{(asym)ab} = \frac{\partial \bar{x}^r}{\partial x^k} \frac{\partial x^i}{\partial \bar{x}^a} \frac{\partial x^j}{\partial \bar{x}^b} \Gamma^k_{(asym)ij} !$$

Definition:  $J^k_{ij} := 2 \Gamma^k_{(asym)ij}$  is the "Torsion tensor"



$$\Gamma^k{}_{ab} = \frac{\partial x^k}{\partial \bar{x}^a} \frac{\partial \bar{x}^b}{\partial x^c} \Gamma^c{}_{ij}$$

Notice: The antisymmetric part of  $\Gamma$  transforms tensorially!

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$$\Rightarrow \bar{\Gamma}_{(asym)ab}^k = \frac{\partial \bar{x}^k}{\partial x^c} \frac{\partial x^i}{\partial \bar{x}^a} \frac{\partial x^j}{\partial \bar{x}^b} \Gamma_{(asym)ij}^c !$$

Definition:  $J^k{}_{ij} := 2 \Gamma_{(asym)ij}^k$  is the "Torsion tensor"

(Notice: Since  $\Gamma$  is not a tensor, but  $\Gamma_{asym}$  is,  $\Gamma_{sym}$  is not a tensor)



In General Relativity: one assumes torsionless  $\nabla$ , i.e.:  $T = 0$ .

Idea: "(Extended) equivalence principle:"

Christoffel  $\Gamma$  will express gravitational and pseudo forces.  
Therefore, we require that around each  $p \in M$  there exists a chart so that  $\Gamma(p) = 0$  (i.e. no such forces in freefall).

This rules out the existence of torsion:

Why? The torsion is a tensor.

$\Rightarrow$  It transforms linearly with invertible jacobian matrices

$$\bar{T}^i_{jk}(\bar{x}) = \frac{\partial \bar{x}^i}{\partial x^a} \frac{\partial x^b}{\partial \bar{x}^j} \frac{\partial x^c}{\partial \bar{x}^k} T^a_{bc}(x)$$

$$\bar{J}_{jk}^i(\bar{x}) = \frac{\partial \bar{x}^i}{\partial x^a} \frac{\partial x^b}{\partial \bar{x}^j} \frac{\partial x^c}{\partial \bar{x}^k} J_{bc}^a(x)$$

$\Rightarrow$  If  $J_{ij}$  vanishes in one cds, it vanishes in all cds.

Proposition:

Vice versa, if  $J_{jk}^i(x) = 0 \forall x \in M$ ,  
then there is for every  $p \in M$  a chart with  $\Gamma_{jk}^i(p) = 0$ .

Recall:

$\xi$  is autoparallel to a path  $\gamma: t \rightarrow x(t)$  if

$\xi \in T'(M)$

$$\nabla_{\dot{\gamma}} \xi = 0$$

Meaning:  $\xi$  is parallel transported along the path  $\gamma$  in  $M$ .

Explicitly:  $\frac{d\xi^k}{dt} + \Gamma_{ij}^k \dot{x}^i \xi^j = 0$



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Explicitly: 
$$\frac{d\xi^k}{dt} + \Gamma^k_{ij} \frac{dx^i}{dt} \xi^j = 0$$

Geodesics:

A curve  $\gamma: t \rightarrow x(t)$  is called a **geodesic** if  $\dot{\gamma}$  is autoparallel along  $\gamma$ , i.e. if

$$\nabla_{\dot{\gamma}} \dot{\gamma} = 0$$

Meaning: The path  $\gamma$  in  $M$  is such that the path's tangent vectors are parallel translates of each other.



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Meaning: The path  $\gamma$  in  $M$  is such that the path's tangent vectors are parallel translates of each other.

□  $\Rightarrow$  In charts, geodesics  $x^r(t)$  obey:

$$\frac{d^2 x^k}{dt^2} + \Gamma^k(x)^{ij} \frac{dx^i}{dt} \frac{dx^j}{dt} = 0 \quad (*)$$

□ Theory of ordinary differential equations:

$\Rightarrow$  Given  $p = \gamma(0)$ , each initial condition  $\xi = \dot{\gamma}(0)$  belongs to a unique geodesic  $\gamma_\xi$  of non zero length.

subscript indicates initial condition vector

□ Notice: If  $\gamma_\xi(t)$  solves  $(*)$  then  $\gamma_\xi(\lambda t)$  also solves  $(*)$  and for  $\lambda \in \mathbb{R}$ :

$$\gamma_\xi(\lambda t) = \gamma_\xi(t)$$

(1)



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$$\gamma_{\lambda\xi}(t) = \gamma_\xi(\lambda t) \quad (G)$$

(Exercise: verify)



## "Exponential map":

□ Consider a fixed point  $p \in M$ .

The exponential map is defined through:

$$\exp_p: T_p(M) \rightarrow M \quad \left( \begin{array}{l} \text{really from a neighborhood} \\ \text{of } 0 \text{ in } T_p(M) \text{ to a neighborhood} \\ \text{of } p \text{ in } M. \end{array} \right)$$

$$\exp_p: \xi \rightarrow \exp_p(\xi) := \gamma_\xi(1)$$

where  $\gamma$  is the geodesic with  $\gamma_\xi(0) = p, \dot{\gamma}_\xi(0) = \xi$ .

□ Observe:

From (G) we obtain:

$$v(\lambda) = v(1) = \exp_p(\lambda \rho) \quad (E)$$

## "Geodesic" or "Riemann normal" coordinates:

□  $\exp_p$  is a diffeomorphism from a neighborhood of  $0 \in T_p(M) \simeq \mathbb{R}^n$  into a neighborhood of the point  $p \in M$ .  
 $\simeq$  isomorphic

$\Rightarrow \exp_p$  provides a chart around  $p$ :

□ Choose a basis, say  $e_1, e_2, \dots, e_n$  of  $T_p(M)$ , then:  
 $\xi = \xi^i e_i$

□ Through  $\exp_p$ , the  $\xi^i$  become the coordinates of points in a neighborhood of  $p \in M$ :

$$(\xi^1, \dots, \xi^n) \rightarrow \exp_p(\xi^i e_i) \in M$$



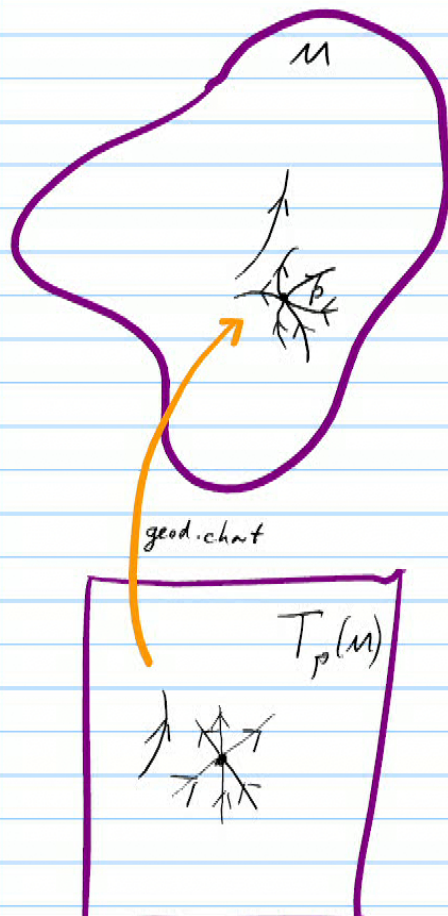
$\Rightarrow$  Geodesics,  $\gamma$ , through  $p$  are straight lines in a normal cds about  $p$ !

□ Recall (E):

$$\gamma_{\xi}(\lambda) = \exp_p(\lambda \xi)$$

for varying  $\lambda$  one moves along the geodesic in  $M$ .

for varying  $\lambda$  one moves on a straight line in the coordinate system of the  $\xi^i$ !



□ Thus: In geodesic cds, geodesics through  $p$  are straight lines of constant velocity  $\xi$ .

□ Does this mean  $\Gamma^k_{ij}(p) = 0$ ? **No!**



Geodesic eqn. at  $p$ :

$$\frac{d^2 x^k}{dt^2} + \Gamma^k_{ij}(p) \frac{dx^i}{dt} \frac{dx^j}{dt} = 0$$

Thus: 
$$\left( \Gamma^k_{sym\ ij}(p) + \Gamma^k_{asym\ ij}(p) \right) \frac{dx^i}{dt} \frac{dx^j}{dt} = 0$$

$\hookrightarrow = 0$  because of type (antisymmetric)<sub>ij</sub> (symmetric)<sup>ij</sup>

$$\Rightarrow \Gamma^k_{sym\ ij}(p) = 0 \text{ in geodesic cds.}$$

$\Rightarrow$  Indeed: If the torsion vanishes,  $\Gamma^k_{asym\ ij}(p) = \frac{1}{2} T^k_{ij}(p) = 0$   
then for each  $p \in M$  there exists a chart in which  
the entire gravity and pseudo force field vanishes at  $p$ :

Note:

Quantum fluctuations  
may induce torsion!

$$\Gamma^k_{ij}(p) = 0$$

Geodesic eqn. at p:

$$\frac{d^2 x^k}{dt^2} + \Gamma^k_{ij}(p) \frac{dx^i}{dt} \frac{dx^j}{dt} = 0$$

Thus:  $\left( \Gamma^k_{sym\ ij}(p) + \Gamma^k_{asym\ ij}(p) \right) \frac{dx^i}{dt} \frac{dx^j}{dt} = 0$

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Note:

Quantum fluctuations  
may induce torsion!  
So, let's nevertheless ask:

$$\Gamma^k_{ij}(p) = 0$$

What would torsion mean, geometrically?



# Abstract definition of Torsion:

□ Assume  $\xi_1$  and  $\xi_2$  are tangent vectors at  $p \in M$ :

Then, the Torsion map is defined as:

$$\mathcal{T} : T_p(M) \times T_p(M) \rightarrow T_p(M)$$

This will be the amount by which an infinitesimal parallelogram spanned by  $\xi_1$  and  $\xi_2$  does not close.

$$\mathcal{T} : \xi_1, \xi_2 \rightarrow \mathcal{T}(\xi_1, \xi_2) := \nabla_{\xi_1} \xi_2 - \nabla_{\xi_2} \xi_1 - [\xi_1, \xi_2]$$

for proof it's a tensor, see Strömmer

□ It is used to define the Torsion tensor,  $J$ ,

$$J \in T_p^1(M)$$

through:

feeding 1 covector & 2 vectors to a (1,2) tensor

$$J(\omega, \xi_1, \xi_2) := \underbrace{\langle \omega, \mathcal{T}(\xi_1, \xi_2) \rangle}_{\text{contraction yields a number}} \in \mathbb{R}$$

we could also write:  $= \omega(\mathcal{T}(\xi_1, \xi_2))$

contraction yields a number

Compare with prior definition:

□ Choose canonical bases  $w := dx^k$ ,  $\xi_1 := \frac{\partial}{\partial x^i}$ ,  $\xi_2 := \frac{\partial}{\partial x^j}$  :

$$\begin{aligned} \square \quad J^k_{ij} &:= dx^k \left( J \left( \frac{\partial}{\partial x^i}, \frac{\partial}{\partial x^j} \right) \right) \\ &= \langle dx^k, J \left( \frac{\partial}{\partial x^i}, \frac{\partial}{\partial x^j} \right) \rangle \quad (\text{more convenient notation}) \end{aligned}$$

$$= \langle dx^k, \nabla_{\frac{\partial}{\partial x^i}} \frac{\partial}{\partial x^j} - \nabla_{\frac{\partial}{\partial x^j}} \frac{\partial}{\partial x^i} - \underbrace{\left[ \frac{\partial}{\partial x^i}, \frac{\partial}{\partial x^j} \right]} \rangle$$

Recall:

$$\nabla_{\frac{\partial}{\partial x^i}} \frac{\partial}{\partial x^j} = \Gamma^k_{ij} \frac{\partial}{\partial x^k}$$

$$\left( \frac{\partial}{\partial x^i} \frac{\partial}{\partial x^j} - \frac{\partial}{\partial x^j} \frac{\partial}{\partial x^i} \right) f = 0 \quad \forall f$$

$$= \langle dx^k, \Gamma^r_{ij} \frac{\partial}{\partial x^r} - \Gamma^r_{ji} \frac{\partial}{\partial x^r} \rangle = \Gamma^r_{ij} \delta^k_r - \Gamma^r_{ji} \delta^k_r$$

$$\square \Rightarrow \boxed{J^k_{ij} = \Gamma^k_{ij} - \Gamma^k_{ji}}$$



Compare with prior definition:

□ Choose canonical bases  $\omega := dx^k$ ,  $\xi_1 := \frac{\partial}{\partial x^i}$ ,  $\xi_2 := \frac{\partial}{\partial x^j} :$

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Recall:

$$\nabla_{\frac{\partial}{\partial x^i}} \frac{\partial}{\partial x^j} = \Gamma^k_{ij} \frac{\partial}{\partial x^k}$$

$$\left( \frac{\partial}{\partial x^i} \frac{\partial}{\partial x^j} - \frac{\partial}{\partial x^j} \frac{\partial}{\partial x^i} \right) f = 0 \quad \forall f$$

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□  $\Rightarrow$

$$\boxed{J^k_{ij} = \Gamma^k_{ij} - \Gamma^k_{ji}}$$

Compare with prior definition:

□ Choose canonical bases  $\omega := dx^k$ ,  $\xi_1 := \frac{\partial}{\partial x^i}$ ,  $\xi_2 := \frac{\partial}{\partial x^j}$  :

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$$= \langle dx^k, \nabla_{\frac{\partial}{\partial x^i}} \frac{\partial}{\partial x^j} - \nabla_{\frac{\partial}{\partial x^j}} \frac{\partial}{\partial x^i} - \underbrace{\left[ \frac{\partial}{\partial x^i}, \frac{\partial}{\partial x^j} \right]} \rangle$$

Recall:

$$\nabla_{\frac{\partial}{\partial x^i}} \frac{\partial}{\partial x^j} = \Gamma^k_{ij} \frac{\partial}{\partial x^k}$$

$$\left( \frac{\partial}{\partial x^i} \frac{\partial}{\partial x^j} - \frac{\partial}{\partial x^j} \frac{\partial}{\partial x^i} \right) f = 0 \quad \forall f$$

$$= \langle dx^k, \Gamma^r_{ij} \frac{\partial}{\partial x^r} - \Gamma^r_{ji} \frac{\partial}{\partial x^r} \rangle = \Gamma^r_{ij} \delta^k_r - \Gamma^r_{ji} \delta^k_r$$

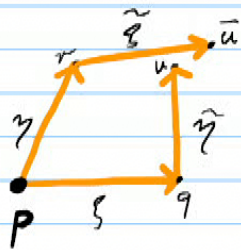
□  $\Rightarrow$

$$\boxed{J^k_{ij} = \Gamma^k_{ij} - \Gamma^k_{ji}}$$



# Geometric meaning of torsion? Parallelograms would not close!

Travel from  $p$  infinitesimally in  $\xi$  and then  $\eta$  direction, and compare with the reverse. (In flat space:  $x^r + \eta^r + \xi^r = x^r + \xi^r + \eta^r$ .)



$$\begin{aligned} \xi, \eta &\in T_p' \\ \tilde{\xi} &\in T_r' \\ \tilde{\eta} &\in T_q' \end{aligned}$$

Recall parallel transport:  $\nabla_j v = 0$

$$\frac{dv^k}{dt} + \Gamma_{ij}^k \frac{dx^i}{dt} v^j = 0$$

$\tilde{\xi}(r) = ?$

$\tilde{\xi}^k(x^i + \eta^i) \approx \xi^k(x^i) + \frac{d\xi^k}{dt}(x^i)$     Now use  $v := \xi, \frac{dx^i}{dt} = \eta^i$ :

$$= \xi^k(x^i) - \Gamma(x)^k_{ij} \eta^i \xi^j$$

$\Rightarrow$  (ds. of  $\bar{u}$ :  $x^a + \eta^a + \xi^a - \Gamma(x)^a_{ij} \eta^i \xi^j$ )

Analogously obtain: Cds. of  $u$ :  $x^a + \xi^a + \eta^a - \Gamma(x)^a_{ij} \xi^i \eta^j$

Torsion!



$\Rightarrow$  Cd. distance from  $u$  to  $\bar{u}$  is:  $(\Gamma(x)^a_{ji} - \Gamma(\bar{u})^a_{ij}) \eta^i \xi^j = T^a_{ji} \eta^i \xi^j$  ✓

Comment: We had:

$$\tilde{\xi}^k(x^i + \eta^i) \approx \xi^k(x^i) + \frac{d\xi^k}{dx^i}(x^i) = \xi^k(x^i) - \Gamma(x)^k_{ij} \eta^i \xi^j$$

this is also:

$$= \xi^k(x^i) - (\eta^i \xi^k_{,i} + \Gamma(x)^k_{ij} \eta^i \xi^j) + \eta^i \xi^k_{,i}$$

$$= \xi^k(x^i) - \eta^i \xi^k_{,j} + \eta^i \xi^k_{,i}$$

Thus: cd distance from  $u$  to  $\bar{u}$  is:

$$(\sqrt{x^a + \eta^a + \xi^a - \eta^i \xi^k_{,i} + \eta^i \xi^k_{,i}}) - (\sqrt{x^a - \xi^a - \eta^a + \eta^i \xi^k_{,i} - \eta^i \xi^k_{,i}}) = T^a_{ij} \eta^i \xi^j$$



$$\tilde{\xi}^k(x^i + \eta^i) \approx \xi^k(x^i) + \frac{d\xi^k}{dx^i}(x^i) = \xi^k(x^i) - \Gamma(x)^k_{ij} \eta^i \xi^j$$

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$$= \xi^k(x^i) - \eta^i \xi^k_{,ji} + \eta^i \xi^k_{,i}$$

Thus: cd distance from  $u$  to  $\bar{u}$  is:

$$(\cancel{x^a} + \eta^a + \xi^a - \eta^i \xi^k_{,ji} + \eta^i \xi^k_{,i}) - (\cancel{x^a} - \xi^a - \eta^a + \xi^i \eta^k_{,ji} - \eta^i \xi^k_{,i}) = J^a_{ji} \eta^i \xi^j$$

Recall that indeed:  $J: \eta, \xi \rightarrow J(\eta, \xi) = \nabla_{\eta} \xi - \nabla_{\xi} \eta - [\eta, \xi]$

## Curvature:

Assume  $\xi, \eta$  and  $\xi, \eta$  are tangent vectors at  $p \in M$ .

# Curvature:

Assume  $\xi_1, \xi_2$  and  $\xi_3$  are tangent vectors at  $p \in M$ .

□ The curvature map,  $R$ , is defined through:

$$R: \xi_1, \xi_2, \xi_3 \rightarrow \underbrace{R(\xi_1, \xi_2)}_{\substack{\text{an operator, or} \\ \text{map, acting on } \xi_3}} \xi_3 = \underbrace{(\nabla_{\xi_1} \nabla_{\xi_2} - \nabla_{\xi_2} \nabla_{\xi_1} - \nabla_{[\xi_1, \xi_2]})}_{\in T_p'(M)} \xi_3$$

□ It defines the curvature tensor,  $R$ ,

← can be fed one covector and 3 vectors to yield a number

$$R \in T_3^1(M)$$

through:

$$R(\omega, \xi_1, \xi_2, \xi_3) := \underbrace{\omega(R(\xi_1, \xi_2)\xi_3)}_{\omega(R(\xi_1, \xi_2)\xi_3)} \in \mathbb{R}$$



In a chart:

$$R^i{}_{jkl} = \left\langle dx^i, R\left(\frac{\partial}{\partial x^k}, \frac{\partial}{\partial x^l}\right) \frac{\partial}{\partial x^j} \right\rangle$$

$$= \left\langle dx^i, \left( \nabla_{\frac{\partial}{\partial x^k}} \nabla_{\frac{\partial}{\partial x^l}} - \nabla_{\frac{\partial}{\partial x^l}} \nabla_{\frac{\partial}{\partial x^k}} - \underbrace{\nabla_{\left[\frac{\partial}{\partial x^k}, \frac{\partial}{\partial x^l}\right]} \right)}_{=0} \right) \frac{\partial}{\partial x^j} \right\rangle$$

$$= \left\langle dx^i, \nabla_{\frac{\partial}{\partial x^k}} \Gamma^s{}_{lj} \frac{\partial}{\partial x^s} - \nabla_{\frac{\partial}{\partial x^l}} \Gamma^s{}_{kj} \frac{\partial}{\partial x^s} \right\rangle$$

$$= \left\langle dx^i, \left( \Gamma^s{}_{lj,k} + \Gamma^r{}_{lj} \Gamma^s{}_{kr} - \Gamma^s{}_{kj,l} - \Gamma^r{}_{kj} \Gamma^s{}_{lr} \right) \frac{\partial}{\partial x^s} \right\rangle$$

$$= \Gamma^i{}_{ljk} - \Gamma^i{}_{kjl} + \Gamma^s{}_{lj} \Gamma^i{}_{ks} - \Gamma^s{}_{kj} \Gamma^i{}_{ls}$$

(at origin of geodesic cds they vanish.)

In a chart:

$$R^i{}_{jkl} = \left\langle dx^i, R\left(\frac{\partial}{\partial x^k}, \frac{\partial}{\partial x^l}\right) \frac{\partial}{\partial x^j} \right\rangle$$

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$$= \Gamma^i{}_{lj,k} - \Gamma^i{}_{kj,l} + \Gamma^s{}_{lj} \Gamma^i{}_{ks} - \Gamma^s{}_{kj} \Gamma^i{}_{ls}$$

(at origin of geodesic cds they vanish.)



## Curvature tensor's meaning?

### Intuition:

- Contains derivatives of  $\Gamma$   $\Rightarrow$
- expresses variation in gravitational forces  $\Rightarrow$
- expresses the strength and direction of "tidal forces".

### Geometry:

- Curvature expresses noncommutativity of two parallel transports.  $m \neq n$ .

Proposition: (Ricci Identity)

Assume the torsion vanishes and  
that  $\xi$  is a vector field. Then:

$$\xi^a{}_{;cd} - \xi^a{}_{;dc} = R^a{}_{cdb} \xi^b$$

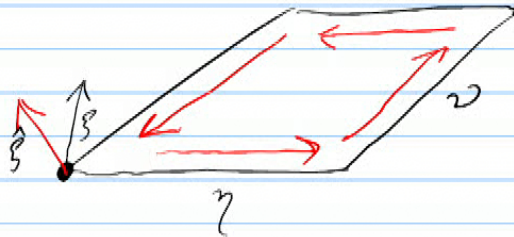
(here:  $\xi^a{}_{;cd} := \xi^a{}_{;c;d}$  etc.)

Remark:

(a bit messy to derive because need Taylor expansion,  
see, e.g., text by Stewart or Einstein)

It implies that for parallel transport  
along infinitesimal parallelogram:

$$(\hat{\xi} - \xi)^a \approx \eta^b \nu^c R^a{}_{bcd} \xi^d$$







# The "Bianchi Identities":

□ They are automatic relations among torsion and curvature, by construction.

□ Preparation: ▽ for maps!

Consider an arbitrary  $F(M)$ -linear map:

$$K: \underbrace{\xi_1 \times \xi_2 \times \dots \times \xi_r}_{\text{tangent vectors}} \rightarrow \underbrace{K(\xi_1, \dots, \xi_r)}_{\text{tangent vector}}$$

(e.g. Torsion or Curvature map)

i.e. at each  $p \in M$ :

$$K: T_p(M)^r \rightarrow T_p(M)'$$



□ Preparation: ▽ for maps!

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i.e. at each  $p \in \mathcal{M}$ :

$$K: T_p(\mathcal{M})^r \rightarrow T_p(\mathcal{M})'$$

□ We can view  $K$  as a tensor  $\tilde{K} \in T_p(\mathcal{M})'_r$ ,

## 1st Bianchi Identity:

$$\sum_{\text{cyclic}} R(\xi, \eta) \nu = \sum_{\text{cyclic}} \left( \mathcal{L}(\mathcal{L}(\xi, \eta), \nu) + (\nabla_{\xi} \mathcal{L})(\eta, \nu) \right)$$

## 2nd Bianchi Identity:

$$\sum_{\text{cyclic}} \left( (\nabla_{\xi} R)(\eta, \nu) + R(\mathcal{L}(\xi, \eta), \nu) \right) = 0$$

with obvious simplification in case  $\mathcal{L} = 0$ .

Note: They are automatically obeyed equations, just like any set of lin. operators  
or like the homogeneous Maxwell equations



$$\underbrace{\sum_{\text{cyclic}} R(\xi, \eta) v = 0}_{\text{cyclic}}$$

Indeed:

$$(\nabla_{\xi} \nabla_{\eta} - \nabla_{\eta} \nabla_{\xi}) v - \nabla_{[\xi, \eta]} v + \text{cyclic}$$

$$= \nabla_{\xi} (\nabla_{\eta} v - \nabla_{\nu} \eta) - \nabla_{[\eta, \nu]} \xi + \text{cyclic}$$

$\downarrow$  skip by 1 cyclically      $\downarrow$  skip by 1 cyclically

Exercise: Prove that:  $\nabla_{\eta} v - \nabla_{\nu} \eta = [\eta, \nu]$  (easy!)  
 without torsion:

$$= \nabla_{\xi} [\eta, \nu] - \nabla_{[\eta, \nu]} \xi + \text{cyclic}$$

// because again  $\nabla_a b - \nabla_b a = [a, b]$

$$= [\xi, [\eta, \nu]] + \text{cyclic}$$

Recall:

Assume  $A, B, C$  are linear maps  $V \rightarrow V$

$$\text{Then: } [A, [B, C]] + [C, [A, B]] + [B, [C, A]] = 0$$

i.e., the Jacobi identity holds.

Proof: Simply spell out the commutators.

Remark: This means that, e.g., in quantum mechanics, all objects that  
that need to be representable as operators on the Hilbert space  
must obey the Jacobi identity, e.g., generators of symmetries.

This is why the Jacobi  
identity is one of the  
axioms of Lie Algebras.



$= 0$  by Jacobi identity for all lin. maps.

Recall:

Assume  $A, B, C$  are linear maps  $V \rightarrow V$

Then:  $[A, [B, C]] + [C, [A, B]] + [B, [C, A]] = 0$

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