

Title: Antiferroquadrupolar correlations in the quantum spin ice candidate Pr₂Zr₂O₇

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URL: <http://pirsa.org/17060053>

Abstract: We present an experimental study of the quantum spin ice candidate pyrochlore compound Pr₂Zr₂O₇ by means of magnetization measurements, specific heat and neutron scattering. We confirm that the spin excitation spectrum is essentially inelastic [1] and consists in a broad flat mode centered at about 0.4 meV with a magnetic structure factor which resembles the spin ice pattern. The new experimental results obtained under an applied magnetic field, interpreted in the light of mean field calculations, draw a new picture where quadrupolar interactions play a major role and overcome the magnetic exchange coupling. We determine a range of acceptable parameters able to account for the observations and propose that the actual ground state of this material is an antiferroquadrupolar liquid with spin-ice like excitations [2]. The influence of disorder is also discussed.

Antiferroquadrupolar correlations in the QSI $\text{Pr}_2\text{Zr}_2\text{O}_7$

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Our point:

1) Role of transverse interactions

PHYSICAL REVIEW B 94, 165153 (2016)

$$\mathcal{H} = \frac{1}{2} \sum_{\langle i,j \rangle} \mathcal{J}^{zz} \sigma_i^z \sigma_j^z + \sum_i (g_{\parallel} \mu_B \vec{z}_i \cdot \vec{h}) \sigma_i^z$$
$$+ \frac{1}{2} \sum_{\langle i,j \rangle} -\mathcal{J}^{\pm} (\sigma_i^+ \sigma_j^- + \sigma_i^- \sigma_j^+)$$
$$+ \frac{1}{2} \sum_{\langle i,j \rangle} \mathcal{J}^{\pm\pm} (\gamma_{ij} \sigma_i^+ \sigma_j^+ + \gamma_{ij}^* \sigma_i^- \sigma_j^-)$$

Jeff's talk
S. Onoda et al,
S.B. Lee et al

2) Role of disorder : « Quantum SI route »

PHYSICAL REVIEW B 94, 134428 (2016)

In preparation

$$\mathcal{H}_{\text{m-el}} = \sum_i v_i \sigma_i^+ + v_i^* \sigma_i^-$$

Spin Freezing in the Pyrochlore Antiferromagnet $\text{Pr}_2\text{Zr}_2\text{O}_7$

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Hinatsu⁴, T. Kitazawa⁵, Y. Kiuchi⁵, Z. Hiroi⁵, S. Takagi¹

ARTICLE

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DOI: 10.1038/ncomms2914

Quantum fluctuations in spin-ice-like $\text{Pr}_2\text{Zr}_2\text{O}_7$

K. Kimura¹, S. Nakatsuji^{1,2}, J.-J. Wen³, C. Broholm^{3,4,5}, M.B. Stone⁵, E. Nishibori⁶ & H. Sawa⁶

PHYSICAL REVIEW B **94**, 134428 (2016)

Magnetic properties and crystal field in $\text{Pr}_2\text{Zr}_2\text{O}_7$

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PHYSICAL REVIEW B 94, 165153 (2016)

Antiferroquadrupolar correlations in the quantum spin ice candidate $\text{Pr}_2\text{Zr}_2\text{O}_7$

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C. Decorse,⁶ M. Ciomaga Hatnean,⁷ and G. Balakrishnan⁷

Materials Research
Express

Structural and magnetic properties of single-crystals
of the geometrically frustrated zirconium pyrochlore,
 $\text{Pr}_2\text{Zr}_2\text{O}_7$

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D S Keeble¹ and G Balakrishnan¹

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PRL 118, 107206 (2017)

PHYSICAL REVIEW LETTERS

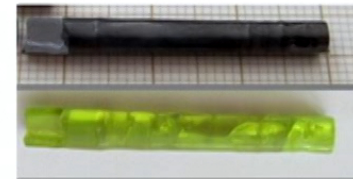
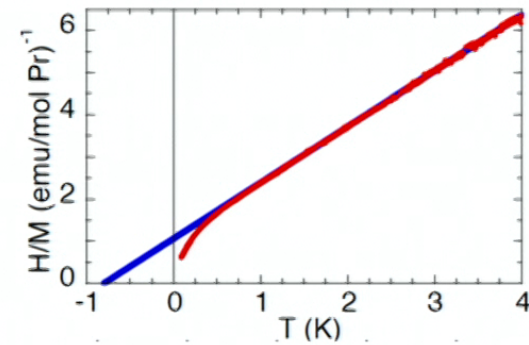
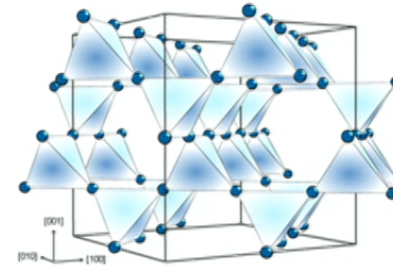
week ending
10 MARCH 2017

Disordered Route to the Coulomb Quantum Spin Liquid: Random Transverse Fields on Spin Ice in $\text{Pr}_2\text{Zr}_2\text{O}_7$

J.-J. Wen,^{1,2,3} S. M. Koohpayeh,¹ K. A. Ross,^{1,4} B. A. Trump,⁵ T. M. McQueen,^{1,5,6} K. Kimura,^{7,8} S. Nakatsuji,^{7,9}
Y. Qiu,⁴ D. M. Pajerowski,⁴ J. R. D. Copley,⁴ and C. L. Broholm^{1,4,6}

Introduction

- A pyrochlore magnet
- Pr^{3+} is a NK ion
- Its CEF ground state is a doublet (spanning the spin $\frac{1}{2}$ states) well protected from the 1st excited state (10 meV)
- No long range ordering down to 60 mK
- AF interactions, as inferred from $\chi(T)$
- Crystal growth is an issue, role of disorder outlined long ago



Introduction

- C_p shows a peak at 2K, followed by an upturn at low T (interpreted in terms of monopoles dynamics)
- The neutron signal is mostly inelastic, with a lorentzian shape

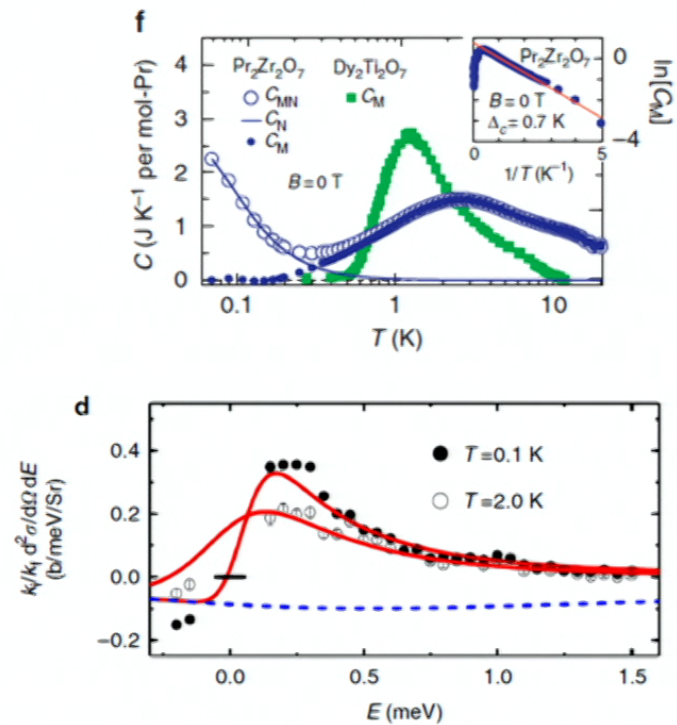
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Quantum fluctuations in spin-ice-like $\text{Pr}_2\text{Zr}_2\text{O}_7$

K. Kimura¹, S. Nakatsuji^{1,2}, J.-J. Wen³, C. Broholm^{3,4,5}, M.B. Stone⁵, E. Nishibori⁶ & H. Sawa⁶



Introduction

- C_p shows a peak at 2K, followed by an upturn at low T (interpreted in terms of monopoles dynamics)
- The neutron signal is mostly inelastic, with a quite **broad** lorentzian shape
- Described as a « dynamical spin ice » (the INS response resembles the spin ice pattern)

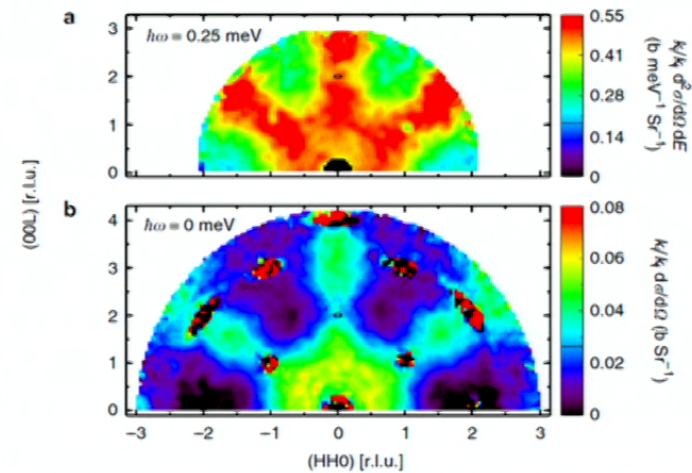
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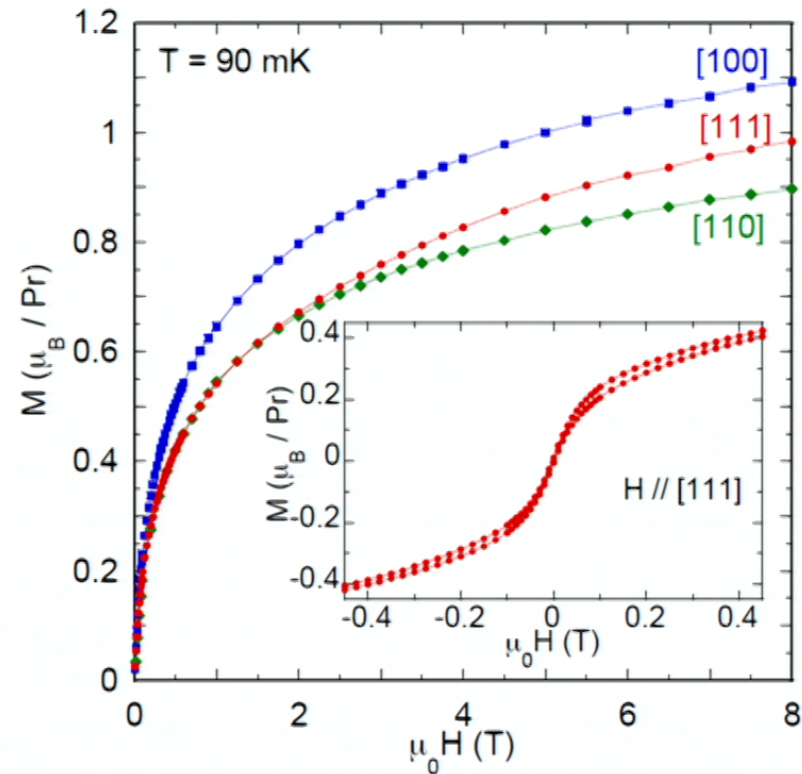
Quantum fluctuations in spin-ice-like $\text{Pr}_2\text{Zr}_2\text{O}_7$

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Magnetization

- $M(H)$ struggles to grow with increasing field
- Different from the expected Ising behaviour
- Something fights against the growth of the magnetization

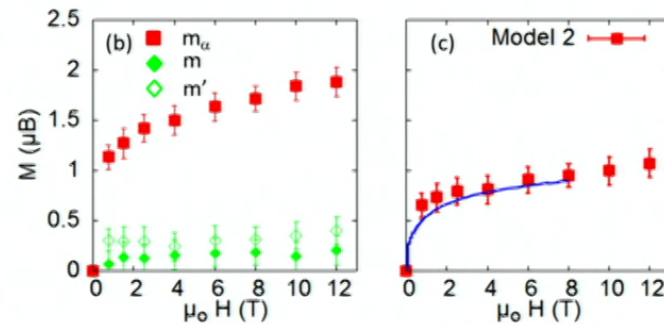
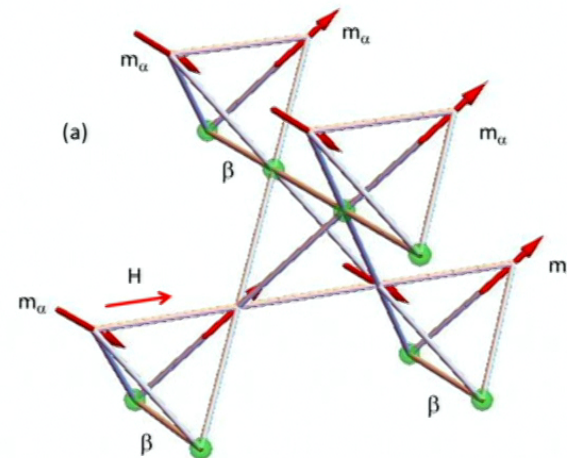


Neutron diffraction

D23@ILL

$H // (1 -1 0)$

- Under applied field, no magnetic moment in the β chains inferred from neutron diffraction
- Smooth increase with H , consistent with $M(H)$

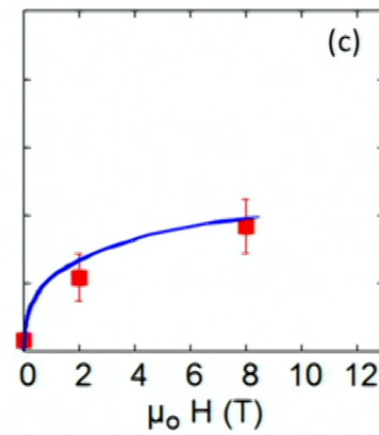
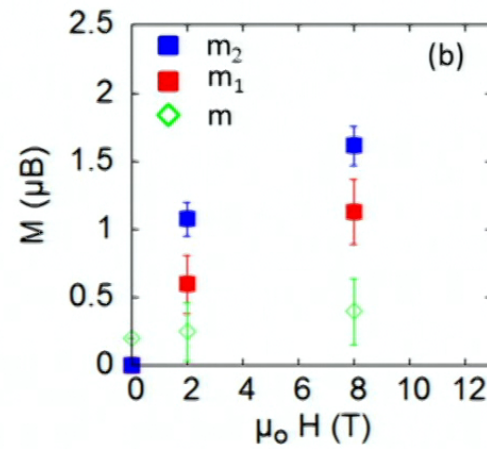
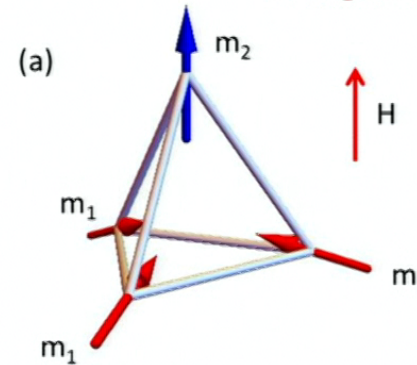


Neutron diffraction

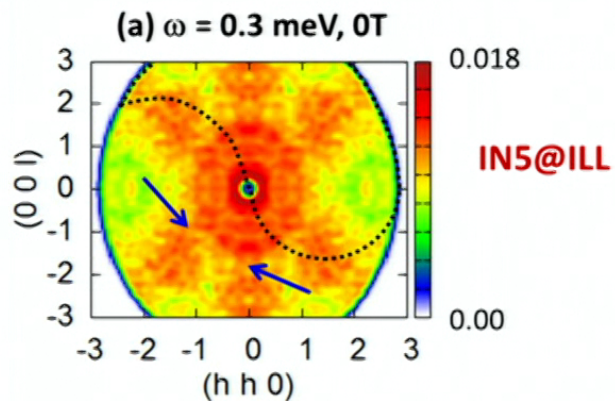
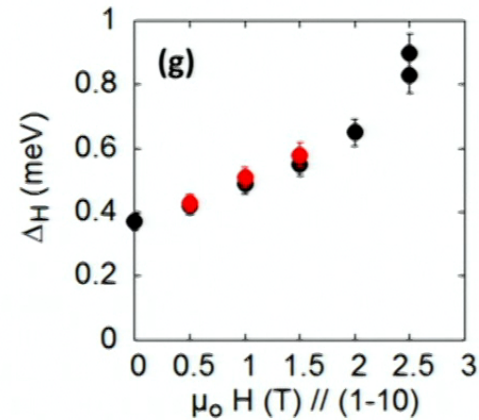
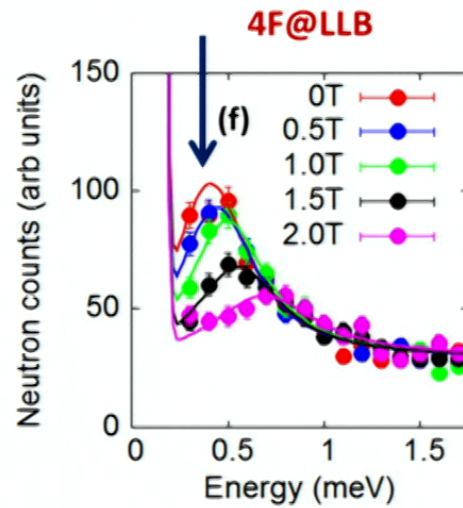
$H // (111)$

Smooth increase with H , consistent with $M(H)$

D23@ILL

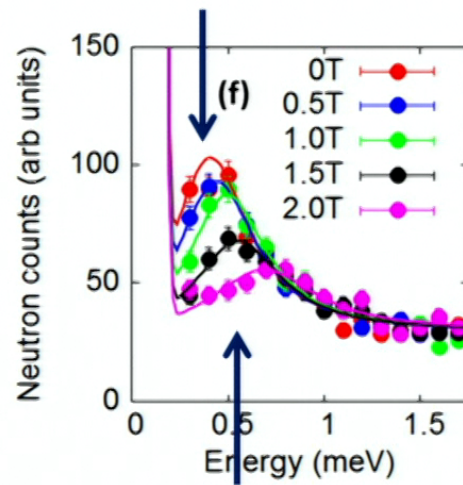


Inelastic neutron scattering

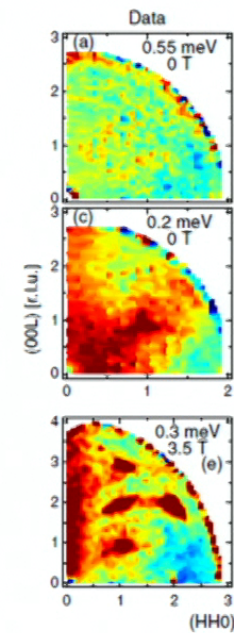
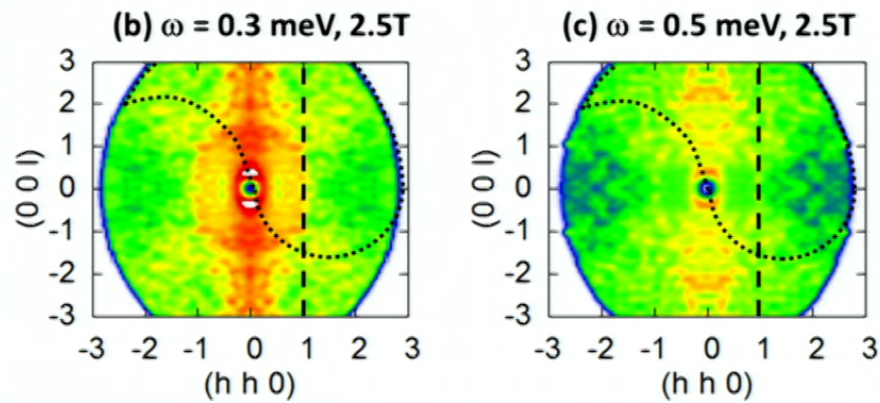


- Under applied field $H // (1-1-0)$, the peak of the lorentzian profile shifts to higher energy
- Confirm the dynamical spin ice pattern

Inelastic neutron scattering



PRL 118, 107206 (2017)



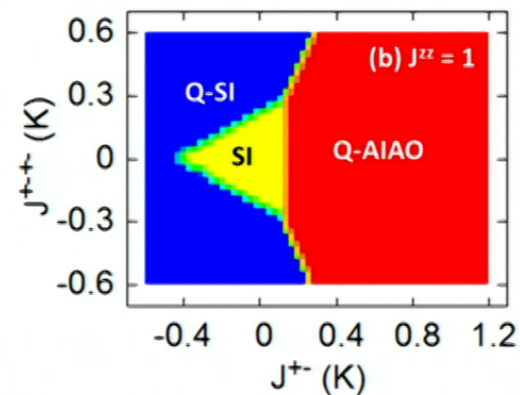
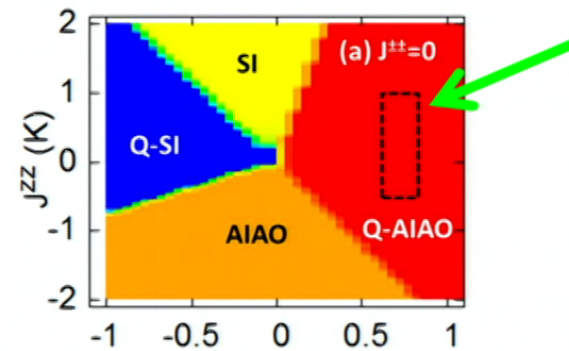
Input of MF

$$\mathcal{H} = \frac{1}{2} \sum_{\langle i,j \rangle} \mathcal{J}^{zz} \sigma_i^z \sigma_j^z + \sum_i (g_{\parallel} \mu_B \vec{z}_i \cdot \vec{h}) \sigma_i^z$$

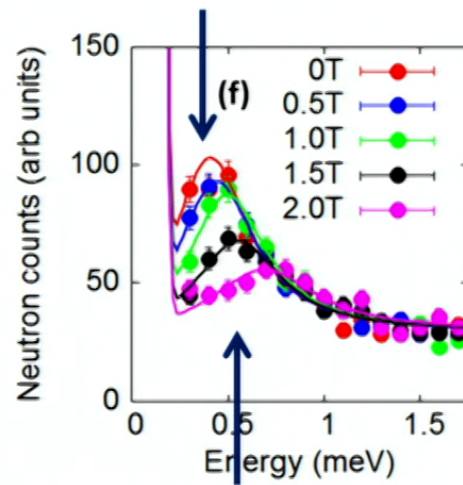
$$+ \frac{1}{2} \sum_{\langle i,j \rangle} -\mathcal{J}^{\pm} (\sigma_i^+ \sigma_j^- + \sigma_i^- \sigma_j^+)$$

$$+ \frac{1}{2} \sum_{\langle i,j \rangle} \mathcal{J}^{\pm\pm} (\gamma_{ij} \sigma_i^+ \sigma_j^+ + \gamma_{ij}^* \sigma_i^- \sigma_j^-)$$

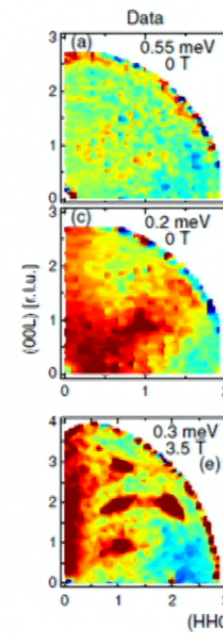
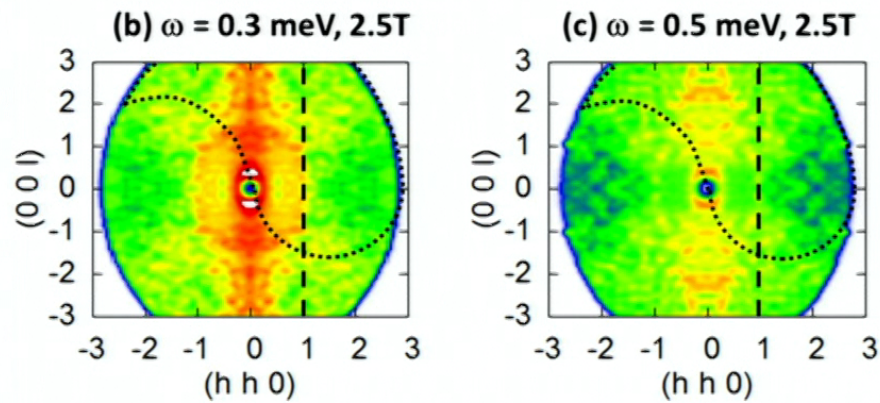
- Mean field, K=0 phases
- Pretty « similar » to the work by Lee et al although by far less sophisticated



Inelastic neutron scattering



PRL 118, 107206 (2017)



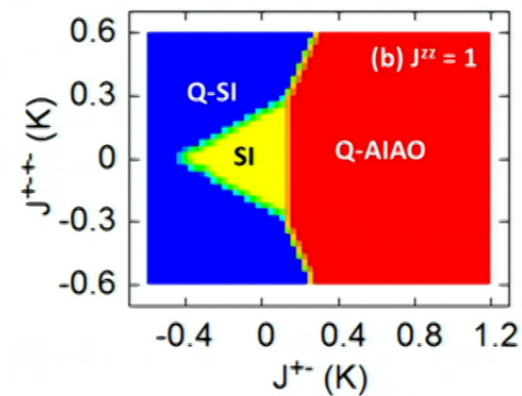
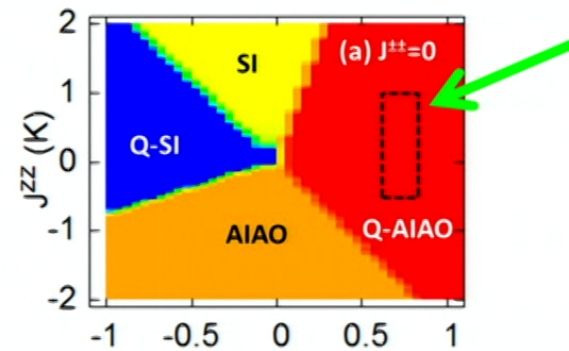
Input of MF

$$\mathcal{H} = \frac{1}{2} \sum_{\langle i,j \rangle} \mathcal{J}^{zz} \sigma_i^z \sigma_j^z + \sum_i (g_{\parallel} \mu_B \vec{z}_i \cdot \vec{h}) \sigma_i^z$$

$$+ \frac{1}{2} \sum_{\langle i,j \rangle} -\mathcal{J}^{\pm} (\sigma_i^+ \sigma_j^- + \sigma_i^- \sigma_j^+)$$

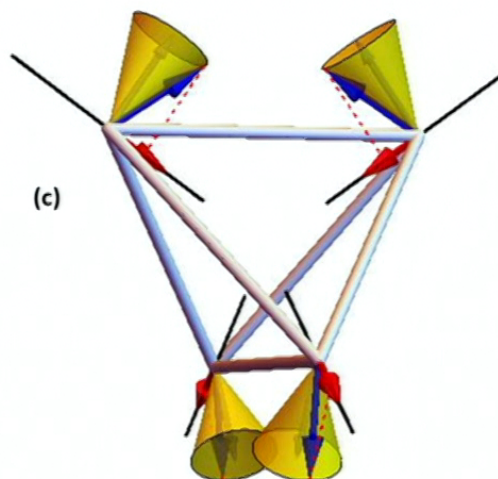
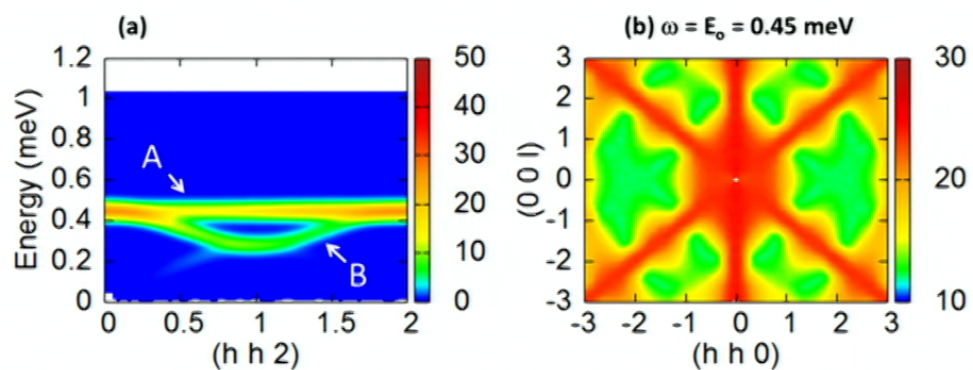
$$+ \frac{1}{2} \sum_{\langle i,j \rangle} \mathcal{J}^{\pm\pm} (\gamma_{ij} \sigma_i^+ \sigma_j^+ + \gamma_{ij}^* \sigma_i^- \sigma_j^-)$$

- Mean field, K=0 phases
- Pretty « similar » to the work by Lee et al although by far less sophisticated

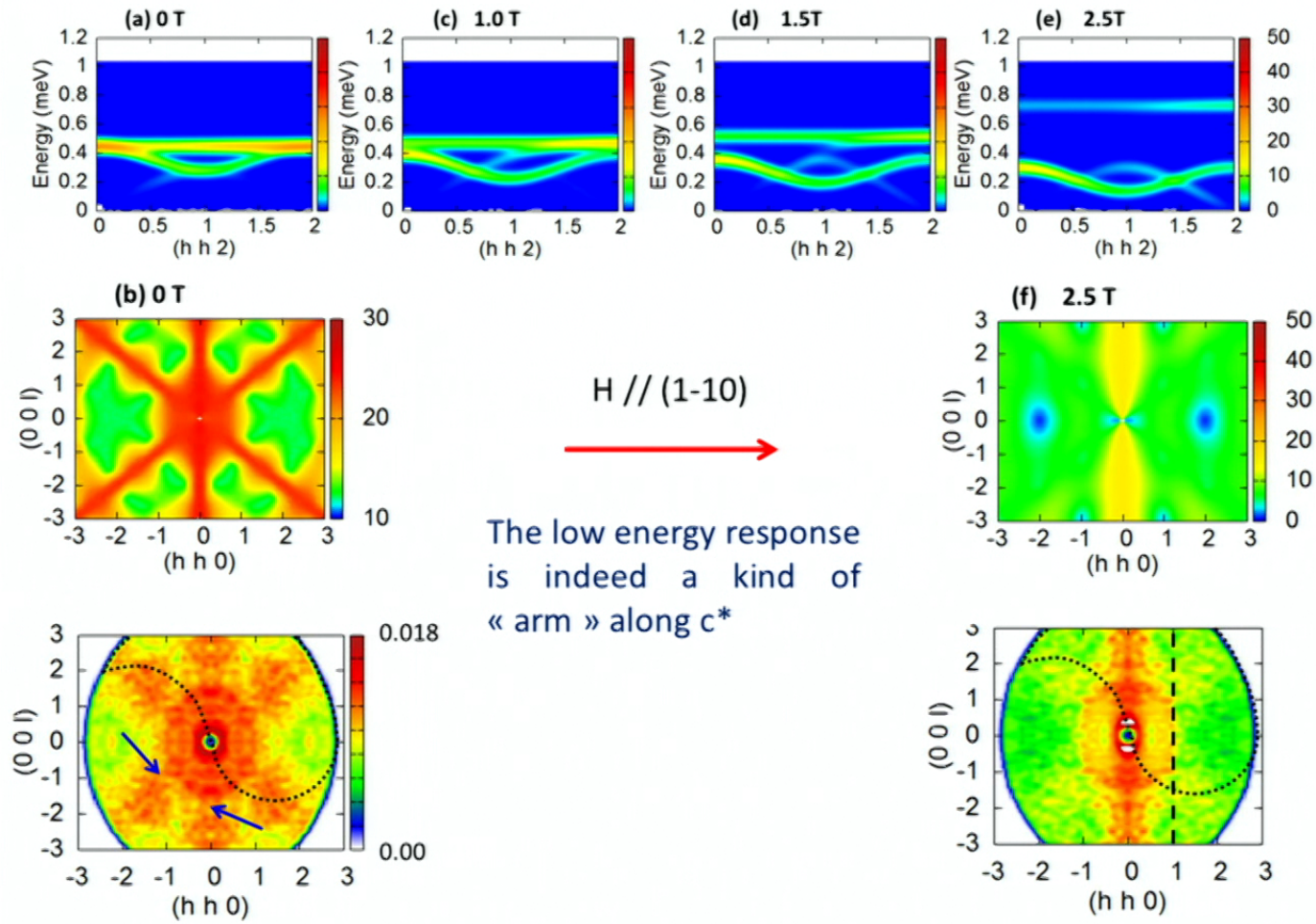


Spin and pseudo spin dynamics

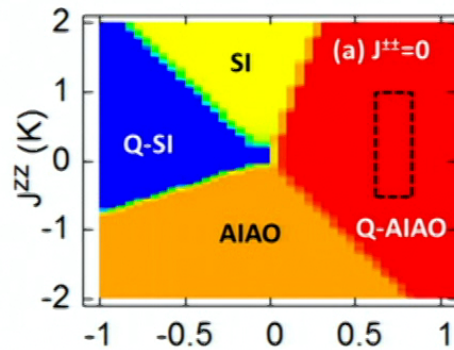
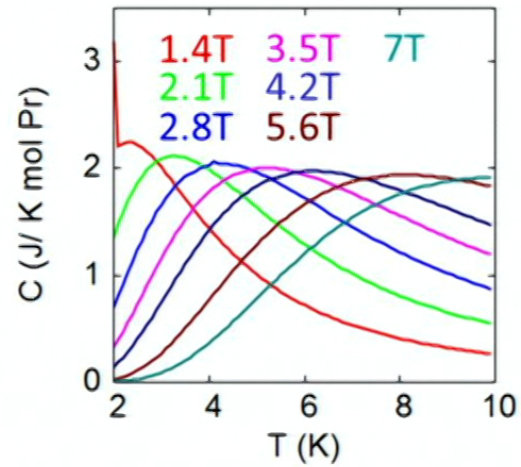
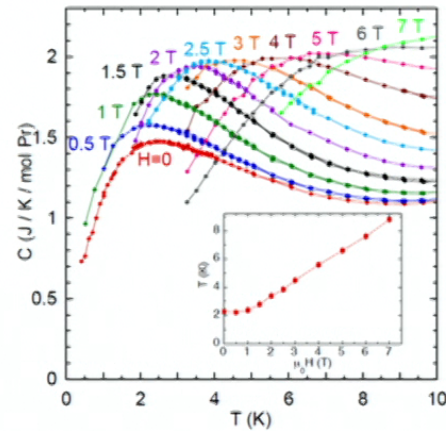
- The RPA pseudo spin dynamics consists of
 - (A) a flat mode
 - (B) a dispersive feature
- (A) is characterized by the spin ice pattern
- The 2-in 2-out **magnetic** nature of excitation (A) arises from the peculiar projection of the corresponding « pseudo-spin » onto the z (magnetic) axis



Spin and pseudo spin dynamics



Specific heat



Based on the analysis of diffraction, INS and C_p , we propose:

$$0.7 \leq J^{\pm} \leq 0.8 \text{ K}$$

$$-0.5 \leq J^{zz} \leq 1 \text{ K}$$

Discussion: input of MF

BUT

- *THE MEAN FIELD APPROACH IS CLEARLY NOT SATISFYING SINCE IT PREDICTS LRO*
- *OTHER INGREDIENT ? → DISORDER, EXTREMELY IMPORTANT IN NK MAGNETS*

PHYSICAL REVIEW B **94**, 134428 (2016)

A. Local imperfections of the crystal field

In order to model local imperfections of the crystal field, we introduced a nontrigonal component of the crystal field (1) of the type $B_2^2 C_2^2$. Here, B_2^2 represents the magnitude of a strain coupled to one quadrupole operator (C_2^2). It would be more

See also

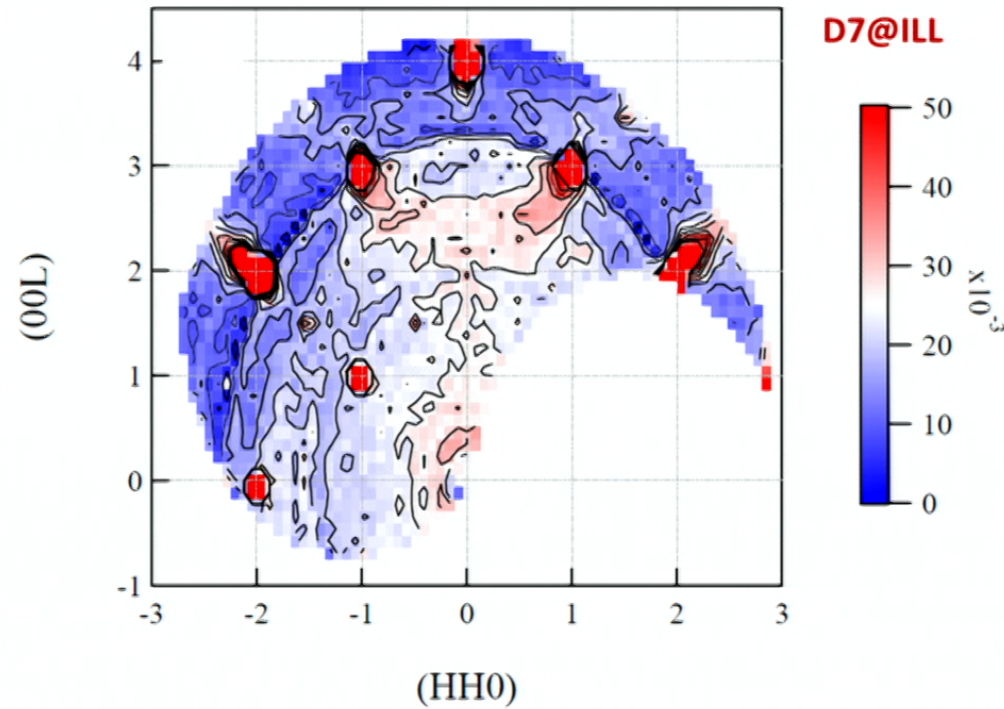
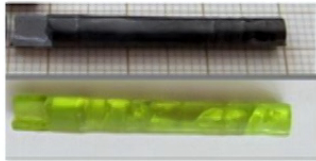
Duijn et al, *PRL* 94 177201 (2005)

Foronda, *PRL* 114, 017602 (2015)

Wen, *PRL* 118, 107206 (2017)

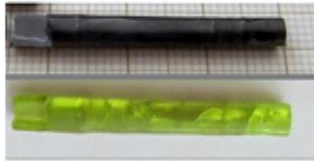
Savary, *PRL* 118, 087203 (2017)

Diffuse scattering reveals disorder



- diffuse scattering measured by polarized neutron diffraction @D7 (ILL) → lattice origin
- Intrinsic defects in the material

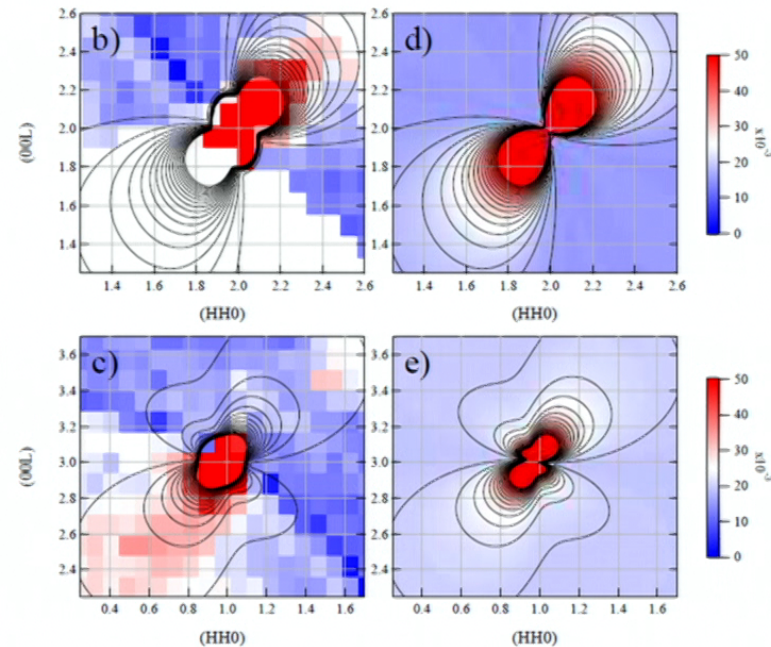
Diffuse scattering reveals disorder



- Well represented by so called Huang scattering with lobes and « pinch points » !
- This would correspond to point defects creating a random strain described by :

$$g(\mathbf{e}) = \frac{3}{4\pi^2} \frac{\gamma}{(e^2 + \gamma^2)^{5/2}},$$

- The parameter γ (reflects the strength of disorder) $\approx 10^{-4}$ to 10^{-5}



Random strain and transverse « field »

By virtue of the magneto-elastic coupling, the random strain perturbs the electronic density:



In the pseudo spin ½ language this introduces a random transverse field v

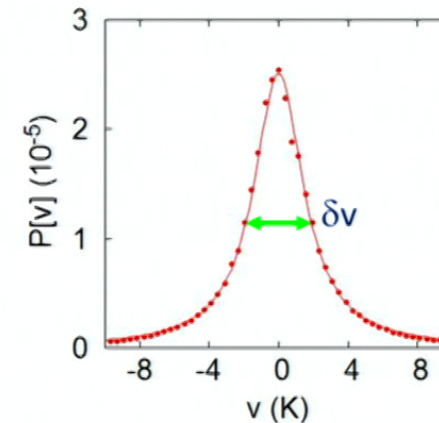
$$\begin{aligned} \mathcal{H}_{m-el} = & e_1(\Gamma_3^1)[B_{21}^{xx} Q_{xz} + B_{22}^{xx} Q_{x^2-y^2}] \\ & + e_2(\Gamma_3^1)[-B_{21}^{xx} Q_{yz} + B_{22}^{xx} Q_{xy}] \\ & + e_1(\Gamma_3^2)[B_{21}^{zx} Q_{xz} + B_{22}^{zx} Q_{x^2-y^2}] \\ & + e_2(\Gamma_3^2)[-B_{21}^{zx} Q_{yz} + B_{22}^{zx} Q_{xy}], \end{aligned}$$

$$g(\mathbf{e}) = \frac{3}{4\pi^2} \frac{\gamma}{(e^2 + \gamma^2)^{5/2}},$$

$$\mathcal{H}_{m-el} = \sum_i v_i \sigma_i^+ + v_i^* \sigma_i^-$$

$$\delta v \approx k\gamma$$

$$\gamma \approx 10^{-5} \text{ to } 10^{-4} \rightarrow \delta v \approx \text{few K}$$

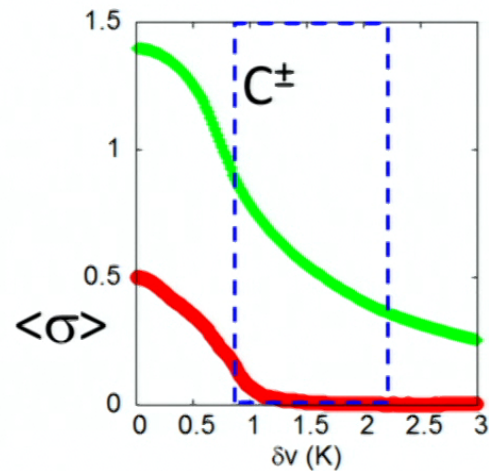


Input of MF

$$\mathcal{H} = \frac{1}{2} \sum_{\langle i,j \rangle} \mathcal{J}^{zz} \sigma_i^z \sigma_j^z + \frac{1}{2} \sum_{\langle i,j \rangle} -\mathcal{J}^{\pm} (\sigma_i^+ \sigma_j^- + \sigma_i^- \sigma_j^+)$$

$$\mathcal{H}_{\text{m-el}} = \sum_i v_i \sigma_i^+ + v_i^* \sigma_i^-$$

For the relevant value of J^{\pm} and J^{zz} ,
 δv progressively kills the AIAO LRO
 yet NN SI and/or Q correlations persist

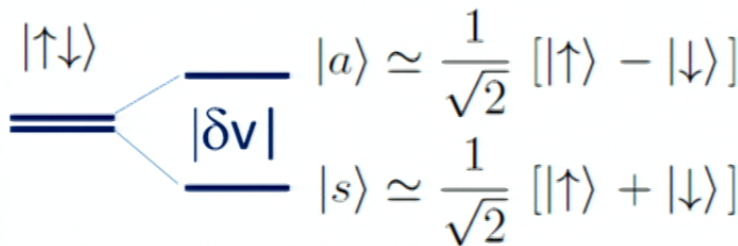


$$C^{zz} = \frac{1}{N} \sum_{i, \langle j \rangle_i} \langle \sigma_i^z \rangle \langle \sigma_j^z \rangle$$

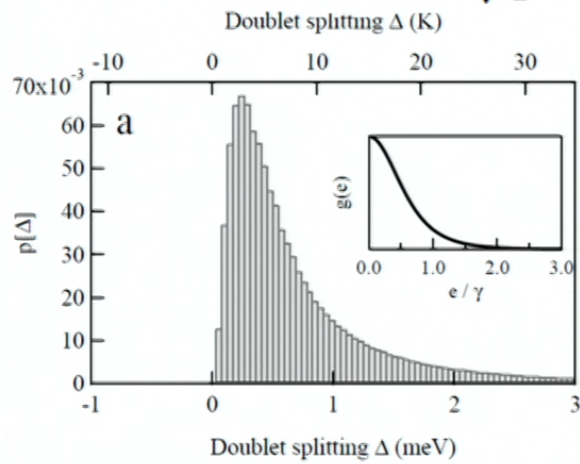
$$C^{\pm} = \frac{1}{N} \sum_{i, \langle j \rangle_i} \langle \sigma_i^+ \rangle \langle \sigma_j^- \rangle$$

Real space mean field calculation
 (7x7x7 unit cells)

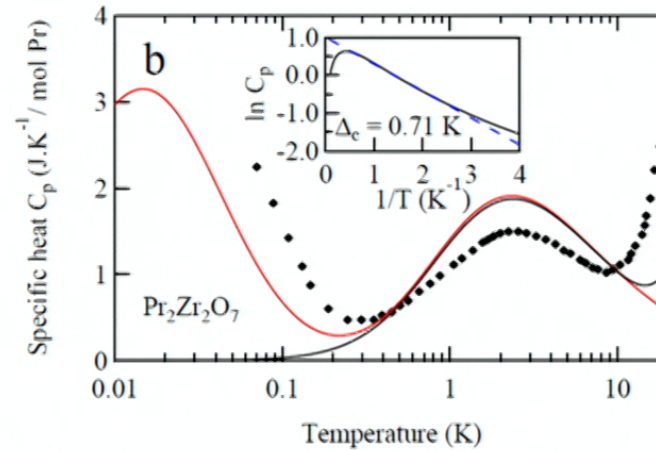
Input of MF



Correlated by J^\pm and J^{zz}



$$\gamma \approx 1.25 \times 10^{-4} \quad \text{or} \quad \delta v \approx 2\text{K}$$



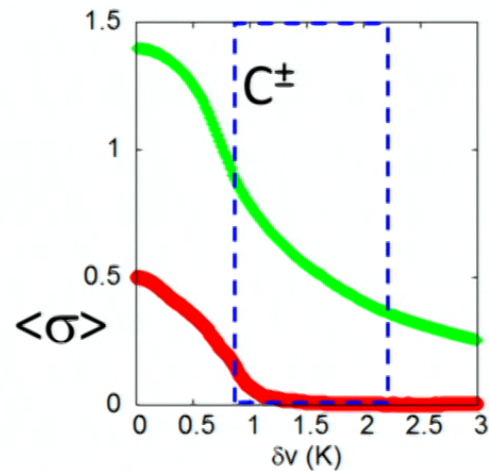
Distribution of Schottky peaks +
Hyperfine interaction at very low T

Input of MF

$$\mathcal{H} = \frac{1}{2} \sum_{\langle i,j \rangle} \mathcal{J}^{zz} \sigma_i^z \sigma_j^z + \frac{1}{2} \sum_{\langle i,j \rangle} -\mathcal{J}^{\pm} (\sigma_i^+ \sigma_j^- + \sigma_i^- \sigma_j^+)$$

$$\mathcal{H}_{\text{m-el}} = \sum_i v_i \sigma_i^+ + v_i^* \sigma_i^-$$

For the relevant value of \mathcal{J}^{\pm} and \mathcal{J}^{zz} ,
 δv progressively kills the AIAO LRO
 yet NN SI and/or Q correlations persist



$$C^{zz} = \frac{1}{N} \sum_{i, \langle j \rangle_i} \langle \sigma_i^z \rangle \langle \sigma_j^z \rangle$$

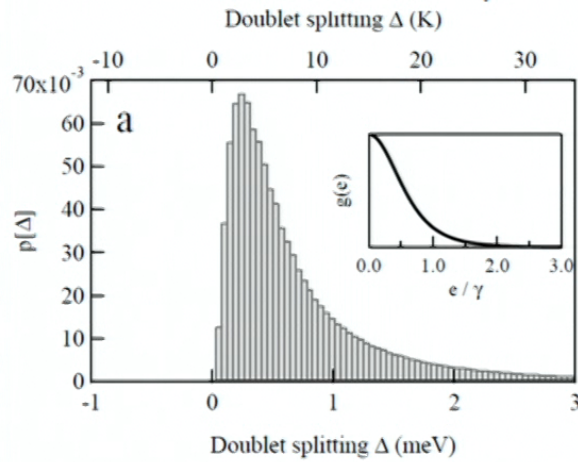
$$C^{\pm} = \frac{1}{N} \sum_{i, \langle j \rangle_i} \langle \sigma_i^+ \rangle \langle \sigma_j^- \rangle$$

Real space mean field calculation
 (7x7x7 unit cells)

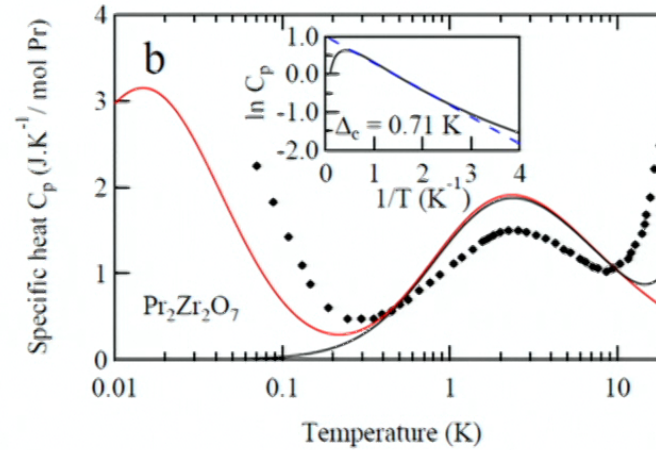
Input of MF

$$\begin{array}{c}
 |\uparrow\downarrow\rangle \\
 \parallel \\
 \text{---} \\
 \text{---} \\
 |\delta v| \\
 \text{---} \\
 \text{---} \\
 |a\rangle \simeq \frac{1}{\sqrt{2}} [|\uparrow\rangle - |\downarrow\rangle] \\
 |s\rangle \simeq \frac{1}{\sqrt{2}} [|\uparrow\rangle + |\downarrow\rangle]
 \end{array}$$

Correlated by J^{\pm} and J^{zz}

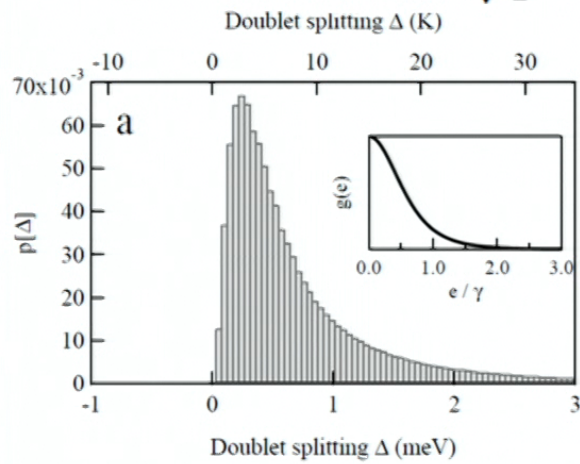
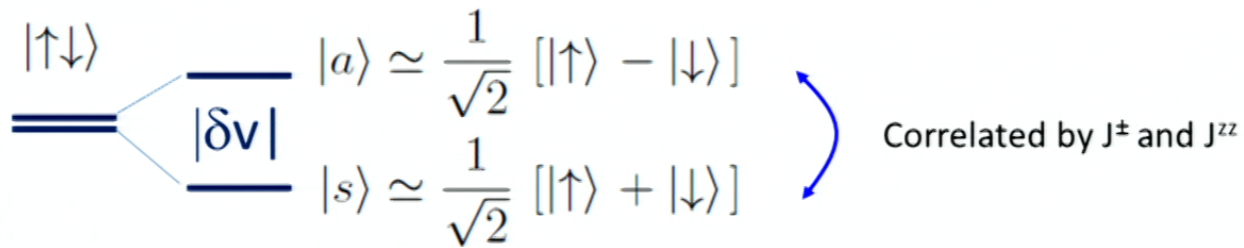


$$\gamma \approx 1.25 \times 10^{-4} \quad \text{or} \quad \delta v \approx 2\text{K}$$

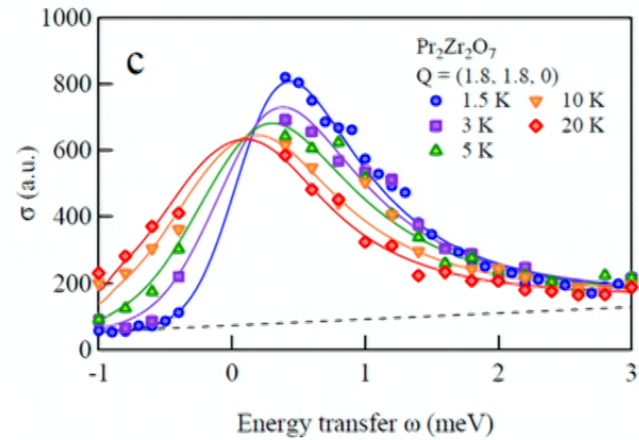


Distribution of Schottky peaks +
Hyperfine interaction at very low T

Input of MF

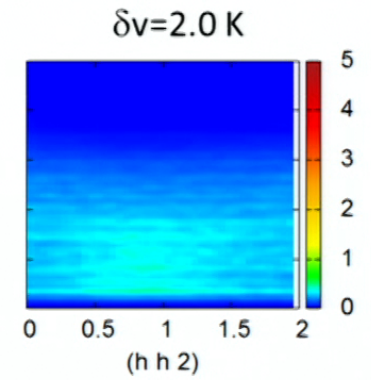
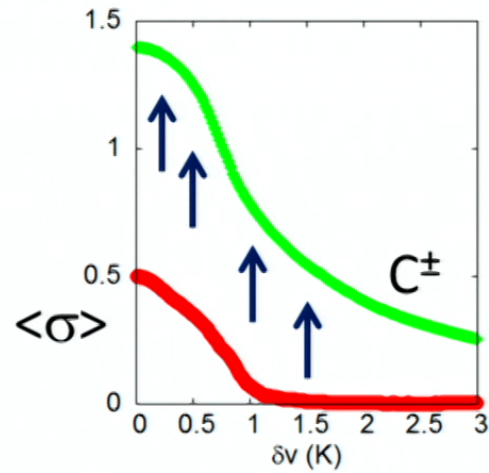
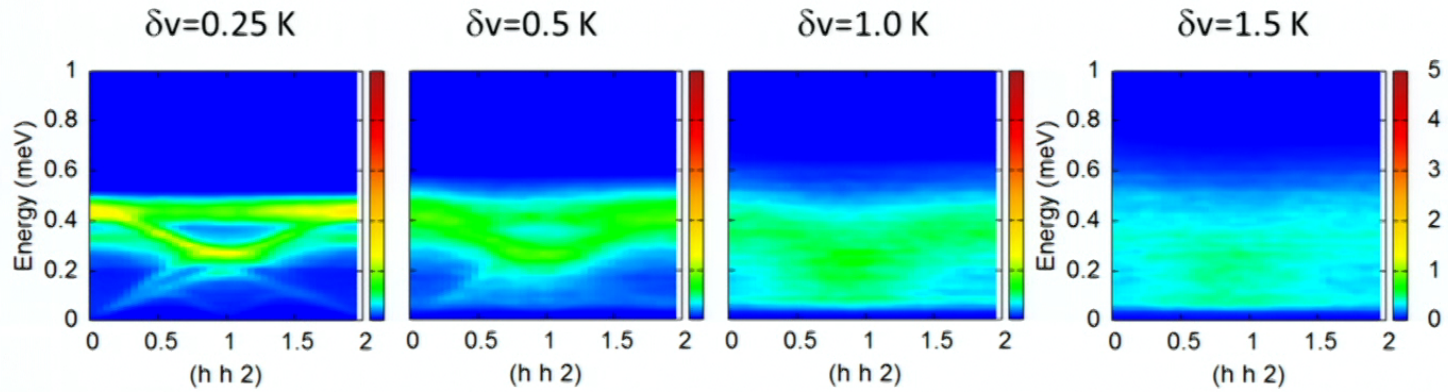


$\gamma \approx 1.25 \times 10^{-4}$ or $\delta v \approx 2\text{K}$



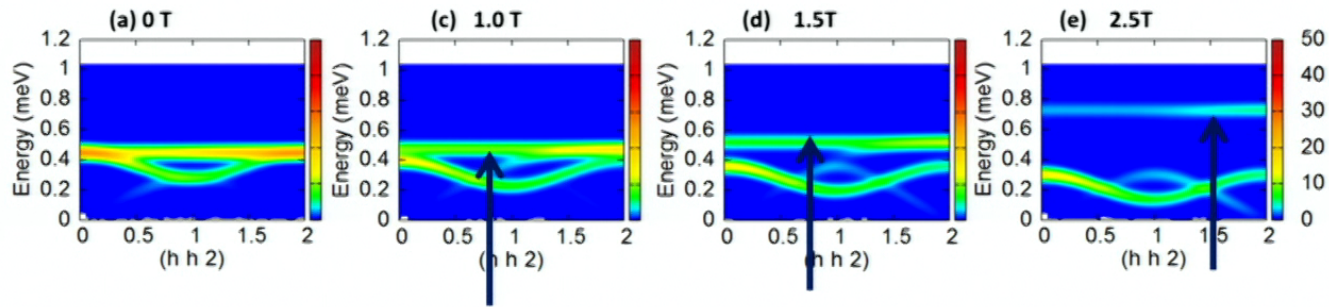
Magnetic transition within the doublets

Input of MF

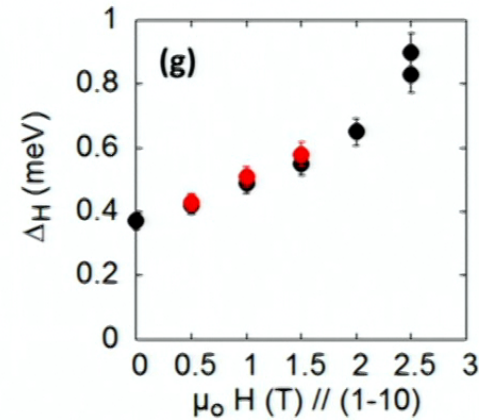
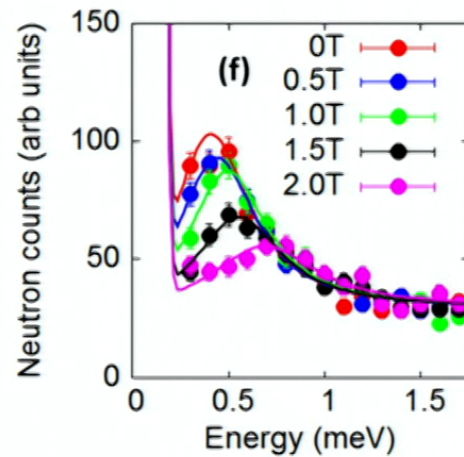


Classical spin dynamics calculation
(real space $7 \times 7 \times 7$ unit cells)

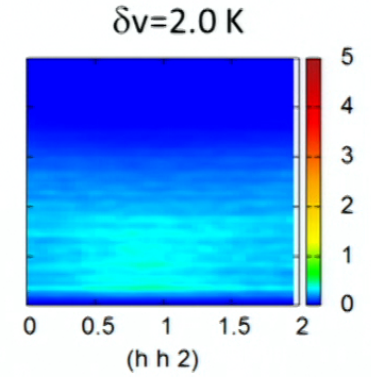
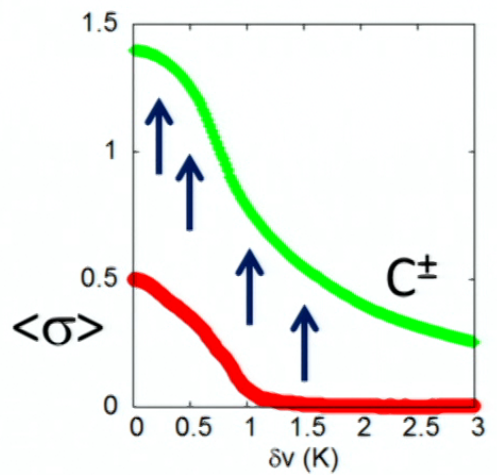
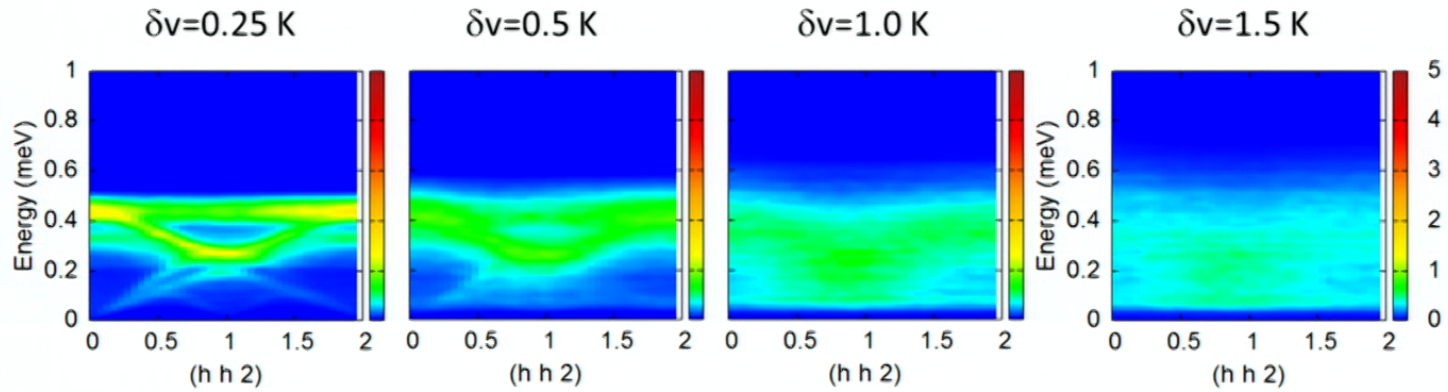
Spin and pseudo spin dynamics



Under applied field, a featureless flat mode detaches from the dynamical spin ice mode and shifts to higher energy



Input of MF



Classical spin dynamics calculation
(real space 7x7x7 unit cells)

Transverse interactions

- AF QUADRUPOLAR INTERACTIONS RELEVANT IN $\text{Pr}_2\text{Zr}_2\text{O}_7$ → FINITE AND POSITIVE J^\pm
- → INS (zero and applied field)
- → Fiel evolution of C_p

Role of disorder : stabilizes a disordered yet correlated state

- DIRECT OBSERVATION OF DIFFUSE SCATTERING THAT DEMONSTRATE THE EXISTENCE OF INTRINSIC DEFECTS WHICH MODIFY THE 4f ELECTRONIC DENSITY
- → No long order
- → C_p (peak at 2K and hyperfine upturn)
- → INS (broad blurred spin ice pattern)

$$\delta v \approx J^\pm \approx J^{zz}$$

Special thanks to

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Thank you all for your attention !