

Title: Quantum singularity resolution in homogeneous cosmology and the implications for shape dynamics

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Abstract: I will present results on the quantization of an FRLW model that utilises a Schrodinger-type evolution equation. In contrast to standard Wheeler--DeWitt-type quantisations, the quantum model resolves the classical singularity, exhibits a quantum bounce, and displays novel early-universe phenomenology. A global scale emerges because of a scale anomaly, and suggests an interesting scenario for quantum shape dynamics. I will give the details of the quantization procedure and show how these techniques can be used more generally for anisotropic models. I will end by speculating about how these techniques might be applicable to a genuine quantum shape model of the universe.

Quantum singularity resolution in homogeneous cosmology and the implications for shape dynamics

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Shape Dynamics Workshop

Perimeter Institute

17 May 2017

Context

A biased, distorted, and selective history of shapes:

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- Symmetry trading paper (Gomes, Koslowski, sg – 2010) \Rightarrow gravity as a 3d Weyl gauge theory.
- Problem: volume preserving condition \Rightarrow a big mess!
- How to deal with the global mode? \Rightarrow dynamical similarity?

Dynamical Similarity

- Perhaps the missing insight?
- A 'new' kind of symmetry.
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- I will **not** quotient by dynamical similarity!

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What goes wrong?

⇒ We believe it is the treatment of time

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- Qualitative features carry over to Bianchi I and IX.
- Natural way to quotient by dynamical similarity at the quantum level.

The classical model

A simple FRWL cosmology

Homogeneous and isotropic mini-superspace model with massless scalar field and vanishing spatial curvature.

Hamiltonian:

$$H = N \left[-\frac{\kappa}{12V_0 a} \pi_a^2 + \frac{1}{2V_0 a^3} \pi_\phi^2 + \frac{V_0 a^3}{\kappa} \Lambda \right]$$

Variables

- a - scale factor
- ϕ - scalar field

Parameters

- Λ - cosmological constant
- π_ϕ - momentum of scalar field

(V_0 - fiducial volume; κ - Newton constant)

Geometry of configuration space

The Hamiltonian can be written as

$$H = N \left(g^{ab} p_a p_b - \frac{\Lambda}{2} \right) \quad (1)$$

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Solutions \Rightarrow geodesics on (\mathcal{C}, g)

Restriction

Geometrically, $a > 0$ puts us in Rindler spacetime.

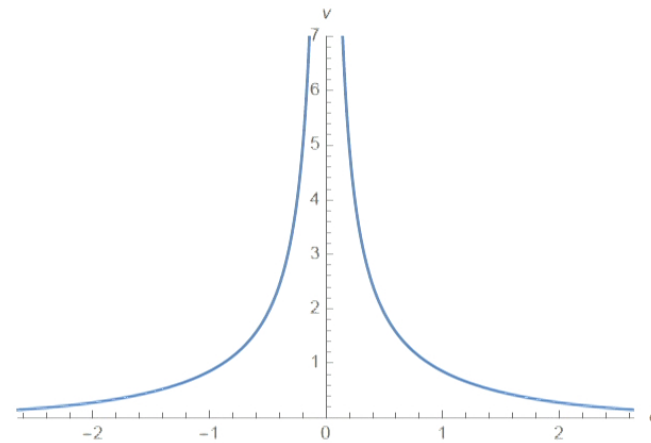
Boundary \Rightarrow geodesic incompleteness on (\mathcal{C}, g) (classical singularity and quantum complications)

Solutions

Boost symmetry suggests Rindler coordinates: $v \propto a^3$, $\varphi \propto \phi$

$$v = v_0 + \sqrt{t - 1}$$

$$\phi = \phi_\infty + \operatorname{arctanh}(t^{-1})$$



Parameters

- v_0 - initial time
- φ_∞ - asymptotic φ
- Λ - total energy
- π_φ - conserved p

Symmetries

- Time translations
- Boost invariance
- Time units
- Space units

Symmetries \Rightarrow no physically relevant parameters.!

Quantization Preliminaries

Hamilton–Jacobi equation:

$$\mathcal{H}\left(q, \frac{\partial S}{\partial p}\right) = \frac{\Lambda}{2}$$

Quantization ambiguity

$$\hat{\mathcal{H}}\Psi = \frac{\Lambda}{2}\Psi \quad \text{or} \quad \hat{\mathcal{H}}\Psi = i\frac{\partial\Psi}{\partial\tau}$$

⇒ Timeless (Wheeler–DeWitt) vs genuine evolution (unimodular gravity-like)

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Relational quantization:

- Take approach with genuine evolution.
- τ - unobservable parameter ordering successive states.

Failure of timeless (Wheeler–DeWitt) approach

Wheeler–DeWitt equation is timeless.

⇒ deparametrize wrt internal time. (e.g., φ)

Problem (example)

$\langle v \rangle$ problematic when $\varphi \rightarrow \infty$ (takes definite values)

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Relation quantization $\Rightarrow \hat{\varphi}$ does not take definite values.

$\langle v \rangle, \langle \varphi \rangle$ protected by uncertainty relations!

Unitarity

Boundary of $\mathcal{C} \Rightarrow \hat{\mathcal{H}}$ not essentially self-adjoint!

$$\langle \Phi, \hat{\mathcal{H}}\Psi \rangle = \langle \hat{\mathcal{H}}\Phi, \Psi \rangle + \text{boundary}$$

(Boundary term \Rightarrow quantum imprint of classical singularity.)

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Self-adjoint Extensions

- Guaranteed by von Neumann. ($\hat{\mathcal{H}}$ - symmetric and real)
- Trick: the equation (boundary = 0) is conformally invariant.
- Use eigenfunctions on conformal completion $\tilde{g} = v^2 g$ (maps $\mathbb{R}^+ \times \mathbb{R} \rightarrow \mathbb{R}^2$) to anchor solutions in Minkowski.
- Put phase between solutions to remove Λ -(scale) dependence.

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Relation quantization $\Rightarrow \hat{\varphi}$ does not take definite values.

$\langle v \rangle, \langle \varphi \rangle$ protected by uncertainty relations!

We will use genuine evolution: $\hat{\mathcal{H}}\Psi = \square\Psi = i\frac{\partial\Psi}{\partial\tau}$.

'Bound' states ($\Lambda < 0$)

Spectrum of Klein–Gordon operator on Rindler for $\Lambda < 0$:

- AdS-like states exist.
- Spectrum is discrete but unbounded: (k - momentum of φ)

$$\Lambda_n = \Lambda_{\text{ref}} e^{2\pi n/k} \quad (n \in \mathbb{Z})$$

- Conformal tower of states with accumulation point at $\Lambda = 0$.
- Analogues of atomic systems: *Efimov effect*.
- Eigenstates behave like $\psi_\Lambda \propto K_{ik}(\sqrt{\Lambda}v) \sim e^{-v}$.

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'Unbound' states ($\Lambda > 0$)

- dS-like states exist and persist to late-time, semi-classical limit.
- Continuous spectrum given by phase-shifted Bessel functions:

$$\psi_\Lambda \sim J_{ik}(\sqrt{\Lambda}v) + \tan \theta I_{ik}(\sqrt{\Lambda}v)$$

$\theta \propto k \log \Lambda/\Lambda_{\text{ref}}$ ($U(1)$ s.a. ext. parameter).

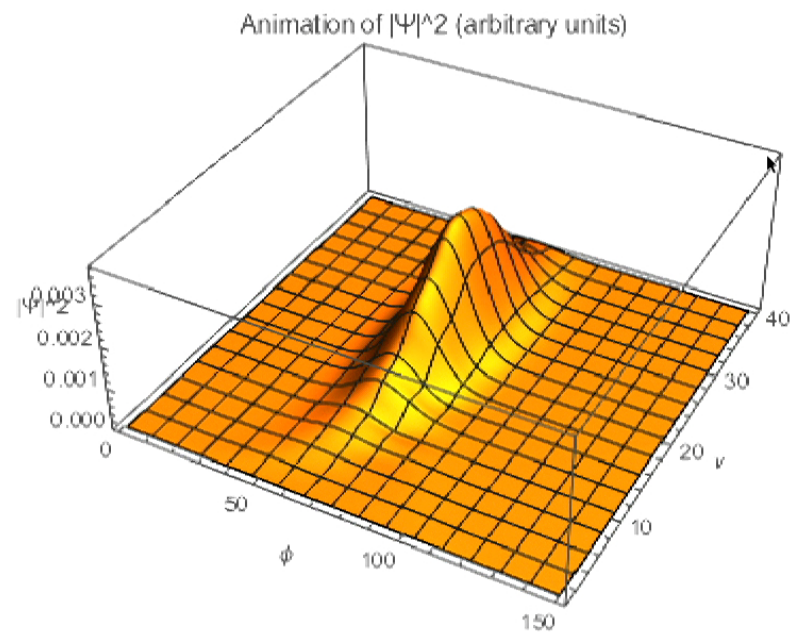
- Gaussian states lead to semi-classical physics.
- New dimensionful parameter $\rightarrow \Lambda_{\text{ref}}$ gives meaning to size of quantum effects (\hbar).
- Self-adjoint extension parameter: phase between in- and out-state.

Analogue model

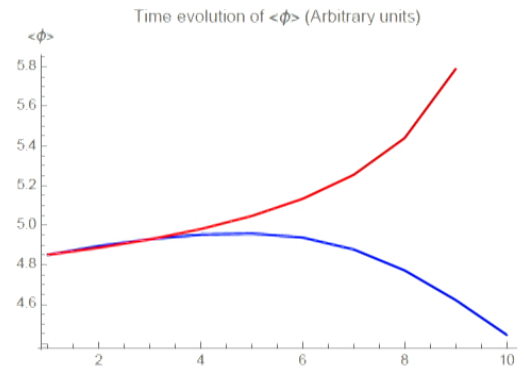
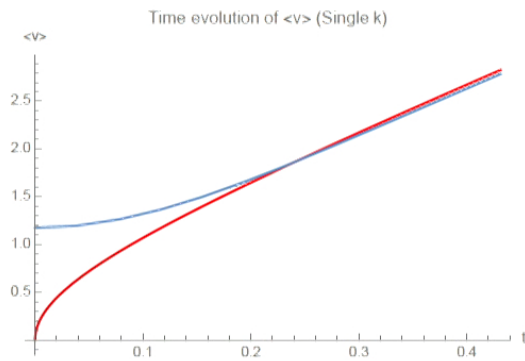
- System is mathematically equivalent to N -particle system with $1/r^2$ potential (in a particular regime).
- Regime in question is a good effective model for some 3-body atomic systems.
- Known as *Efimov effect* \Rightarrow vast literature.
- Unbound states can be scattered off bound states.
- Self-adjoint extension parameter gives scattering length via phase shift, θ .
- θ is macro parameter encoding micro-physics.

How seriously can we take the analogue model or simulator of the early universe?

Bounce Solution



Expectation values

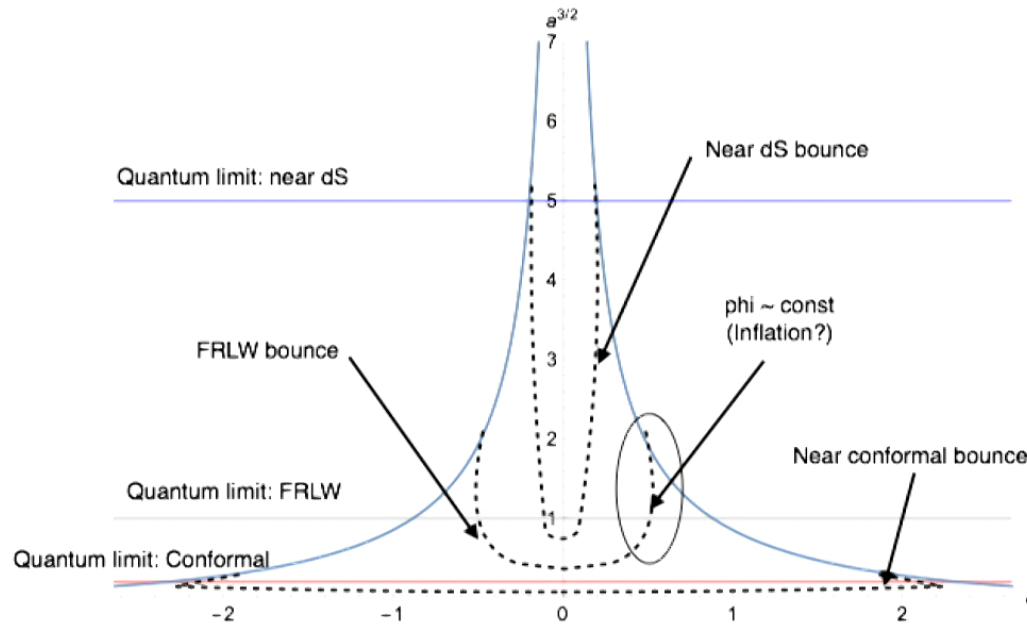


$\langle v \rangle$ has a min when $v_{\text{class}} = 0$.

$\langle \phi \rangle$ is finite when $\phi_{\text{class}} \rightarrow \infty$.

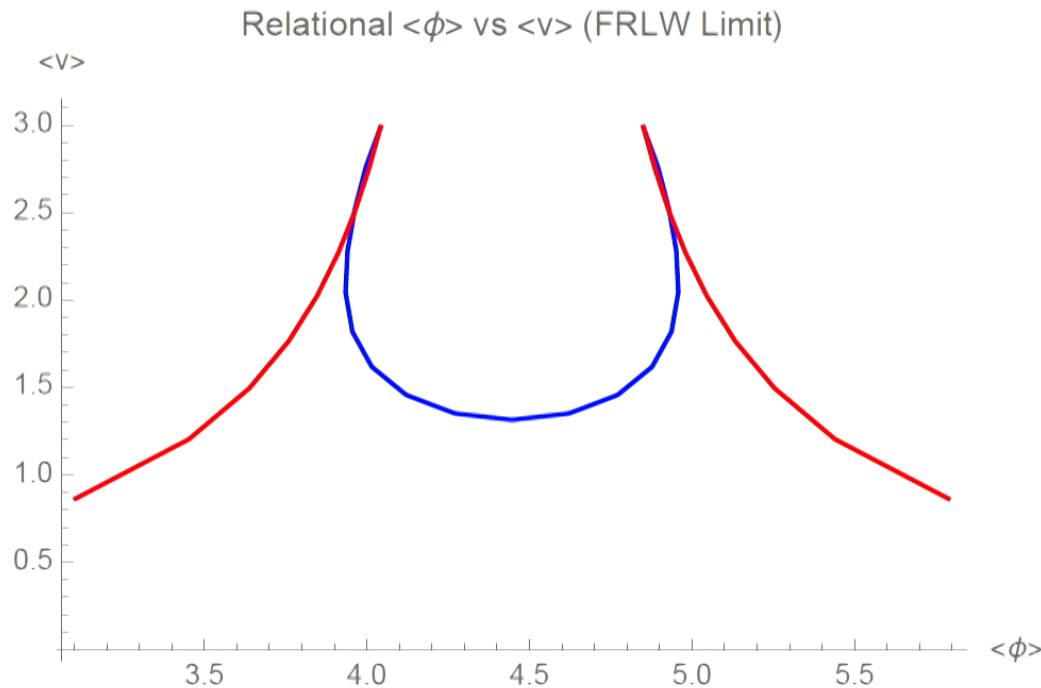
Singularity resolved and classical limit recovered!

General Classical/Quantum Regimes (Cartoon)



New parameter, Λ_{ref} , gives meaning to \hbar and relative size of quantum effects.

General classical/quantum regimes (simulation)



- Simple example in $\hbar \sim 1$ regime.
- Note inflationary period where $\frac{\partial\langle\phi\rangle}{\partial\langle v\rangle} \approx 0$.

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- Resolves singularity (improvement over WDW).
- Recovers semi-classical dS limit.
- Suggest novel phenomenology in quantum regime.
- Provides unique macro-parameter, Λ_{ref} , encoding micro-physics.
- New platform for condensed matter quantum simulation of the early universe.

Generalizations

- Bianchi I: v wavefunction is identical.
- Bianchi IX: Bessel functions \Rightarrow non-analytic with same asymptotic properties.
- BKL \Rightarrow more general mechanism for singularity resolution via quantum effects.

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- Bianchi I: ν wavefunction is identical.
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- BKL \Rightarrow more general mechanism for singularity resolution via quantum effects.
- Dynamical similarity quotient \Rightarrow integrate over all $\Lambda_{\text{ref}}??$
- Conformal anomaly and holographic renormalization.
- Absorption/Decay: AdS-‘bound’ state \Rightarrow dS-‘emission’ (need non-linear corrections).
- Self-adjoint extension parameter \Rightarrow UV-completion required (conformal?).