Title: Weak measurements, decoherence and cosmology

Date: Jun 01, 2017 04:50 PM

URL: http://pirsa.org/17060029

Abstract: In this work we consider a recent proposal in which gravitational interactions are mediated via the exchange of classical information and apply it to a quantized Friedman-Robertson-Walker (FRW) universe with the assumption that any test particles must feel a classical metric. We show that such a model results in decoherence in the FRW state that manifests itself as a dark energy fluid that fills the spacetime. Motivated by quantum-classical interactions this model is yet another example of theories with violation of energy-momentum conservation whose signature could have significant consequences for the observable universe.

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Weak measurements, decoherence and cosmology



(CQG-34,11 (2017) - arXiv:1605.05980)

Natacha Altamirano

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PI Day June 2017

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Sunyaev-Zeldovich effect

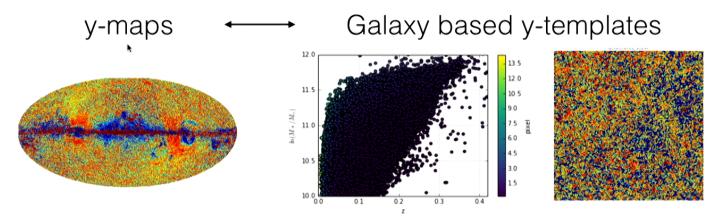
With Chiamaka Okoli and Niayesh Afshordi

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Sunyaev-Zeldovich effect

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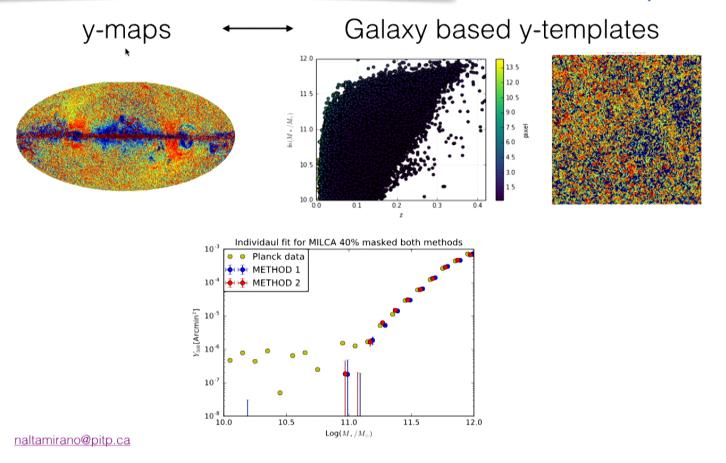


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Sunyaev-Zeldovich effect

With Chiamaka Okoli and Niayesh Afshordi



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QUANTUM INTERACTIONS

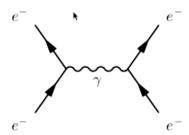
- Mediated by virtual particles
- Produce entanglement

CLASSICAL INTERACTIONS

- Ehrenfest theorem
- → Do not produce entanglement

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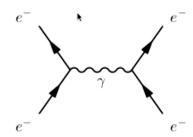
Pirsa: 17060029 Page 6/43



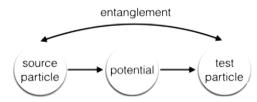
As an example of a quantum interaction we consider electrodynamic theory that is **dominated** by **local interactions** and **long range** forces arise as fluctuations of **gauge field**.

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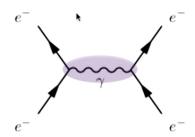


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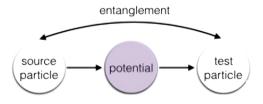


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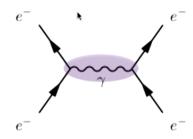
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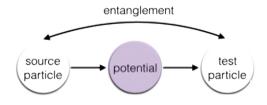
The quantum interactions described by quantum field theory are governed by unitary evolution

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Pirsa: 17060029 Page 9/43



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The quantum interactions described by quantum field theory are governed by unitary evolution

Gravitation remains stubbornly resistant to quantization

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General Relativity and Gravitation, Vol. 28, No. 5, 1996

On Gravity's Role in Quantum State Reduction

N.

Roger Penrose^{1,2}

Received August 22, 1995. Rev. version December 12, 1995

The stability of a quantum superposition of two different stationary mass distributions is examined, where the perturbing effect of each distribution on the space-time structure is taken into account, in accordance with the principles of general relativity. It is argued that the definition of the time-translation operator for the superposed space-times involves an inherent ill-definedness, leading to an essential uncertainty in the energy of the superposed state which, in the Newtonian limit, is proportional to the gravitational self-energy E_{Δ} of the difference between the two mass distributions. This is consistent with a suggested finite lifetime of the order of \hbar/E_{Δ} for the superposed state, in agreement with a certain proposal made by the author for a gravitationally induced spontaneous quantum state reduction, and with closely related earlier suggestions by Diósi and by Ghirardi et al.

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On Gravity's Role in Quantum State Reduction

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VOLUME 40, NUMBER 3

AUGUST 1, 1989

Models for universal reduction of macroscopic quantum fluctuations

L. Diósi

Central Research Institute for Physics, H-1525 Budapest 114, P.O. Box 49, Hungary (Received 17 October 1988; revised manuscript received 21 March 1989)

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use principles of relativity to **limit the lifetime of spatial quantum superpositions** and, as a result, **breaking the unitary evolution of the wavefunction**

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(Classical channel model

for Gravitational decoherence)

(Kafri *et.al.* arXiv:1401.0946 NJP - 2013)

A classical channel model for gravitational decoherence.

D. Kafri, J.M. Taylor¹ and G. J. Milburn²

Abstract

We show that, by treating the gravitational interaction between two mechanical resonators as a classical measurement channel, a gravitational decoherence model results that is equivalent to a model first proposed by Diosi. The resulting decoherence model implies that the classically mediated gravitational interaction between two gravitationally coupled resonators cannot create entanglement. The gravitational decoherence rate (and the complementary heating rate) is of the order of the gravitationally induced normal mode splitting of the two resonators.

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This approach is motivated by the fact that **gravity cannot be shielded** and therefore any observer can in principle **gain information about the quantum state sourcing gravity**

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(Classical channel model

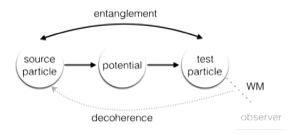
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Pirsa: 17060029 Page 17/43

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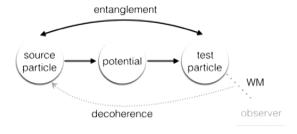
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This decoherence is **present in general** and not limited to gravity. It is also perfectly **compatible with unitary evolution**.

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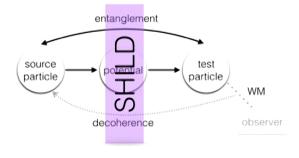
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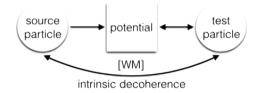
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Gravity can not be shielded



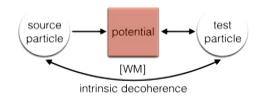
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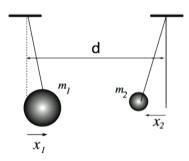
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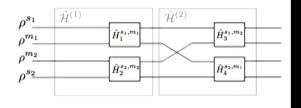
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$$\frac{d\hat{\rho}_{s_1s_2}}{dt} = -\frac{i}{\hbar}[\hat{H}_0 + \sum_i \Omega_i + K\hat{x}_1\hat{x}_2, \hat{\rho}_{s_1s_2}] - \left(\frac{1}{4D} + \frac{K^2D}{4\hbar^2}\right)\sum_i [\hat{x}_i, [\hat{x}_i, \hat{\rho}_{s_1s_2}]]$$

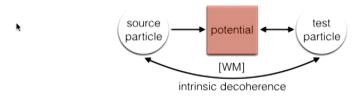
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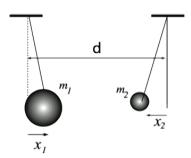
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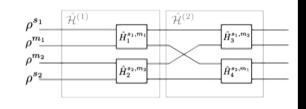
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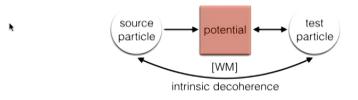
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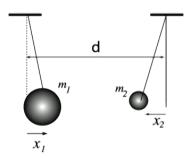
UNITARY

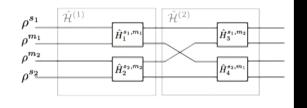
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$$-\left(\frac{1}{4D} + \frac{K^2D}{4\hbar^2}\right) \sum_{i} [\hat{x}_i, [\hat{x}_i, \hat{\rho}_{s_1 s_2}]]$$

UNITARY

DECOHERENCE

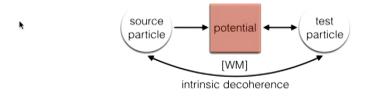
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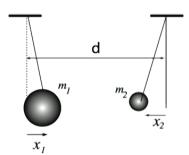
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IOP Institute of Physics
of Physics
of Physics

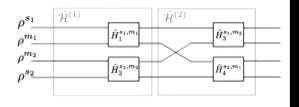
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PAPER

Unitarity, feedback, interactions—dynamics emergent from repeated measurements

Natacha Altamirano (20), Paulina Corona-Ugalde (1), Robert B Mann and Magdalena Zych





$$\frac{d\hat{\rho}_{s_1 s_2}}{dt} = -\frac{i}{\hbar} [\hat{H}_0 + \sum_i \Omega_i + K\hat{x}_1 \hat{x}_2, \hat{\rho}_{s_1 s_2}]$$

$$-\left(\frac{1}{4D} + \frac{K^2D}{4\hbar^2}\right) \sum_{i} [\hat{x}_i, [\hat{x}_i, \hat{\rho}_{s_1s_2}]]$$

UNITARY

DECOHERENCE

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[Review]

EINSTEIN EQUATIONS

$$R_{ab} - \frac{1}{2}Rg_{ab} = 8\pi GT_{ab}$$

METRIC

$$ds^{2} = -dt^{2} + a^{2}(t) \left(\frac{1}{1 - kr^{2}} dr^{2} + r^{2} d\Omega^{2} \right)$$

ENERGY-MOMENTUM TENSOR $T_{ab} = (\rho + P)u_au_b + Pg_{ab}$

$$T_{ab} = (\rho + P)u_a u_b + P g_{ab}$$

$$H^2 + \frac{k}{a^2} = \frac{8\pi G}{3}\rho$$

$$H^{2} + \frac{k}{a^{2}} = \frac{8\pi G}{3}\rho \qquad \qquad \frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3P) \qquad H = \frac{\dot{a}}{a}$$

$$H = \frac{\dot{a}}{a}$$

• EQUATION OF STATE $w(t) \longrightarrow P = w(t)\rho$

$$v(t) \longrightarrow P = w(t)$$

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}\rho(1+3w(t))$$

$$w>-1/3$$
 \longrightarrow decelerating

$$w<-1/3$$
 \longrightarrow accelerating

Class. Quantum Grav. 34 (2017) 115007 (19pp)

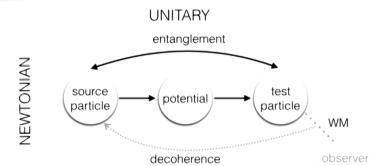
Emergent dark energy via decoherence in quantum interactions

Natacha Altamirano^{1,2}, Paulina Corona-Ugalde^{2,3}, Kiran E Khosla^{4,5}, Gerard J Milburn^{4,5} and Robert B Mann^{1,2}

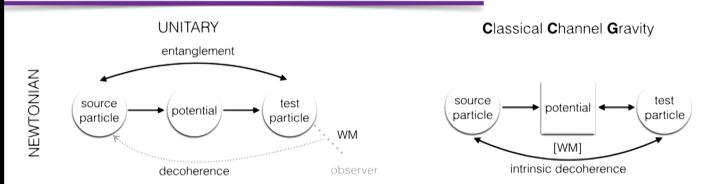
Abstract

In this work we consider a recent proposal that gravitational interactions are mediated via classical information and apply it to a relativistic context. We study a toy model of a quantized Friedman-Robertson-Walker (FRW) universe with the assumption that any test particles must feel a classical metric. We show that such a model results in decoherence in the FRW state that manifests itself as a dark energy fluid that fills the spacetime. Analysis of the resulting fluid, shows the equation of state asymptotically oscillates around the value w = -1/3, regardless of the spatial curvature, which provides the bound between accelerating and decelerating expanding FRW cosmologies. Motivated with quantum-classical interactions this model is yet another example of theories with violation of energy-momentum conservation whose signature could have significant consequences for the observable universe.

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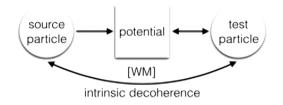
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UNITARY entanglement NEWTONIAN source test potential particle particle WM observer decoherence FRW scale metric test factor particle function entanglement

Classical Channel Gravity

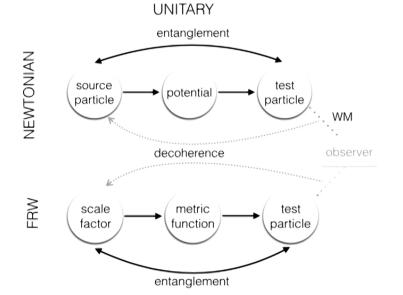


Wheeler De-Witt equation

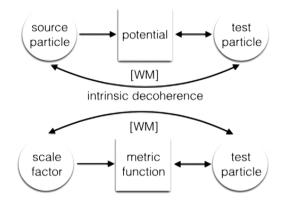
$$\hat{H}|\psi>=0$$

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WODEL



Classical Channel Gravity



Wheeler De-Witt equation

$$\hat{H}|\psi>=0$$

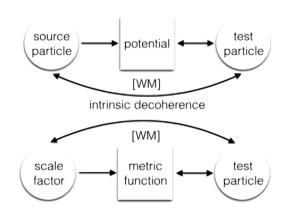
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UNITARY entanglement NEWTONIAN source test potential particle particle WM decoherence observer FRW scale metric test particle factor function entanglement

Wheeler De-Witt equation

$$\hat{H}|\psi>=0$$

Classical Channel Gravity



Master equation

$$d\hat{\rho} = -\frac{i}{\hbar}[\hat{H}, \hat{\rho}] - \frac{\gamma}{8\hbar}[\hat{a}^2, [\hat{a}^2, \hat{\rho}]]$$

Emergent dark energy fluid from quantum decoherence

(Altamirano *et.al.* arXiv:1605.05980 CQG-34,11 (2017))

Hilbert Space



CCG

$$d\hat{\rho} = -\frac{i}{\hbar}[\hat{H}, \hat{\rho}] - \frac{\gamma}{8\hbar}[\hat{a}^2, [\hat{a}^2, \hat{\rho}]]$$

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HAMILTONIAN:

• If no interaction is performed then we have an empty FRW

$$\hat{H} = -\frac{\hat{\pi}^2}{4} - k\hat{a}^2$$

CLASSICAL METRIC:

$$ds^{2} = <\hat{a}^{2}> \left[d\tau^{2} + \frac{1}{1 - kr^{2}}dr^{2} + r^{2}d\Omega^{2}\right]$$

Emergent dark energy fluid from quantum decoherence

(Altamirano et.al. CQG-34,11 (2017) arXiv:1605.05980)





WCQM

$$d\hat{\rho} = -\frac{i}{\hbar}[\hat{H}, \hat{\rho}] - \frac{\gamma}{8\hbar}[\hat{a}^2, [\hat{a}^2, \hat{\rho}]]$$

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Emergent dark energy fluid from quantum decoherence

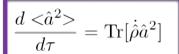
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WCQM

$$d\hat{\rho} = -\frac{i}{\hbar}[\hat{H}, \hat{\rho}] - \frac{\gamma}{8\hbar}[\hat{a}^2, [\hat{a}^2, \hat{\rho}]]$$



■DYNAMICS

$$\hat{H} = -\frac{\hat{\pi}^2}{4} - k\hat{a}^2$$

$$ds^2 = <\hat{a}^2 > [d\tau^2 + \frac{1}{1 - kr^2}dr^2 + r^2d\Omega^2]$$

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(Altamirano *et.al. CQG-34,11 (2017)* arXiv:1605.05980)

Hilbert Space



WCQM

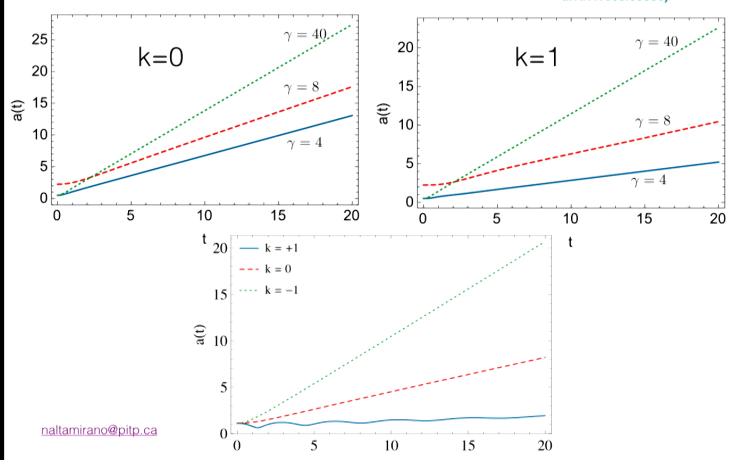
$$d\hat{\rho} = -\frac{i}{\hbar}[\hat{H}, \hat{\rho}] - \frac{\gamma}{8\hbar}[\hat{a}^2, [\hat{a}^2, \hat{\rho}]]$$

 $\frac{d < \hat{a}^2 >}{d\tau} = - < \hat{a}\hat{\pi} + \hat{\pi}\hat{a} > /2 ,$ $\frac{d < \hat{\pi}^2 >}{d\tau} = 2k < \hat{a}\hat{\pi} + \hat{\pi}\hat{a} > + \gamma < \hat{a}^2 > ,$ $\frac{d < \hat{a}\hat{\pi} + \hat{\pi}\hat{a} >}{d\tau} = - < \hat{\pi}^2 > + 4k < \hat{a}^2 > .$ DYNAMICS

 $\hat{H} = -\frac{\hat{\pi}^2}{4} - k\hat{a}^2$ $ds^2 = <\hat{a}^2 > [d\tau^2 + \frac{1}{1 - kr^2}dr^2 + r^2d\Omega^2]$

Emergent dark energy fluid from quantum decoherence

(Altamirano *et.al.* CQG-34,11 (2017) arXiv:1605.05980)



Pirsa: 17060029

Emergent dark energy fluid from quantum decoherence

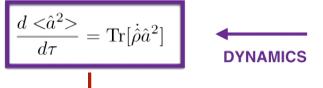
(Altamirano et.al. CQG-34,11 (2017) arXiv:1605.05980)





WCQM

$$d\hat{\rho} = -\frac{i}{\hbar}[\hat{H}, \hat{\rho}] - \frac{\gamma}{8\hbar}[\hat{a}^2, [\hat{a}^2, \hat{\rho}]]$$



$$\hat{H} = -\frac{\hat{\pi}^2}{4} - k\hat{a}^2$$

$$ds^{2} = <\hat{a}^{2}> \left[d\tau^{2} + \frac{1}{1-kr^{2}}dr^{2} + r^{2}d\Omega^{2}\right]$$

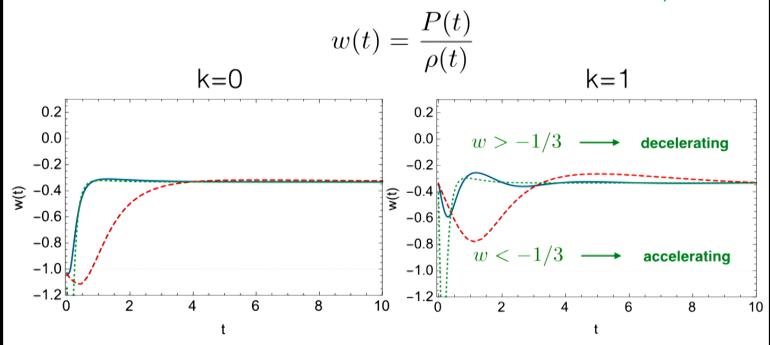
$$\frac{d\hat{H}}{d\tau} = -\frac{\gamma}{4} < \hat{a}^2 >$$

$$G_{ab}(\langle \hat{a}^2 \rangle) = 8\pi T_{ab}$$

Emergent dark energy fluid from

quantum decoherence

(Altamirano *et.al.* CQG-34,11 (2017) arXiv:1605.05980)

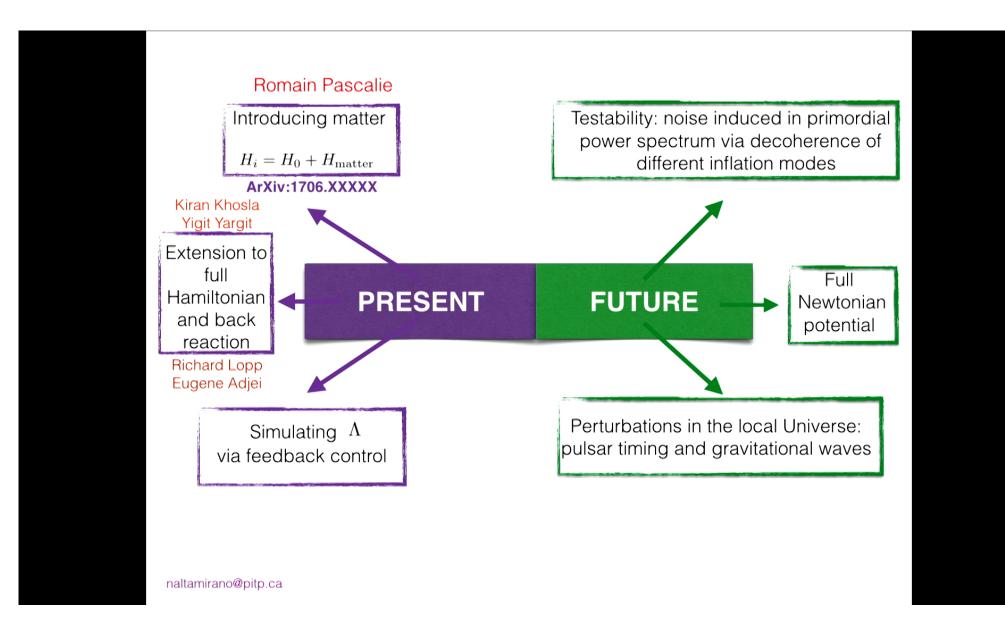




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REMARKS

$$H_i = H_0 + H_{\text{matter}}$$

Introducing matter $\dot{
ho}=rac{i}{\hbar}[H_0+rac{\Lambda}{3}a^4,
ho]+\Big(rac{1}{4D}+rac{D\Lambda^2}{9\hbar^2}\Big)[a^2,[a^2,
ho]]$

Romain Pascalie

Kiran Khosla

PRESENT

Simulating Λ via feedback control

$$H_{\Lambda} = -ka^2 - \frac{p^2}{4} + \frac{1}{3}\Lambda a^4$$

$$\begin{split} \dot{\rho}(t) = & -\frac{i}{\hbar} [\hat{S}_0, \rho(t)] - \frac{i}{\hbar} \frac{\bar{g}_2}{4} [\hat{S}_2, \hat{S}_1 \rho(t) + \rho(t) \hat{S}_1] \\ & - \frac{1}{8D} [\hat{S}_1, [\hat{S}_1, \rho(t)]] - \frac{D}{\hbar^2} \frac{\bar{g}_2^2}{8} [\hat{S}_2, [\hat{S}_2, \rho(t)]] \,. \end{split}$$

Emergent dark energy fluid from

quantum decoherence

(Altamirano *et.al.* CQG-34,11 (2017) arXiv:1605.05980)

