

Title: Approaching the ground state of classical dipolar spin ice Dy₂Ti₂O₇: a renewed study with neutron scattering

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URL: <http://pirsa.org/17060043>

Abstract: The true magnetic ground state in thermally equilibrated classical spin ice compounds such as Dy₂Ti₂O₇ remains an important but yet to be settled issue. In this talk, I will present our recent neutron scattering study of static and dynamic magnetic correlations in isotope-enriched ¹⁶²Dy₂Ti₂O₇ single-crystal samples. Implications within the context of possible quantum effects in dipolar spin ice based on our neutron results will be discussed as well.

Approaching the ground state of classical dipolar spin ice $\text{Dy}_2\text{Ti}_2\text{O}_7$

Yixi Su

**Juelich Centre for Neutron Science JCNS at MLZ
Forschungszentrum Jueich
Garching, Germany**

MLZ is a cooperation between:



7-9th June 2017, Waterloo,
Canada

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(JCNS-MLZ, Garching)

Michel Gingras, Jeff Rau (Univ. of Waterloo)

G. Balakrishnan, M.C. Hatnean, M.R. Lees (Warwick Univ.)

Th. Brueckel (JCNS)

A. Sazonov, V. Hutanu (RWTH Aachen and MLZ)

J. Ollivier (ILL)

F. Demmel (ISIS)

V. Pomjakushin (PSI)

and many more ...

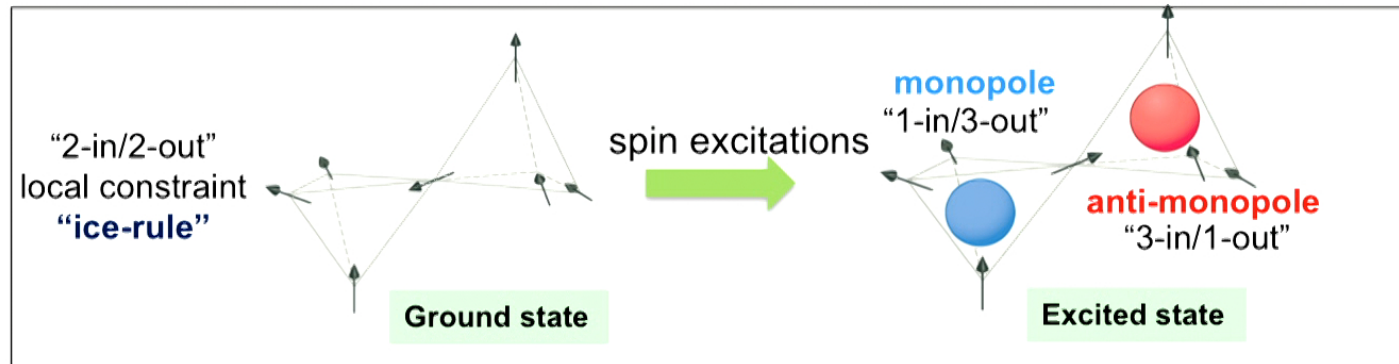


1. Introduction: spin ice – from classical to quantum

➤ Huge success with the dipolar spin ice physics

$$H_{DSI} = -J \sum_{\langle ij \rangle} \mathbf{S}_i \cdot \mathbf{S}_j + Da^3 \sum_{j>i} \left[\frac{\mathbf{S}_i \cdot \mathbf{S}_j}{|\mathbf{r}_{ij}|^3} - \frac{3(\mathbf{S}_i \cdot \mathbf{r}_{ij})(\mathbf{S}_j \cdot \mathbf{r}_{ij})}{|\mathbf{r}_{ij}|^5} \right]$$

- Realization of emergent magnetic monopole and gauge structure
- Highly entropic ground state
- Spin freezing and non-equilibrium physics
- Monopole dynamics and possible Wien effect

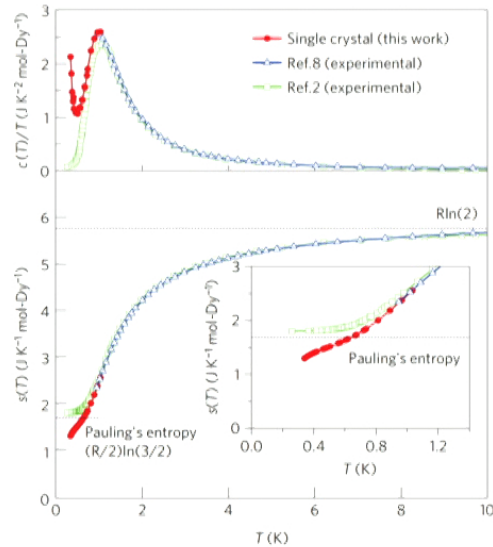


C. Catelnovo, *et al.*, Nature **451**, 42 (2008)
D.J.P. Morris, *et al.*, Science **326**, 411 (2009)

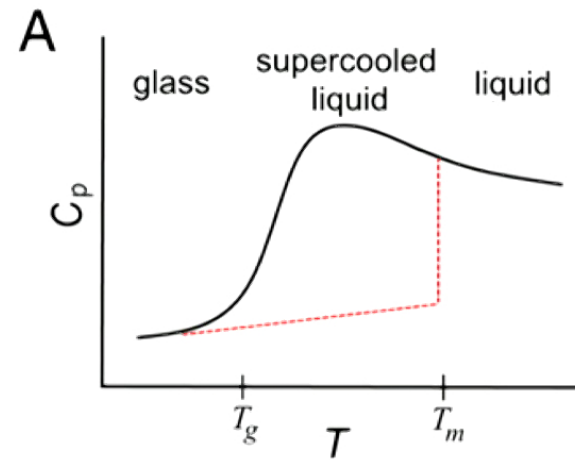
➤ However, controversies still exist

Absence of Pauling's entropy in thermally equilibrated $\text{Dy}_2\text{Ti}_2\text{O}_7$ - possible quantum effect?

"Glassiness" in spin ice - supercooled spin liquid?



D. Pomaranski, *et al.*, Nat. Phys. **9**, 353 (2013)



E.R. Kassner, *et al.*, PNAS **112**, 8549 (2015)

➤ Add quantum fluctuation in spin ice: *anisotropic exchange Hamiltonian*

$$H = \sum_{\langle ij \rangle} \{ J_{zz} S_i^z S_j^z - J_{\pm} (S_i^+ S_j^- + S_i^- S_j^+) \\ + J_{\pm\pm} (\gamma_{ij} S_i^+ S_j^+ + \gamma_{ij}^* S_i^- S_j^-) \\ + J_{z\pm} [S_i^z (\zeta_{ij} S_j^+ + \zeta_{ij}^* S_j^-) + (\zeta_{ij} S_i^+ + \zeta_{ij}^* S_i^-) S_j^z] \}$$

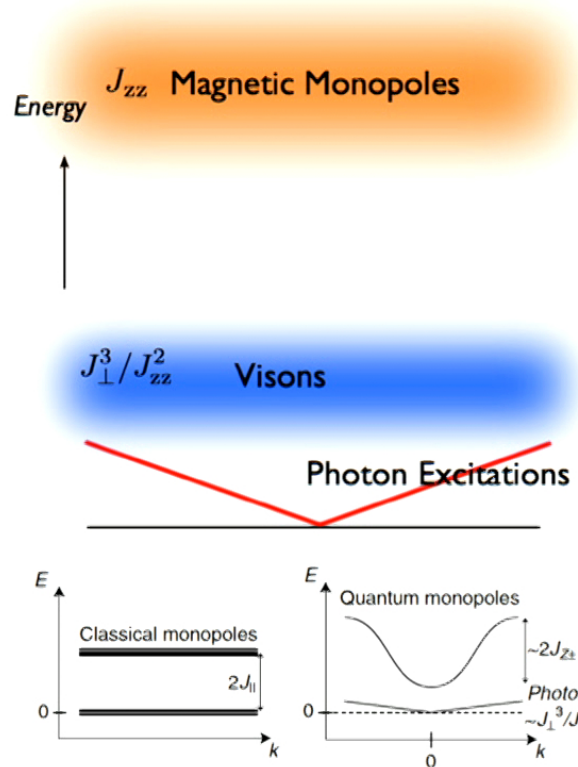
Transverse exchange interaction terms: J_{\pm} , $J_{\pm\pm}$ and $J_{z\pm}$

	$\text{Ho}_2\text{Ti}_2\text{O}_7$	$\text{Yb}_2\text{Ti}_2\text{O}_7$
CEF: 1 st excited state	240 K	680 K
Ground state	Non-Kramers	Kramers doublet: effective $S=1/2$
g-tensor	$g_{\parallel} \approx 19, g_{\perp} = 0$	$g_{\parallel} \approx 1.78, g_{\perp} \approx 4.28$
anisotropy	Ising easy-axis anisotropy	XY easy-plane anisotropy
Localized moment	$\sim 10 \mu\text{B}$	$\sim 3 \mu\text{B}$

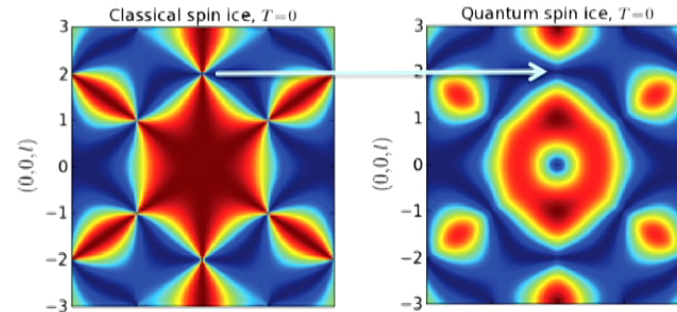
In real materials: Yb-, Pr-based pyrochlore compounds are strong candidates for QSI

➤ QSI: emergent quantum electrodynamics (QED)

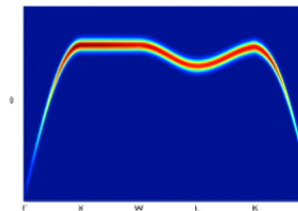
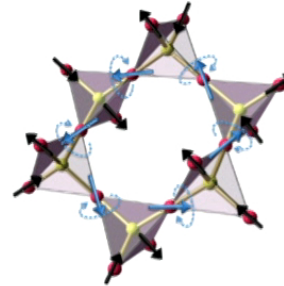
Energy landscape of QSI



Melting of pinch-point



Emergence of artificial photons: collective spin fluctuations over hexagon-loop
 -> could be observed as gapless excitations by inelastic neutron scattering



M.J.P. Gingras and P.A. McClarty, Rep. Prog. Phys. **77**, 056501 (2014)

PHYSICAL REVIEW X **1**, 021002 (2011)

Quantum Excitations in Quantum Spin Ice

Yb₂Ti₂O₇

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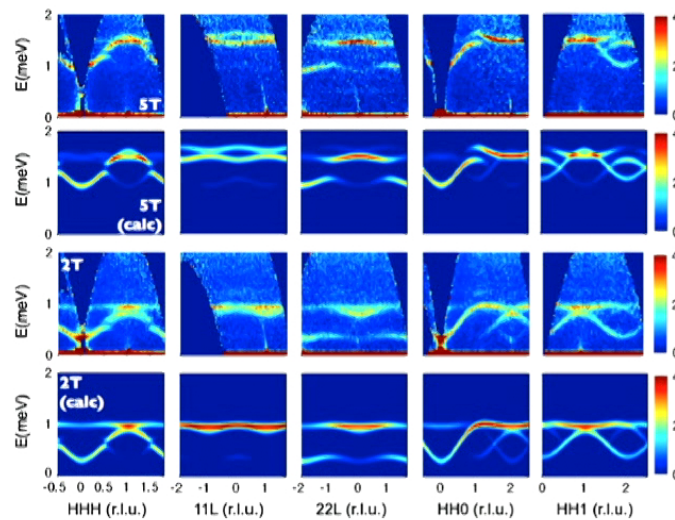
²Ecole Normale Supérieure de Lyon, 46, allée d'Italie, 69364 Lyon Cedex 07, France

³Canadian Institute for Advanced Research, 180 Dundas St. W., Toronto, Ontario, M5G 1Z8, Canada

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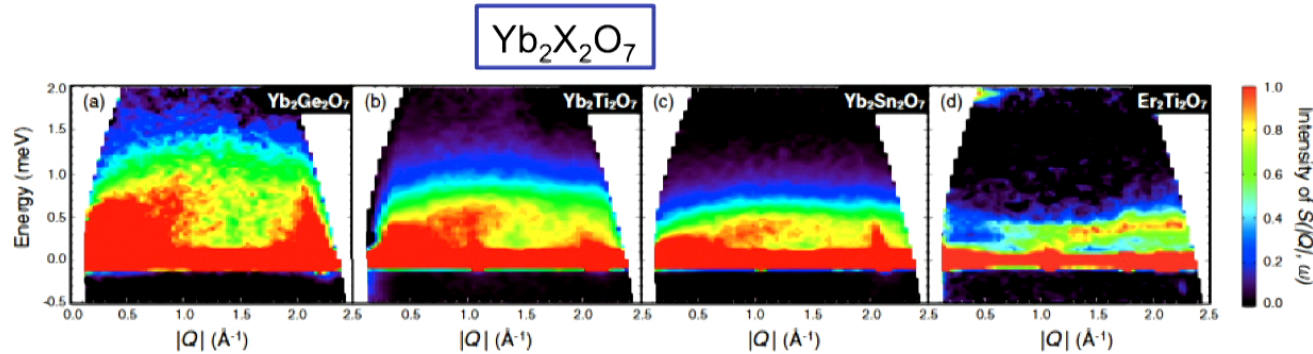
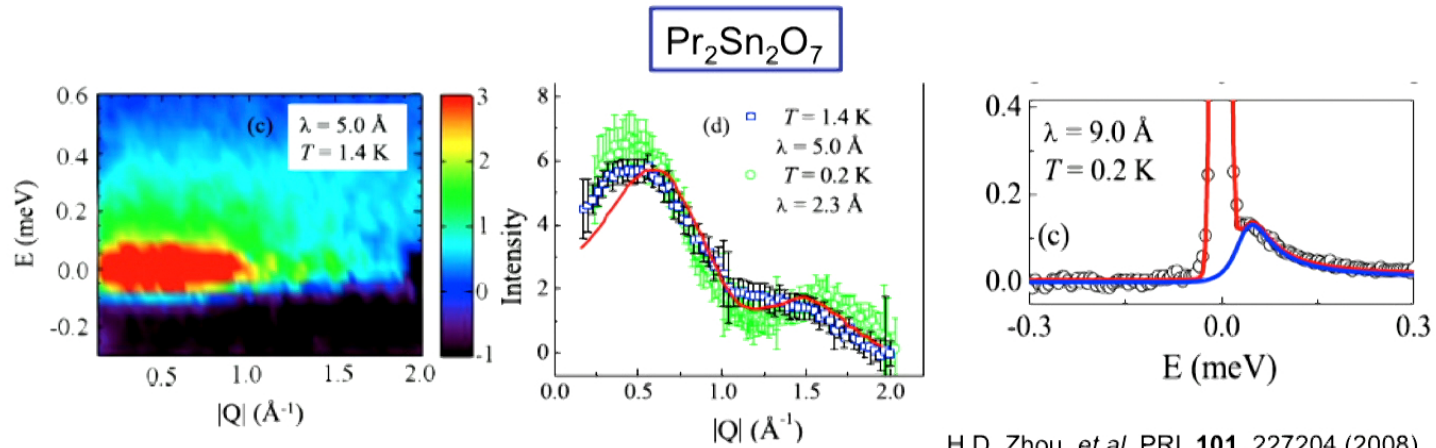
⁵Kavli Institute for Theoretical Physics, University of California, Santa Barbara, California, 93106-4030, USA

(Received 22 July 2011; published 3 October 2011)



$$\begin{aligned}
 J_{zz} &= 0.17 \pm 0.04, & J_{\pm} &= 0.05 \pm 0.01, \\
 J_{\pm\pm} &= 0.05 \pm 0.01, & J_{z\pm} &= -0.14 \pm 0.01,
 \end{aligned}$$

→ Presence of significant transverse exchange terms!



A.M. Hallas, *et al.* Phys. Rev. B **93**, 100403(R) (2016)

2. *Quasi-collinear* ferromagnetic structure and gapped magnetic excitations in $\text{Yb}_2\text{Ti}_2\text{O}_7$

Are Multiphase Competition and Order by Disorder the Keys to Understanding Yb₂Ti₂O₇?

L. D. C. Jaubert,¹ Owen Benton,¹ Jeffrey G. Rau,² J. Oitmaa,³ R. R. P. Singh,⁴ Nic Shannon,¹ and Michel J. P. Gingras^{2,5,6}

¹Okinawa Institute of Science and Technology Graduate University, Onna-son, Okinawa 904-0495, Japan

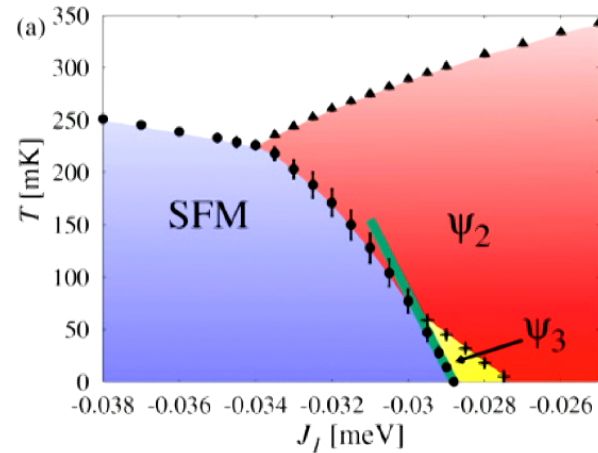
²Department of Physics and Astronomy, University of Waterloo, 200 University Avenue West, Waterloo, Ontario N2L 3G1, Canada

³School of Physics, The University of New South Wales, Sydney 2052, Australia

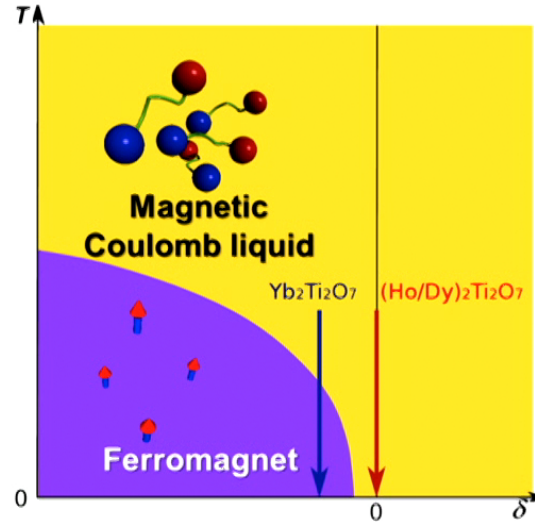
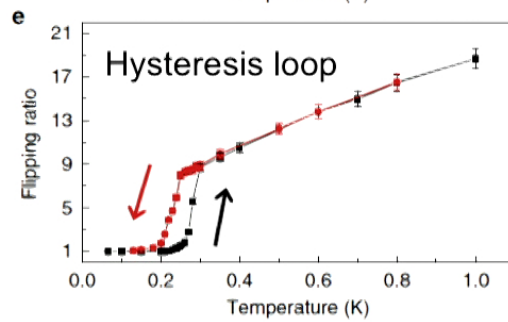
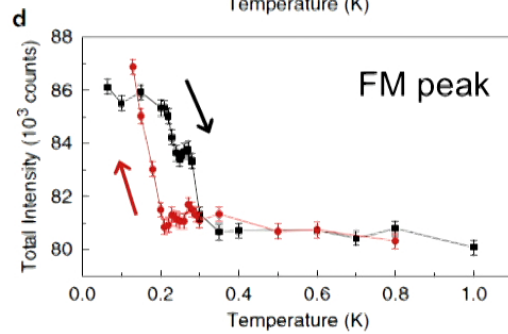
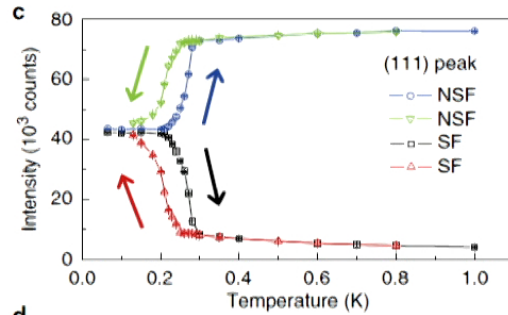
⁴Department of Physics, University of California, Davis, California 95616, USA

⁵Perimeter Institute for Theoretical Physics, 31 Caroline North, Waterloo, Ontario N2L 2Y5, Canada

⁶Canadian Institute for Advanced Research, Toronto, Ontario M5G 1Z8, Canada



Strong sample dependence!



Single crystals grown by Y. Yasui
(Nagoya University, Japan)

L.J. Chang, *et al.*, Nat. Commun. **3**, 992 (2012)

Neutron powder diffraction refinement

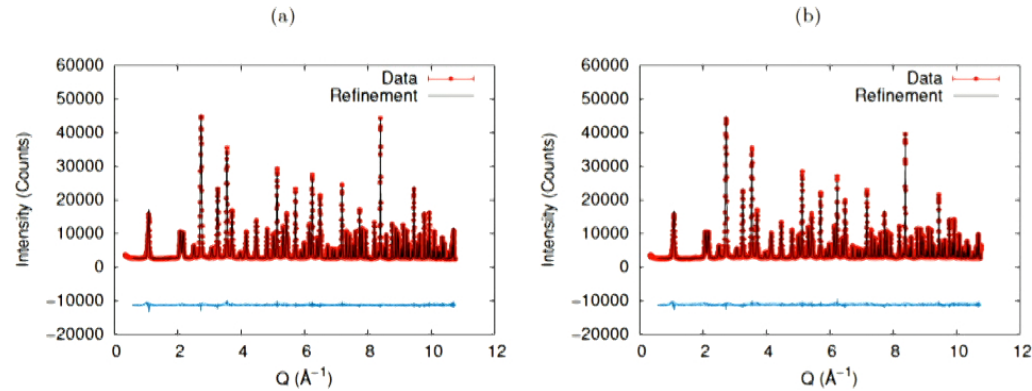


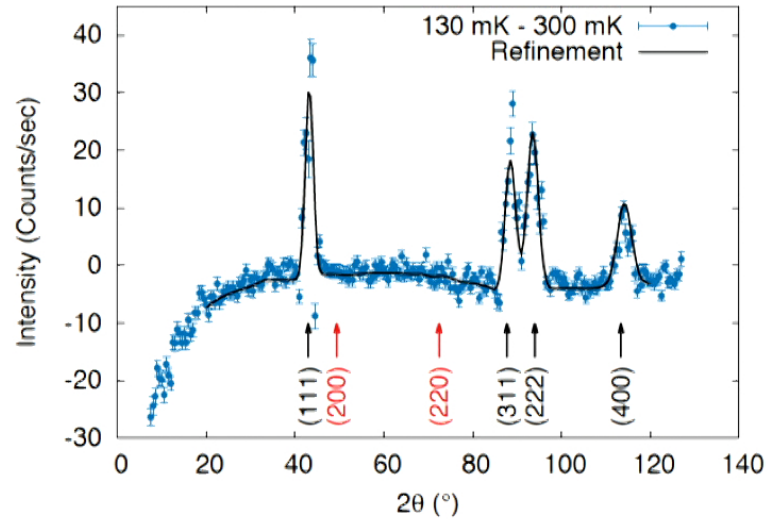
FIG. 1. Data (red dots) and best refinement (black lines) for (a) 2 K and (b) 150 K. The difference between calculated and measured intensities is given by the blue line below each pattern.

Temperature	Occupancy of Yb^{3+} at 16c site	Lattice parameter (\AA)	x	R_{wp}
Stoichiometric model				
2 K	-	10.01275(2)	0.33140(5)	0.0671
150 K	-	10.02061(2)	0.33105(6)	0.0787
1.5 K (Yaouanc <i>et al.</i>)	-	10.0220(5)	0.332(1)	-
150 K (Ross <i>et al.</i>)	-	10.01111(3)	0.33122(3)	0.0414
Stuffed model				
2 K	0.0009(1)	10.01273(2)	0.33126(5)	0.0670
150 K	0.0010(1)	10.02059(2)	0.33095(6)	0.0785
150 K (Ross <i>et al.</i>)	0.002(1)	10.01111(3)	0.33121(3)	0.0415

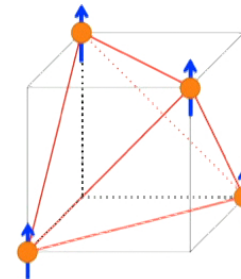
TABLE I. Refined lattice parameters and oxygen 48f Wyckoff position of the space group for the stoichiometric and stuffed models. For comparison, results from Ross *et al.* [18] and Yaouanc *et al.* [2] are also shown. The refinement uncertainties are not mentioned in Ref. [2].

Viviane Pecanha-Antonio, *et al.* (in preparation)

Determination of LT magnetic structure

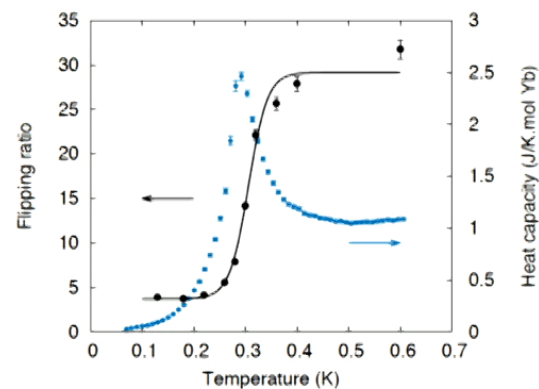
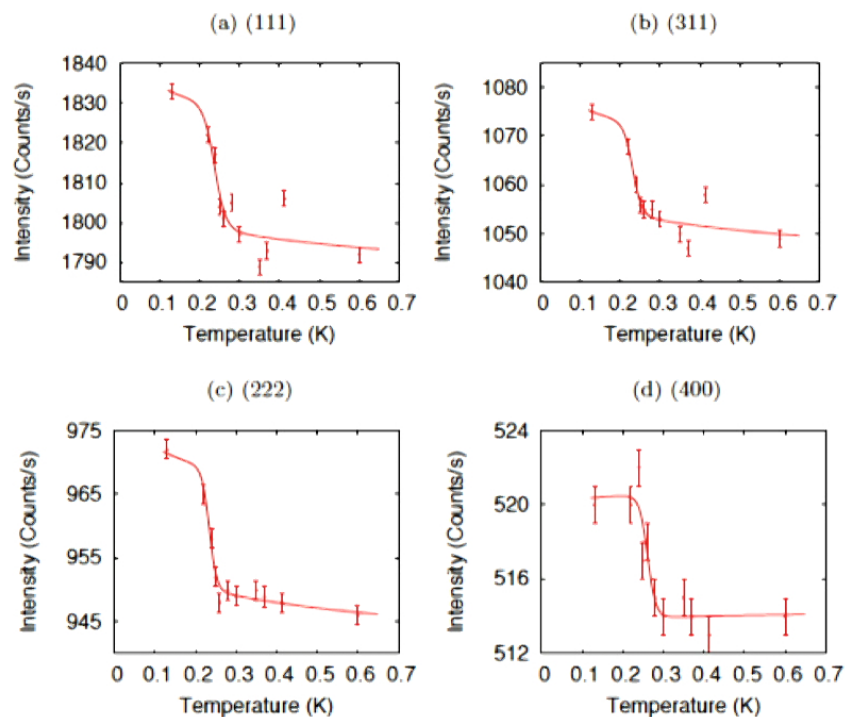


Site	X	Y	Z	$m^a(\mu_B)$	$m^b(\mu_B)$	$m^c(\mu_B)$
1	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	0.00(2)	0.00(2)	0.87(1)
2	$\frac{1}{4}$	$\frac{1}{4}$	$\frac{1}{2}$	-0.00(2)	-0.00(2)	0.87(1)
3	$\frac{1}{2}$	$\frac{1}{4}$	$\frac{1}{4}$	-0.00(2)	0.00(2)	0.87(1)
4	$\frac{1}{4}$	$\frac{1}{2}$	$\frac{1}{4}$	0.00(2)	-0.00(2)	0.87(1)



Viviane Pecanha-Antonio, *et al.* (in preparation)

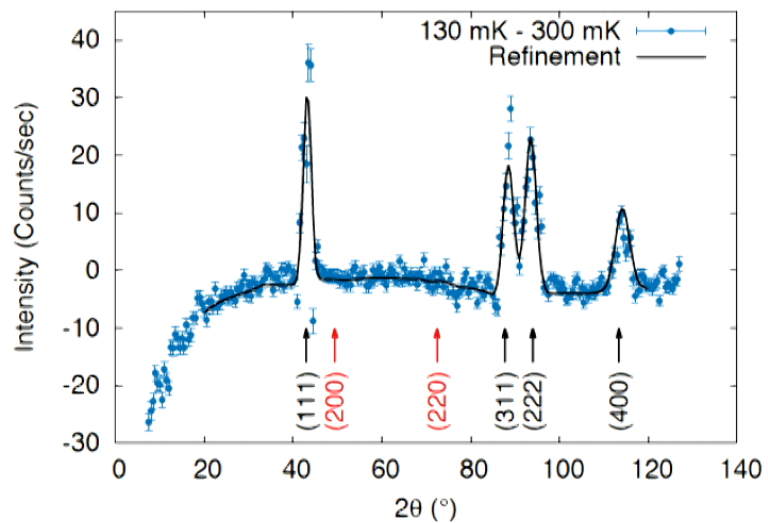
Well-resolvable LT phase transition in $\text{Yb}_2\text{Ti}_2\text{O}_7$



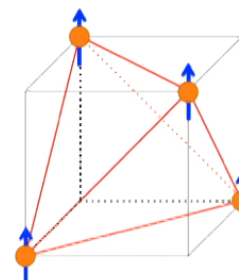
Heat capacity
Neutron depolarisation

Neutron diffraction

Determination of LT magnetic structure



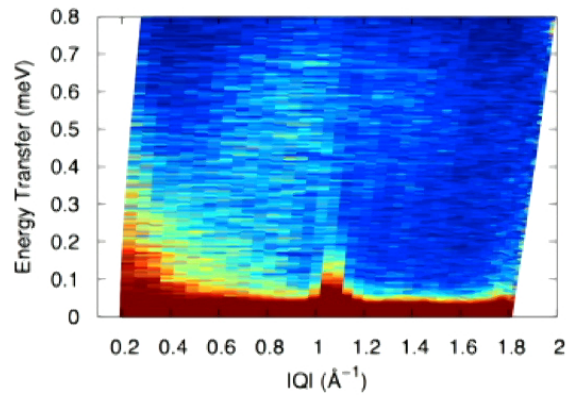
Site	X	Y	Z	$m^a (\mu_B)$	$m^b (\mu_B)$	$m^c (\mu_B)$
1	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	0.00(2)	0.00(2)	0.87(1)
2	$\frac{1}{4}$	$\frac{1}{4}$	$\frac{1}{2}$	-0.00(2)	-0.00(2)	0.87(1)
3	$\frac{1}{2}$	$\frac{1}{4}$	$\frac{1}{4}$	-0.00(2)	0.00(2)	0.87(1)
4	$\frac{1}{4}$	$\frac{1}{2}$	$\frac{1}{4}$	0.00(2)	-0.00(2)	0.87(1)



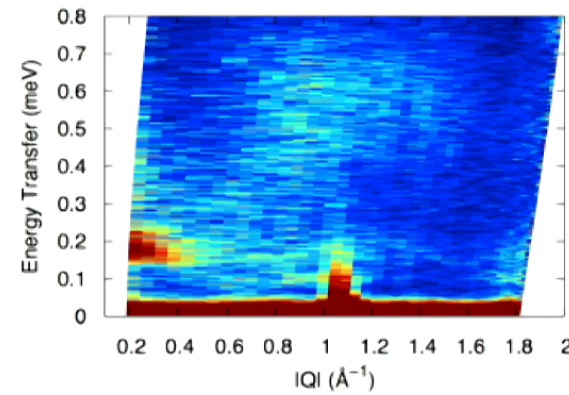
Viviane Pecanha-Antonio, *et al.* (in preparation)

Inelastic neutron scattering on polycrystalline sample

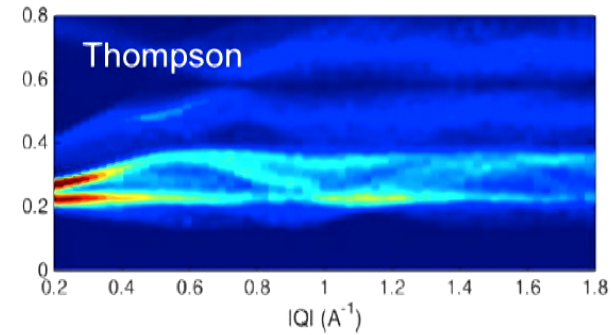
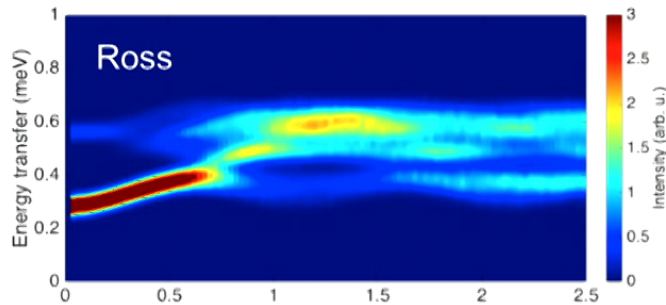
300 mK



50 mK



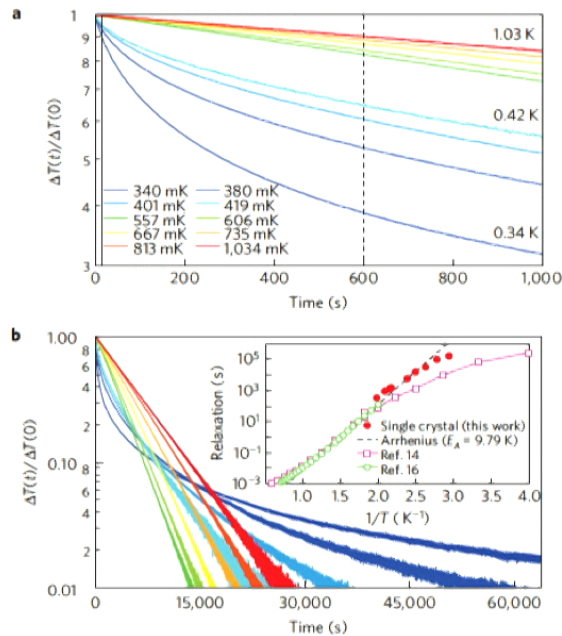
For comparison: LSWT simulations based on the known models



Viviane Pecanha-Antonio, *et al.* (in preparation)

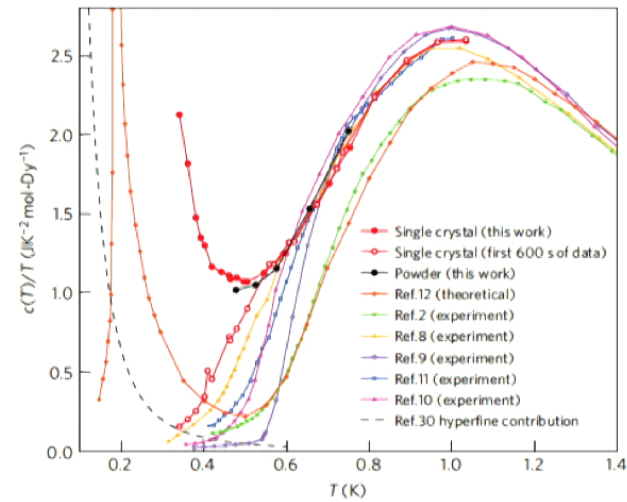
3. Approaching the ground state of classical dipolar spin ice $\text{Dy}_2\text{Ti}_2\text{O}_7$

Zero-point entropy in spin ice $\text{Dy}_2\text{Ti}_2\text{O}_7$



Estimated thermal relaxation time:

- 800mK: $\sim 10^{-1}$ s
- 450mK: $\sim 10^3$ - 10^4 s
- 350mK: $\sim 10^5$ s



- Extremely slow thermal relaxations
➔ out-of-equilibrium state ➔ Wien effect and monopole dynamics
- True ground state of spin ice?

D. Pomaranski, *et al.*, Nat. Phys. **9**, 353 (2013)
V. Kaiser, *et al.*, Nat. Mat. **12**, 1033 (2013)

PHYSICAL REVIEW B 92, 094418 (2015)

Chain-based order and quantum spin liquids in dipolar spin ice

P. A. McClarty,^{1,2} O. Sikora,^{3,4,5} R. Moessner,² K. Penc,⁶ F. Pollmann,² and N. Shannon^{4,5}

PHYSICAL REVIEW B 93, 024402 (2016)

Refrustration and competing orders in the prototypical $\text{Dy}_2\text{Ti}_2\text{O}_7$ spin ice material

P. Henelius,^{1,*} T. Lin,² M. Enjalran,^{3,4} Z. Hao,² J. G. Rau,² J. Altosaar,^{2,5} F. Flicker,⁶ T. Yavors'kii,⁷ and M. J. P. Gingras^{2,8,9}

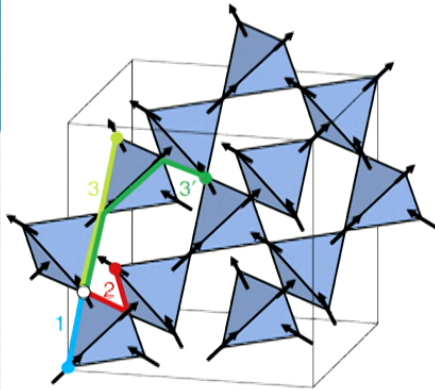
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Received 18 Dec 2015 | Accepted 15 Jul 2016 | Published 25 Aug 2016

DOI: 10.1038/ncomms12592 OPEN

Intermediate magnetization state and competing orders in $\text{Dy}_2\text{Ti}_2\text{O}_7$ and $\text{Ho}_2\text{Ti}_2\text{O}_7$

R.A. Borzi^{1,2}, F.A. Gómez Albarracín^{2,3}, H.D. Rosales^{2,3}, G.L. Rossini^{2,3}, A. Steppke^{4,5}, D. Prabhakaran⁶, A.P. Mackenzie^{4,5}, D.C. Cabra^{2,3} & S.A. Grigera^{1,2,4}



Model	D	J_1	J_2	J_{3a}	$(J_2/3)+J_{3a}$	J_{3b}
Upper	1.3224	3.30	-0.0949	0.07	0.038	-0.0167
Lower	1.3224	3.44	-0.1649	-0.02	-0.075	0.0433
Boundary	1.3224	3.38	-0.1349	0.02	-0.025	0.0167
g-DSM*	1.3224	3.41	-0.14	0.025	-0.0217	0.025
d-DSM	1.3224	3.41	0	-0.02	-0.02	0.07
s-DSM	1.41	1.56	0	0	0	0

1. important role of further-neighbor exchange interactions such as J_2, J_3
2. competing orders leading to re-frustration in the spin-ice regime
3. quantum effect may be enhanced due to dipole-octopole nature of ground state doublet of Dy^{3+}

P. Henelius *et al.*, Phys. Rev. B **93**, 024402 (2016).
P.A. McClarty *et al.*, Phys. Rev. B **92**, 094418 (2015).
J.G. Rau & M.J.P. Gingras, Phys. Rev. B **92**, 144417 (2015).
R.A. Borzi, *et al.*, Nat. Commun. **7**, 12592 (2016).

Estimation of possible quantum effects in spin ice

- Estimation by Jeff Rau et al. (PRB **92**, 144417 (2015)):

$$J_{\pm} \leq 56 \text{ mK}$$

$$g = 12J_{\pm}^3/J_{ZZ}^2 \sim 2 \text{ mK (if taking } J_{ZZ} = 1 \text{ K)}$$

- Estimation by P. Henelius et al. (PRB **93**, 024042 (2016)):

$$J_{\pm} \sim 3 \text{ mK}$$

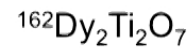
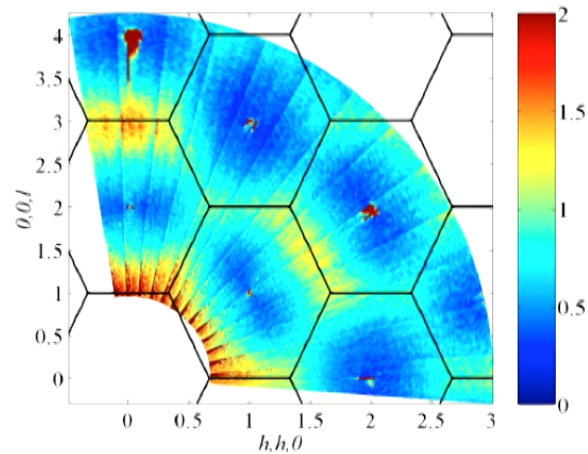
$$g = 12J_{\pm}^3/J_{ZZ}^2 \sim 0.05 \text{ mK (if taking } J_{ZZ} = 4 \text{ K)}$$

- Estimation by P.A. McClarty et al. (PRB **92**, 094418 (2016)):

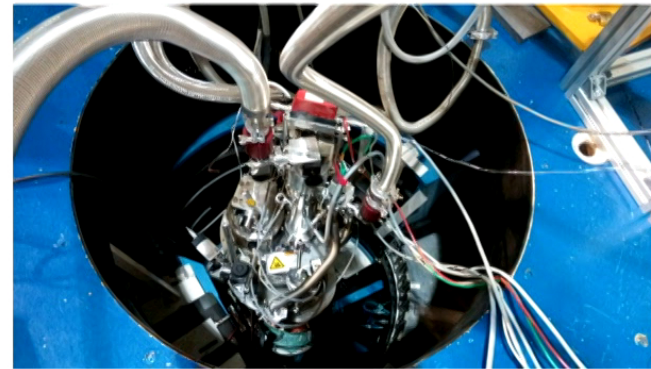
$$g = 12J_{\pm}^3/J_{ZZ}^2 \sim 70 \text{ mK could stabilize QSI for DTO}$$

Experimental challenges with neutrons

1. strong neutron absorption of Dy^{3+} ($\sigma_{\text{abs}} > 4600$ barn at 4.2 Å)
2. large incoherent cross-section of Dy^{3+} (~54 barn/Dy) -> high background
3. extremely slow thermal relaxation to equilibrium below ~600 mK

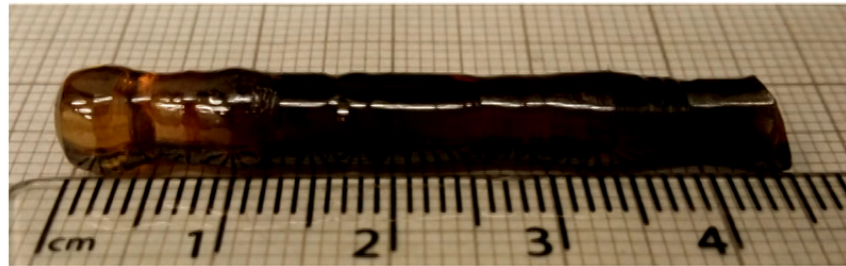


T. Fennell, *et al.* Phys. Rev. B **70**, 134408 (2004);
 T. Yavors'kii, *et al.* Phys. Rev. Lett. **101**, 037204 (2008)



Cryogen-free dilution- and ^3He -
 inserts at MLZ
 -> long time and stable operation

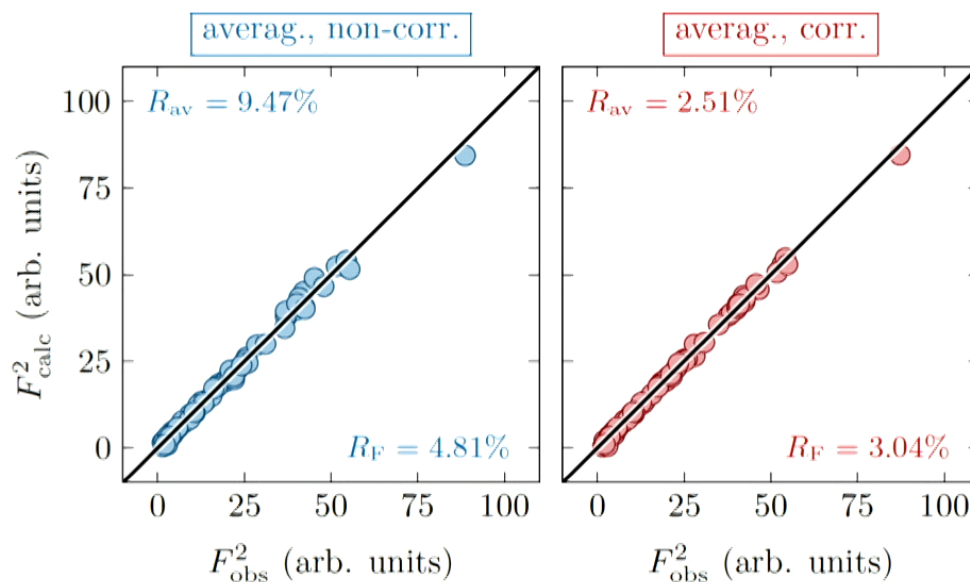
New generation of isotope-enriched $^{162}\text{Dy}_2\text{Ti}_2\text{O}_7$ single crystal



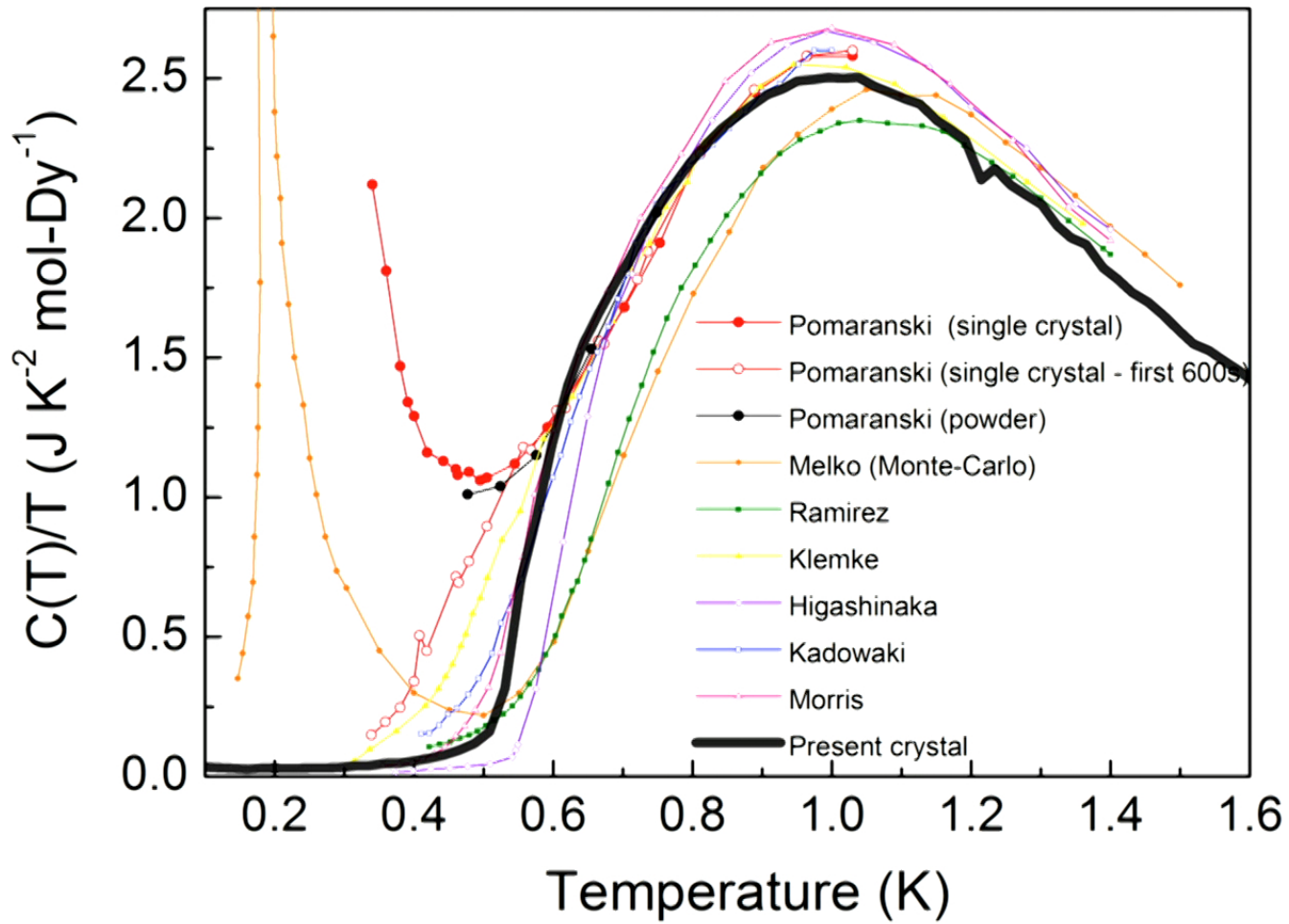
1. much reduced neutron absorption: $\sigma_{\text{abs}} < 1000$ barn at 4.2 Å
2. much reduced incoherent scattering background: 0 barn/ ^{162}Dy
3. grown under optimal conditions to reduce possible stuffing and oxygen defects

M. Ciomaga Hatnean
G. Balakrishnan
(University of Warwick)

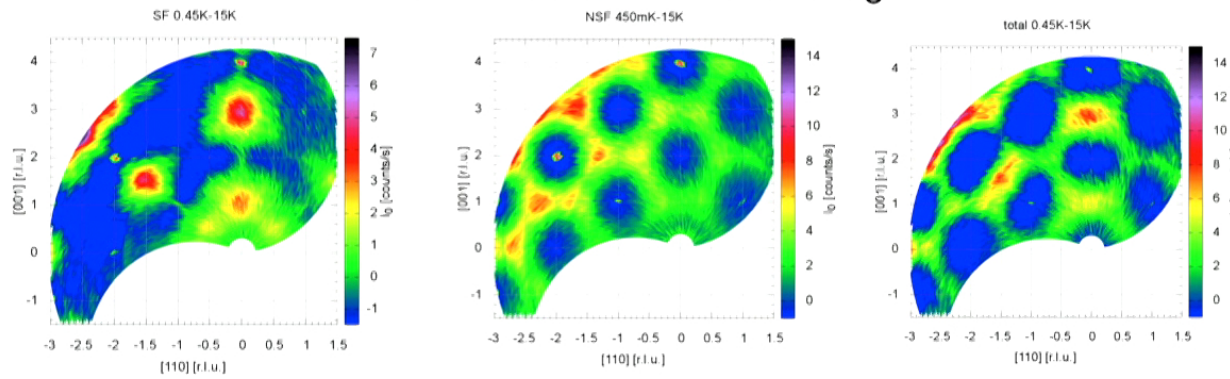
- hot-neutron diffractometer POLI@MLZ: wavelength $\lambda=0.7 \text{ \AA}$
- 296 reflections measured on a $\sim 200 \text{ mg}$ single crystal
- absorption correction via the CCSL package



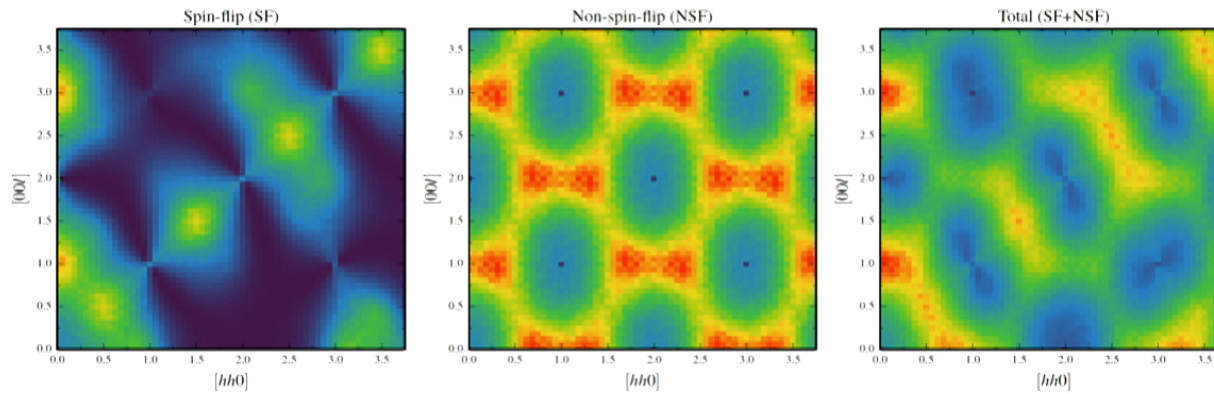
	162-Dy	natural-Dy	Ti	O(1)	O(2)
Occupancy	0.91(2)	0.09(2)	1.00(2)	0.99(2)	1.00(2)



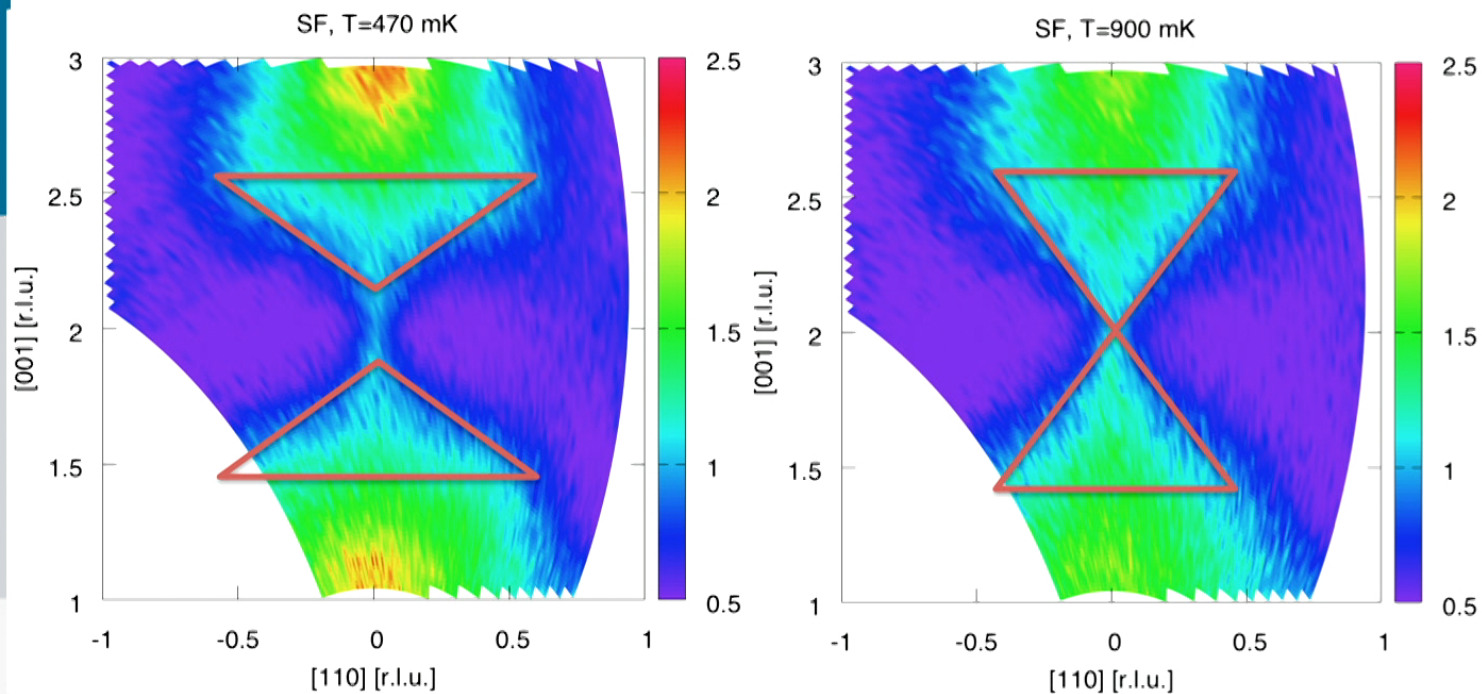
Polarized neutron scattering



Monte-Carlo simulations based on g-DSM* (J.G. Rau and M. Gingras)

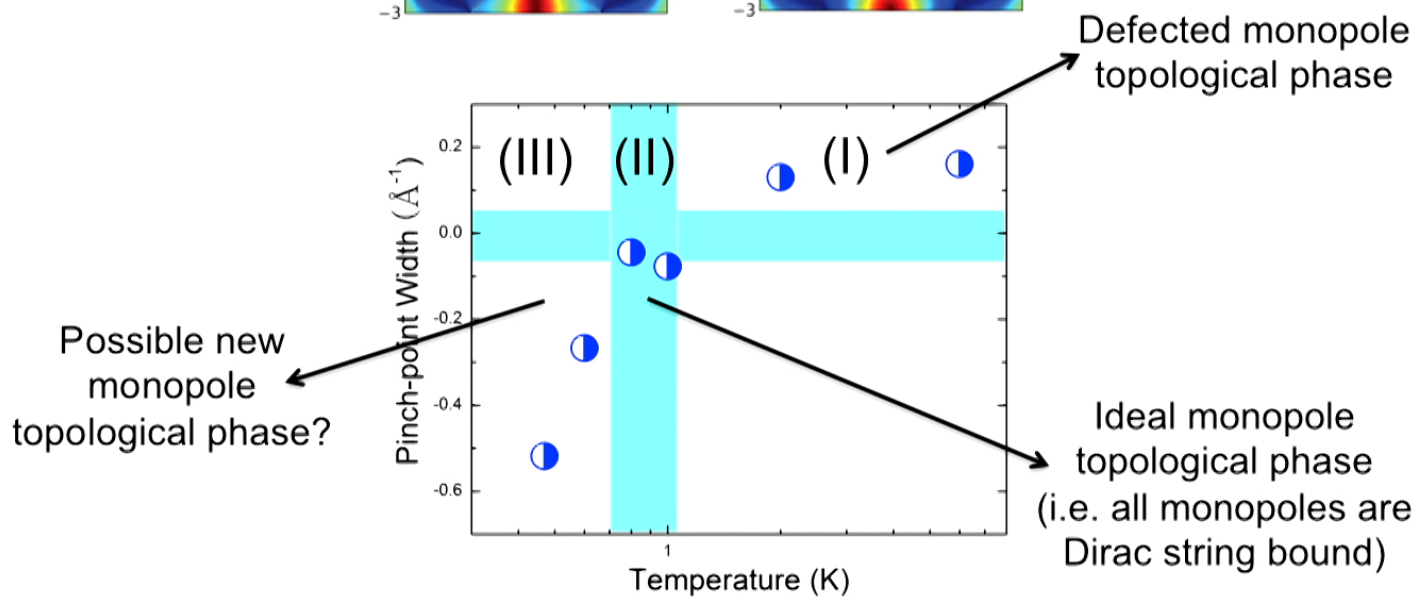
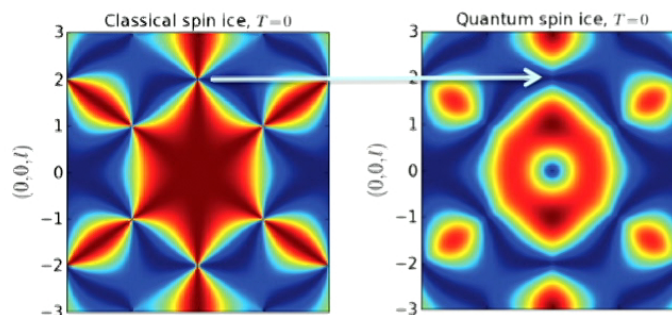


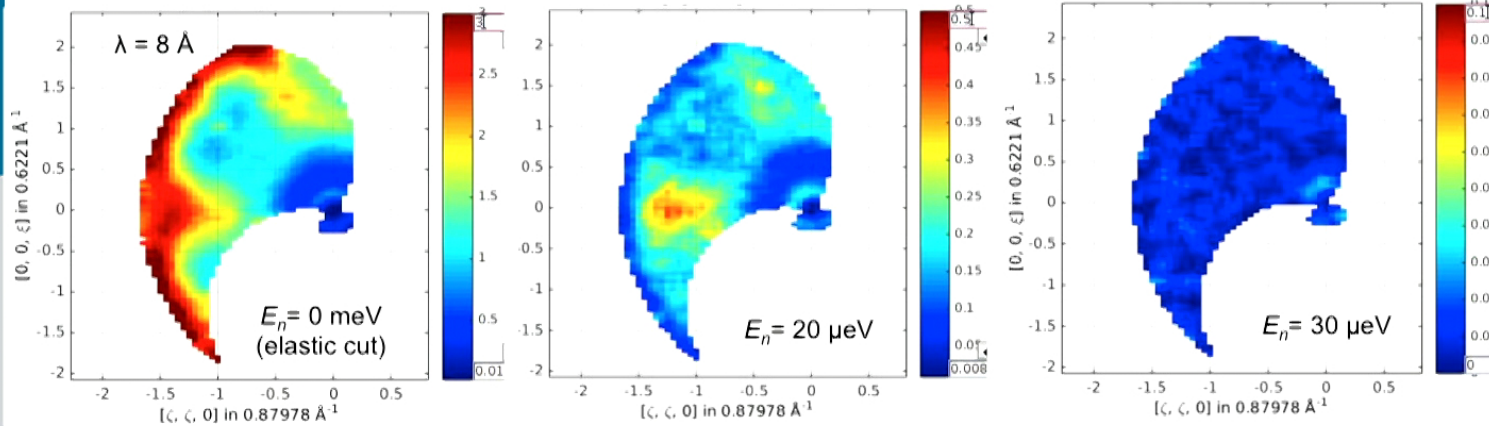
Fate of pinch point



Fate of pinch point

Melting of pinch-point (Owen Benton)



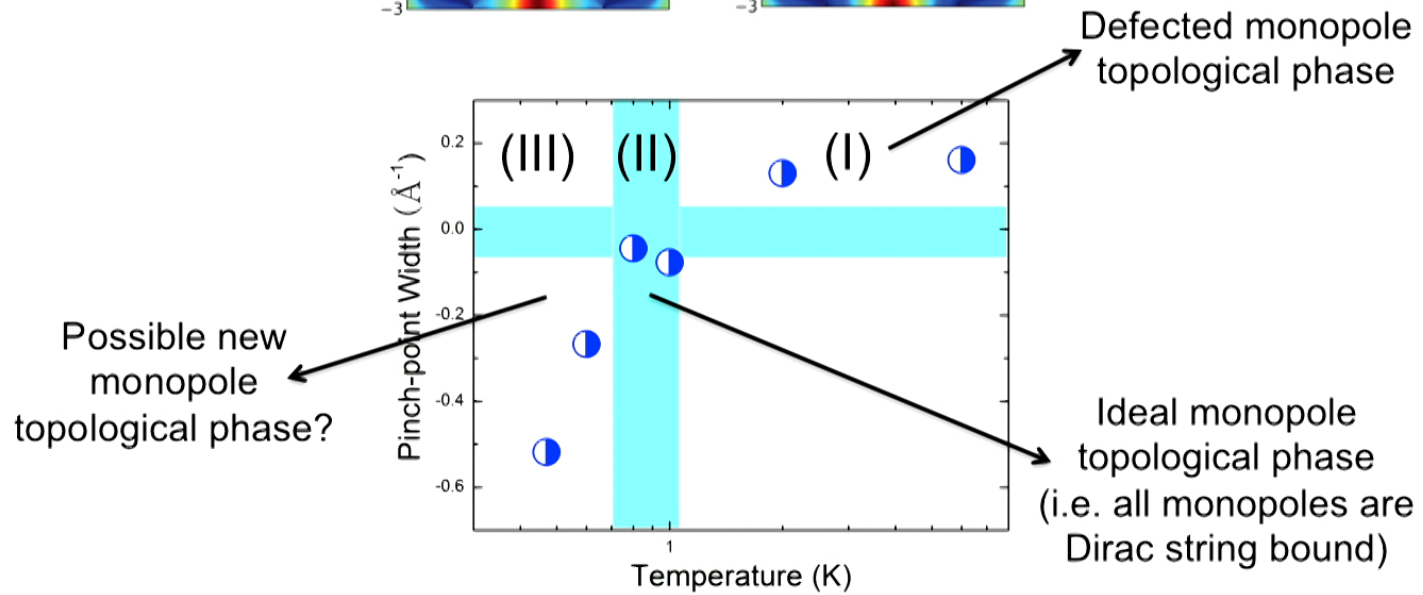
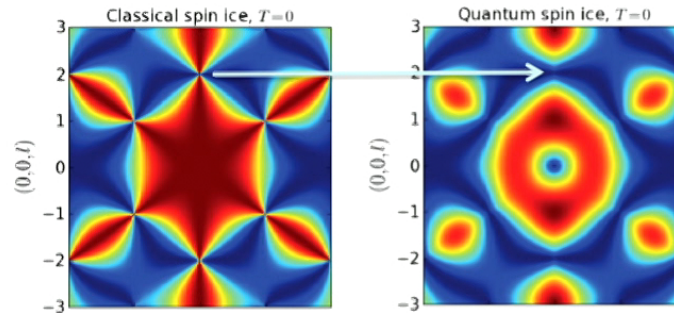


- magnetic scattering almost entirely elastic
- however, indication exists for an extended quasi-elastic component with energy transfer up to $\sim 20 \text{ \mu eV}$
 - upper limit for any possible quantum effects
 - possibly due to competing orders?

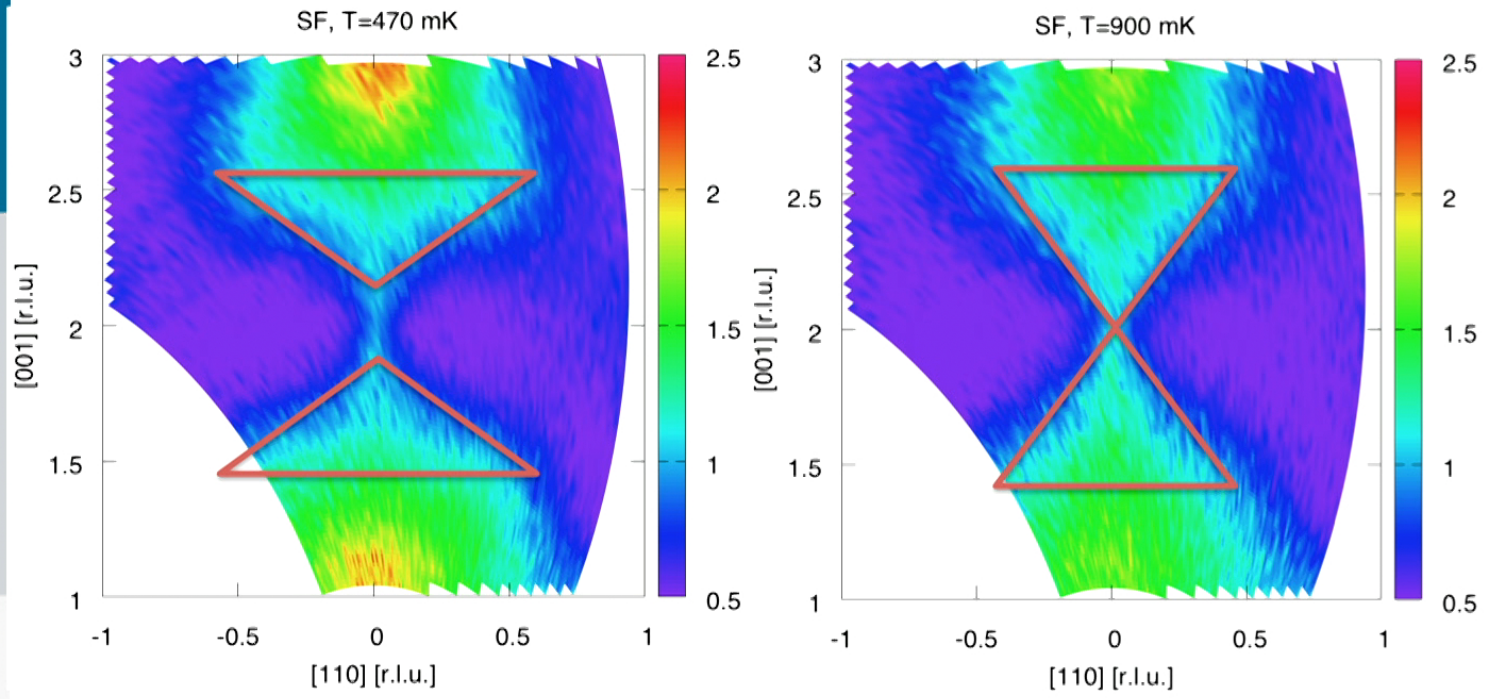
IN5@ILL

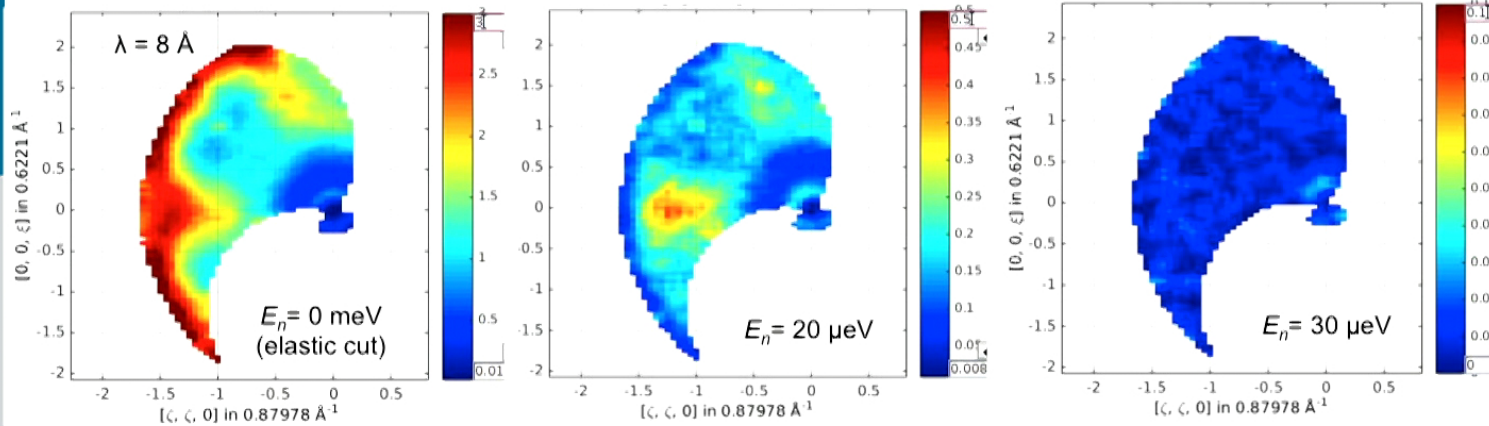
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IN5@ILL

1. Magnetic ground state of $\text{Yb}_2\text{Ti}_2\text{O}_7$
 - Quasi-collinear ferromagnetic structure
 - Gapped magnetic excitations
2. Magnetic ground state of dipolar spin ice $\text{Dy}_2\text{Ti}_2\text{O}_7$
 - Polarized diffuse scattering data \Leftrightarrow various DSM models
 - Evidence of possible melting of pinch point
 - Competing orders and re-frustration



MLZ is a cooperation between: