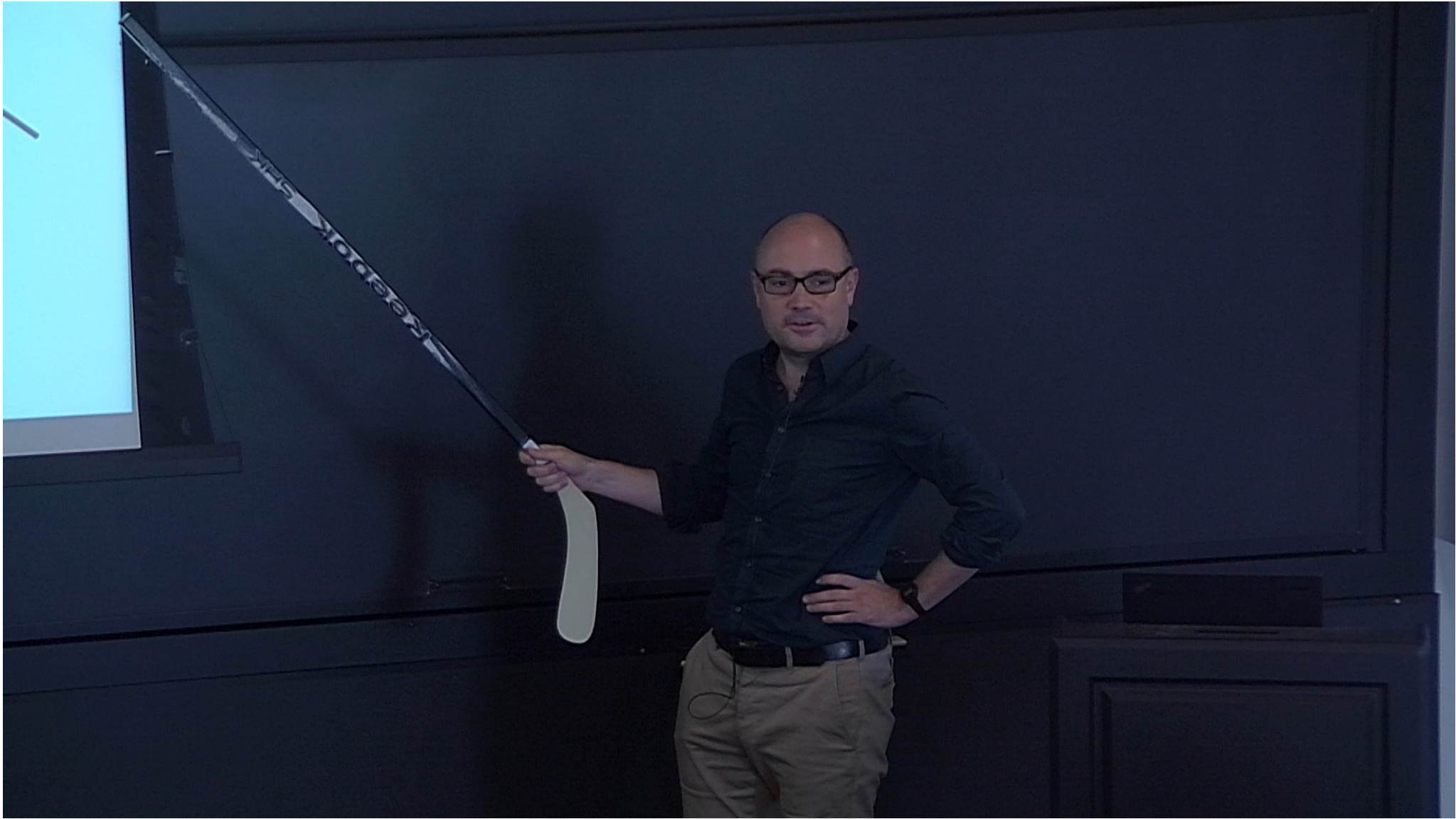


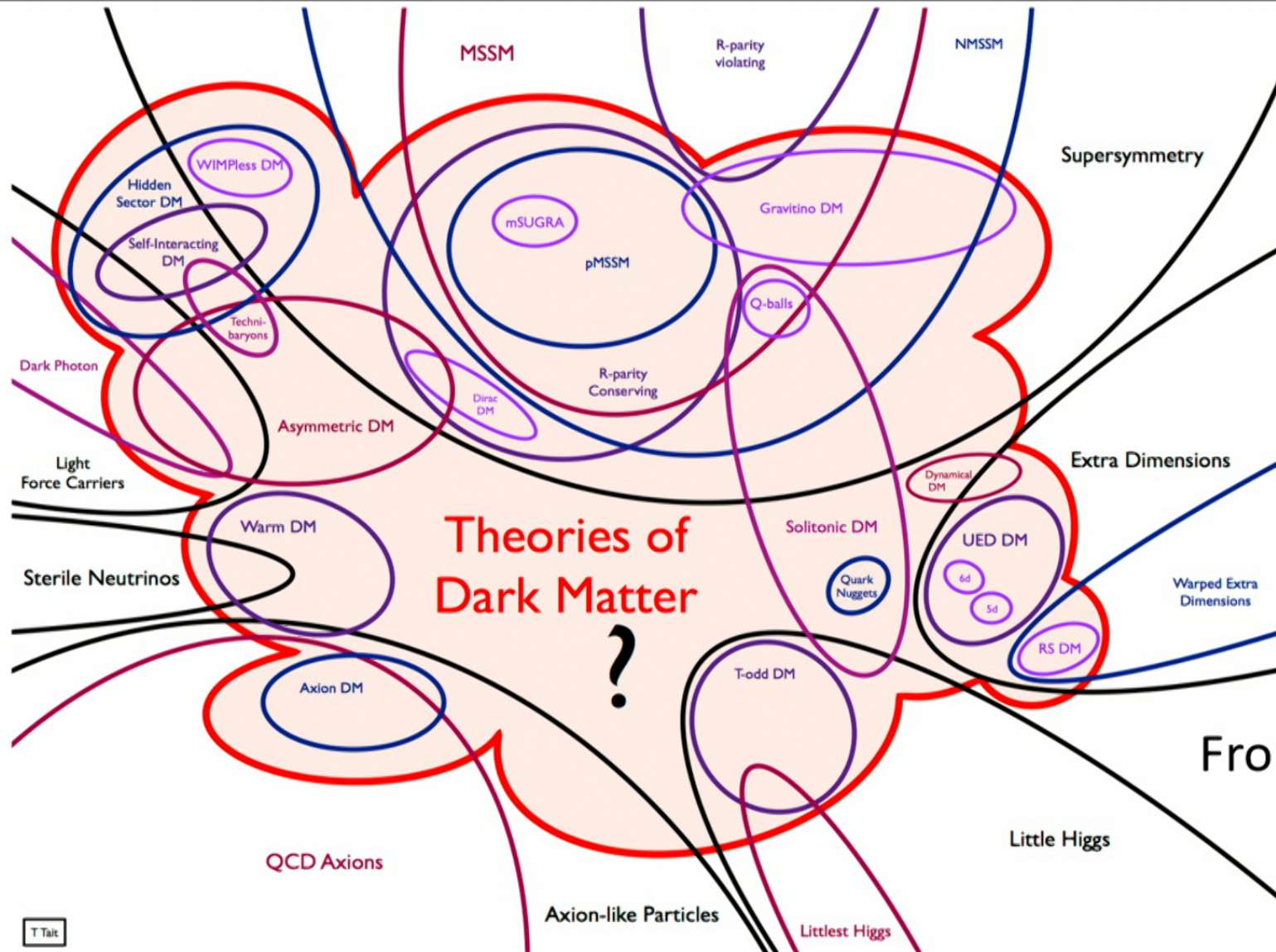
Title: Neutrinos as Signal and Background - detecting and avoiding neutrinos at dark matter detectors

Date: Jul 21, 2017 09:50 AM

URL: <http://pirsa.org/17070017>

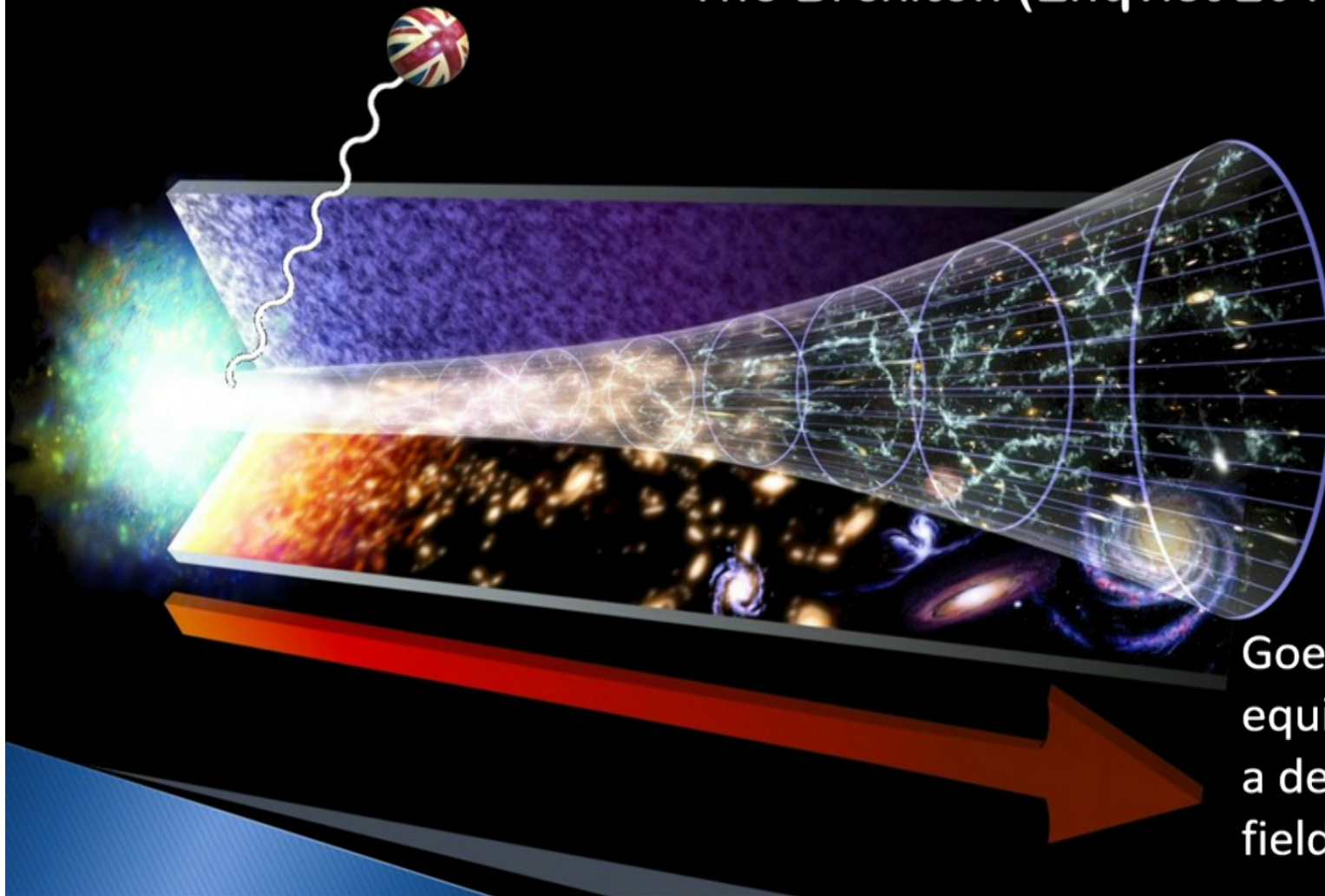
Abstract:





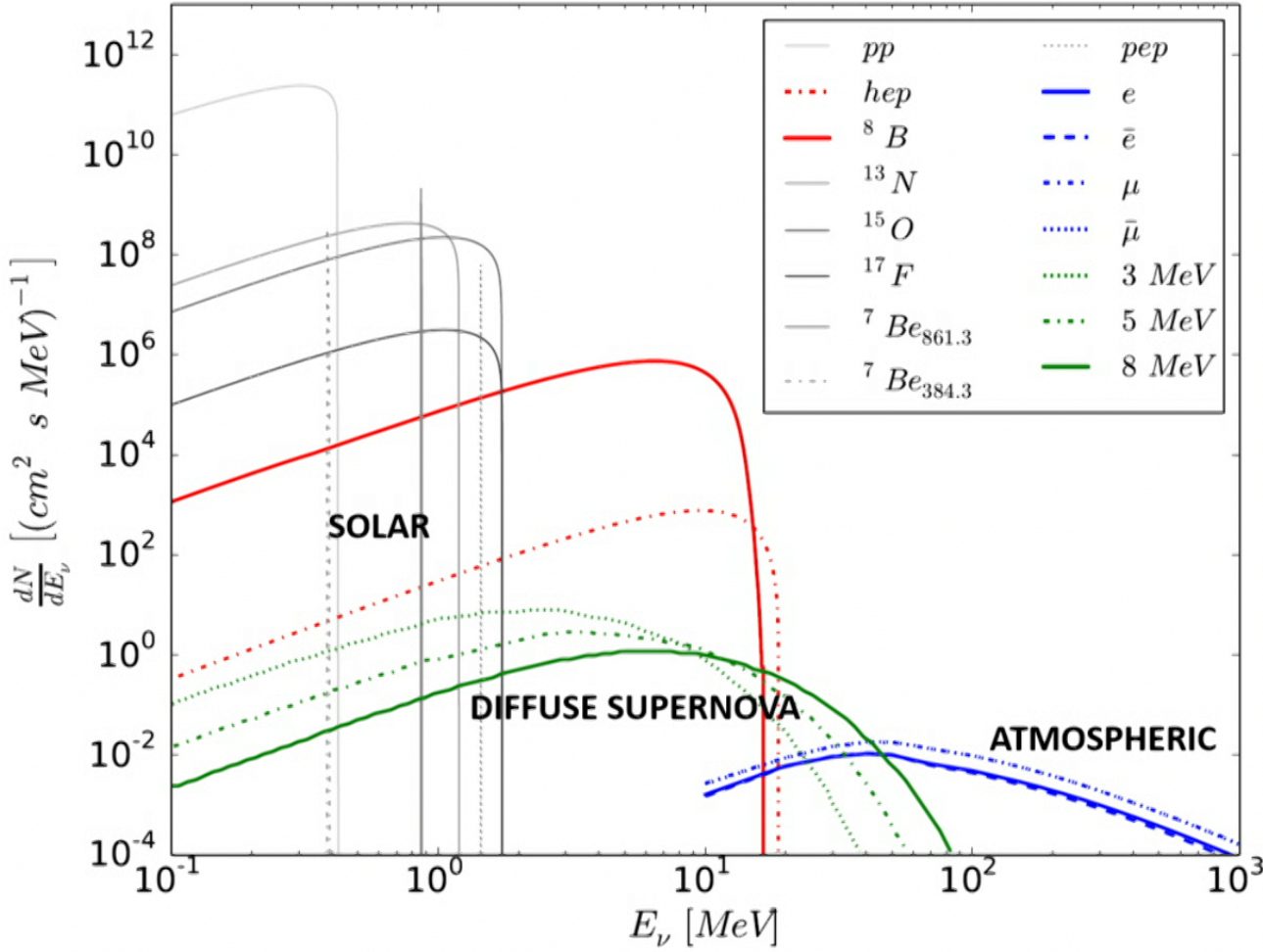
From Tim Tait

# The Brexiton (Enqvist 2017, Rajantie 2017)



Goes out of thermal equilibrium then acts as a decoupled spectator field while it decays

# Neutrino Background



# Coherent Neutrino-Nucleon Interactions

$$\frac{d\sigma}{d(\cos\theta)} = \frac{G_F^2}{8\pi} Q_W^2 E_\nu^2 (1 + \cos\theta) F(Q^2)^2$$

- Enhanced by factor  $N^2$ :

$$Q_W = N - (1 - 4 \sin^2 \theta_W)Z \approx N - 0.08 \times Z \approx N$$

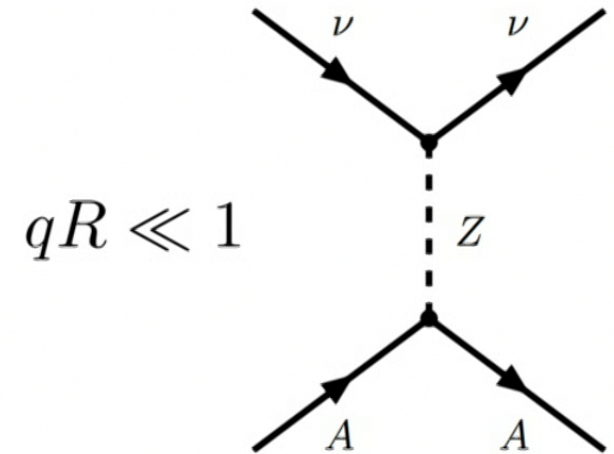
- $\cos\theta$ : angle between in- and outgoing neutrino direction

- $2m_T E_r = q^2 = 2E_\nu^2(1 - \cos\theta)$

$$\Rightarrow \frac{d\sigma}{dE_r} = \frac{G_F^2}{4\pi} Q_W^2 m_T \left(1 - \frac{m_T E_r}{2E_\nu^2}\right) F(Q^2)^2.$$

$$\frac{dR_\nu}{dE_r} = n_T \int_{t_0}^{t_1} \int_{E_\nu^{\min}}^{\infty} \frac{dN(t)}{dE_\nu} \frac{d\sigma(E_\nu, E_r)}{dE_r} dE_\nu dt$$

$$R_\nu = \int_{E_{\text{thr}}}^{E_{\text{up}}} \frac{dR_\nu}{dE_r} dE_r$$



# Coherent Neutrino-Nucleon Interactions

$$\frac{d\sigma}{d(\cos\theta)} = \frac{G_F^2}{8\pi} Q_W^2 E_\nu^2 (1 + \cos\theta) F(Q^2)^2$$

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$$Q_W = N - (1 - 4 \sin^2 \theta_w) Z \approx N - 0.08 \times Z \approx N$$

- $\cos\theta$ : angle between in-

- $2m_T E_r = q^2 = 2E_\nu^2 (1 - \cos\theta)$

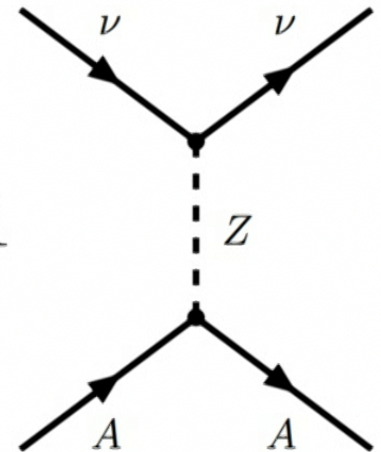
$$\Rightarrow \frac{d\sigma}{dE_r} = \frac{G_F^2}{4\pi} Q_W^2 m_T \left(1 - \frac{m_T E_r}{2E_\nu^2}\right) F(Q^2)^2.$$

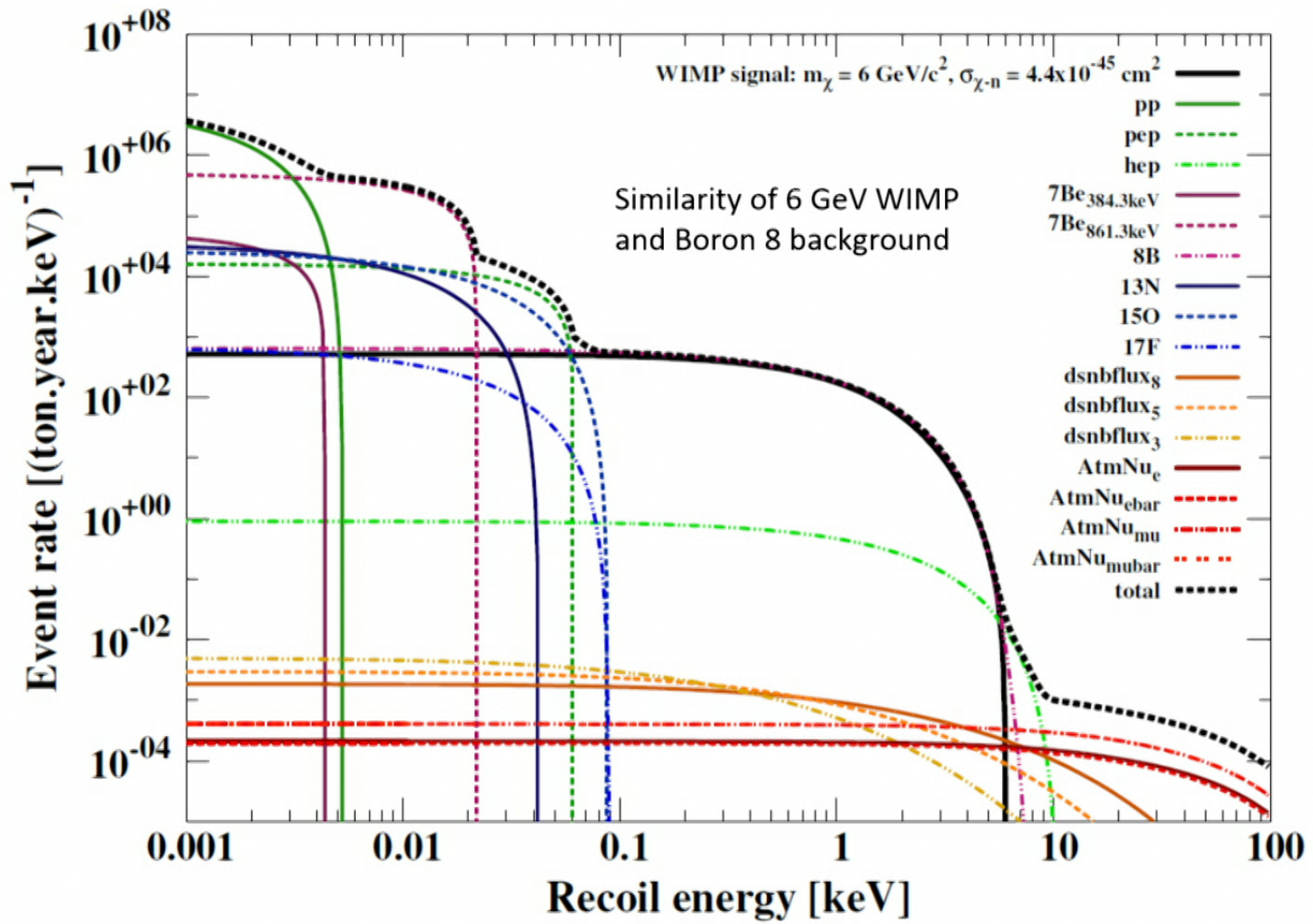
$$\frac{dR_\nu}{dE_r} = n_T \int_{t_0}^{t_1} \int_{E_\nu^{\min}}^{\infty} \frac{dN(t)}{dE_\nu} \frac{d\sigma(E_\nu, E_r)}{dE_r} dE_\nu dt$$

$$R_\nu = \int_{E_{\text{thr}}}^{E_{\text{up}}} \frac{dR_\nu}{dE_r} dE_r$$

**STILL NOT OBSERVED  
IN STANDARD MODEL**

$< 1$





Ruppin et al 1408.3581

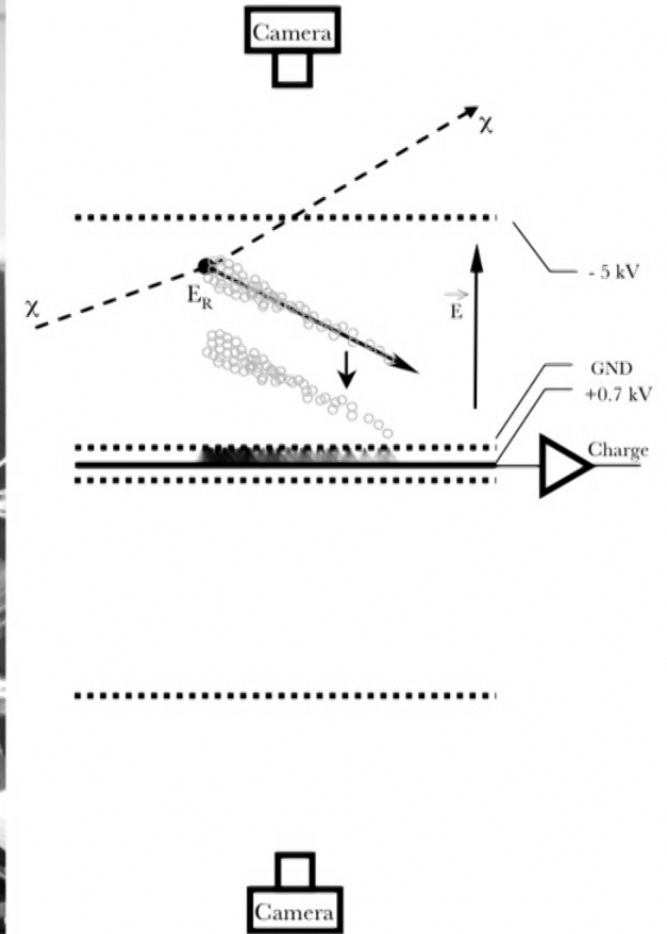


What if we can Tell which direction the dark matter is coming from?

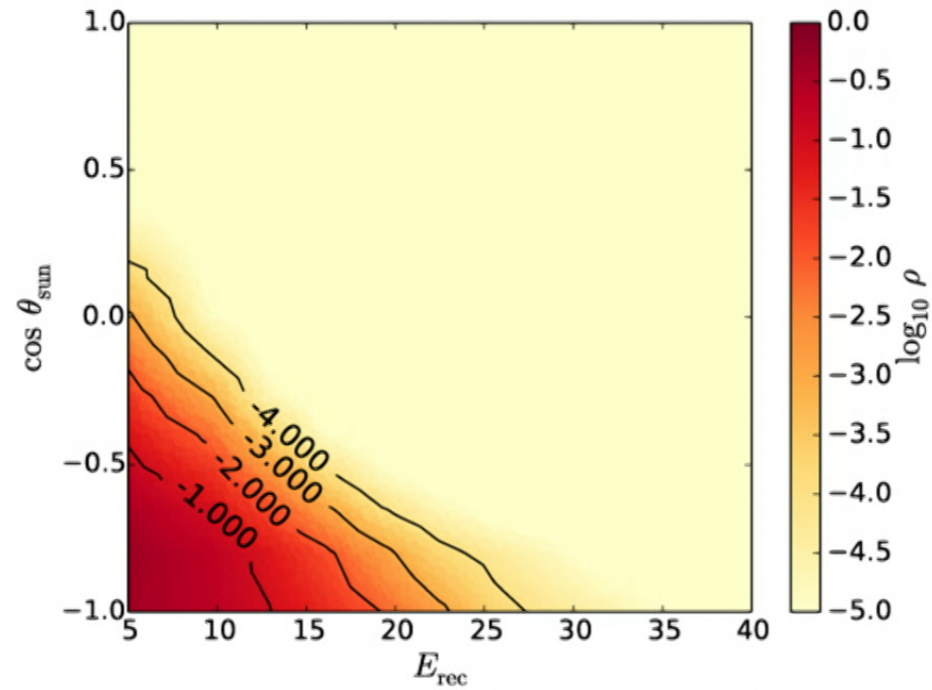
**DIRECTIONAL DARK MATTER DETECTION**

e.g. DMTPC

Faces major technical challenges to scale up but lets carry on regardless...



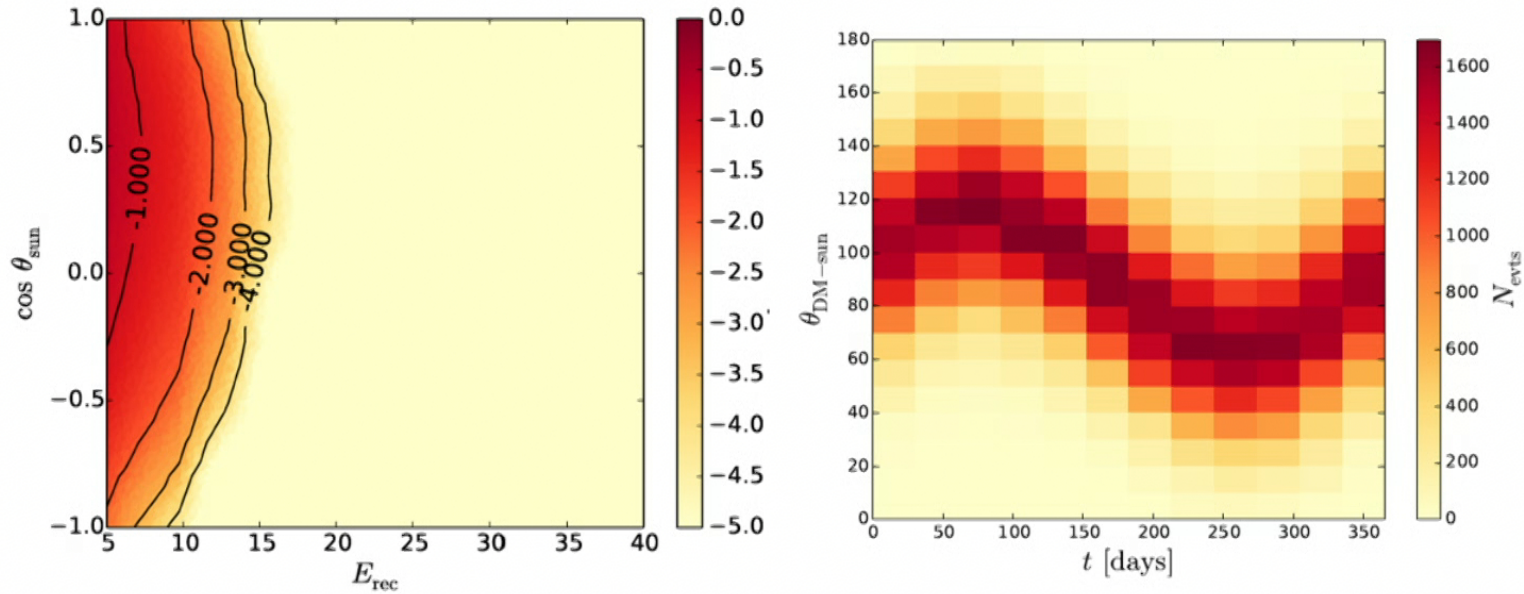
## angle between recoil from Solar neutrino and sun



$$\cos \theta' = \frac{E_\nu + m_T}{E_\nu} \sqrt{\frac{E_r}{2m_T}}$$

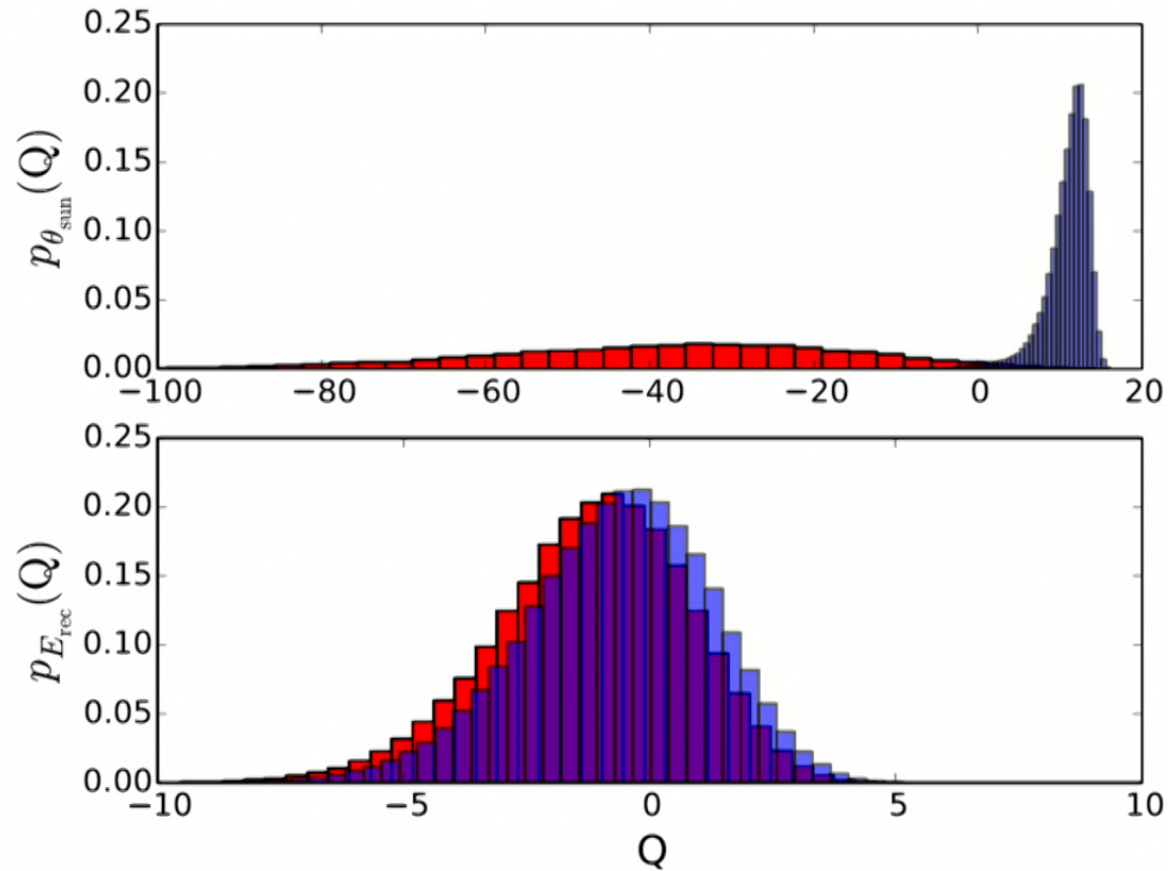
arXiv:1406.5047

## angle between recoil from Dark Matter and sun



- Preferred arrival direction roughly from Cygnus A
- This changes during the year
- Lighter (heavier) dark matter more (less) directional above a given threshold

arXiv:1406.5047



The normalised background only distribution  $p_B(Q_B)$  (blue) and signal plus background distribution  $p_{SB}(Q_{SB})$  (red) including angular information (top) and excluding angular information (bottom) for  $s=10$  and  $b=500$  for a 6 GeV dark matter particle in a  $\text{CF}_4$  detector.

arXiv:1406.5047

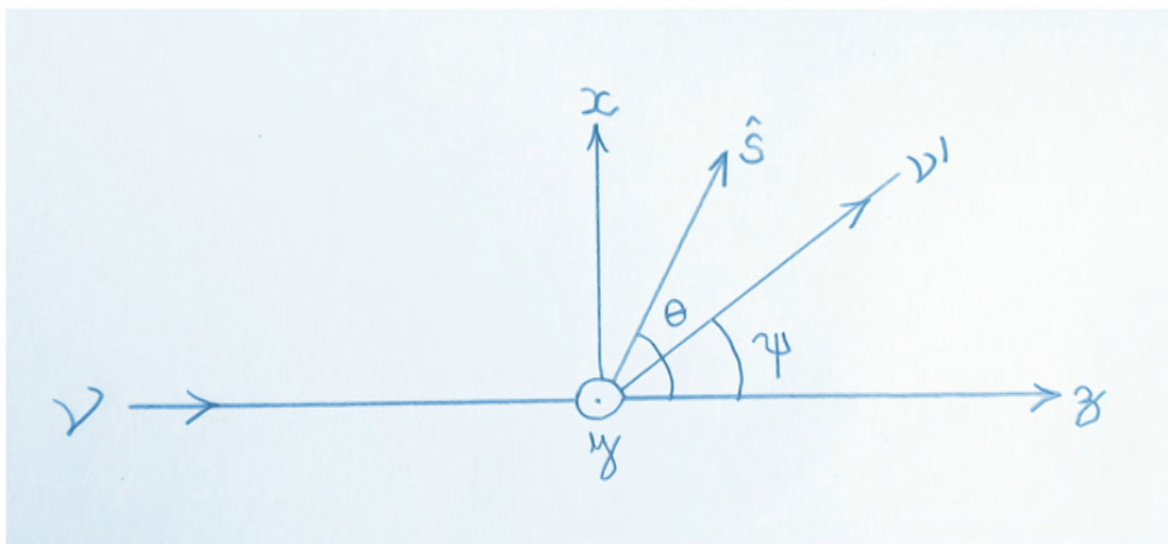
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## Interesting Possibility – Polarised targets

- Polarised targets not very directional for dark matter  
(effect is suppressed when no preferred helicity)
- Polarised targets with unpaired neutrons ARE directional to axial coupling of neutrinos
- Effect usually dwarfed by vector coupling due to coherent enhancement
- Notable exception is Helium-3

arXiv:1605.08727

see also  
“Dark Matter Detection with Polarized Detectors”  
Chiang, Kamionkowski & Krnjaic, arXiv:1202.1807



if  $N=1$  and  $c_A$  due to unpaired neutron

cancellation between  $V$  and  $A$  for particular orientations of the spin and the arrival direction of the neutrino

$$\frac{d\sigma}{d\Omega} = \frac{G_F^2 E_\nu^2}{16\pi^2} \left\{ \underbrace{c_V^2 - 3c_A^2 + (c_V^2 - c_A^2)\cos\psi}_{\text{SI}} + \underbrace{2c_A[(c_V - c_A)\hat{\nu} \cdot \hat{s} + (c_V + c_A)\hat{\nu}' \cdot \hat{s}]}_{\text{SD}} \right\}$$

**SI**

**SD**

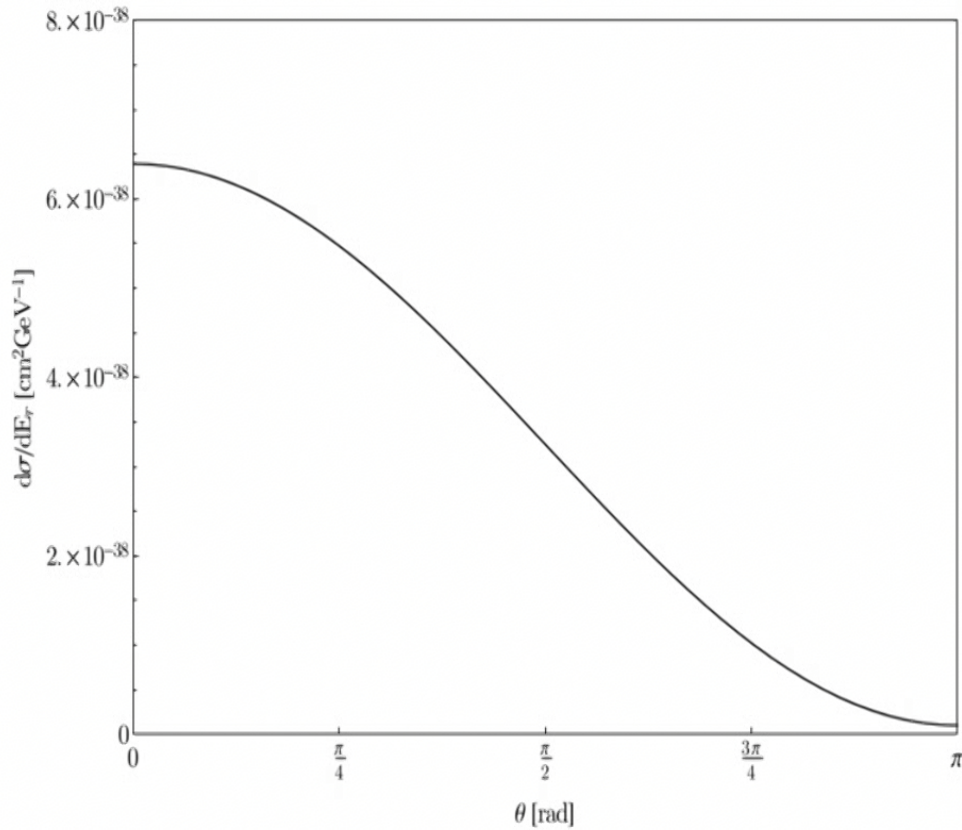
$$c_V^{\text{nucleus}} = Zc_V^p + Nc_V^n$$

$$c_A^{\text{nucleus}} = c_A^{\text{unpaired nucleon}}$$

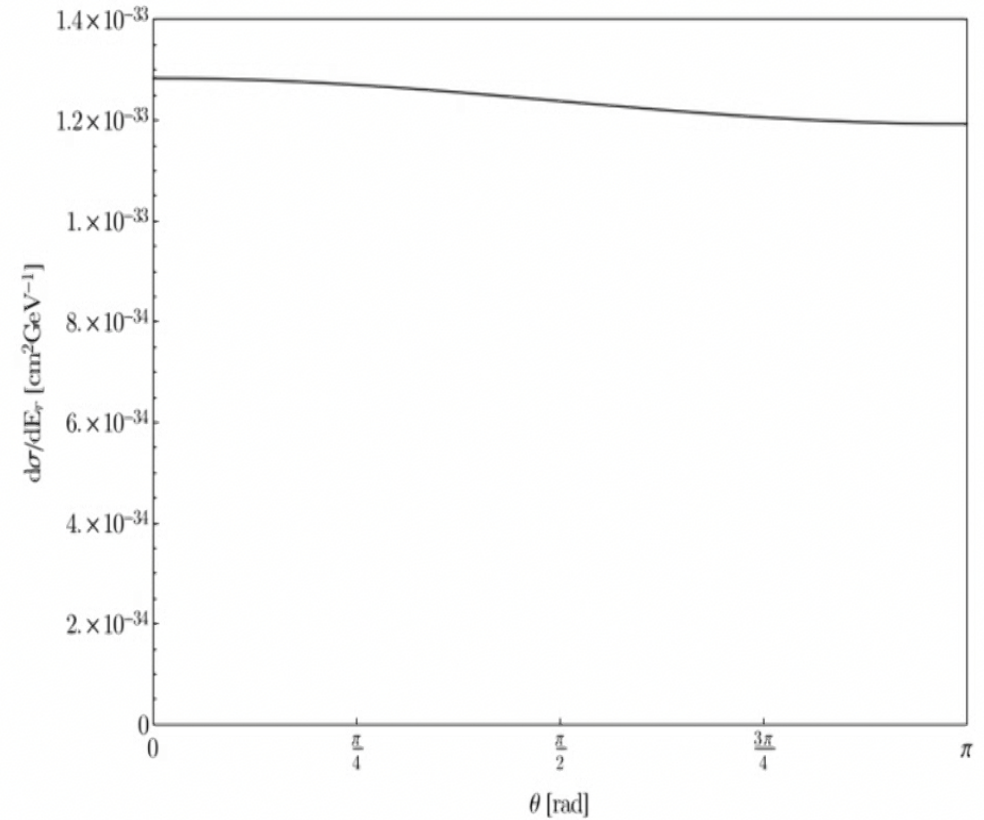
	$c_V$	$c_A$
Proton	$1 - 4\sin^2\theta_W$	1.26
Neutron	-1	-1.26

# 6.4 MeV Neutrino-nucleon cross section as function of angle

For Xenon there is a small effect while for Helium-3 there is almost a complete cancellation.



$^3\text{He}$

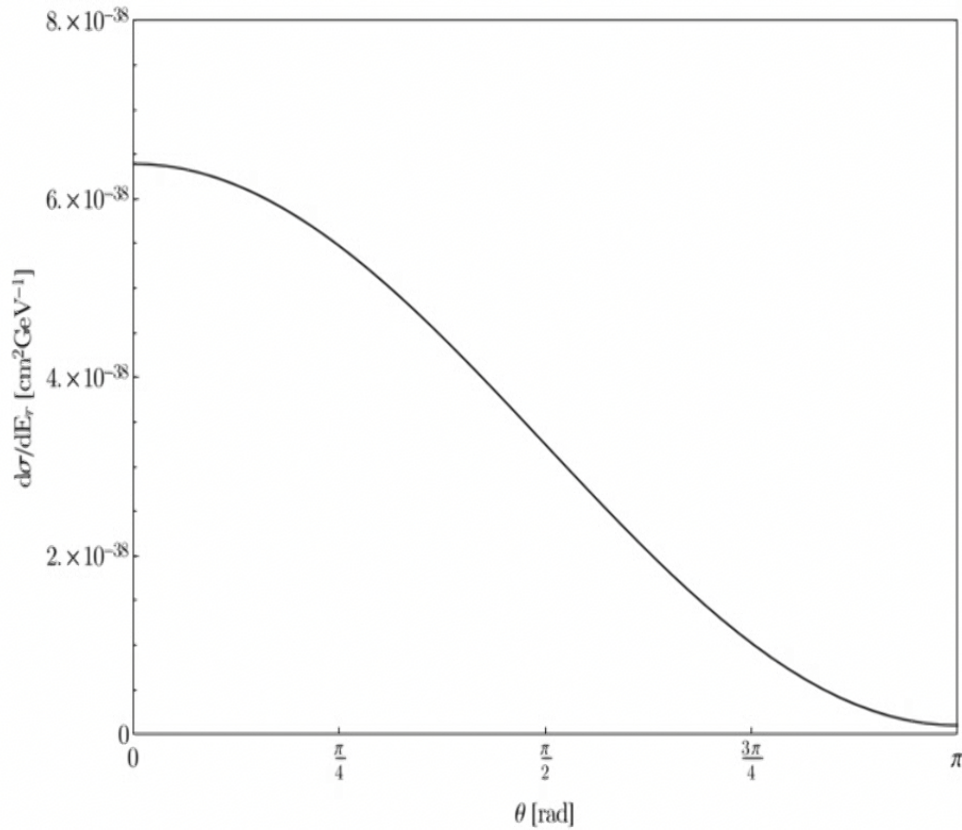


$^{129}\text{Xe}$

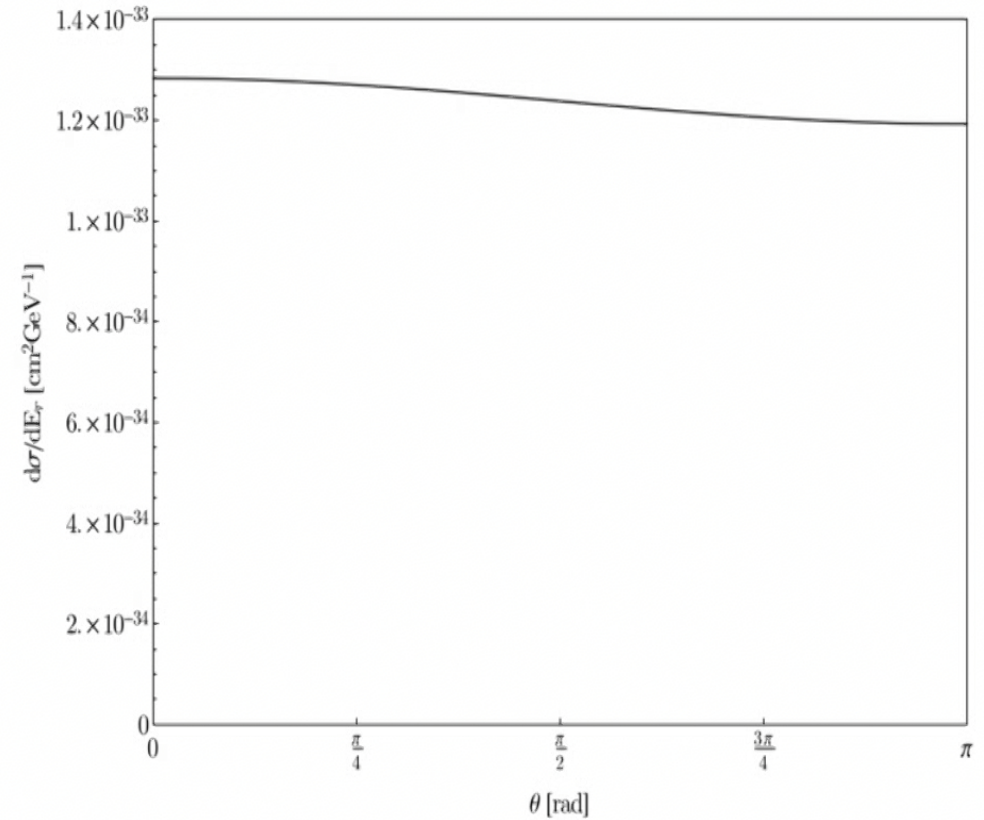
arXiv:1605.08727

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$^{129}\text{Xe}$

arXiv:1605.08727



## Some obvious problems with Helium-3

- Tritium contamination would be a major background
- Simplest Polarisation scheme for He-3 for NMR uses potassium and/or rubidium, both of which are potential contaminants
- Helium-3 makes Xenon look as cheap as water

$$\alpha = \frac{1}{2} \left| \frac{\frac{d\sigma}{dE_r}(0) - \frac{d\sigma}{dE_r}(\pi)}{\frac{d\sigma}{dE_r}(\pi/2)} \right|$$

arXiv:1605.08727

	$\alpha$
$^3\text{He}$	0.97
$^{13}\text{C}$	0.41
$^{15}\text{N}$	0.36
$^{19}\text{F}$	0.22
$^{129}\text{Xe}$	0.04

## We expect to detect Neutrinos. What could we do with this information?

Experiment	$\epsilon$ (ton-year)	$E_{th,n}$ (keV)	$E_{th,o}$ (keV)	$E_{max}$ (keV)	$R(pp)$	$R(^8\text{B})$
G2-Ge	0.25	0.35	0.05	50	–	[62 – 85]
G2-Si	0.025	0.35	0.05	50	–	[3 – 3]
G2-Xe	25	3.0	2.0	30	[2104 – 2167]	[0 – 64]
Future-Xe	200	2.0	1.0	30	[17339 – 17846]	[520 – 10094]
Future-Ar	150	2.0	1.0	30	[14232 – 14649]	[6638 – 12354]
Future-Ne	10	0.15	0.1	30	[1141 – 1143]	[898 – 910]

arXiv:1604.01025

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arXiv:1604.01025

**We expect to detect Neutrinos. What could we do with this information?**

Measure Boron-8 flux using nuclear recoils and pp flux using electron recoils

Can measure the Weinberg angle at very low energies

Exp.	$\phi_\nu^{8\text{B}}$	$\phi_\nu^{pp}$	$\sin^2 \theta_W$
Measured	2.0% <sup>a</sup>	10.6 % <sup>b</sup>	
G2	1.9% (1.9%)	2.5 % (2.5%)	4.6% (4.5%)
Future-Xe	1.8% (0.9%)	0.7% (0.7%)	1.7% (1.7%)
Future-Ar	1.0% (0.6%)	0.6% (0.5%)	1.5% (1.4%)
HyperK <sup>c</sup>	1.43%	—	—

arXiv:1604.01025

# We expect to detect Neutrinos. What could we do with this information?

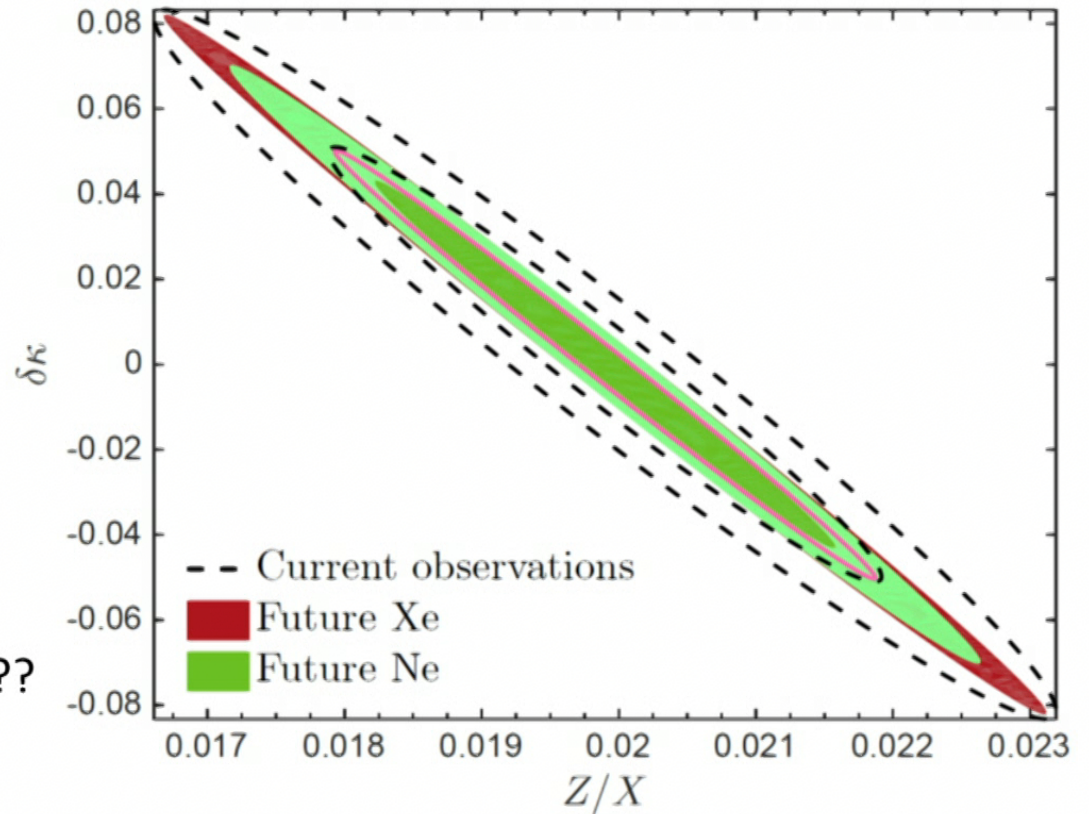
Limits average opacity vs. metallicity

Narrows line but still huge degeneracy

Needs to be broken by observation of CNO neutrinos –

SNO+ ???

Future direct detection experiments ???



arXiv:1604.01025

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## Tests of BSM Physics

Momentum exchanged for pp-neutrino electron events is around 10 keV

Momentum exchanged for neutrino-nucleon events is about MeV scale

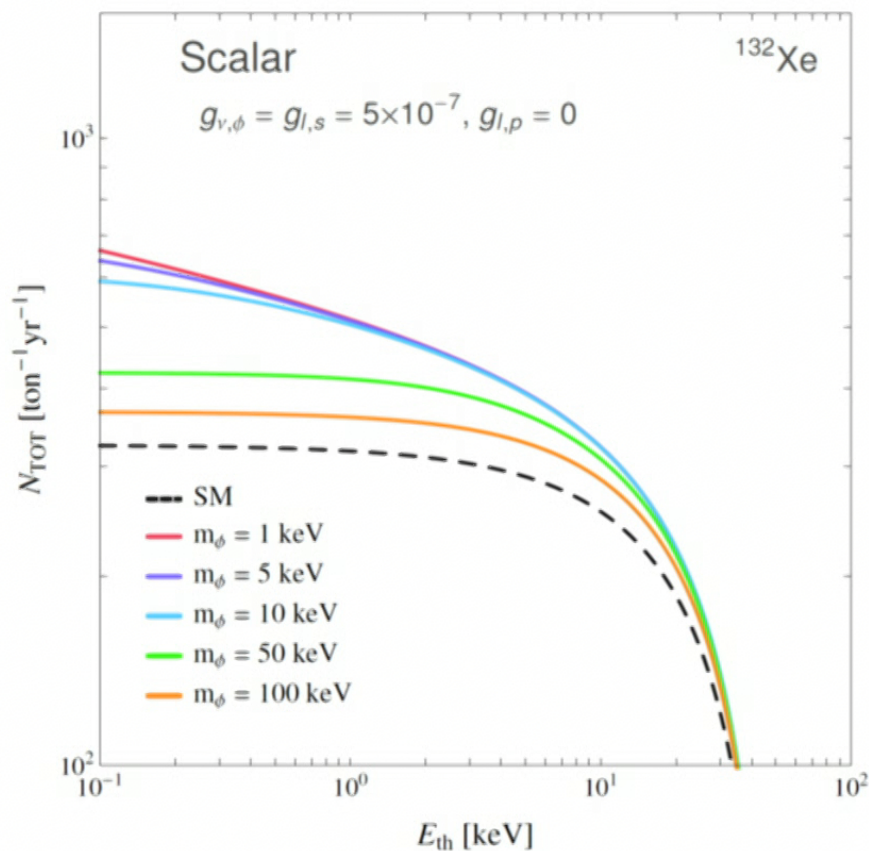
Both  $Q^2$  unstudied in those settings, can probe new interactions.

arXiv:1604.01025

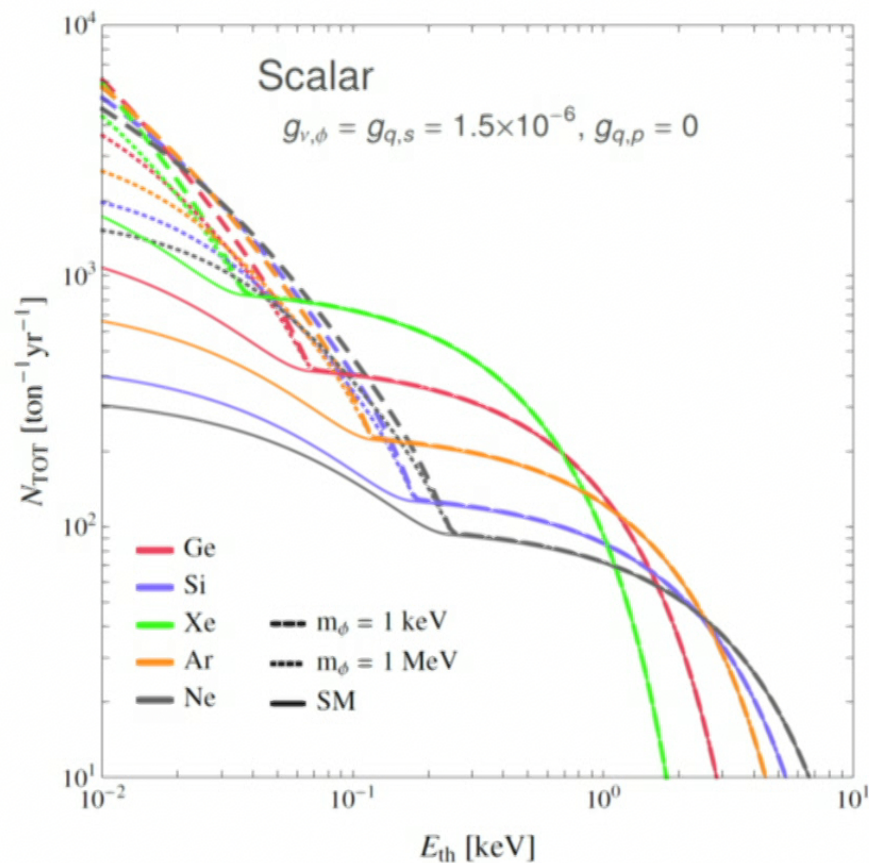
# Tests of BSM Physics

arXiv:1604.01025

$$(g_{\nu,\phi} \phi \bar{\nu}_R \nu_L + h.c.) + \phi \bar{l} g_{e,s} l + \phi \bar{q} g_{q,s} q$$



electron recoils



nuclear recoils

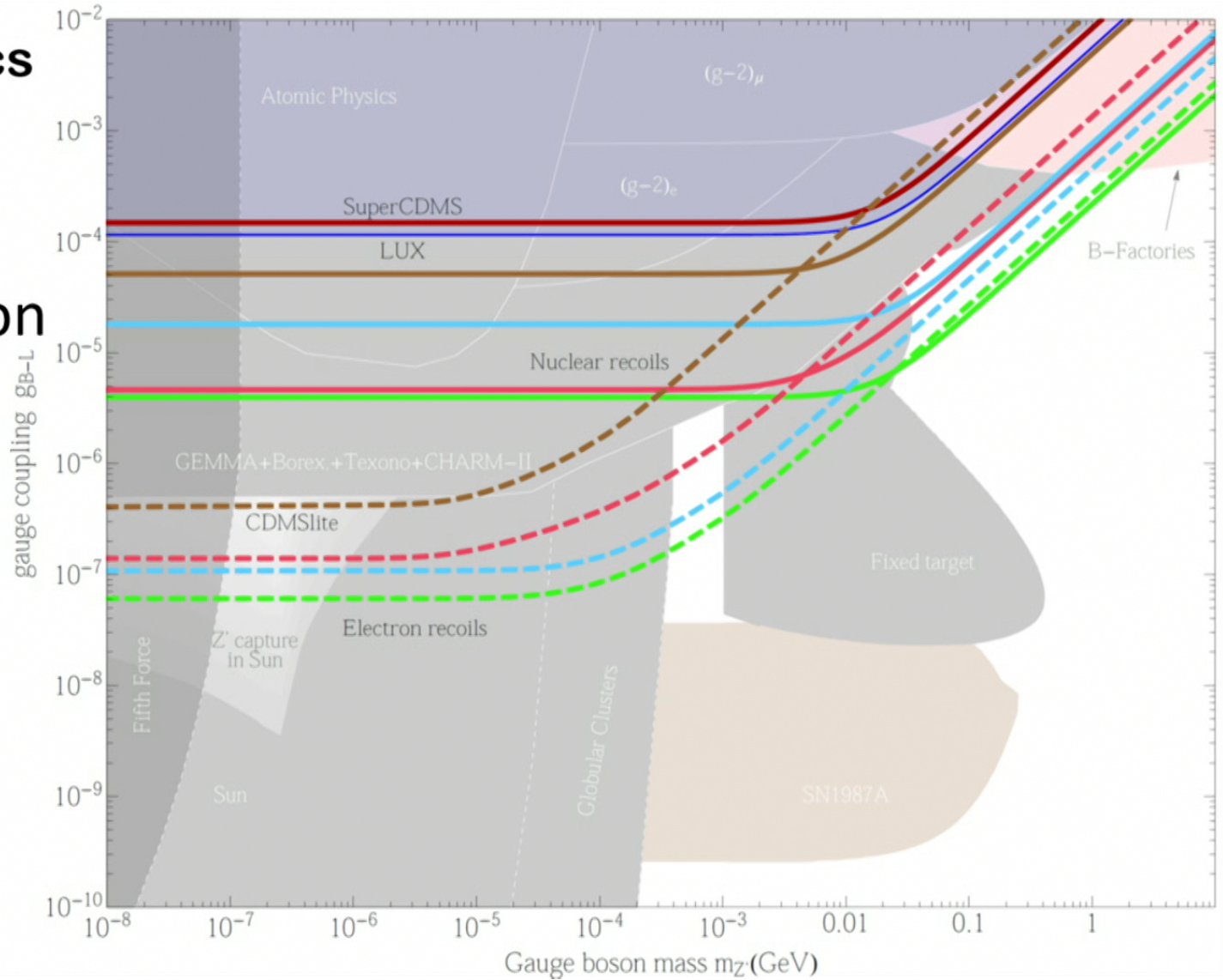
# Tests of BSM Physics

arXiv:1604.01025

## $U(1)_{B-L}$ gauge boson couples to B-L charge of SM particles

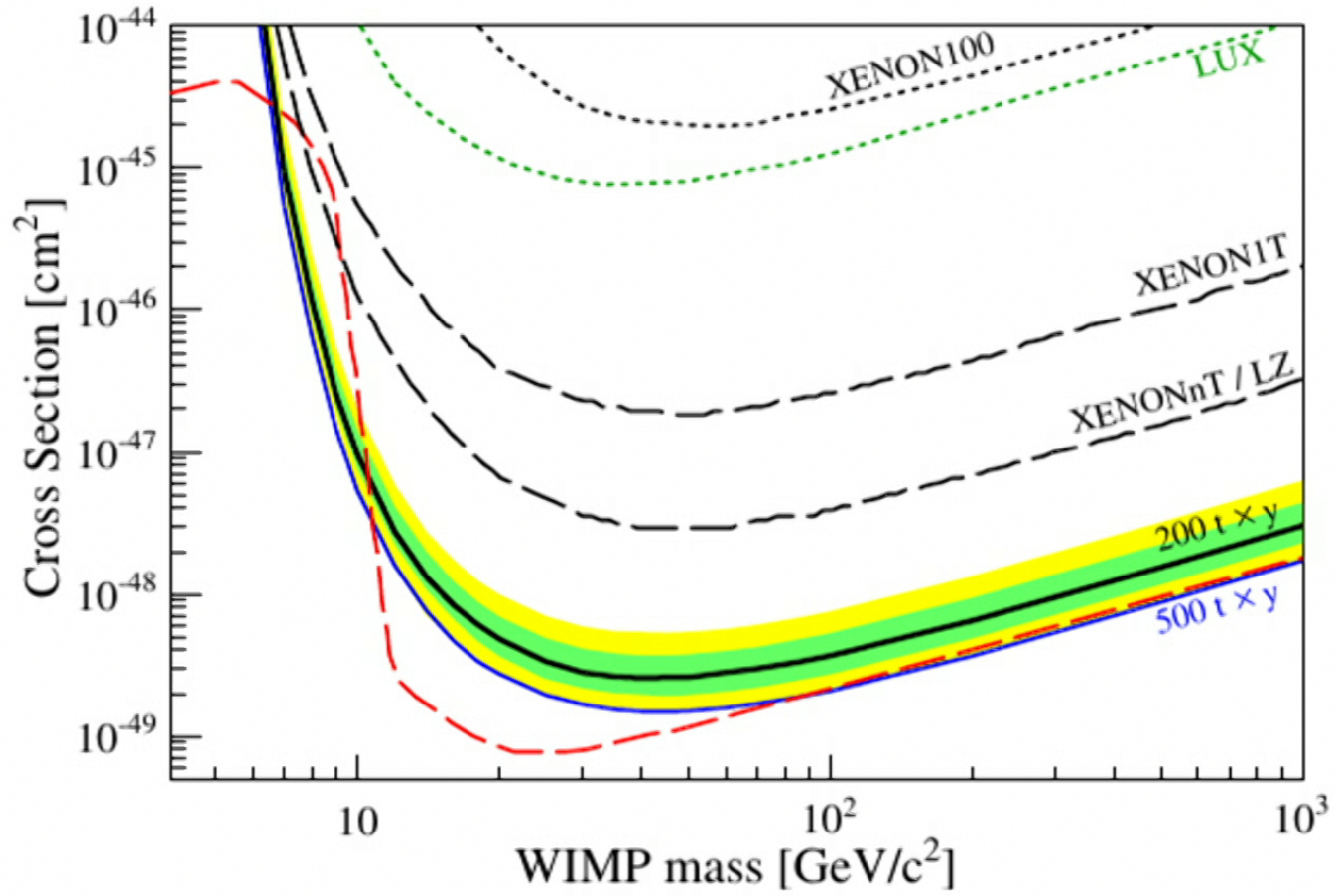
Dashed electron, solid nucleon.

Green future xenon  
 Blue G2 xenon  
 Red G2 germanium

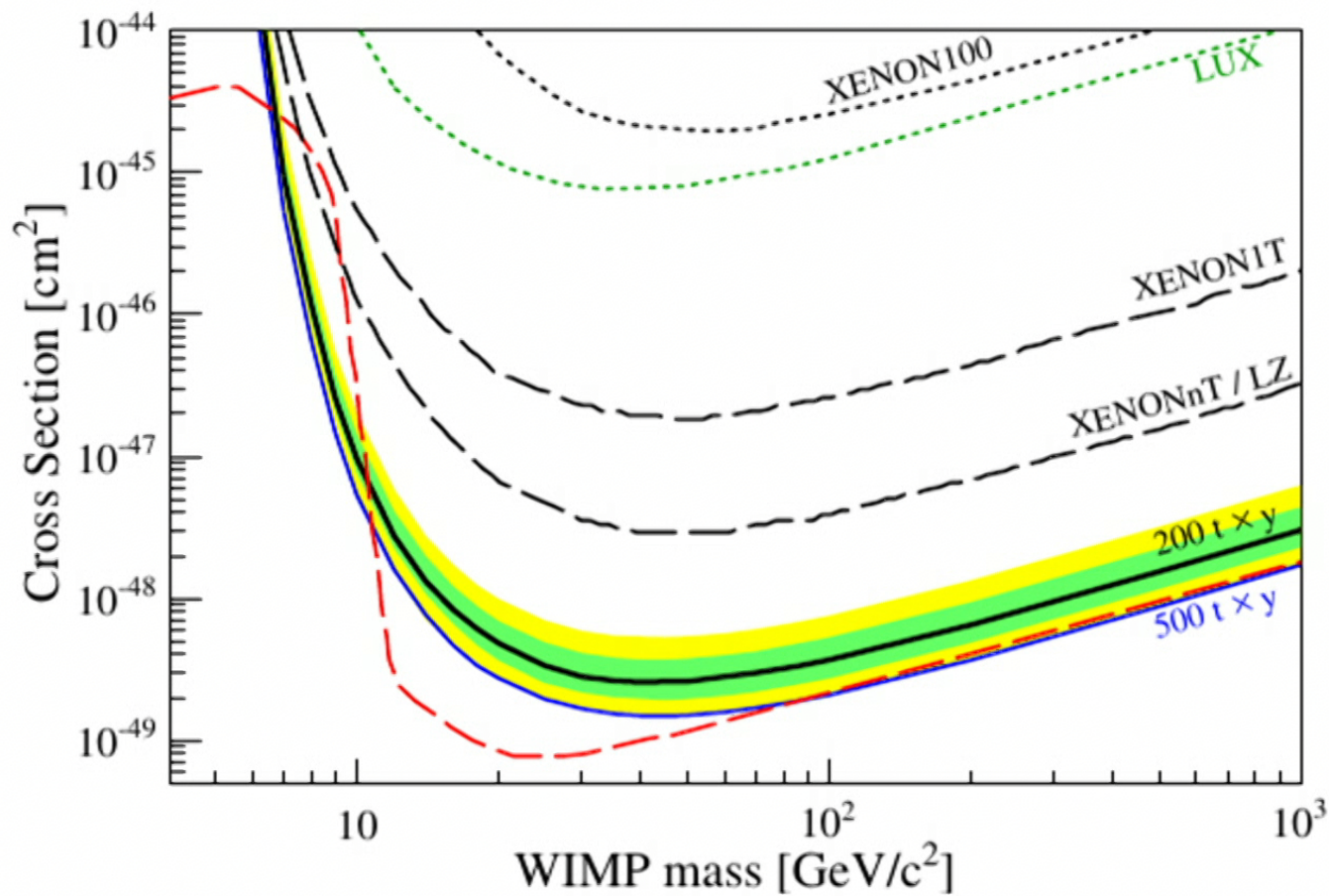




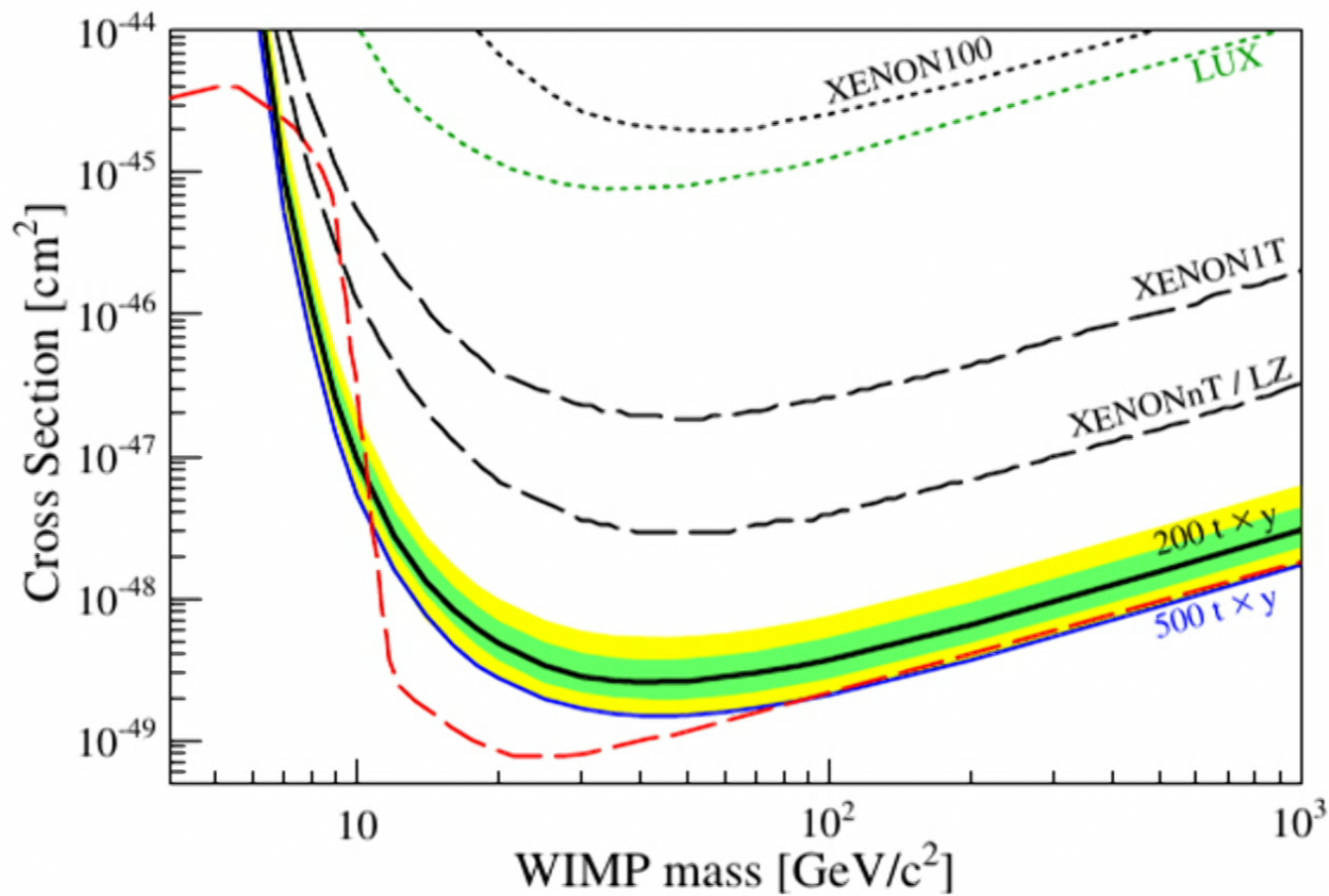
Doesn't seem possible to measure Diffuse Supernova Neutrino Background even with Darwin



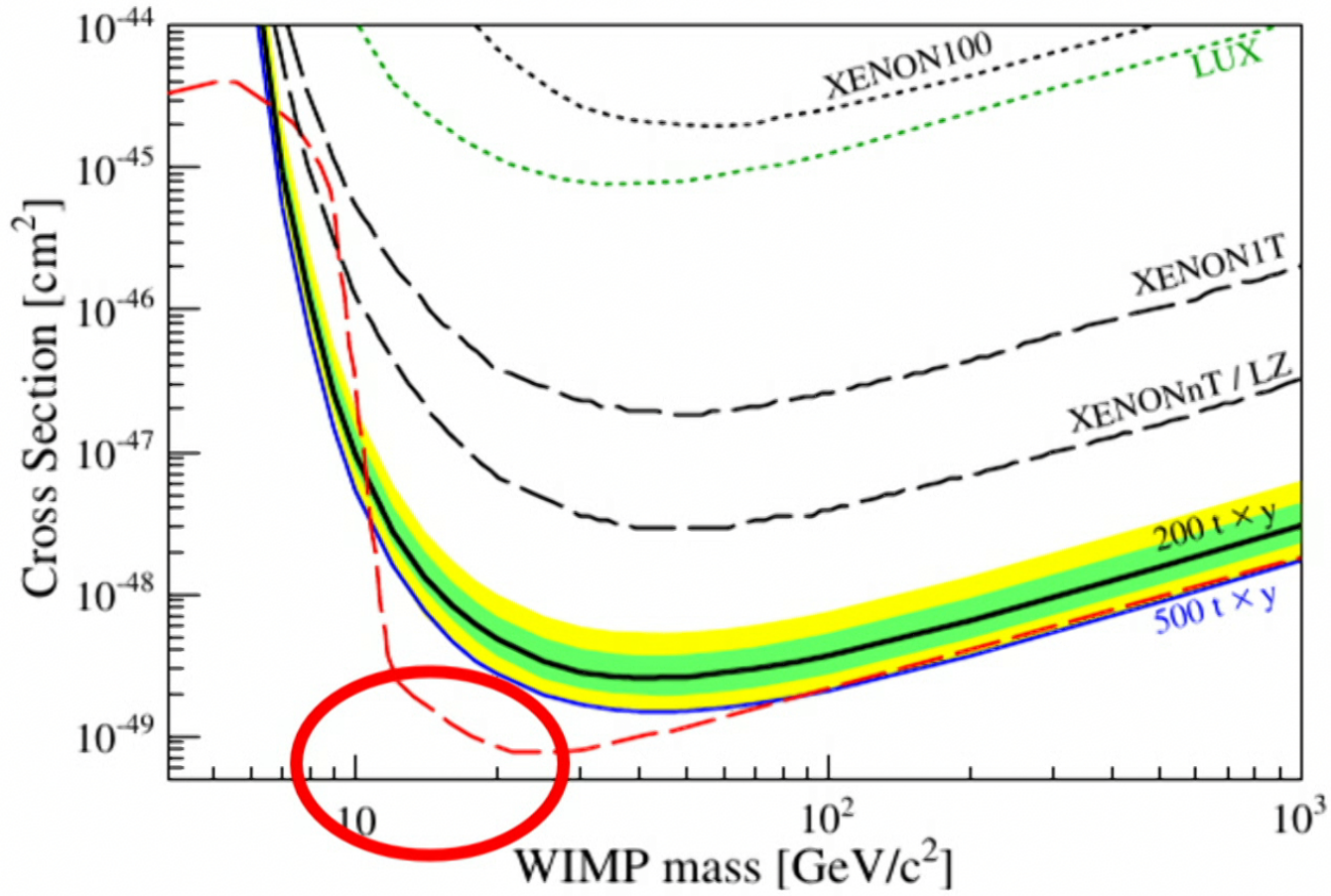
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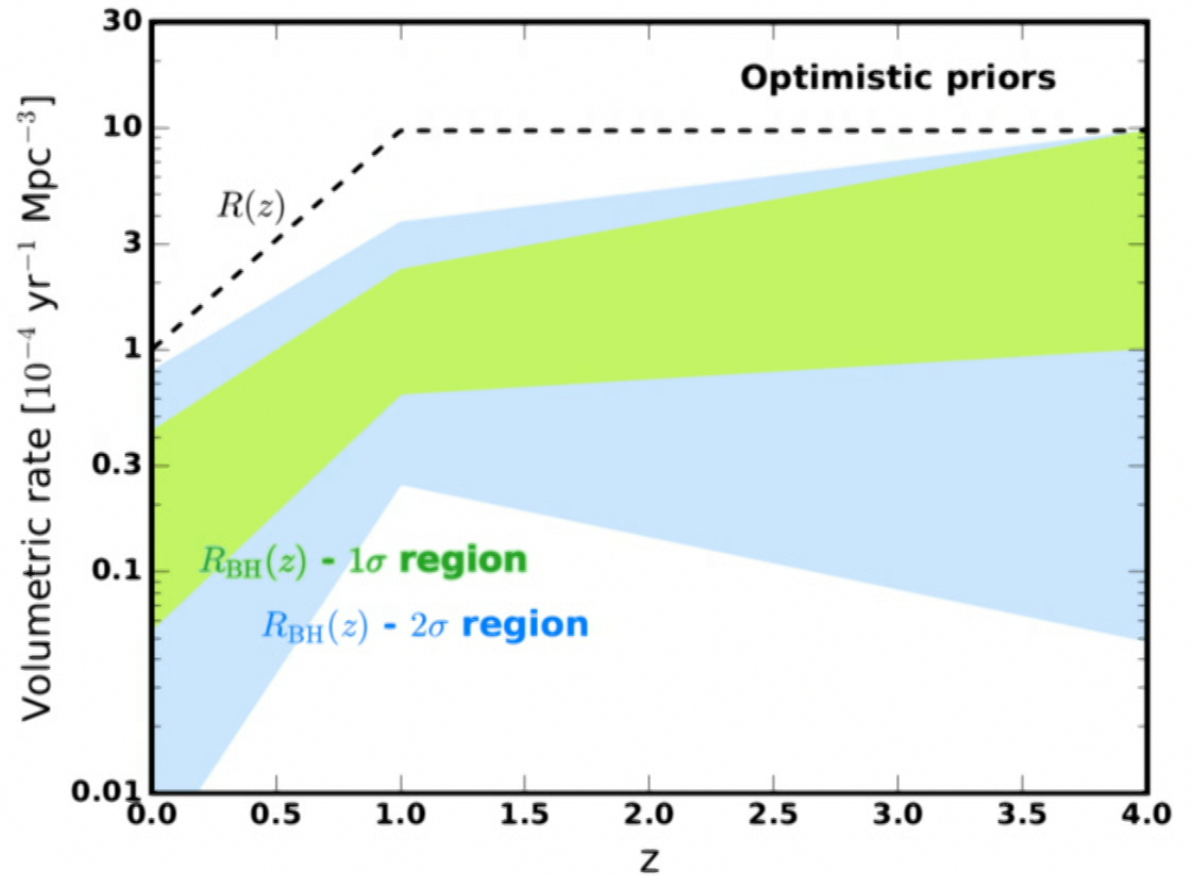


# Diffuse Supernova Neutrino Background

Can hopefully be seen with Super-K (gadzoos) and Hyper-K and probably DUNE.

**Difficult** with Darwin Xenon or Darkside because neutrinos too low energy.

BH Formation Rate from DNSB at HyperK Davis and Fairbairn 1704.05073

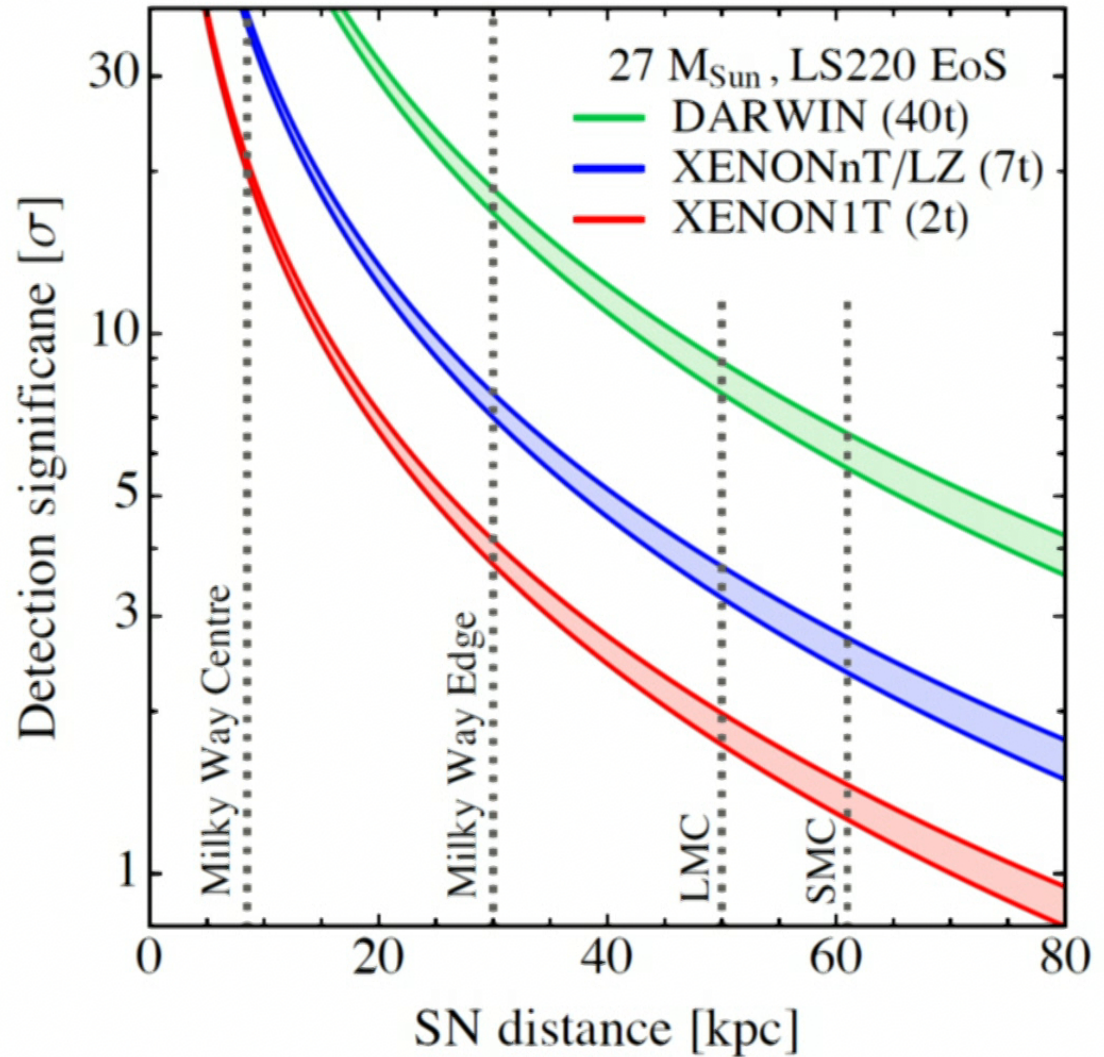


$$\Phi(E) = \frac{c}{H_0} \int_0^{z_{\max}} \frac{dz}{\sqrt{\Omega_m(1+z)^3 + \Omega_\Lambda}} \left[ R_{\text{NS}}(z) F_{\text{NS}}(E(1+z); \bar{E}_{e\text{NS}}, \bar{E}_{x\text{NS}}, L_{e\text{NS}}, L_{x\text{NS}}) \right. \\ \left. + R_{\text{BH}}(z) F_{\text{BH}}(E(1+z); \bar{E}_{e\text{BH}}, \bar{E}_{x\text{BH}}, L_{e\text{BH}}, L_{x\text{BH}}) \right]$$

## Observing a supernova with DM detectors

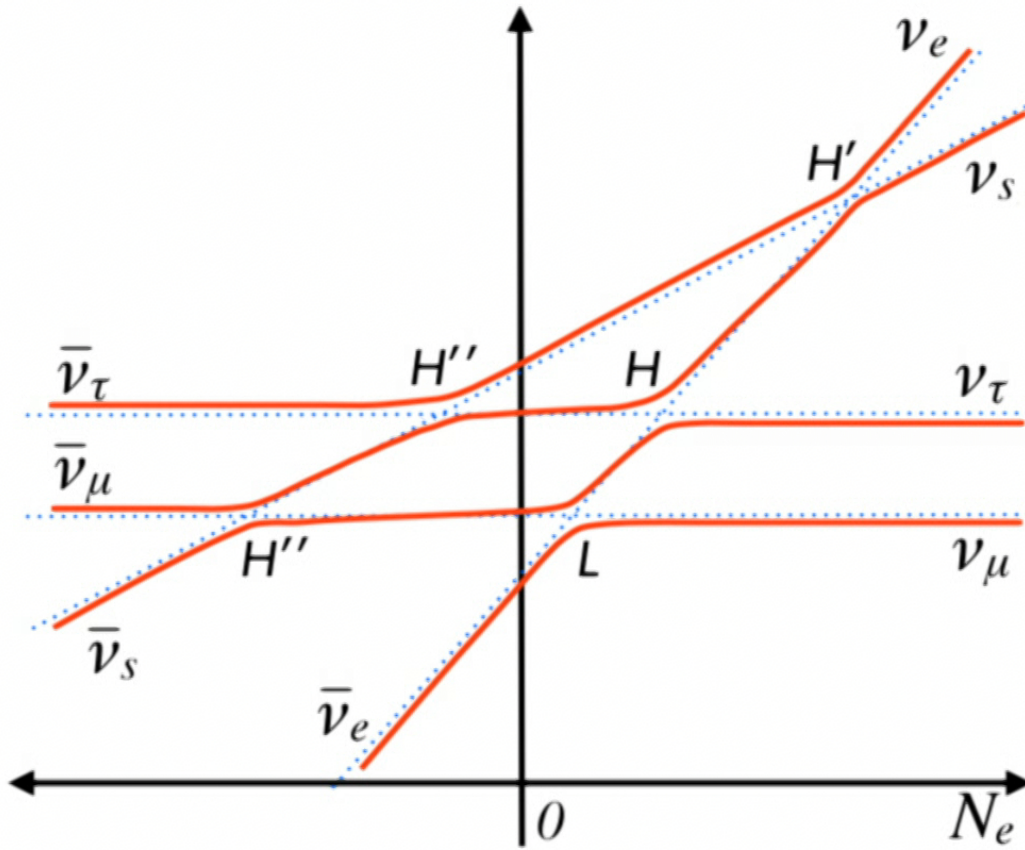
Should be easily possible  
(see e.g. Lang, McCabe et al 1606.09243)

Question is, can we do new  
fundamental physics with  
such observations?



## Mixing with Sterile Neutrino during Supernova Explosion

At the first resonance:



Normal hierarchy

$$\begin{pmatrix} F'_1 \\ F'_2 \\ F'_3 \\ F'_4 \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 - p_{H'} & p_{H'} \\ 0 & 0 & p_{H'} & 1 - p_{H'} \end{pmatrix} \begin{pmatrix} F_{\bar{\nu}_\mu}^0 \\ F_{\bar{\nu}_\tau}^0 \\ F_{\bar{\nu}_s}^0 \\ F_{\bar{\nu}_e}^0 \end{pmatrix}$$

Then:

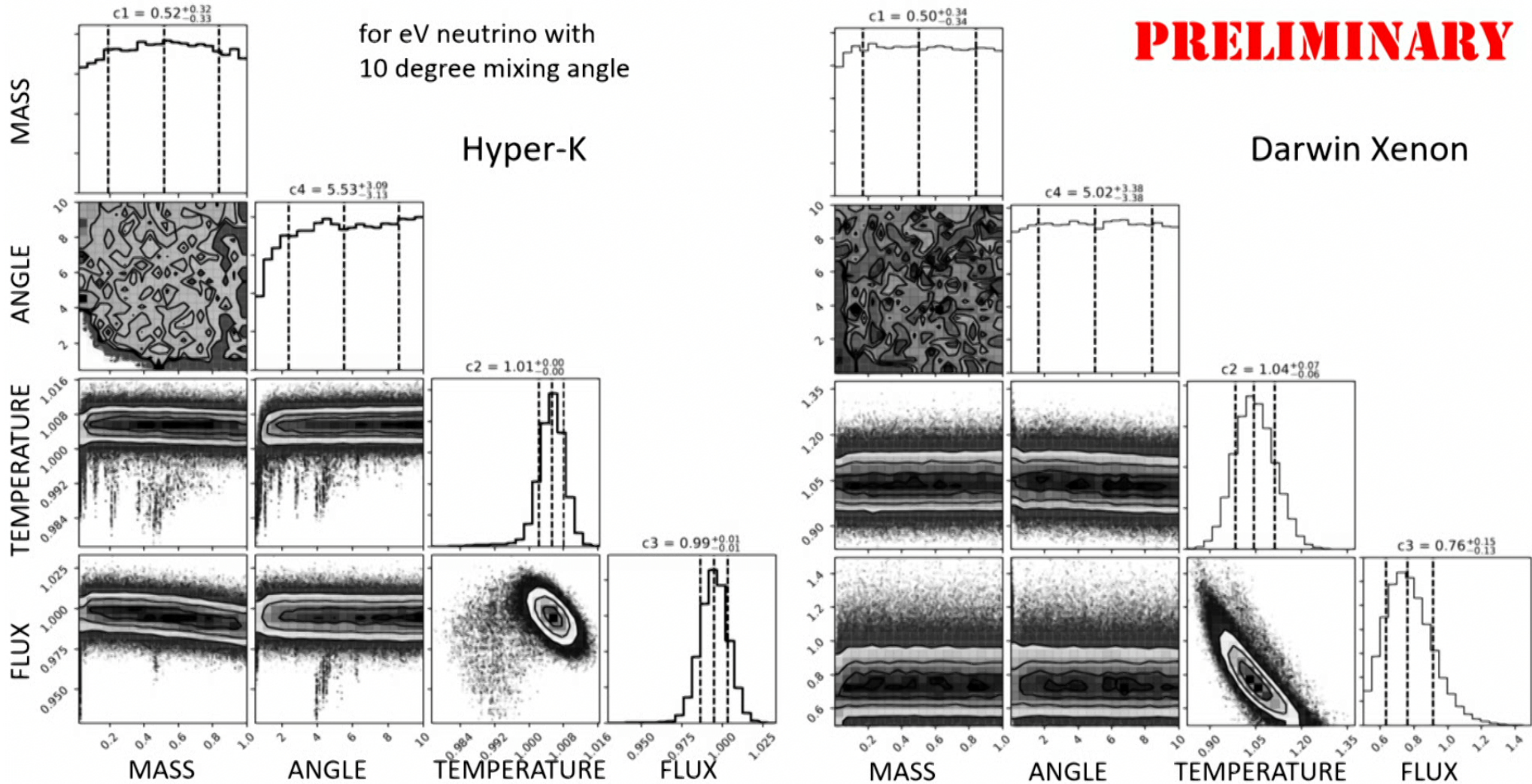
$$\begin{pmatrix} F''_1 \\ F''_2 \\ F''_3 \\ F''_4 \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 - p_H & p_H & 0 \\ 0 & p_H & 1 - p_H & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} F'_1 \\ F'_2 \\ F'_3 \\ F'_4 \end{pmatrix}$$

Finally:

$$\begin{pmatrix} F_1 \\ F_2 \\ F_3 \\ F_4 \end{pmatrix} = \begin{pmatrix} 1 - p_L & p_L & 0 & 0 \\ p_L & 1 - p_L & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} F''_1 \\ F''_2 \\ F''_3 \\ F''_4 \end{pmatrix}$$

# Sensitivity of Hyper-K and Darwin to Sterile mixing from Supernova at 10 kpc

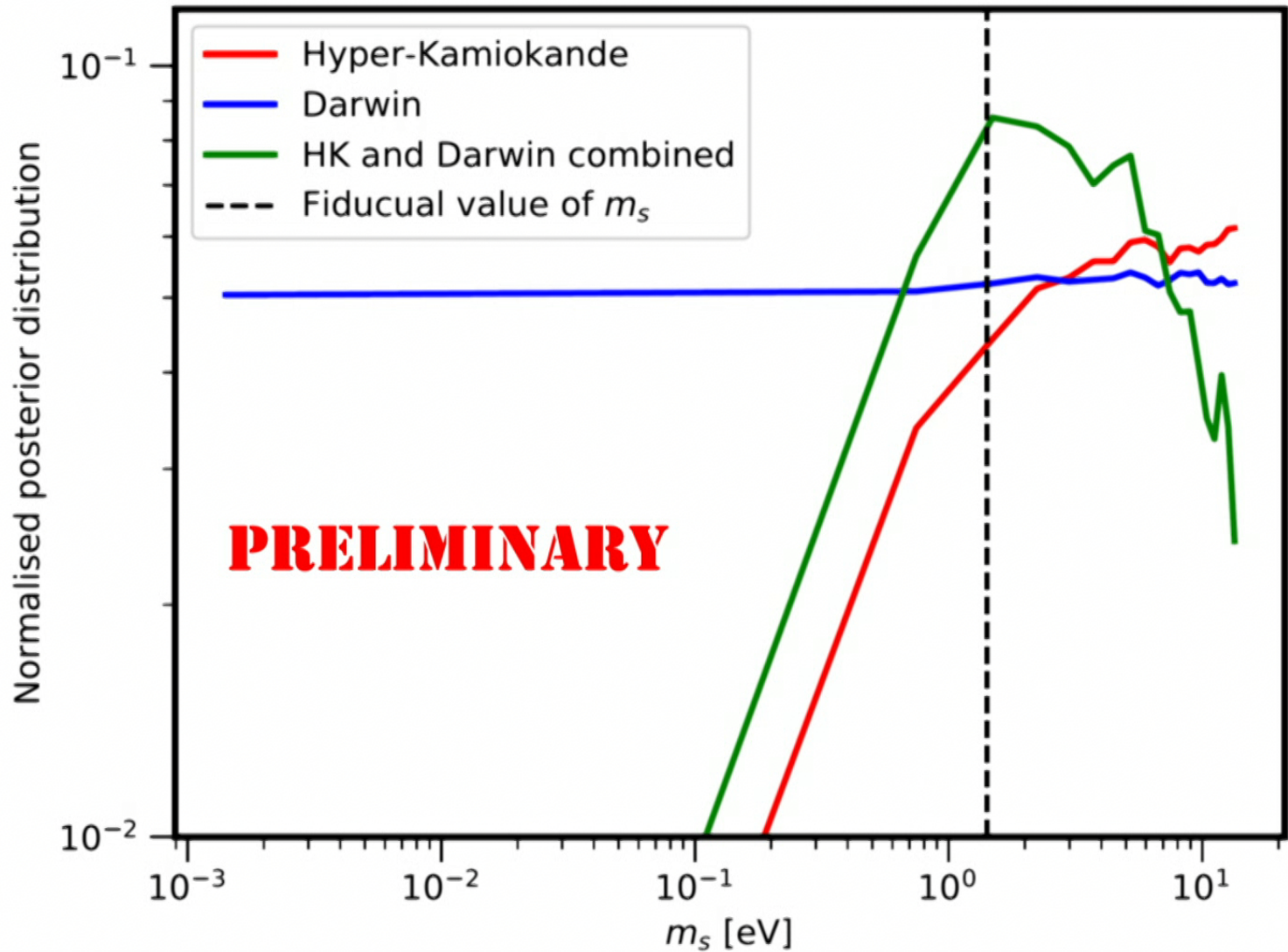
**PRELIMINARY**





Combined  
power of  
Hyper-K  
and  
Darwin Xenon

Added power  
of nuclear  
recoils very  
clear





**King's College London**, most central London University

Currently looking for strong fellowship candidates for five year royal society and STFC fellowships some of which are suitable for flexible working conditions

In addition, recruiting one or two regular postdocs in the Autumn



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Dark Matter detectors are about to detect neutrinos

This is very exciting, and will lead to new tests of solar physics and BSM physics.

This new subject area will continue to grow in the next years.

