

Title: Impact of neutrino interaction uncertainties

Date: Jun 13, 2017 11:30 AM

URL: <http://pirsa.org/17060009>

Abstract:

# Why are neutrino interactions important for neutrino physics?

*Radiative Corrections at the Intensity Frontier of Particle Physics*

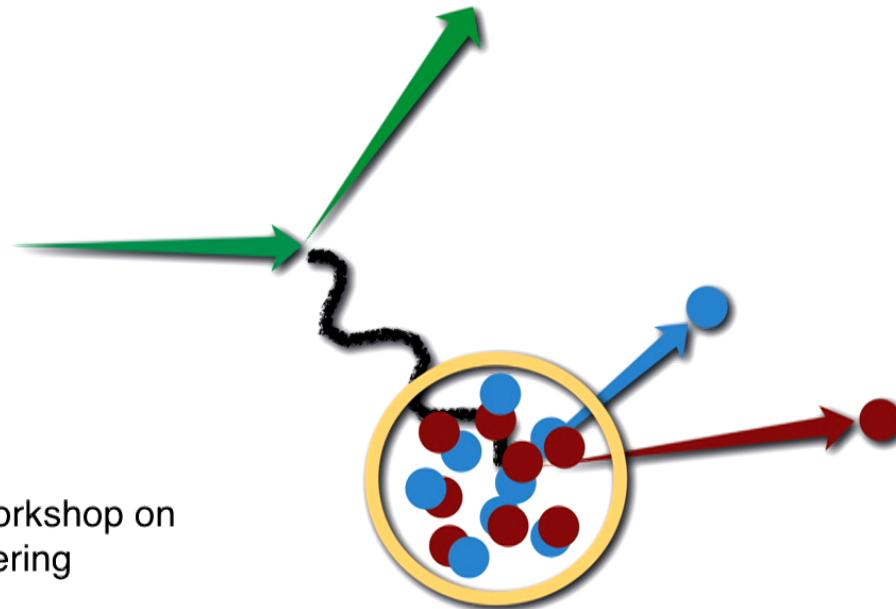


Image from workshop on  
neutrino scattering  
(NuInt2014)

Kendall Mahn  
Michigan State University

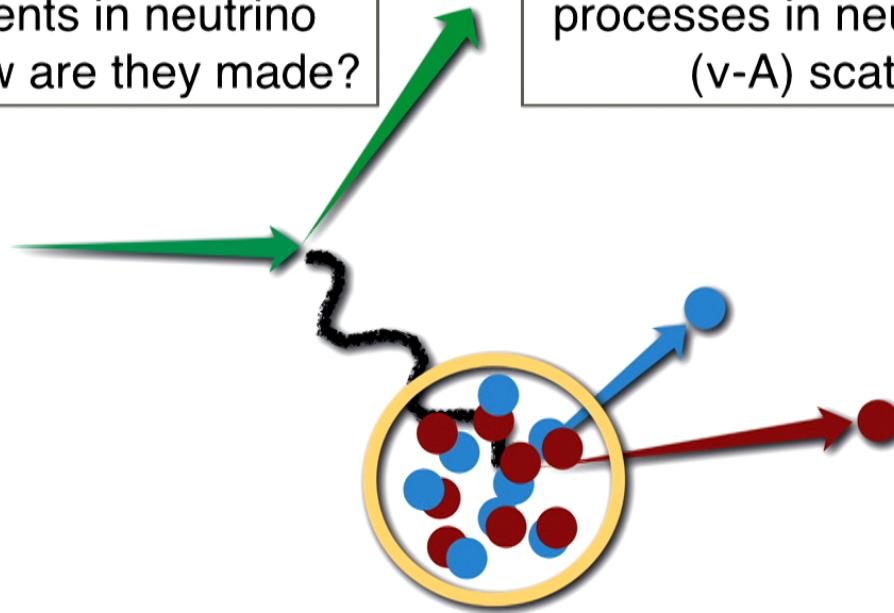


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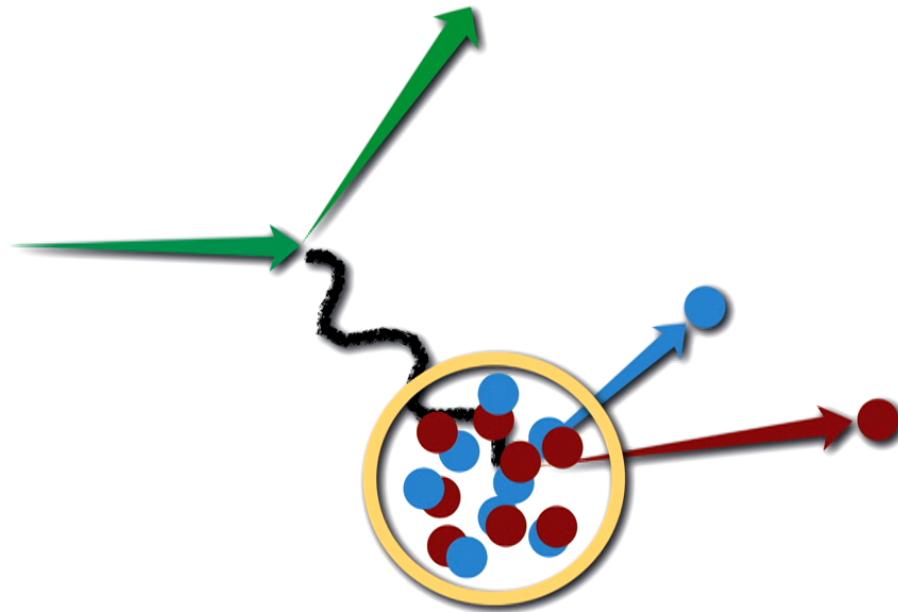
What are the interesting measurements in neutrino physics? How are they made?

What are the relevant processes in neutrino nuclear ( $\nu$ -A) scattering?



Why is  $\nu$ -A scattering important to neutrino oscillation and cross section expt's?

# Neutrino oscillation

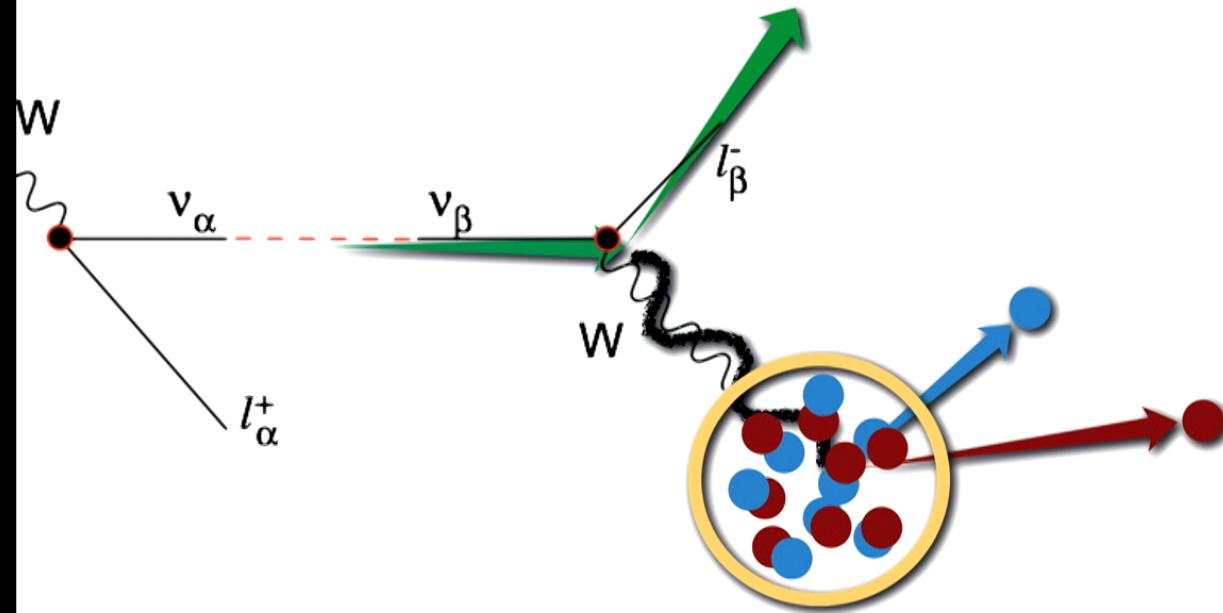


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# Neutrino oscillation

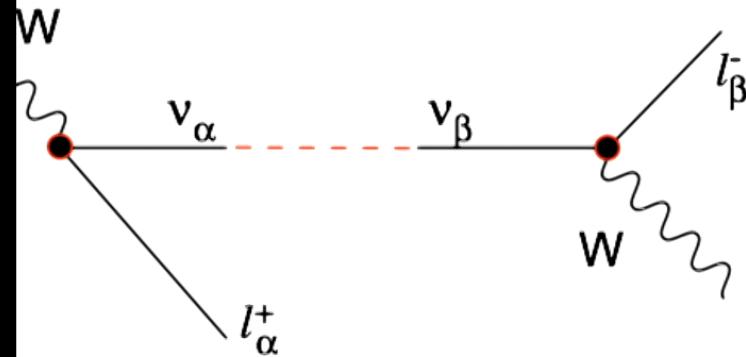


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# Neutrino oscillation



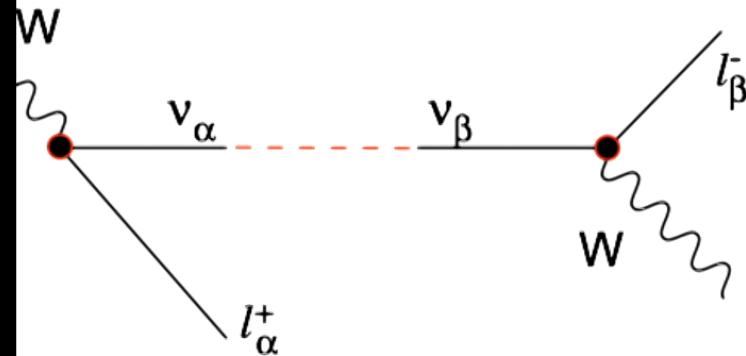
- Interference between mass and flavor eigenstates
- Appearance ( $\beta \neq \alpha$ ) or disappearance ( $\beta = \alpha$ )

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# Neutrino oscillation



$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

- Interference between mass and flavor eigenstates
- Appearance ( $\beta \neq \alpha$ ) or disappearance ( $\beta = \alpha$ )

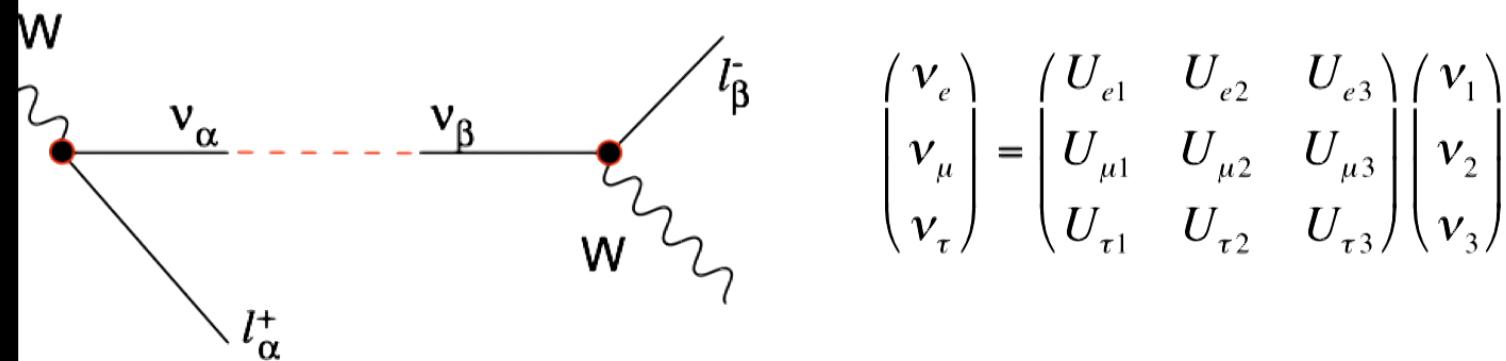
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# Neutrino oscillation

$$P_{\alpha\beta} = \delta_{\alpha\beta} - 4 \sum_{i>j} \text{Re} \left[ U_{\beta i} U_{\alpha i}^* U_{\beta j}^* U_{\alpha j} \right] \sin^2 \left( \frac{1.27 \Delta m_{ij}^2 L}{E} \right) + 2 \sum_{i>j} \text{Im} \left[ U_{\beta i} U_{\alpha i}^* U_{\beta j}^* U_{\alpha j} \right] \sin \left( \frac{2.54 \Delta m_{ij}^2 L}{E} \right)$$



- Interference between mass and flavor eigenstates
- Appearance ( $\beta \neq \alpha$ ) or disappearance ( $\beta = \alpha$ )

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# Neutrino oscillation

- Three rotation angles ( $\theta_{12}$ ,  $\theta_{13}$ ,  $\theta_{23}$ )
  - $\theta_{12}$ : solar and reactor experiments
  - $\theta_{13}$ : reactor and long-baseline (accelerator) experiments
  - $\theta_{23}$ : long-baseline and atmospheric experiments
- One complex phase  $\delta_{\text{CP}}$ 
  - additional phases if neutrinos are “Majorana”, not measured in oscillation experiments

Standard factorization of  $U$

# Neutrino oscillation

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$$U = \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{+i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\alpha_1/2} & 0 \\ 0 & 0 & e^{i\alpha_2/2} \end{pmatrix}$$

Standard factorization of  $U$

# Current knowledge and open questions

Parameter	best-fit	$3\sigma$
$\Delta m_{21}^2$ [10 <sup>-5</sup> eV <sup>2</sup> ]	7.37	6.93 – 7.97
$ \Delta m^2 $ [10 <sup>-3</sup> eV <sup>2</sup> ]	2.50 (2.46)	2.37 – 2.63 (2.33 – 2.60)
$\sin^2 \theta_{12}$	0.297	0.250 – 0.354
$\sin^2 \theta_{23}, \Delta m^2 > 0$	0.437	0.379 – 0.616
$\sin^2 \theta_{23}, \Delta m^2 < 0$	0.569	0.383 – 0.637
$\sin^2 \theta_{13}, \Delta m^2 > 0$	0.0214	0.0185 – 0.0246
$\sin^2 \theta_{13}, \Delta m^2 < 0$	0.0218	0.0186 – 0.0248
$\delta/\pi$	1.35 (1.32)	(0.92 – 1.99) ((0.83 – 1.99))

Known: separation of solar  
and atmospheric splittings

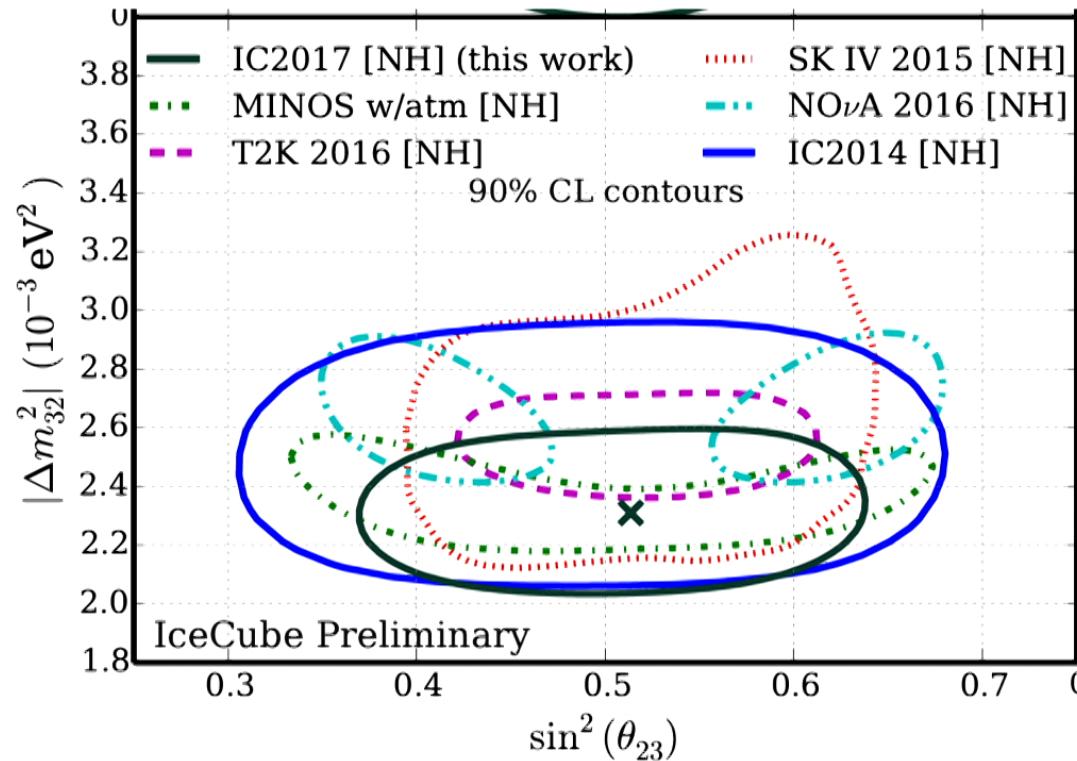
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Known: separation of solar  
and atmospheric splittings

Unknown: ordering of mass  
splittings (mass hierarchy)

# Current knowledge and open questions



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# The experiments

Current generation:

Accelerator-based: Tokai-to-Kamioka (T2K), NOvA

Atmospheric: IceCube, Super-Kamiokande

Reactor: Daya Bay, RENO, Double Chooz

Future:

Accelerator-based (+ atmospheric): DUNE, Hyper-Kamiokande

Atmospheric: PINGU

Reactor: JUNO

# Oscillation probabilities

$$P(\nu_\mu \rightarrow \nu_\mu) \sim 1 - (\cos^4 \theta_{13} \sin^2 2\theta_{23} + \sin^2 2\theta_{13} \sin^2 \theta_{23}) \sin^2 \Delta m_{31}^2 \frac{L}{4E}$$

- Precision measurement of  $2\theta_{23}$  and  $\Delta m_{31}^2$
- CPT tests with antineutrino mode ( $\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu$ )

$$\begin{aligned} P(\nu_\mu \rightarrow \nu_e) &\sim \boxed{\sin^2 2\theta_{13}} \times \boxed{\sin^2 \theta_{23}} \times \boxed{\frac{\sin^2[(1-x)\Delta]}{(1-x)^2}} \\ &\quad \boxed{-\alpha \sin \delta} \times \sin 2\theta_{12} \sin 2\theta_{13} \sin 2\theta_{23} \times \sin \Delta \frac{\sin[x\Delta]}{x} \frac{\sin[(1-x)\Delta]}{(1-x)} \\ &\quad + \alpha \cos \delta \times \sin 2\theta_{12} \sin 2\theta_{13} \sin 2\theta_{23} \times \cos \Delta \frac{\sin[x\Delta]}{x} \frac{\sin[(1-x)\Delta]}{(1-x)} \\ &\quad + \mathcal{O}(\alpha^2) \end{aligned}$$

M. Freund, Phys.Rev. D64 (2001) 053003

$$\alpha = \left| \frac{\Delta m_{21}^2}{\Delta m_{31}^2} \right| \sim \frac{1}{30} \quad \Delta \equiv \frac{\Delta m_{31}^2 L}{4E} \quad x \equiv \frac{2\sqrt{2}G_F N_e E}{\Delta m_{31}^2}$$

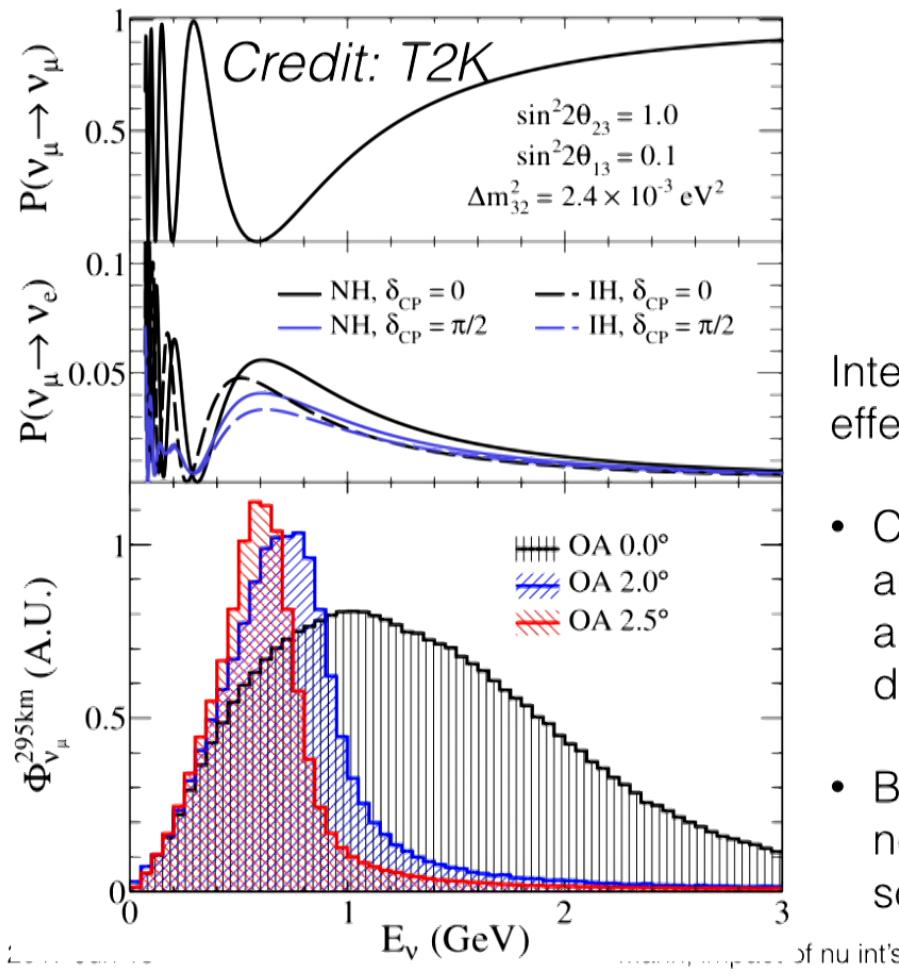
- $\sin^2 2\theta_{13}$  dependence of leading term
- $\theta_{23}$  dependence of leading term: "octant" dependence ( $\theta_{23} = />/<45^\circ?$ )
- CP odd phase  $\delta$ : asymmetry of probabilities  $P(\nu_\mu \rightarrow \nu_e) \neq P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e)$  if  $\sin \delta \neq 0$
- Matter effect through  $x$ :  $\nu_e (\bar{\nu}_e)$  enhanced in normal (inverted) hierarchy

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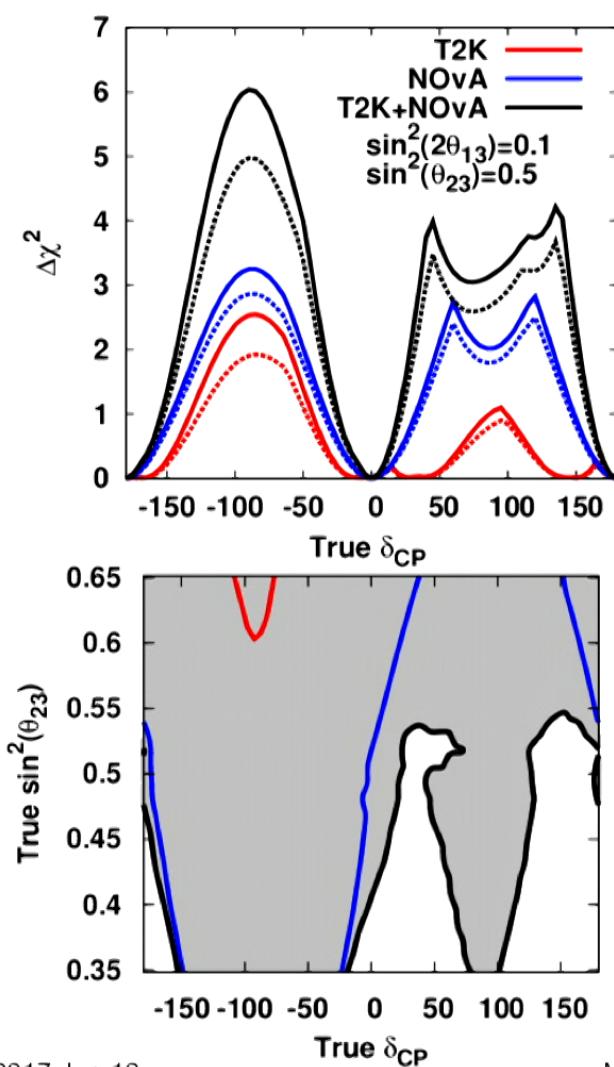
# Details of oscillation experiments



Interdependence of different effects in oscillation probability

- Combined analysis of neutrino and antineutrino data, appearance and disappearance samples
- But, fluxes are different for neutrino vs. antineutrino sources

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- Combinations of different baselines/energies give enhanced physics reach

# Aside: What do we know about CP violation?

*Slides from Tanaka, Neutrino 2016*

- Observe

- more  $\nu_e$  candidates than predicted
- fewer  $\bar{\nu}_e$  candidates than predicted

in the case of NH,  $\delta_{CP} = -\pi/2$  that induces the largest asymmetry

observed vs. expected number of  $\nu_e$  and  $\bar{\nu}_e$  candidates

	OBS.	$\delta_{CP}=-\pi/2$	$\delta_{CP}=0$	$\delta_{CP}=+\pi/2$	$\delta_{CP}=\pi$
$\nu_e$	<b>32</b>	27.0	22.7	18.5	22.7
$\bar{\nu}_e$	<b>4</b>	6.0	6.9	7.7	6.8

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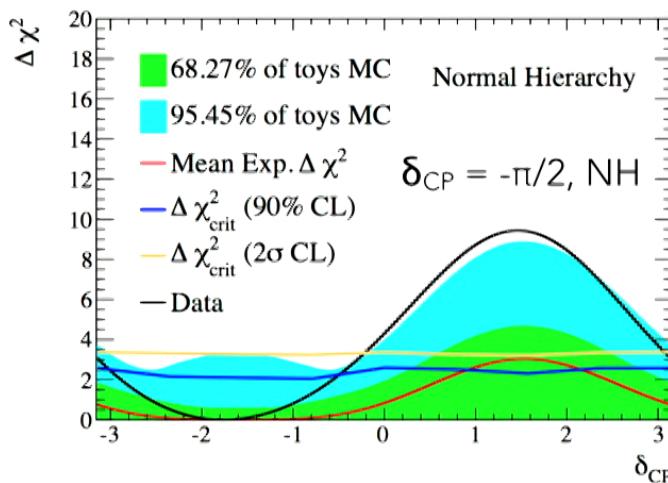
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	OBS.	EXPECTED (NH, $\sin^2 \Theta_{23}=0.528$ )			
		$\delta_{CP}=-\pi/2$	$\delta_{CP}=0$	$\delta_{CP}=+\pi/2$	$\delta_{CP}=\pi$
$\nu_e$	<b>32</b>	27.0	22.7	18.5	22.7
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- Toy MC run to assess probability of outcome given a set of "true" parameters
- Below: fraction where  $\delta_{CP} = 0$  excluded at 90% or 2  $\sigma$  CL for NH,  $\delta_{CP} = -\pi/2, 0$

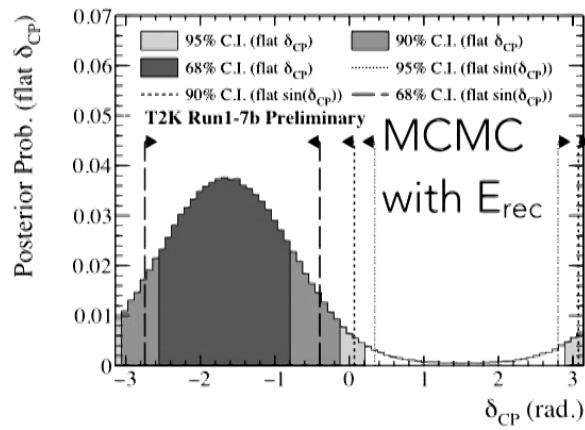
	TRUE PARAMETERS	
	$\delta_{CP}=-\pi/2, \text{NH}$	$\delta_{CP}=0, \text{NH}$
90%	0.187	0.102
2 $\sigma$	0.089	0.047

Impact of nu int's

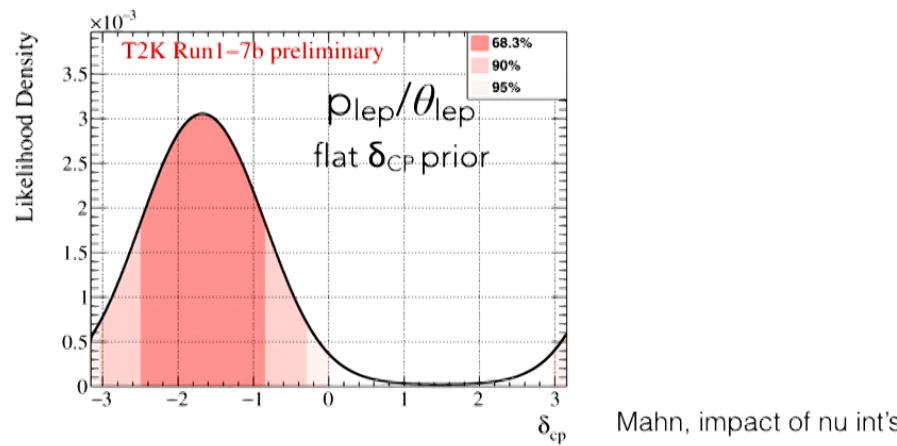
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## Aside: What do we know about CP violation?

*Slides from Tanaka, Neutrino 2016*



- Left: posterior probability distribution in  $\delta_{CP}$  marginalizing over all other parameters
  - negligible dependence on priors except for  $\delta_{CP}$ 
    - (flat in  $\delta_{CP}$  vs.  $\sin \delta_{CP}$ )

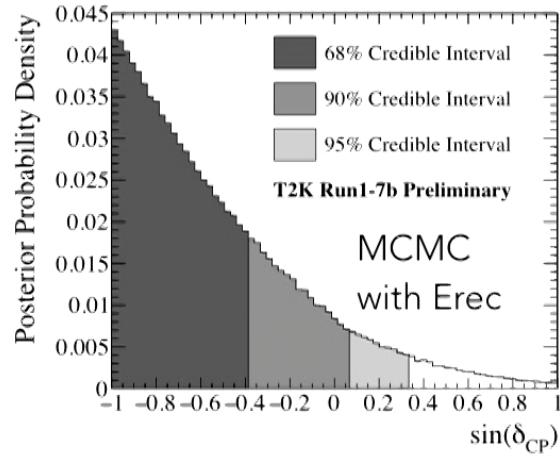


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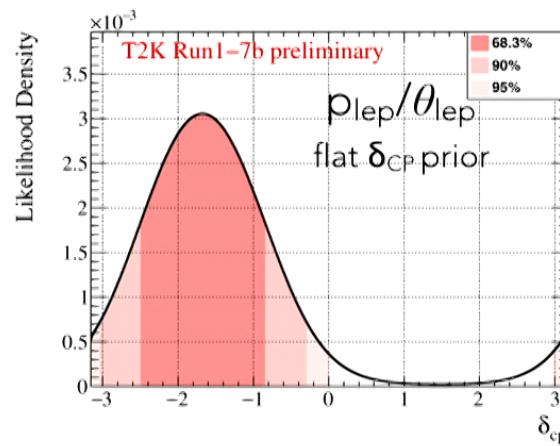
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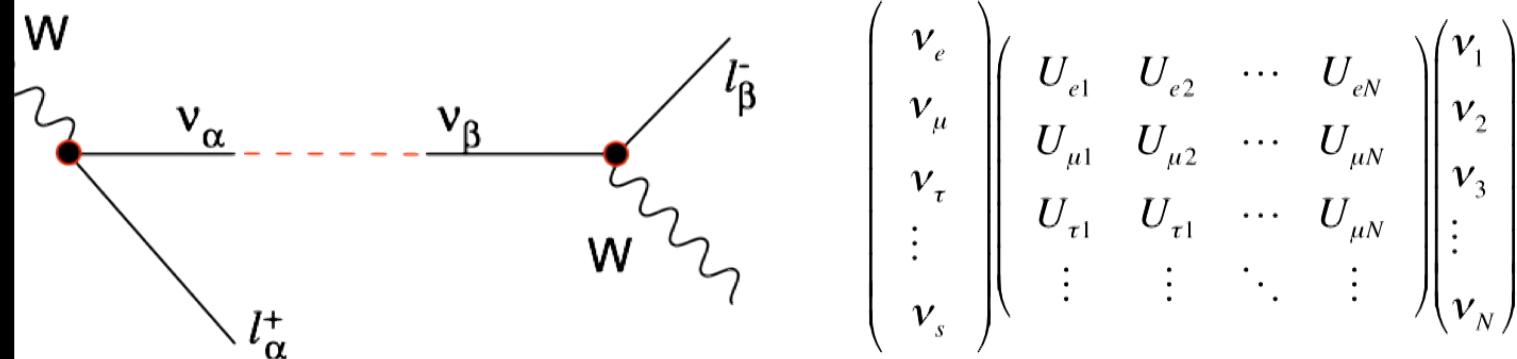
*Slides from Tanaka, Neutrino 2016*

- Posterior probability distributions for  $\theta_{23}$  octant and hierarchy with MCMC analysis
  - mild preference for  $\theta_{23} > \pi/4$  and normal hierarchy

	NH	IH	SUM
$\sin^2\theta_{23} \leq 0.5$	0.218	0.072	0.290
$\sin^2\theta_{23} > 0.5$	0.529	0.181	0.710
SUM	0.747	0.253	1.000

# What about non-standard effects?

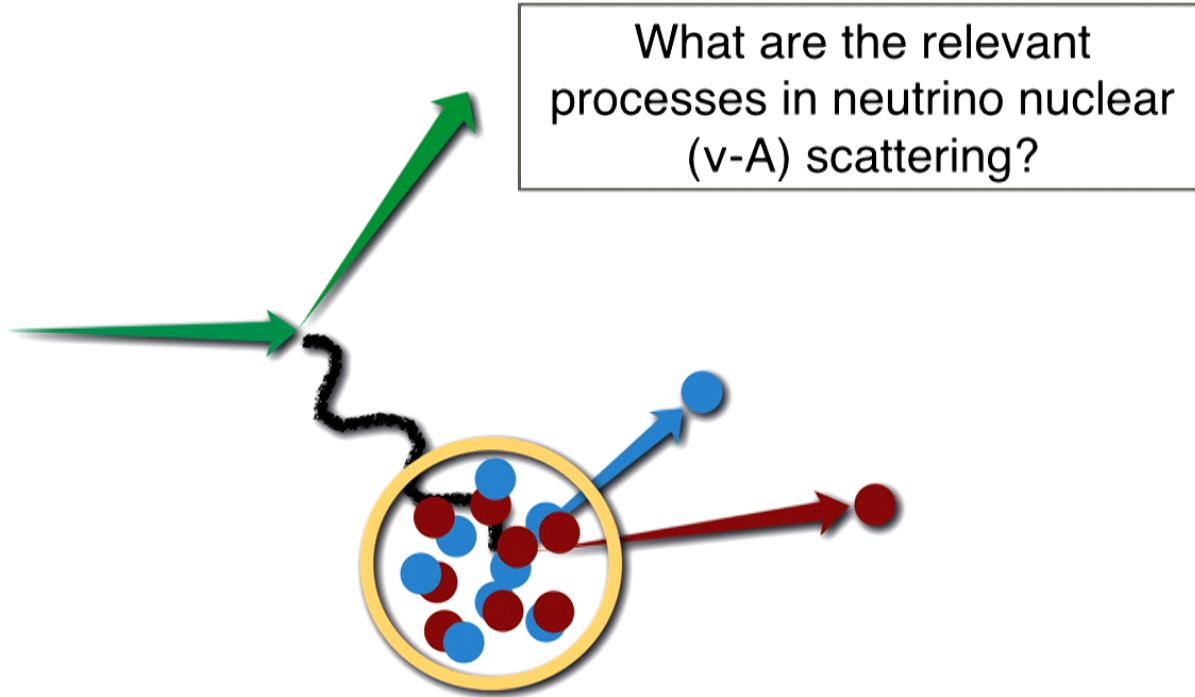
*See Andre's talk*



- Light sterile neutrinos: anomalous appearance or disappearance through sterile. Similar setup to “standard oscillation” except shorter distances
  - Short baseline neutrino program at Fermilab
  - Long baseline neutral current disappearance
- Non-Standard Interactions (at source, or detector via matter effects), often constrained by charged lepton data
- Tests of unitarity of mixing matrix

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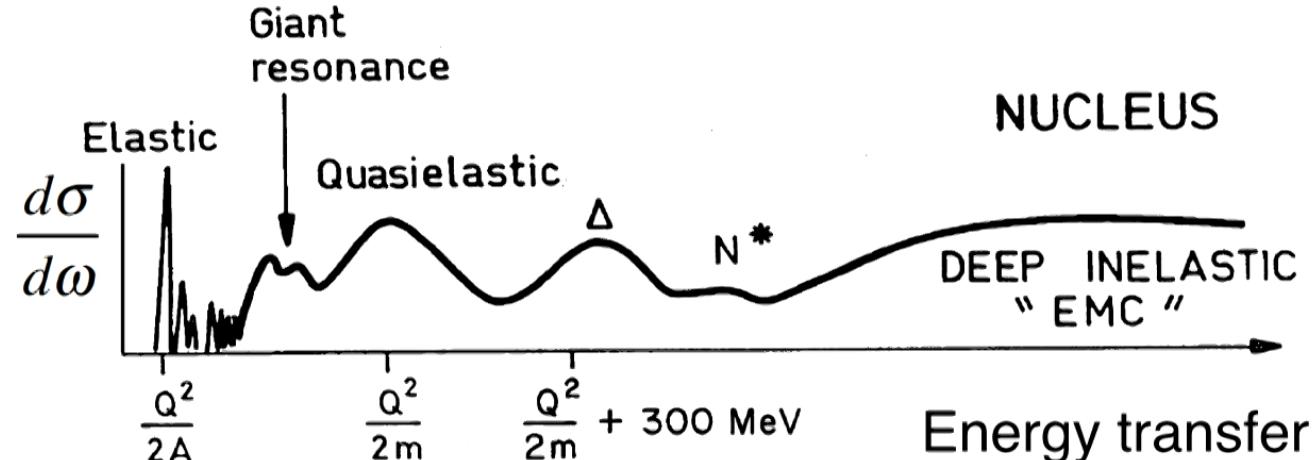
For these details and more, (Neutrino Scattering Theory Experiment Collaboration (NuSTEC) White Paper:

<https://arxiv.org/abs/1706.03621>

Subscribe to newsletter (backups)

# Processes in Neutrino Scattering

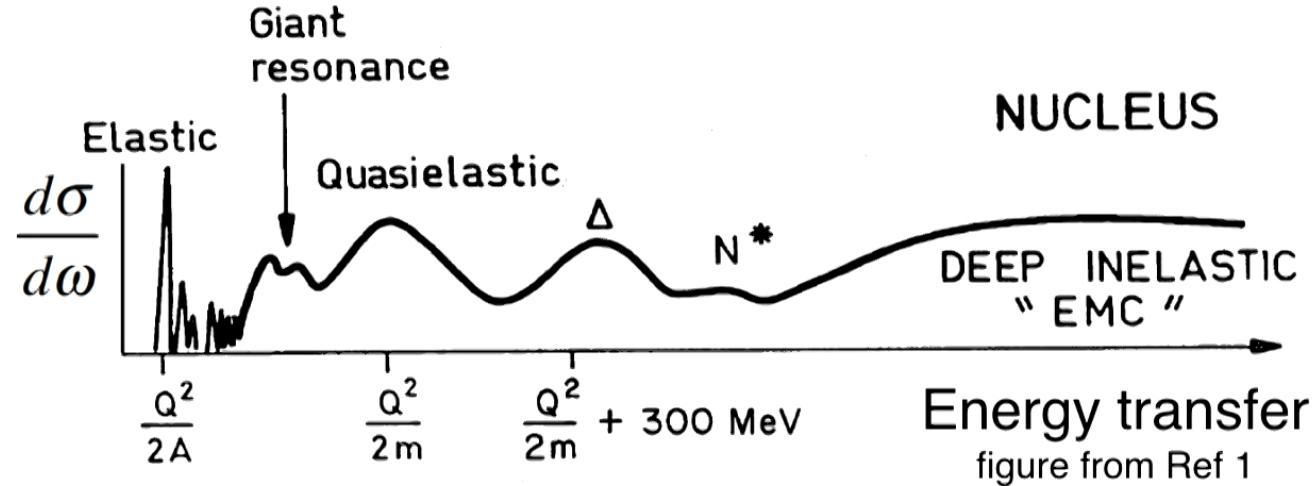
B Foris, and C N Papanicolas, Electron Scattering and Nuclear Structure



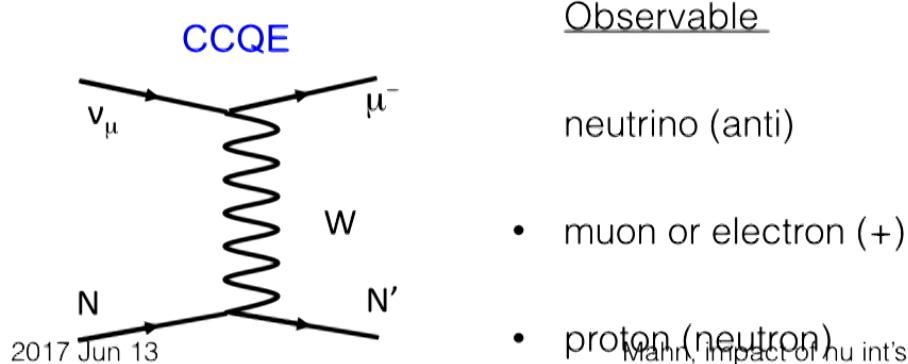
- Reactor energies: cross section well known
- Need charged current (CC) interactions to detect flavor
- Neutral current (NC) also relevant (backgrounds, non-standard oscillation)

# Processes in Neutrino Scattering

B Foris, and C N Papanicolas, Electron Scattering and Nuclear Structure



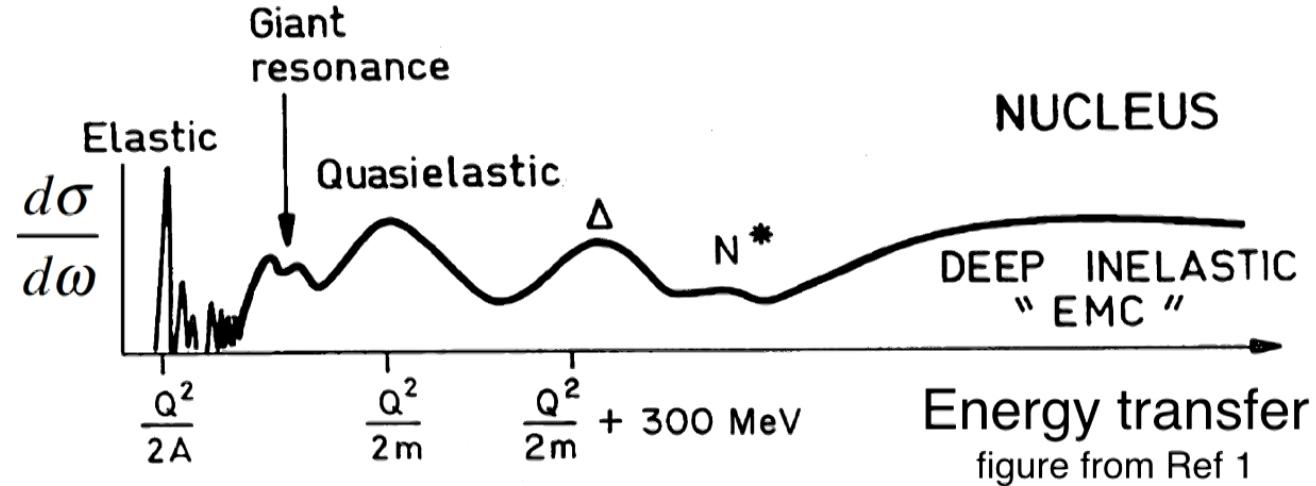
- Charged Current Quasi Elastic (CCQE) and multinucleon processes (2p2h)



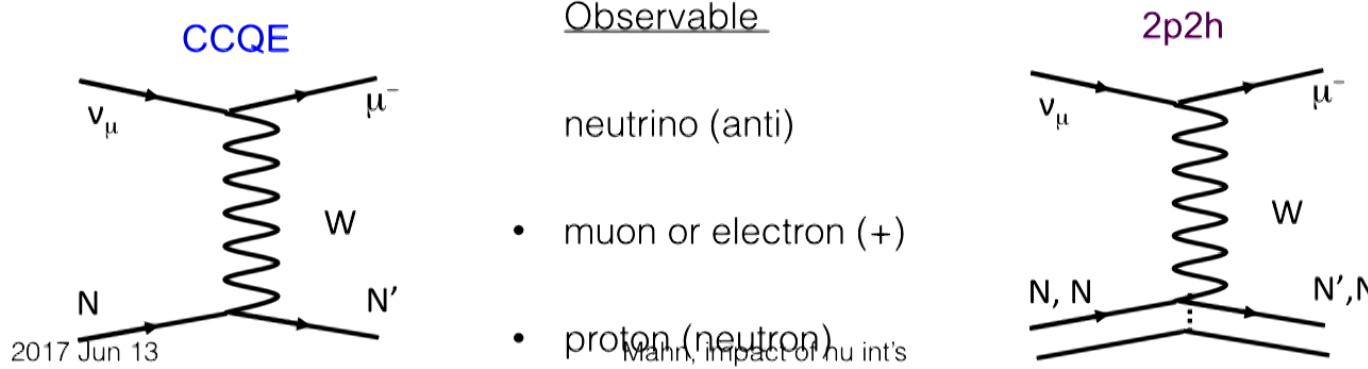
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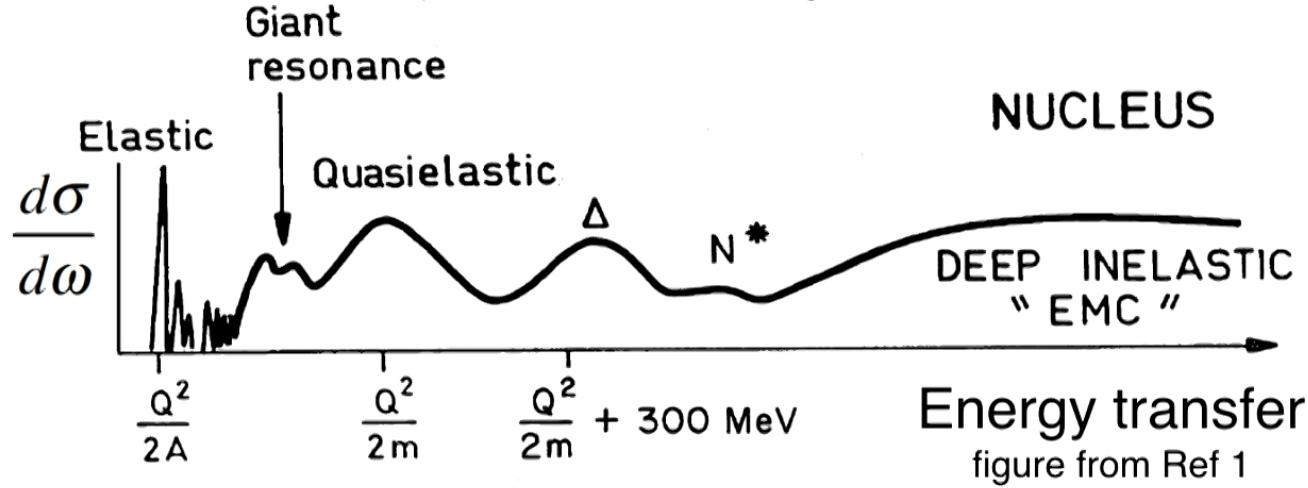


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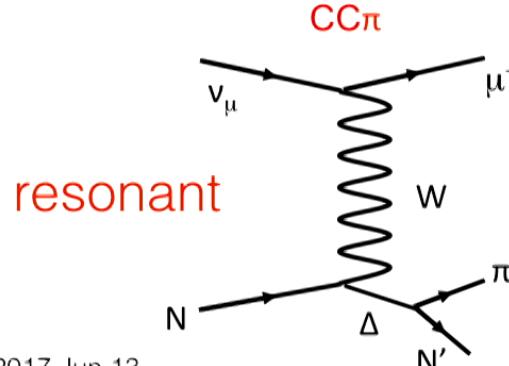


# Processes in Neutrino Scattering

B Foris, and C N Papanicolas, Electron Scattering and Nuclear Structure



- Production of pions, CC1 $\pi^{+/0/-}$  and NC1 $\pi^{+/0/-}$



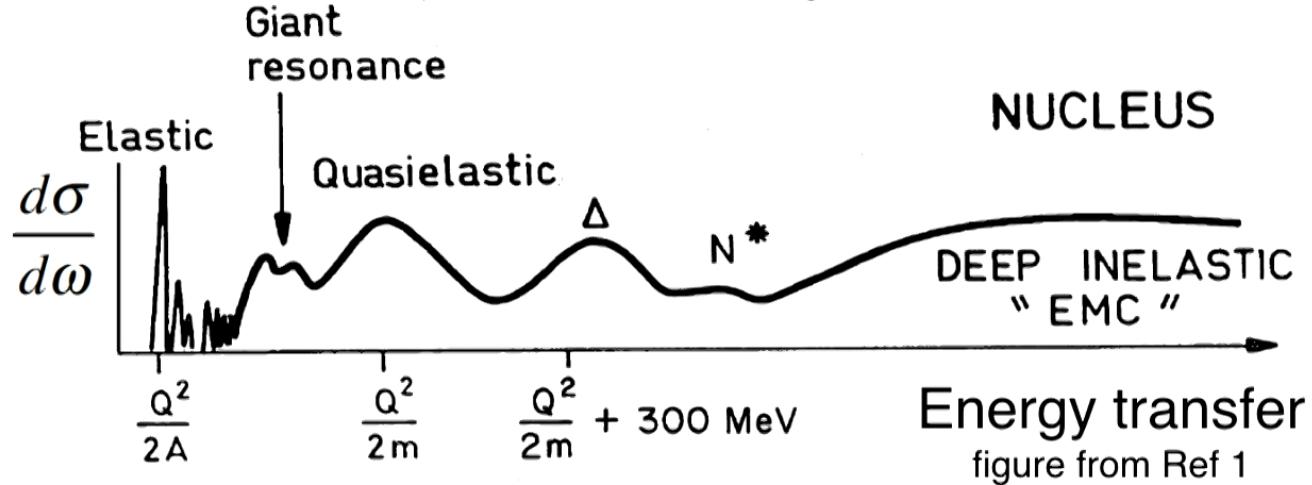
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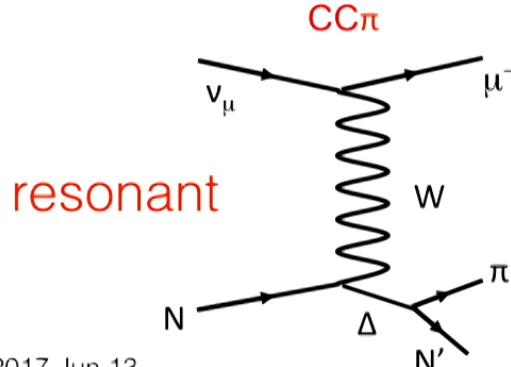
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# Processes in Neutrino Scattering

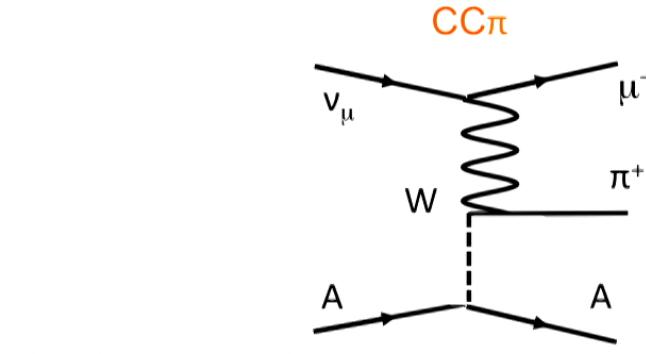
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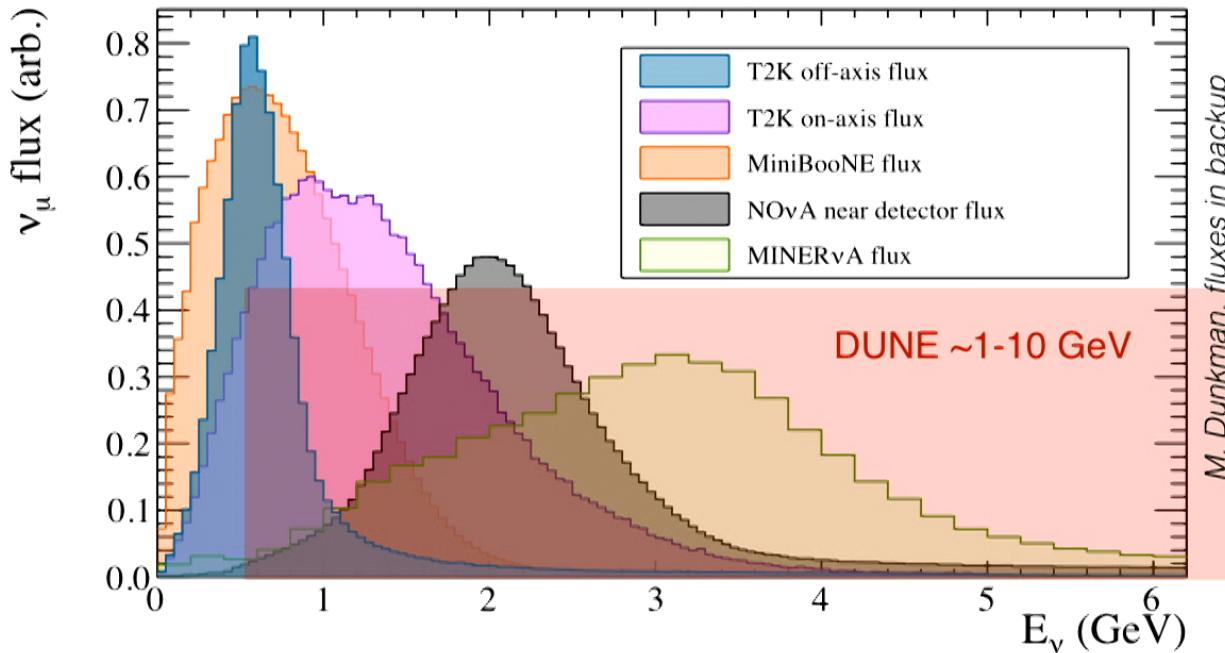
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# Neutrino Sources and Nuclear Effects



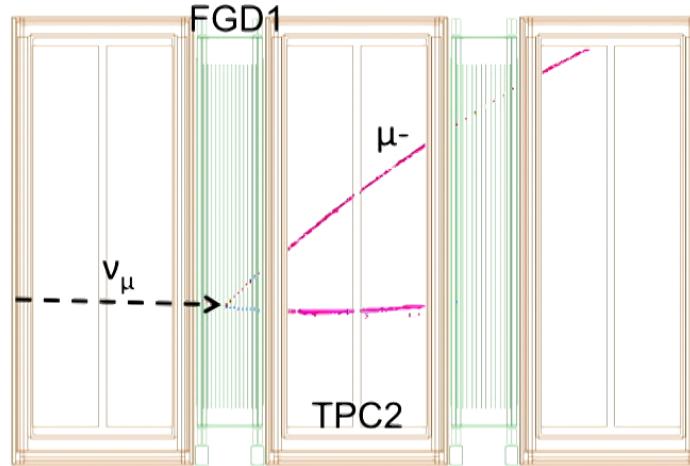
- Neutrino decay-in-flight beams are **not mono-energetic**
  - Spread of beam is larger than nuclear effects
- **Nuclear target** needed for experimental statistical sample size

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# Nuclear effects: process vs. topology



T2K experiment example

- CC0 $\pi$  “topology”: 1 muon, no pion
- Includes CCQE, 2p2h, CC1 $\pi$  (pion absorbed in nucleus)

# An overly generic oscillation analysis

$$N_{\text{FD}}^{\alpha \rightarrow \beta}(\mathbf{p}_{\text{reco}}) = \sum \phi_\alpha(E_{\text{true}}) \times P_{\alpha\beta}(E_{\text{true}}) \times \sigma_\beta^i(\mathbf{p}_{\text{true}}) \times \epsilon_\beta(\mathbf{p}_{\text{true}}) \times R_i(\mathbf{p}_{\text{true}}; \mathbf{p}_{\text{reco}})$$

Rate is a combination of:

- Flux
- Oscillation probability
- All processes which contribute to the topology
- Efficiency
- Relationship (R) between inferred and true energy of the process from kinematics

# An overly generic oscillation analysis

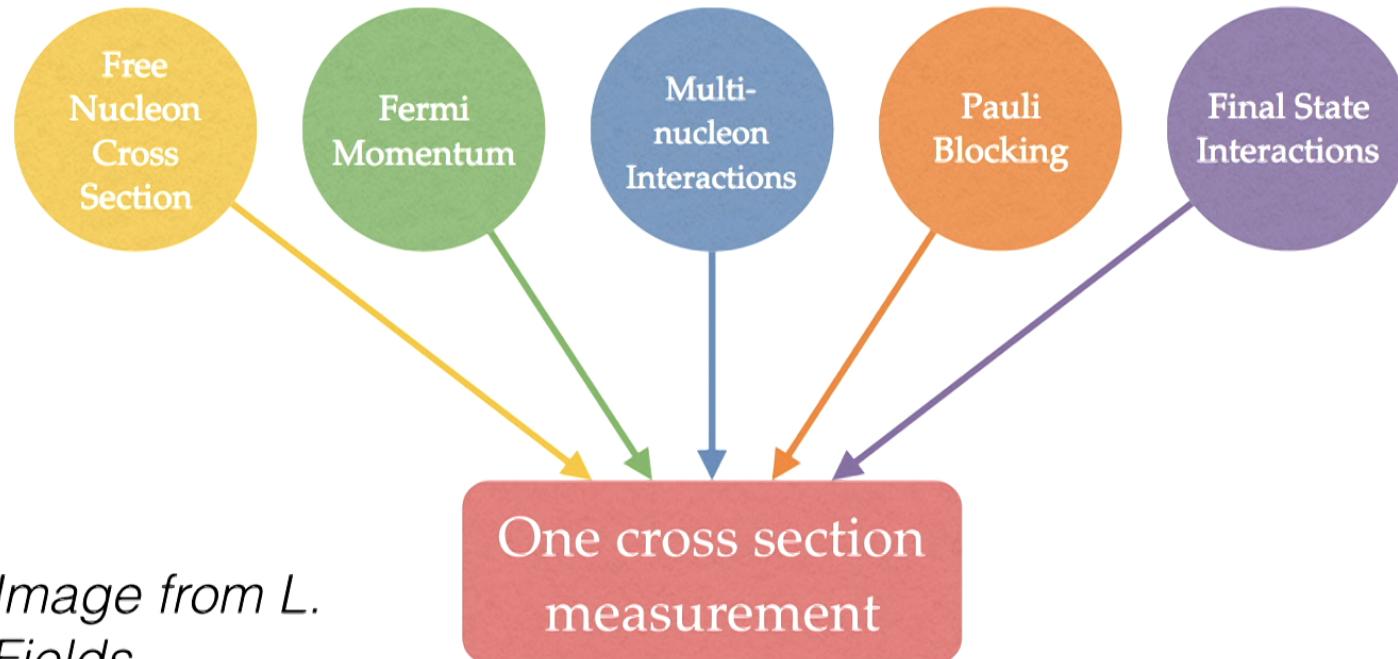
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$$N_{\text{ND}}^\alpha(\mathbf{p}_{\text{reco}}) = \sum \phi_\alpha(E_{\text{true}}) \times \sigma_\alpha^i(\mathbf{p}_{\text{true}}) \times \epsilon_\alpha(\mathbf{p}_{\text{true}}) \times R_i(\mathbf{p}_{\text{true}}; \mathbf{p}_{\text{reco}})$$

Near detector provides event rate (constrains flux, cross section and (some) of the detector response). Inherent difficulties:

1. Energy dependance enters in integral over energy due to  $P$
2.  $\nu_e$  appearance (but ND measures  $\nu_\mu$  rate)
3. Not pure flavor beam (neutrino and antineutrino contributions)
4. Wide flux spectrum (not possible to isolate cross section processes)
5. Small differences in flux and detector response of ND and FD
6. Correct association between true and reconstructed variables

# The problem



*Image from L.  
Fields*

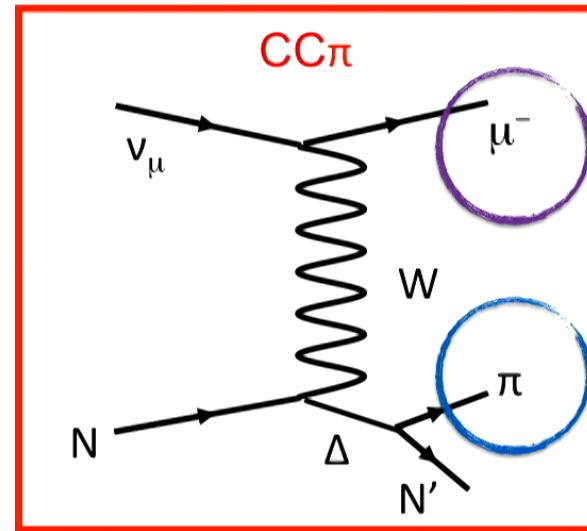
- Measurements are sensitive to multiple physics effects on signal and background processes

# Example: Single Pion Production Puzzles

- Important background (and signal) process for oscillation (non-standard) physics
- **Challenge to model!**
  - non-resonant backgrounds?
  - pion re-interaction (final state interaction) model?
  - Busted single nucleon model?
- Example from NUISANCE software framework in references thanks to C. Wilkinson - *work in progress*
- *Related work in Gabe's talk*

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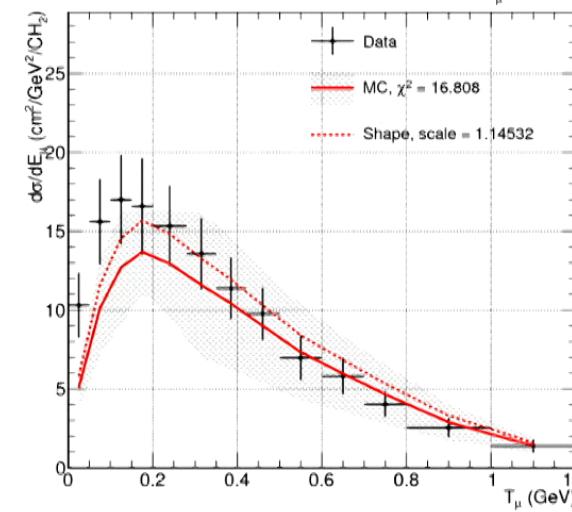
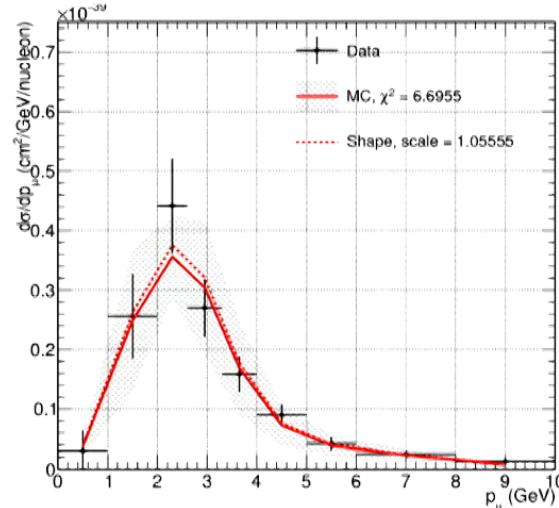
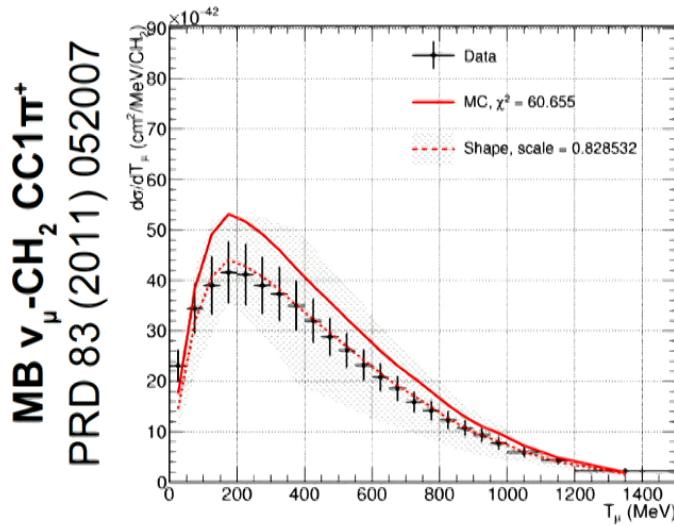
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# Single Pion Production Puzzles

- Reasonable agreement in outgoing muon spectrum
- Terrible agreement in outgoing pion spectrum
- Model development essential



# TENSIONS2016 Workshop

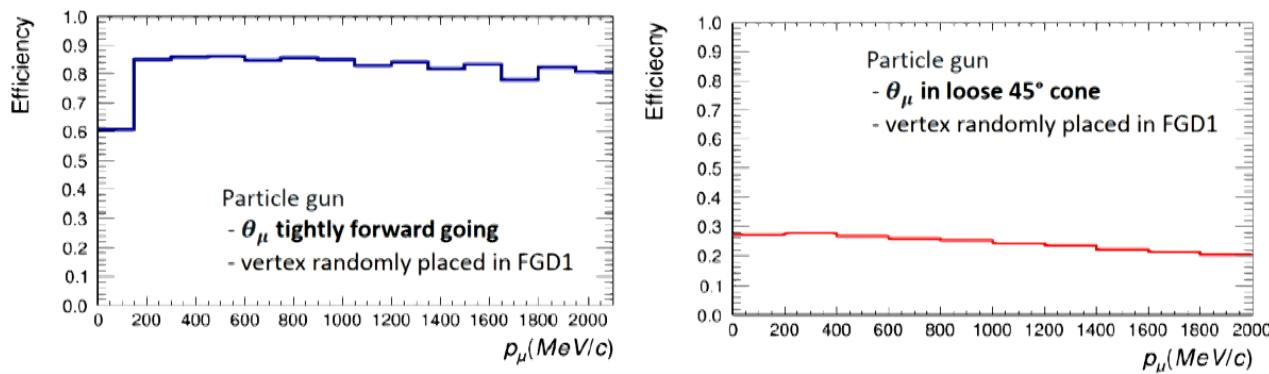
Inability to reconcile MiniBooNE, MINERvA QE-like, and single pion measurements within a single model (“data tensions”)

- T2K attempts to fit [QE: PRD93 no.7, 072010 (2016)]
- Single pion pain

Motivated a workshop (writeup underway) to understand:

- Differences in signal definition? Selection? Extraction?
- **Role of efficiency in cross section measurements**

- **Example 1** – want to measure  $p_\mu$  for single muons using TPC.
  - The efficiency is very dependent on the underlying  $\theta_\mu$  distribution.
  - The underlying  $\theta_\mu$  distribution depends on the neutrino scattering model



*Credit: S. Dolan, T2K-XSEC workshop and State of Nu-tion speaker*

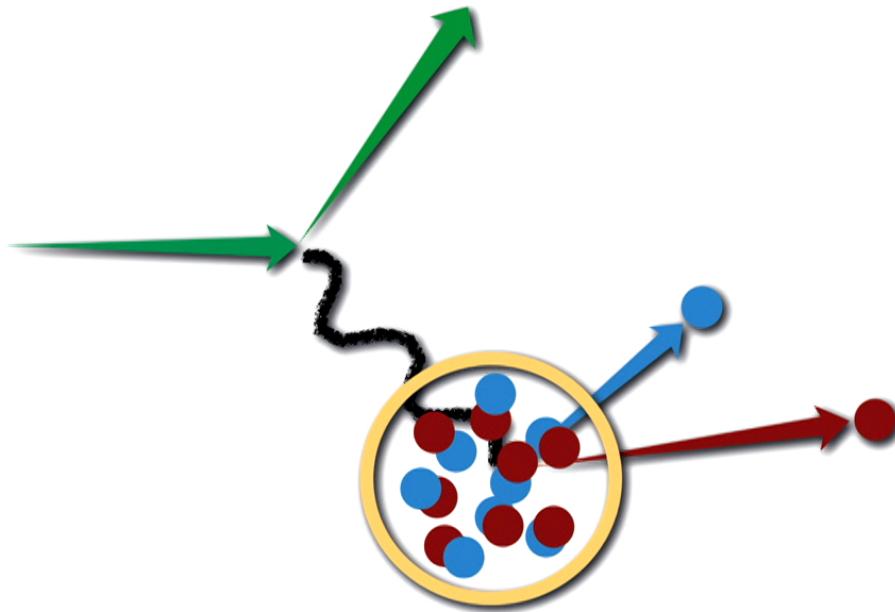
Need correct angular and momentum distributions for signal/background processes!

- Getting latest models (and appropriate uncertainty) into software to use

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Why is  $\nu$ -A scattering important to neutrino oscillation and cross section expt's?

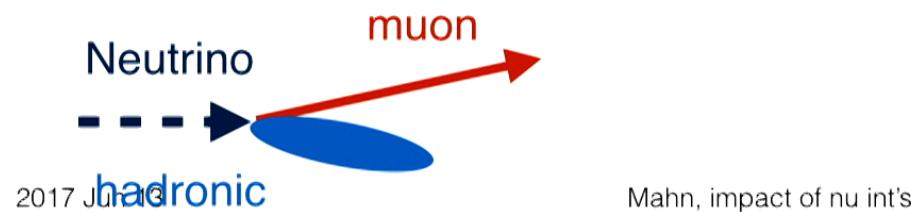
Impact of alternate models on oscillation physics

# Why is v-A important for oscillation expts?

- Oscillation depends on energy
  - Estimate from hadronic and/or leptonic information

$$E_\nu^{QE} = \frac{m_p^2 - m'_n{}^2 - m_\mu^2 + 2m'_n E_\mu}{2(m'_n - E_\mu + p_\mu \cos \theta_\mu)} \quad E_\nu = E_\mu + \sum E_{hadronic}$$

- Nuclear effects bias true and estimated neutrino energy



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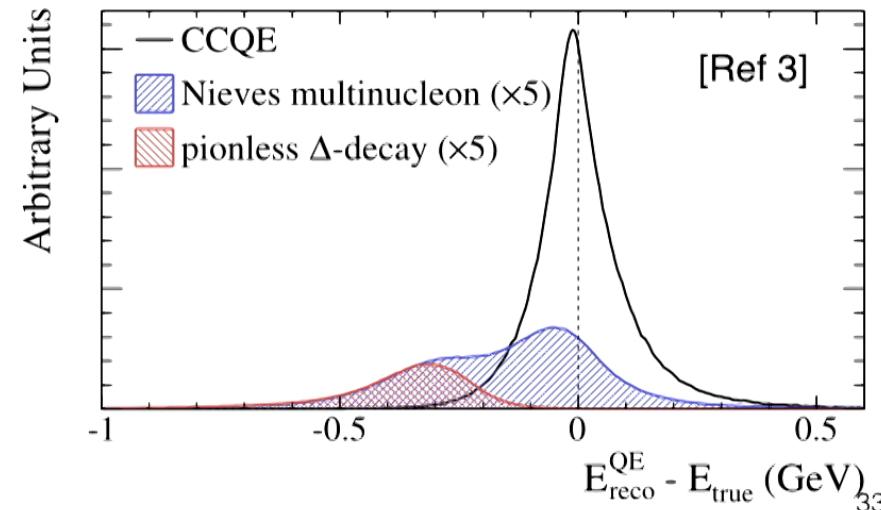
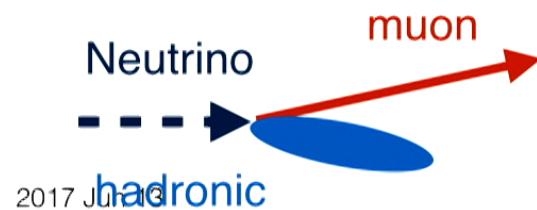
33

# Why is v-A important for oscillation expts?

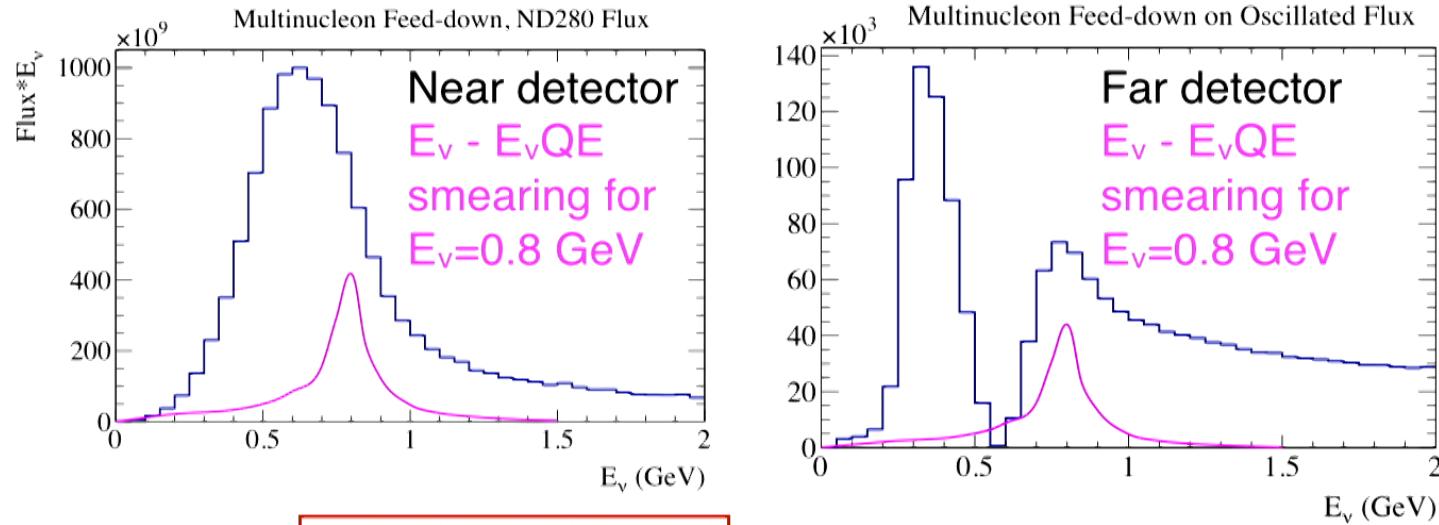
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- Nuclear effects bias true and estimated neutrino energy



# Why is ν-A important for oscillation expts?



$$ND(\nu_\mu) = \boxed{\Phi(E_\nu) \times \sigma(E_\nu, A)} \times \epsilon_{ND}$$

$$FD(\nu_\mu) = \boxed{\Phi(E_\nu) \times \sigma(E_\nu, A)} \times \epsilon_{FD} \times P(\nu_\mu \rightarrow \nu_e)$$

- Even with a near detector, **critical reliance on model**
  - QE method assumptions: 2p2h feed-down to oscillation peak from [Ref 4]
  - Hadronic method assumptions: particle multiplicity, detection threshold

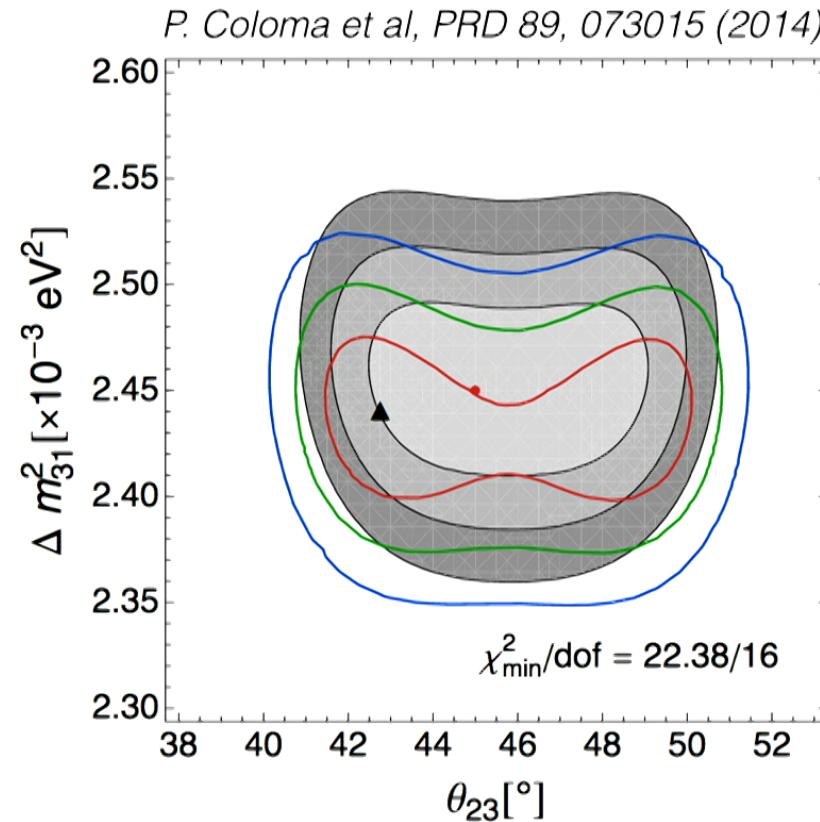
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# Effect of v-A on oscillation physics

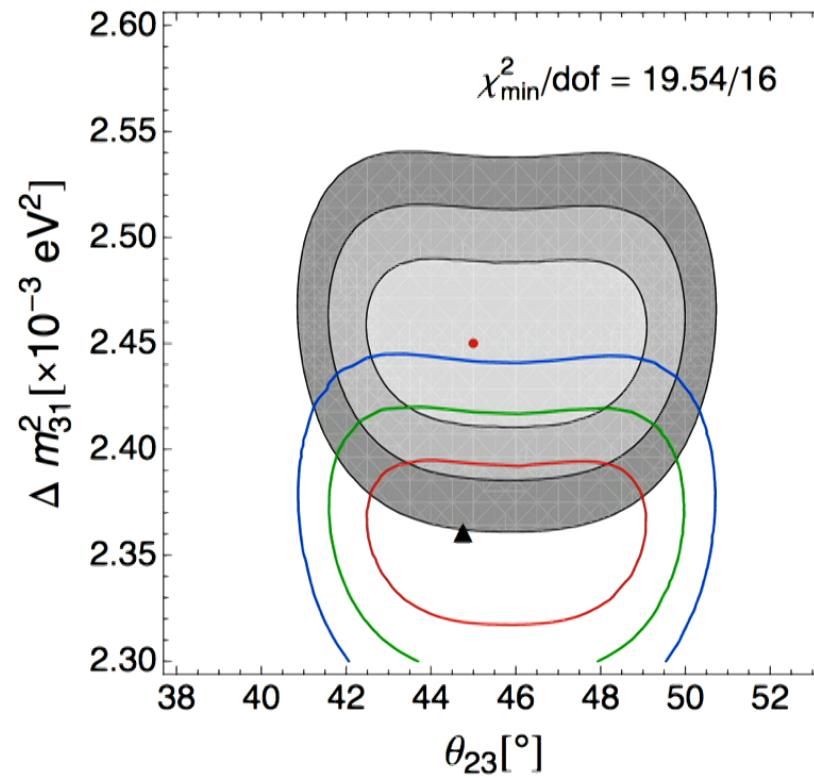
- **Model choice determines:**
- Size of effect on oscillation physics
- Parameterization of the v-A uncertainties



# Effect of v-A on oscillation physics

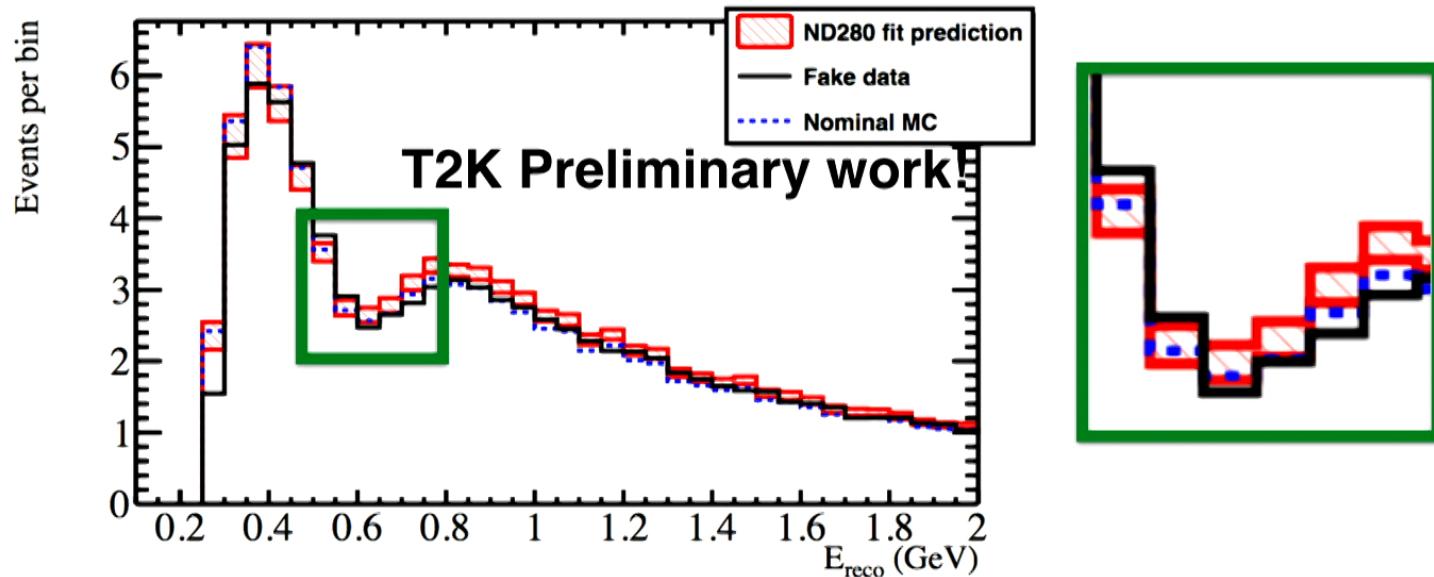
- **Model choice determines:**
- Size of effect on oscillation physics
- Parameterization of the v-A uncertainties
- Which oscillation parameter(s) are affected
- Which method (QE or calorimetric) is used

P. Coloma et al, PRD 89, 073015 (2014)



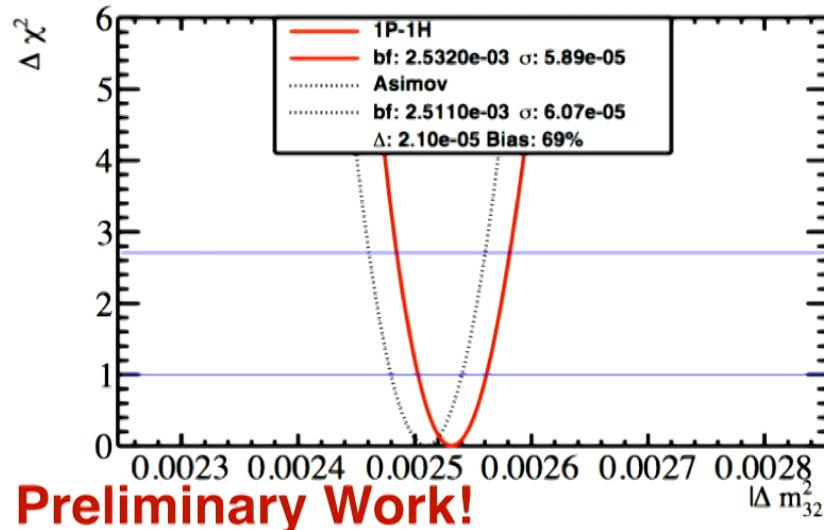
# An example of the pitfalls

- Take a QE model of interest (full simulation or likely, smeared simulation)
- Run entire T2K oscillation analysis chain (fit ND, propagate uncertainties, fit neutrino, antineutrino samples) and evaluate effect on oscillation parameters

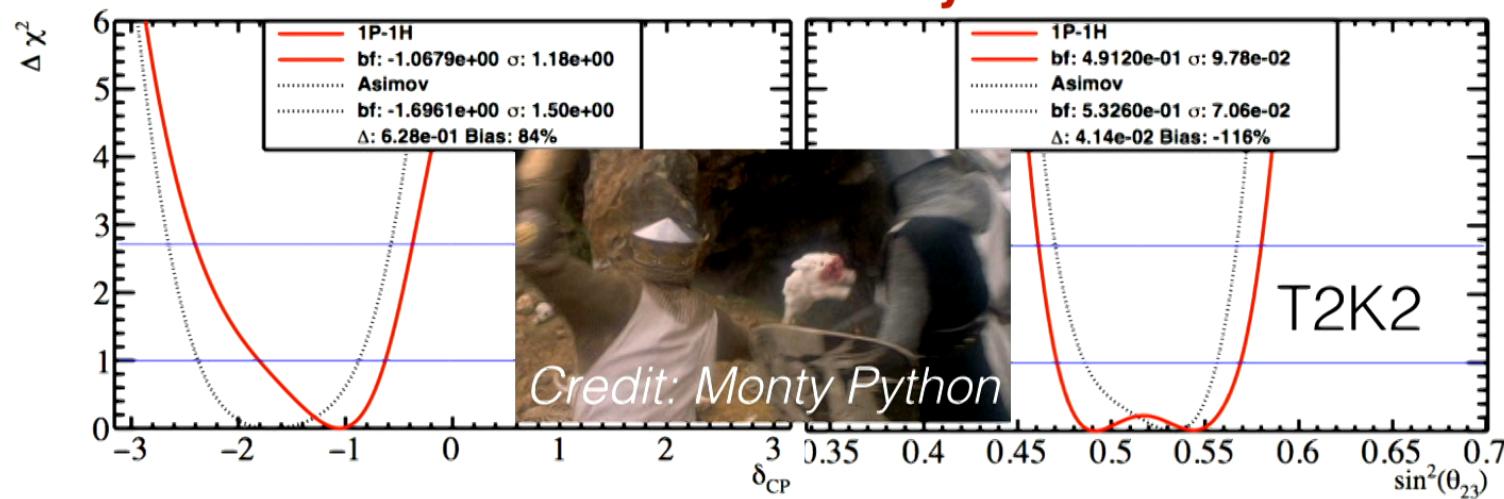


# An example of the pitfalls

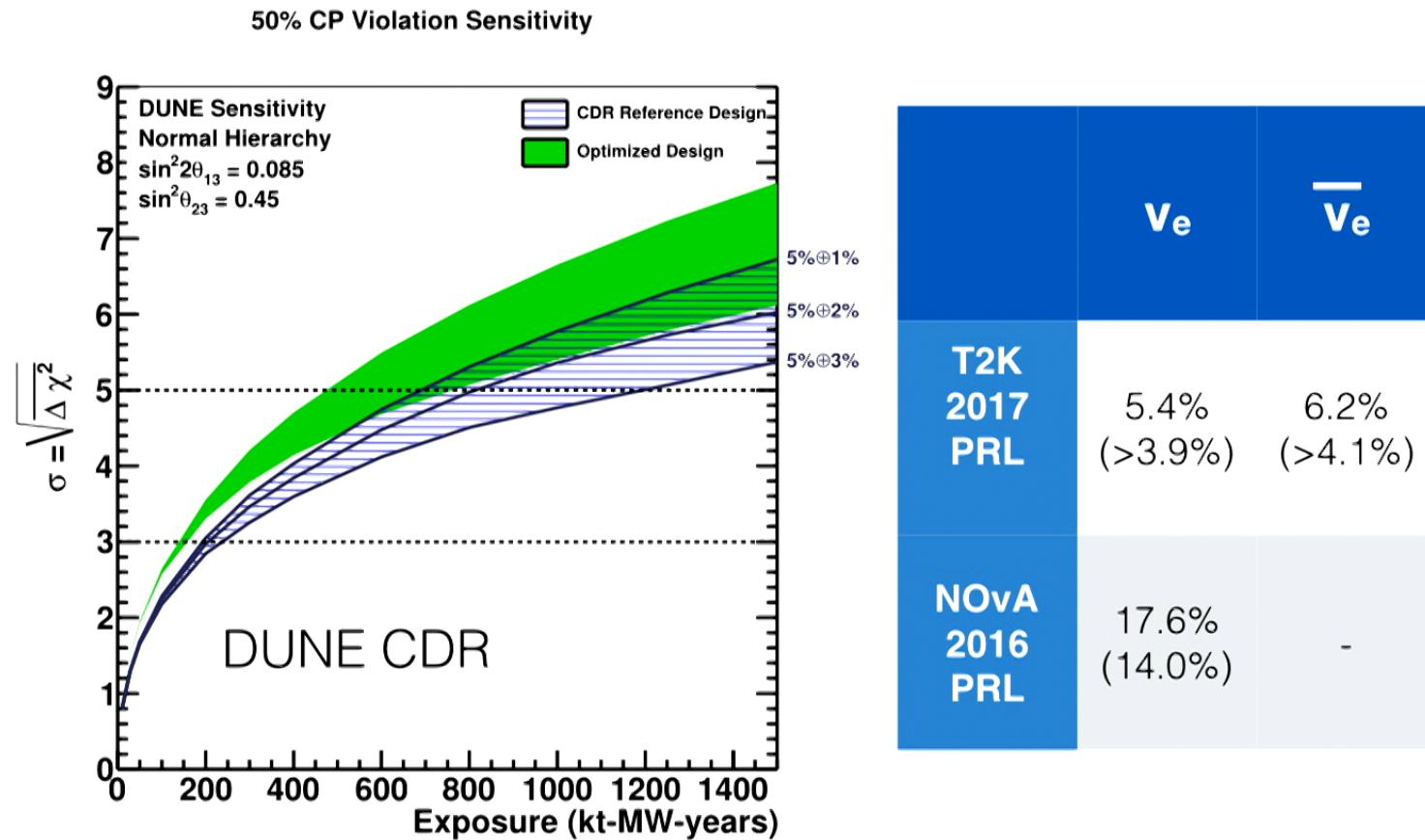
- For an extended run of T2K, significant bias possible.
- Creation of multiple minima
- We mustn't run away!



T2K Preliminary Work!



# How well do we need to know $\nu$ -A?

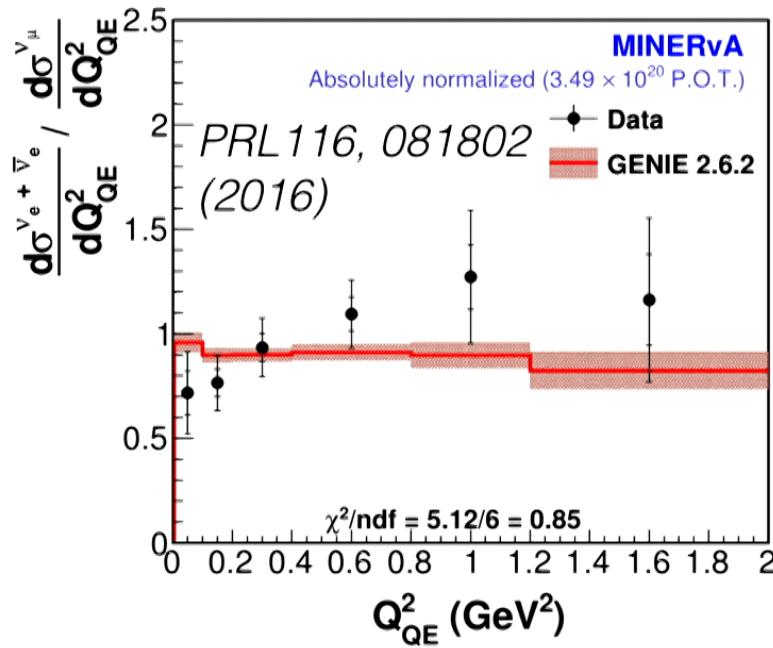


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## Known challenge: electron neutrino cross sections



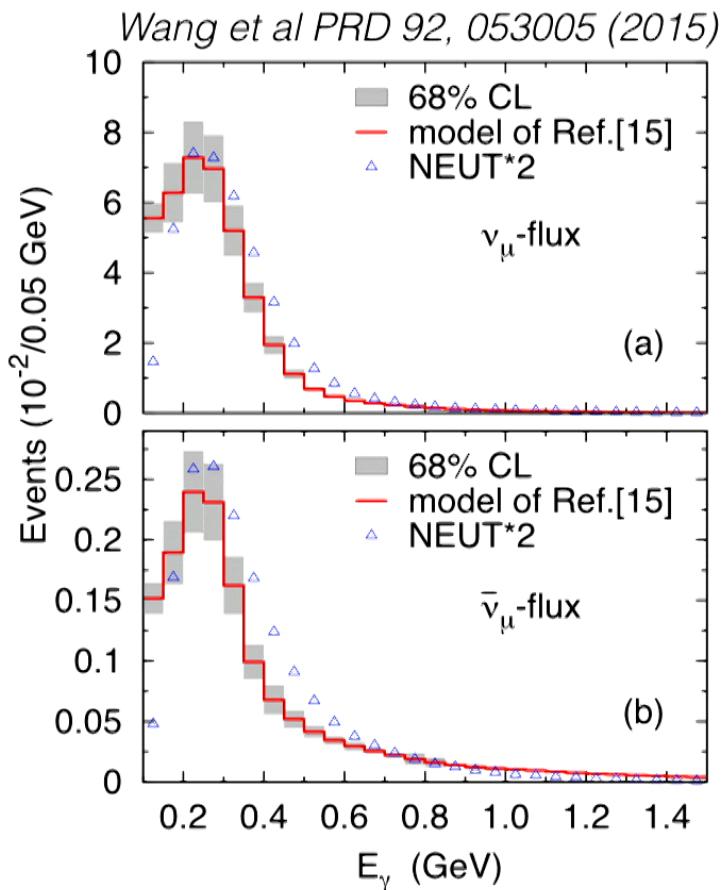
- Measurements from MINERvA, NOvA, T2K of  $\nu_e$  interactions
- Sets scale of different  $\nu_e/\nu_\mu$  nuclear effects of <15-30%
- Tough measurement: small sample size, separation from backgrounds.
- Effects are ~few percent. **Please help calculate!**

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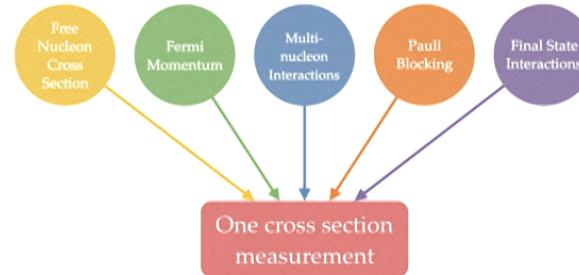
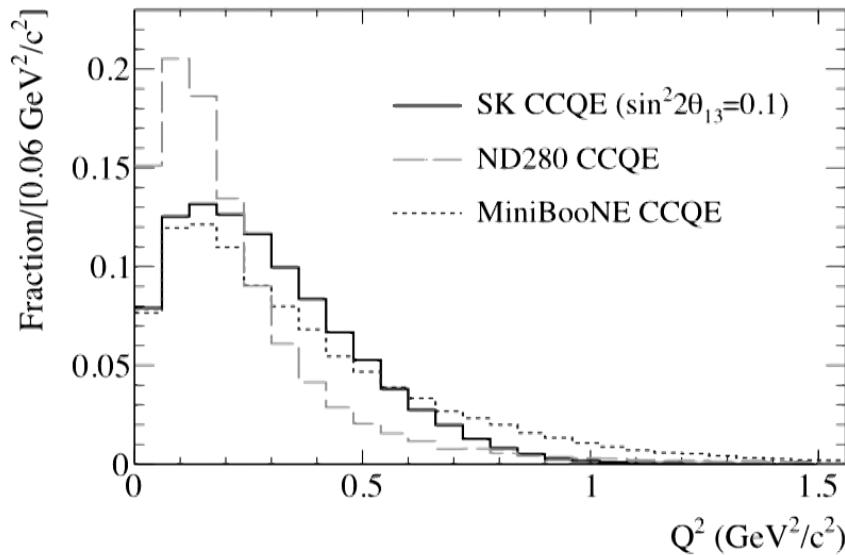
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# Known challenge: neutral current processes



- NC single photon processes ( $\Delta \rightarrow N\gamma$ ) are small (1% of event rate) but difficult to distinguish from signal. Also, photonuclear absorption of 1 photon from NC  $\pi^0$  production.
- Measurements are challenging (like  $\nu_e$ ) though in principle shared with ND
- **Help:** Are there any anomalous contributions? Differences between antineutrino production rates?

# Impact of single nucleon form factors



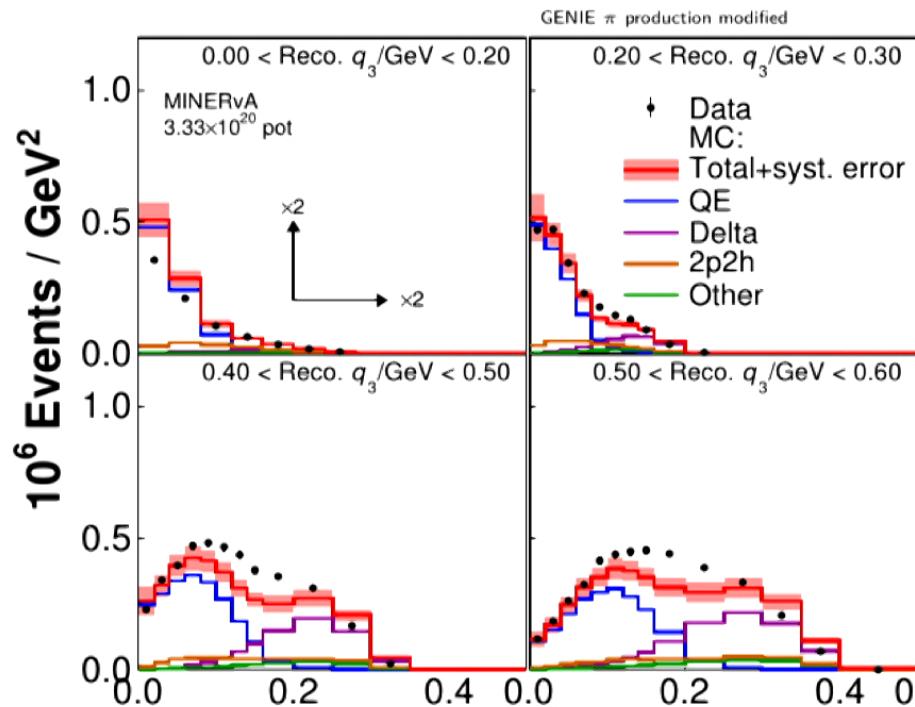
Reliance in T2K analysis on parametrization extrapolation from low to high  $Q^2$

- Testing this now in T2K analyses. How important is this?

Complicated by other processes measured by ND  
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# Novel MINERvA Test of Nuclear Effects

Phys. Rev. Lett. 116, 071802  
(2016), plot from NuInt2016



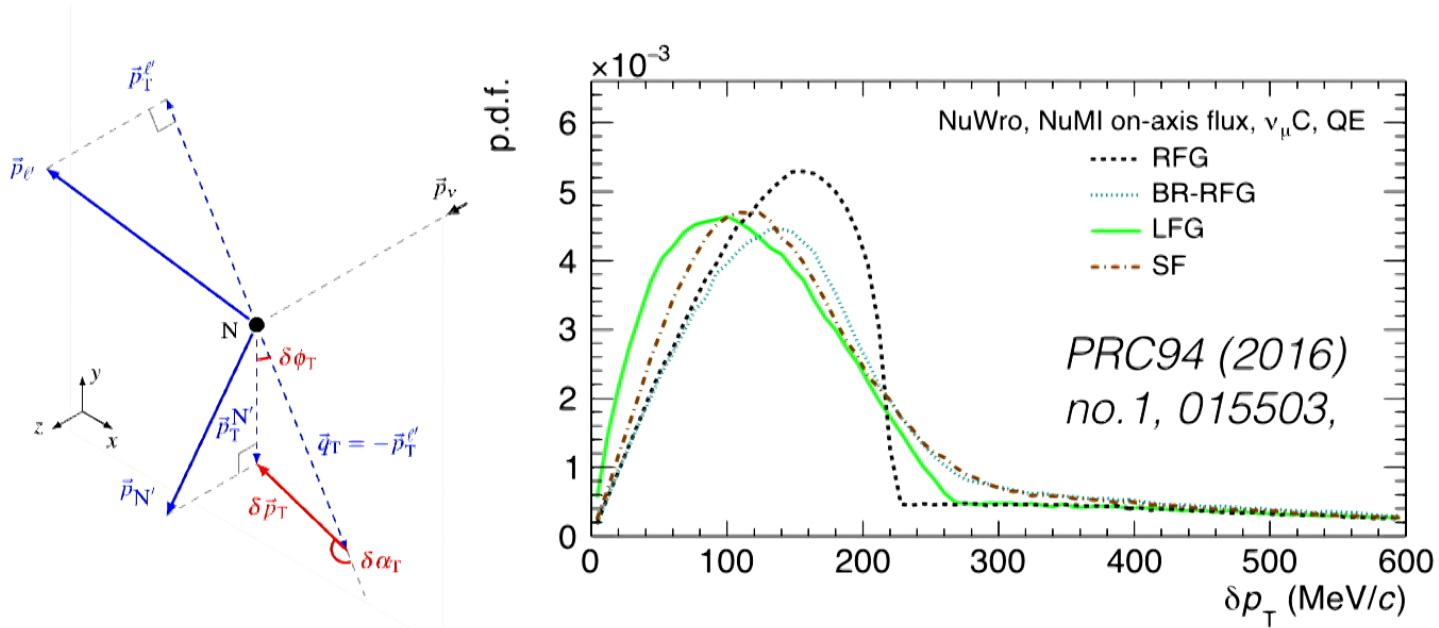
- CC inclusive selection, use estimate of available energy from all charged particle tracks (proton, pion, muon)
- 2p2h model not “enough”: **new experimental way to test model**

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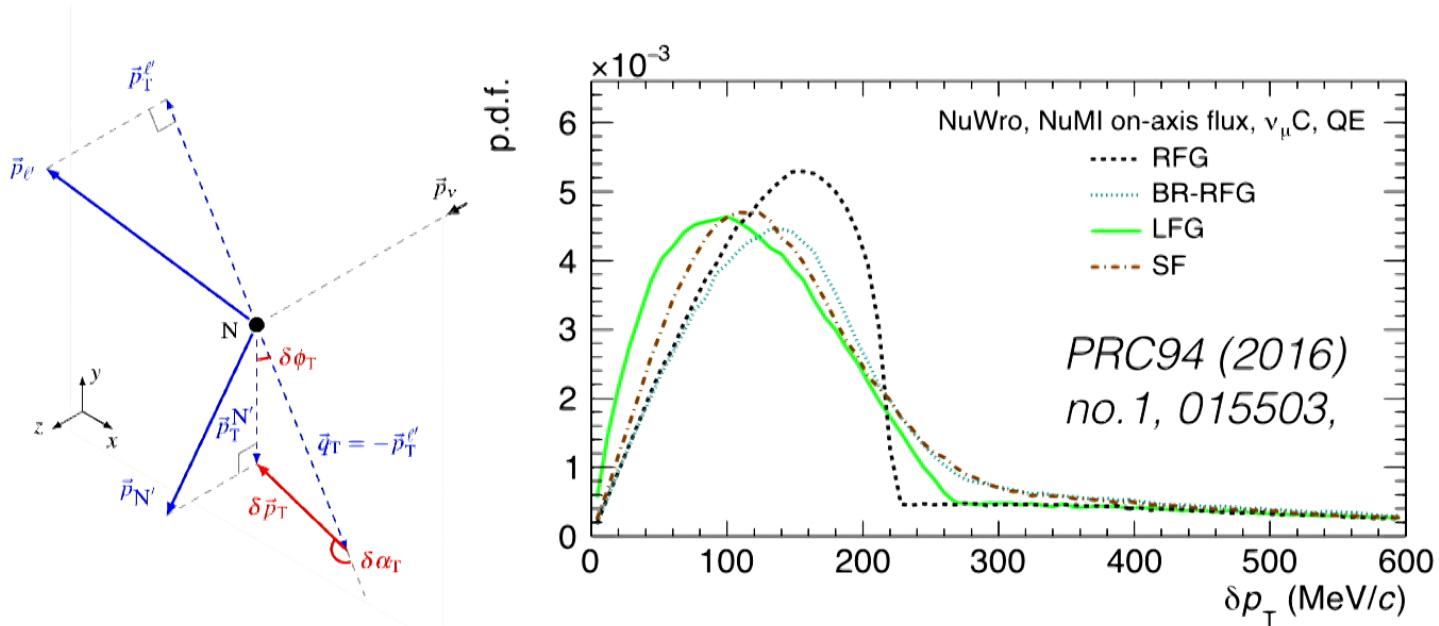
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# Novel approach: Transversity



- Kinematic variables relative to the beam direction are sensitive to nuclear effects

# Novel approach: Transversity



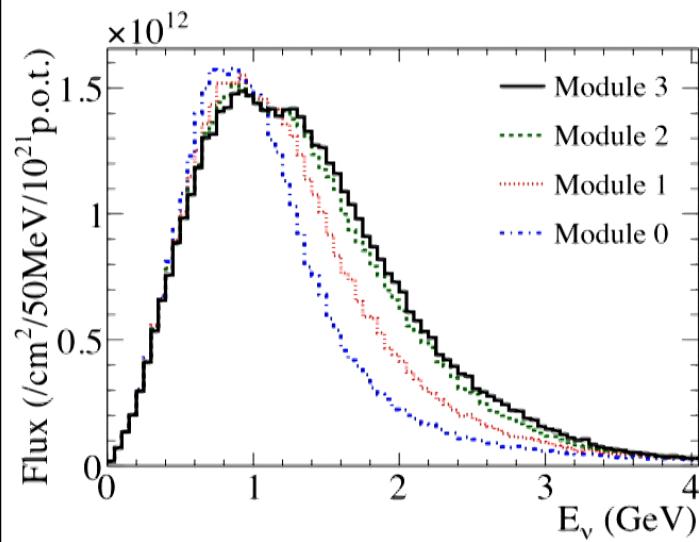
- Kinematic variables relative to the beam direction are sensitive to nuclear effects
- **Challenges:** Requires robust modeling of final state particle composition multiplicity and kinematics => Semi-inclusive or better theory. Care with uncertainties of neutrino source direction

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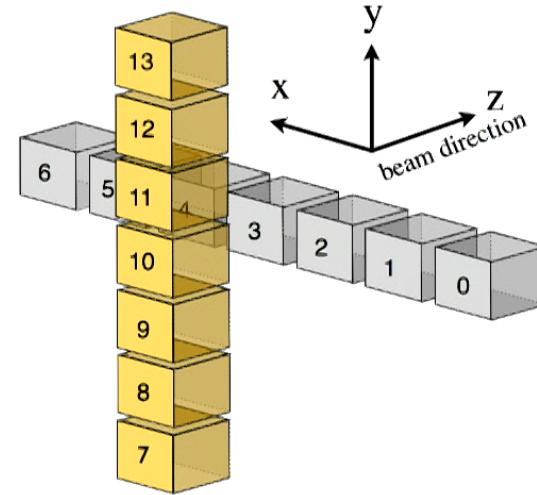
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# Novel: use of flux techniques

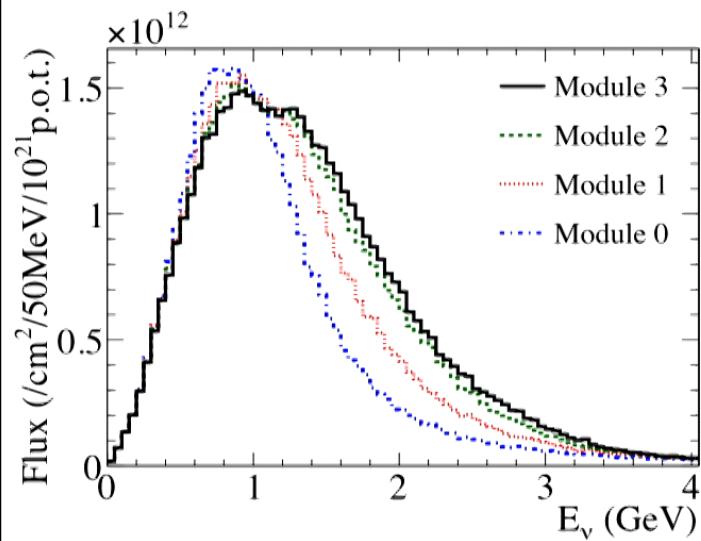


Phys. Rev. D 93,  
072002 (2016)

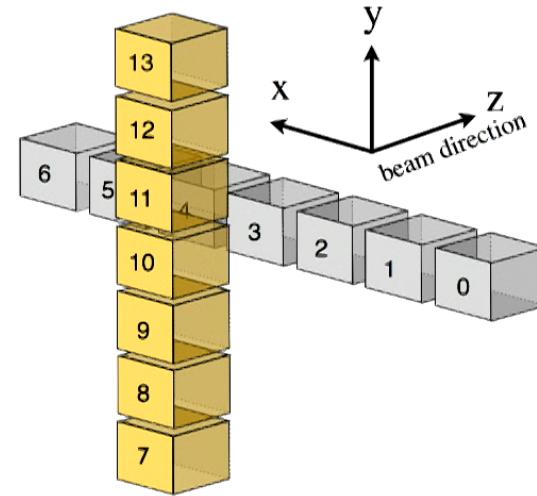


- Flux changes across on-axis T2K detector due to off-axis effect

# Novel: use of flux techniques

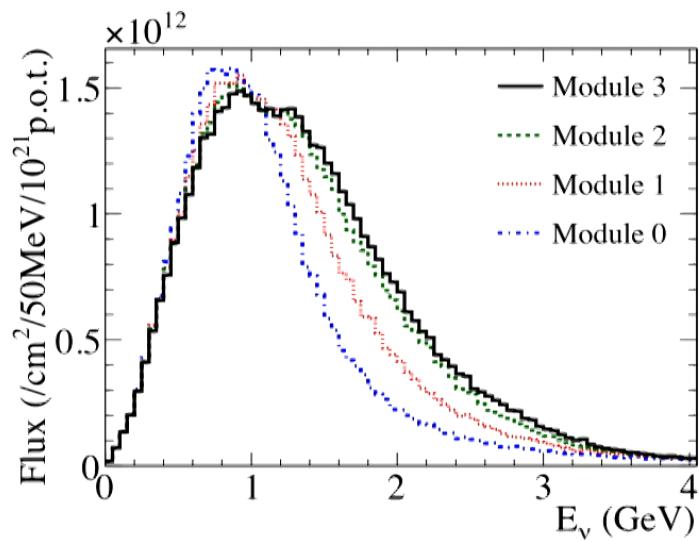


Phys. Rev. D 93,  
072002 (2016)

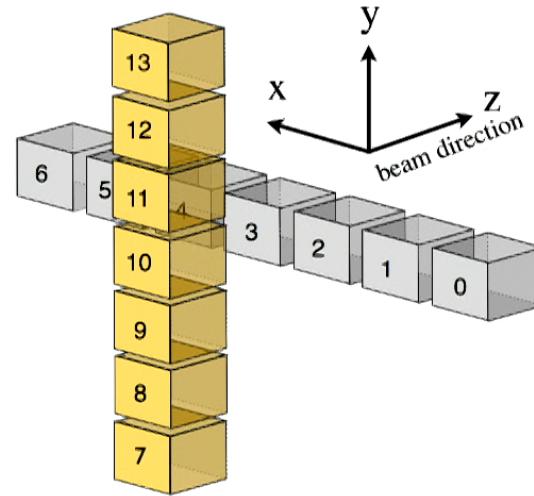


- Flux changes across on-axis T2K detector due to off-axis effect
- Compare nearby detectors to infer cross section using known flux properties

# Novel: use of flux techniques

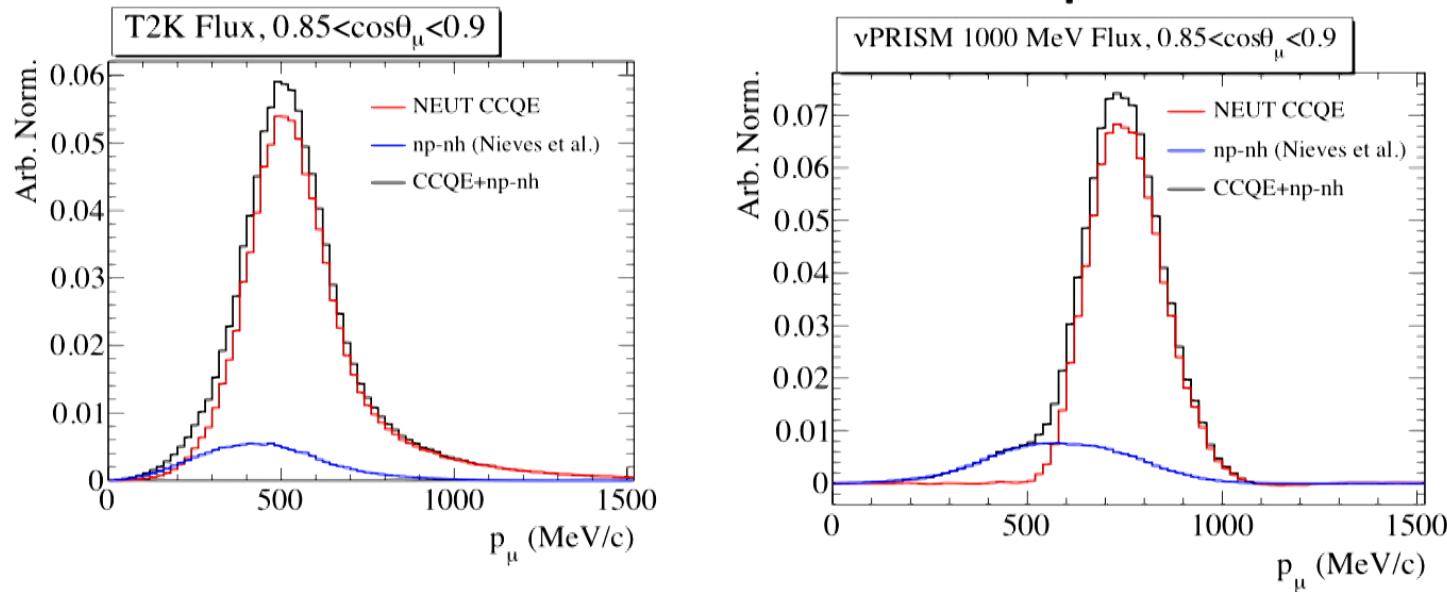


Phys. Rev. D 93,  
072002 (2016)



- Flux changes across on-axis T2K detector due to off-axis effect
- Compare nearby detectors to infer cross section using known flux properties
  - Similar to breaking cross section-flux interplay with nu-e scattering

# Novel: use of flux techniques



- Probe of nuclear effects from data at multiple positions in beam
- First phase: ~1% statistical uncertainty on electron neutrino cross section
  - Dominant uncertainty for Hyper-Kamiokande experiment (3%)

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# Summary

- Exciting era of precision in oscillation programs comes at the price of lots of work. Nuclear effects in neutrino interactions are important for the current and future oscillation physics program.
- Help welcome in robust model building:
  - Tools to exchange theory/determine what effects are significant
  - Single nucleon knowledge (IQCD/z-expansion). Would nuclear effects change this?
  - Leveraging nu-e scattering
  - Differences between electron and muon neutrinos
  - Rare but irritating backgrounds (NC photon production)