

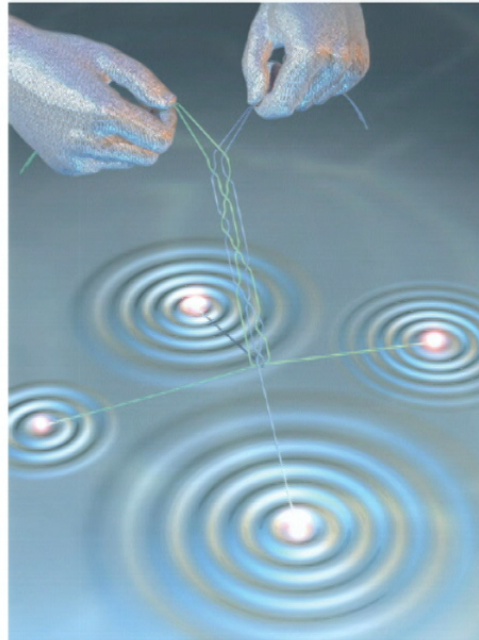
Title: Topological Quantum Computation

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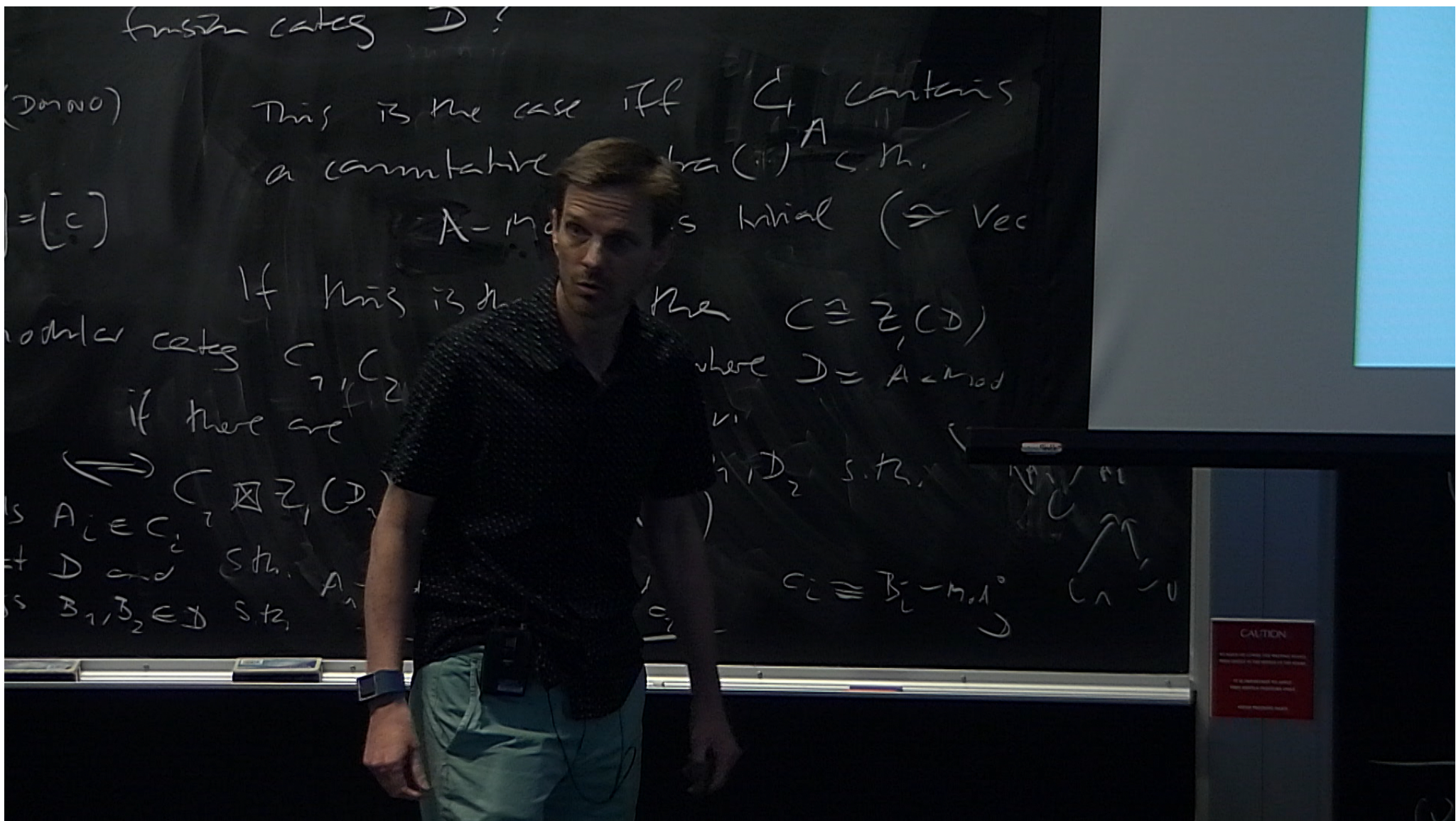
Abstract: The (Freedman-Kitaev) topological model for quantum computation is an inherently fault-tolerant computation scheme, storing information in topological (rather than local) degrees of freedom with quantum gates typically realized by braiding quasi-particles in two dimensional media. I will give an overview of this model, emphasizing the mathematical aspects.

# Topological Quantum Computation



Eric Rowell, PI, July 2017





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## Exchange Statistics: Anyons

How does the wave function  $\psi(z_1, z_2)$  of point-like particles change under  $z_1 \leftrightarrow z_2$ ?



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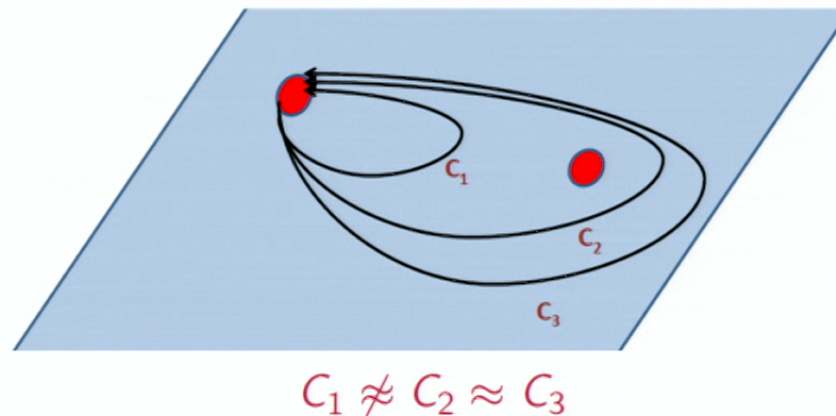
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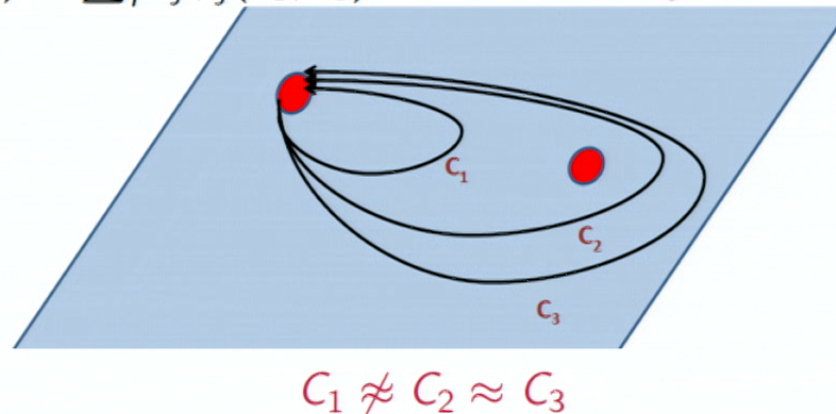




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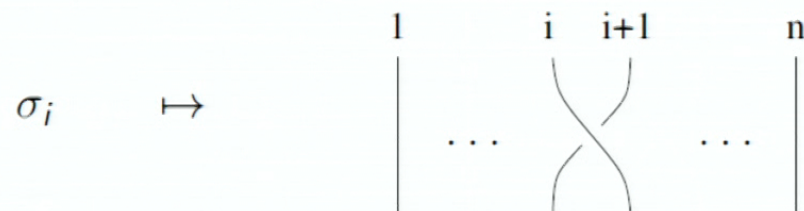
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- ▶ Particle exchange  $\rightsquigarrow$  reps. of braid group  $\mathcal{B}_n$
- ▶ Why?  $\pi_1(\mathbb{R}^3 \setminus \{z_i\}) = 1$  but  $\pi_1(\mathbb{R}^2 \setminus \{z_i\}) = F_n$  Free group.

The hero is the **Braid Group**  $\mathcal{B}_n$ :  $\sigma_1, \dots, \sigma_{n-1}$  with

(R1)  $\sigma_i \sigma_{i+1} \sigma_i = \sigma_{i+1} \sigma_i \sigma_{i+1}$

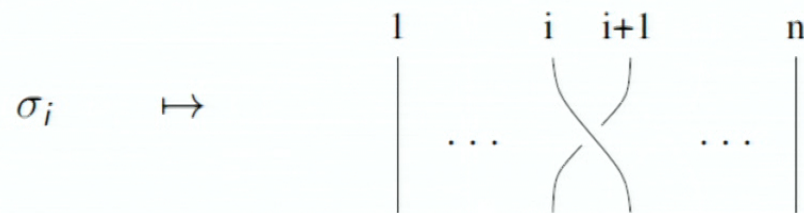
(R2)  $\sigma_i \sigma_j = \sigma_j \sigma_i$  if  $|i - j| > 1$



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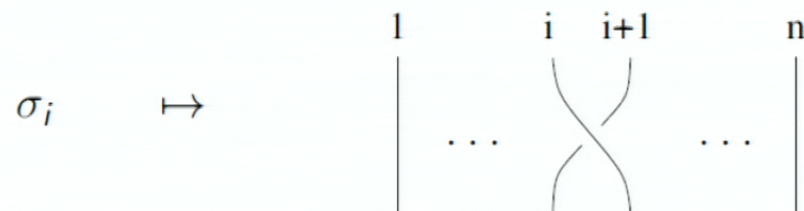
Motions of  $n$  points in a disk/Mapping Class Group.



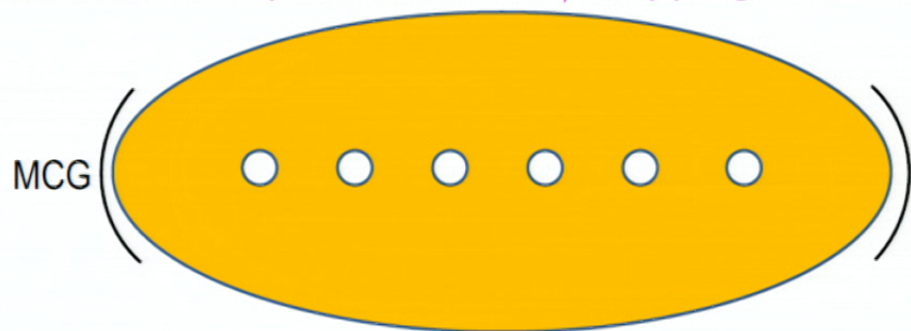
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# Topological Model

## Computation

output

apply gates

initialize

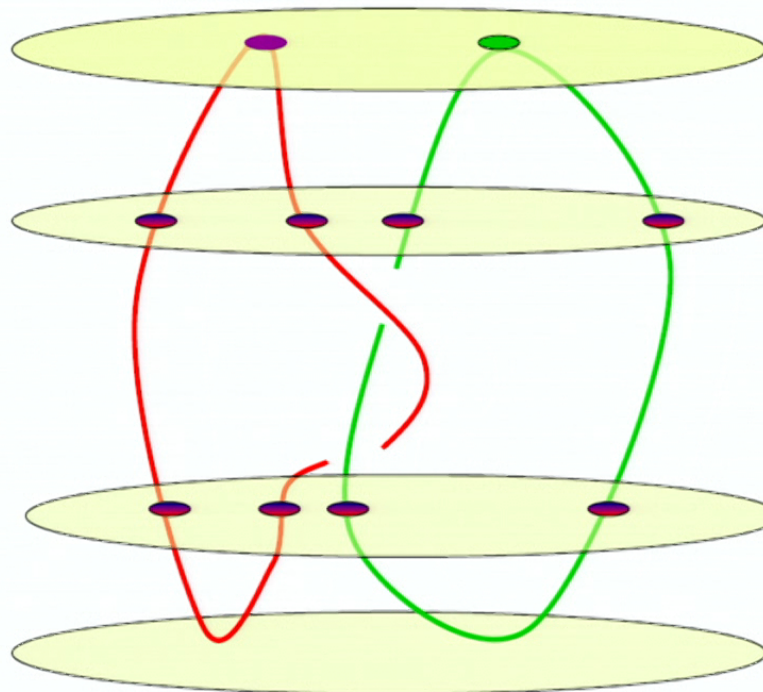
vacuum

## Physics

measure  
(fusion)

braid  
anyons

create  
anyons





# Foundational Questions

1. How to model Anyons on Surfaces? State Spaces?  
Topological Quantum Circuits?
2. Why is TQC Fault-tolerant?

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1. How to model Anyons on Surfaces? State Spaces?  
Topological Quantum Circuits?
2. Why is TQC Fault-tolerant?
3. How Powerful are TQCs?

# Modeling Anyons on Surfaces

Definition (Nayak, et al '08)

a (bosonic) system is in a **topological phase** if its low-energy effective field theory is a **topological quantum field theory** (TQFT).



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A (2+1)D **TQFT** assigns to any (surface, boundary data)  $(M, \ell)$  a **Hilbert** space:

$$(M, \ell) \rightarrow \mathcal{H}(M, \ell).$$

$\begin{matrix} & & x_1, y_0 z & & x_1, y & & x_1, z \\ & \diagdown & & \diagup & & \diagdown & & \diagup \\ & x & & y & & x & & y \end{matrix} = \parallel$ 
 $c_{x,y} \cdot c_{y,x} = id_{y \otimes x} \quad \forall x, y$

$x \rightarrow \frac{x}{c_{x,y}} + \frac{x}{(c_{y,x})^{-1}} = \frac{x}{c_{x,y}}$

is a modular category  $\subset$  the  $\mathbb{Z}_2$  of some fusion category  $\mathcal{D}$ ?

(Deno) This is the case iff  $\mathcal{C}_1$  contains a commutative  $A = (1) \subset \mathcal{C}_1$ .  
 $A$ -mod  $\cong$  trivial  $(\cong \text{Vec})$

$\mathcal{C} = [\mathcal{C}]$

If this is the case  $\mathcal{C} \cong \mathbb{Z}_2(\mathcal{D})$   
 where  $\mathcal{D} \cong A$ -mod

modular category  $\mathcal{C}_1, \mathcal{C}_2$  are 1  
 if there are fusion

$\mathcal{C}_1 \boxtimes \mathbb{Z}_2(\mathcal{D}_1) \cong \mathcal{C}_2$   
 as  $A_i \in \mathcal{C}_i$  s.t.  $A$ -mod  $\cong \mathcal{C}_1$   
 at  $\mathcal{D}$  and  $\mathcal{D}_1$  s.t.  $A$ -mod  $\cong \mathcal{C}_1$   
 as  $B_1, B_2 \in \mathcal{D}$  s.t.

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Boundary  $\bigcirc$  labeled by  $i \in \mathcal{L}$ : finite set of colors  $\leftrightarrow$  (anyons).

$(x, r_x) \otimes (y, r_y) = (x \otimes y, r_{x \otimes y})$  where  $r_{x \otimes y} = r_x \otimes r_y$

$1 = (1, r_1)$  where  $r_1(1) = 1$

$c_{(x, r_x), (y, r_y)} = c_x(r_y)$

$\mathcal{C} = \mathbb{Z}_2(\mathcal{C})$  trivial? Yes, for action  $\mathcal{C}$   
 ↑ labeled fusion category  $\mathcal{C}$  with dim  $\mathcal{C} \neq 0$



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## Basic pieces/Local Axioms

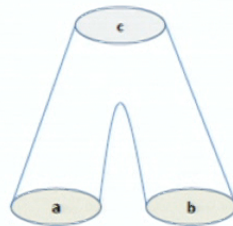
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► pants:

$P :=$



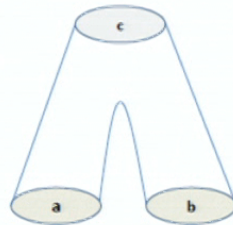
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$$\mathcal{H}(P; a, b, c) = \mathbb{C}^{N(a,b,c)} \swarrow \text{choices!}$$

## Compatibility Axioms

Axiom (Disjoint Union)

$$\mathcal{H}[(M_1, \ell_1) \amalg (M_2, \ell_2)] = \mathcal{H}(M_1, \ell_1) \otimes \mathcal{H}(M_2, \ell_2)$$



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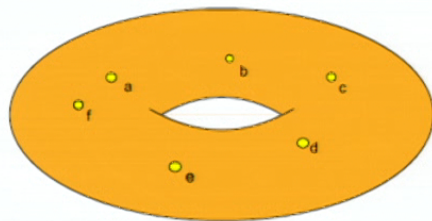
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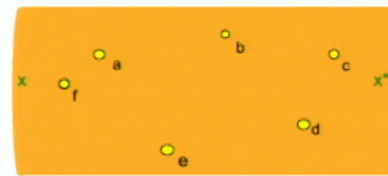
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$(M, \ell)$

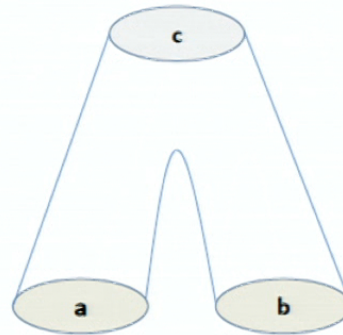


$(M_g, \ell, x, x^*)$



## Fusion Channels

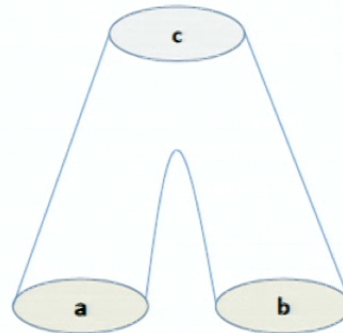
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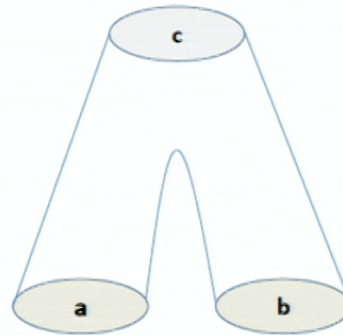


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### Principle

The **Computational Space**  $\mathcal{H}_n := \mathcal{H}(D^2; a, \dots, a)$ : the state space of  $n$  identical type  $a$  anyons in a disk.



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$$\mathfrak{g} \rightsquigarrow U\mathfrak{g} \rightsquigarrow U_q\mathfrak{g} \xrightarrow{q=e^{\pi i/\ell}} \text{Rep}(U_q\mathfrak{g}) \xrightarrow{\langle \text{Ann}(Tr) \rangle} \mathcal{C}(\mathfrak{g}, \ell),$$



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$\text{Rep}(D^\omega G)$  Drinfeld doubles/centers ...

Theorem (Bruillard, Ng, R, Wang)

For fixed  $k$ , *finitely many* models with  $|\mathcal{L}| = k$

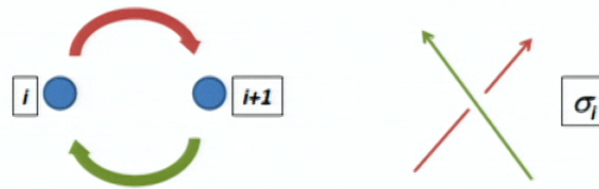
## Anyon Model $\leftrightarrow$ Modular Category

anyonic system	Modular Category
anyon types $x \in \mathcal{L}$	simple $X$
vacuum $0 \in \mathcal{L}$	<b>1</b>
$x^*$ antiparticle	dual $X^*$
$\mathcal{H}(P; x, y, z)$ state space	$\text{Hom}(X \otimes Y, Z)$
particle exchange	braiding $c_{X,X}$
Locality	Gluing Axiom
Entanglement	Disjoint Union Axiom
anyon types distinguishable	$\det(S) \neq 0$
topological spin	$\theta_X$
$n$ anyon state space	$\text{End}(X^{\otimes n})$

# Topological Circuits and Fault-Tolerance

Fix **anyon**  $a$

- ▶  $\sigma_i \in \mathcal{B}_n$  acts on  $\mathcal{H}_n = \text{End}(X_a^{\otimes n})$  by particle exchange



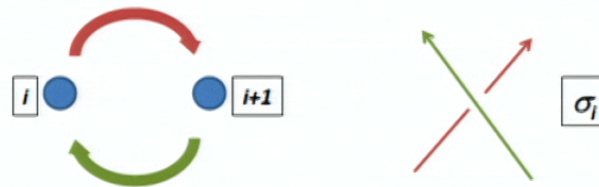
- ▶ **Braid group** representation  $\rho_a : \mathcal{B}_n \rightarrow U(\mathcal{H}_n)$
- ▶ Quantum Gates:  $\rho_a(\sigma_i)$ ,



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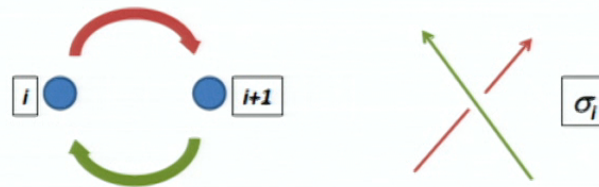
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Information is **de-localized** → **Fault-Tolerance** as errors are **local**.

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## Principle

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quantum gates  $\rho_a(\sigma_i)$  are non-local, topological operations.

## Simulating TQCs on QCM

Freedman, Kitaev & Wang showed TQCs have **hidden locality**: Let  $U(\beta) \in \mathbf{U}(\mathcal{H}_n)$  be a braiding quantum circuit.



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Goal: simulate  $U$  on  $V^{\otimes k(n)}$  by **local** gates for some v.s.  $V$ .

- ▶ Set  $V = \bigoplus_{(a,b,c) \in \mathcal{L}^3} \mathcal{H}(P; a, b, c)$

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### Remark

$V$  can be quite large and  $U(\beta)$  only acts on the subspace  $\mathcal{H}_n$ , non-computational space  $\mathcal{H}_n^\perp$  can be large



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**Idea:** Push braiding gates inside a braided QCM.

## Example Ising $\mathcal{C}(\mathfrak{sl}_2, 4)$

$$\text{Let } R = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 0 & 0 & 1 \\ 0 & 1 & -1 & 0 \\ 0 & 1 & 1 & 0 \\ -1 & 0 & 0 & 1 \end{pmatrix}$$



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*Fibonacci not localizable.*

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KEY:  $\overline{\rho_\tau(\mathcal{B}_n)} \supset SU(f_n) \times SU(f_{n-1})$ , where  $\mathcal{H}_n \cong \mathbb{C}^{f_n} \oplus \mathbb{C}^{f_{n-1}}$ ,

$$f_n = \frac{(1 + \sqrt{5})^n - (1 - \sqrt{5})^2}{2^n \sqrt{5}}$$



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$\dim(\sigma) = \sqrt{2}$  (Ising) while  $\dim(\tau) = \frac{1+\sqrt{5}}{2}$  (Fibonacci).



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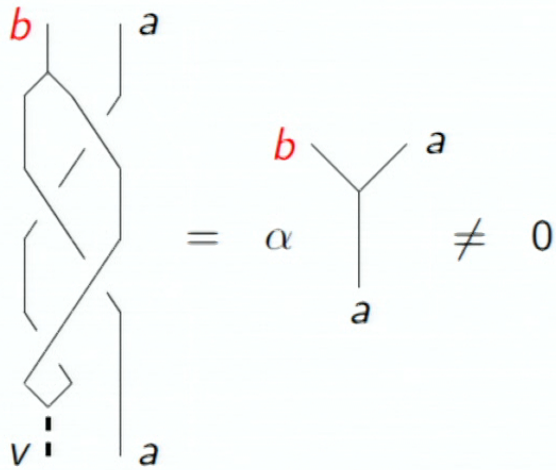
True for “fundamental” anyons in **all** quantum group models.

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If  $\dim(a) > 1$  there is a  $b \neq v$  ( $v = \text{vacuum}$ ) with  $N(a, a, b) \neq 0$ .

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Diagrammatic equation showing a braid of  $a$  and  $b$  lines with a vacuum line  $v$ , equal to a fusion coefficient  $\alpha$  times a tree diagram, which is non-zero:

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THANK YOU!

see: [arXiv:1705.06206](https://arxiv.org/abs/1705.06206)

