

Title: Tensor Networks and Holography

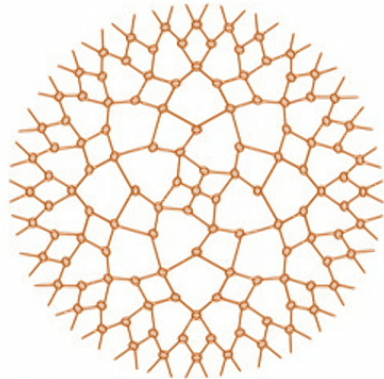
Date: Apr 20, 2017 09:30 AM

URL: <http://pirsa.org/17040041>

Abstract:



McGill

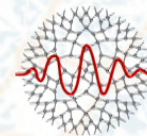


# Tensor Networks and Holography



James Sully

TENSOR NETWORKS FOR QUANTUM FIELD THEORIES II



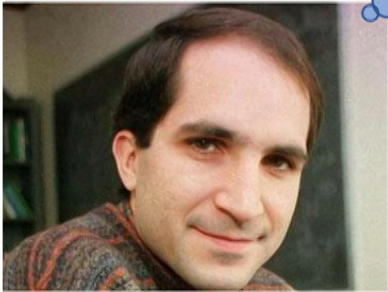
**WELCOME TO  
DISNEYLAND!**



# Introduction

Like many Disney© tales, we begin with young hero, and a dream...

If only we understood quantum gravity...

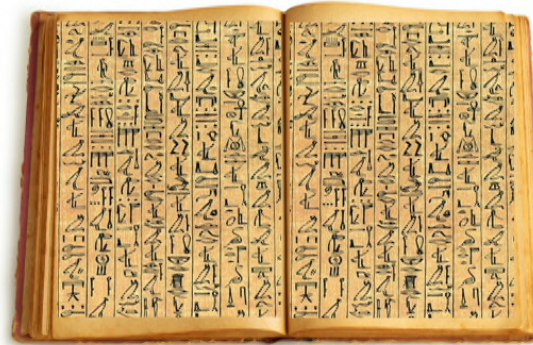
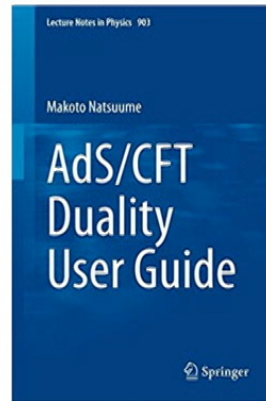


Thankfully his wish was granted:  
**AdS/CFT Duality**



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We have a **UV complete theory of quantum gravity** in terms of a **dual CFT...**



but we haven't been given the tools to completely translate between them.

**How do we see classical geometry and gravitational physics emerge from CFT description?**

**Perhaps we need simpler toy models of holography to make progress...**

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# Outline

1. Review of Gauge/Gravity Duality
2. MERA as Emergent Geometry
3. Gravitational Physics from Toy Tensor Networks

This Talk: pedagogical review

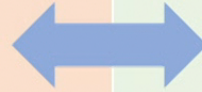
Based on: [Swingle]  
[Pastawski, Yoshida, Harlow, Preskill]  
[Hayden, Nezami, Qi, Thomas, Walter, Yang]

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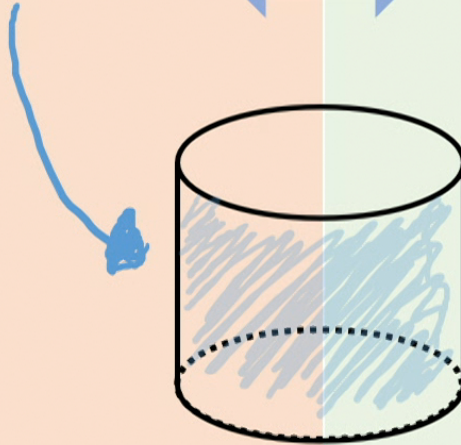
# Gauge/Gravity Duality

# AdS/CFT Duality

**Conformal Field theory**  
d-dimensions

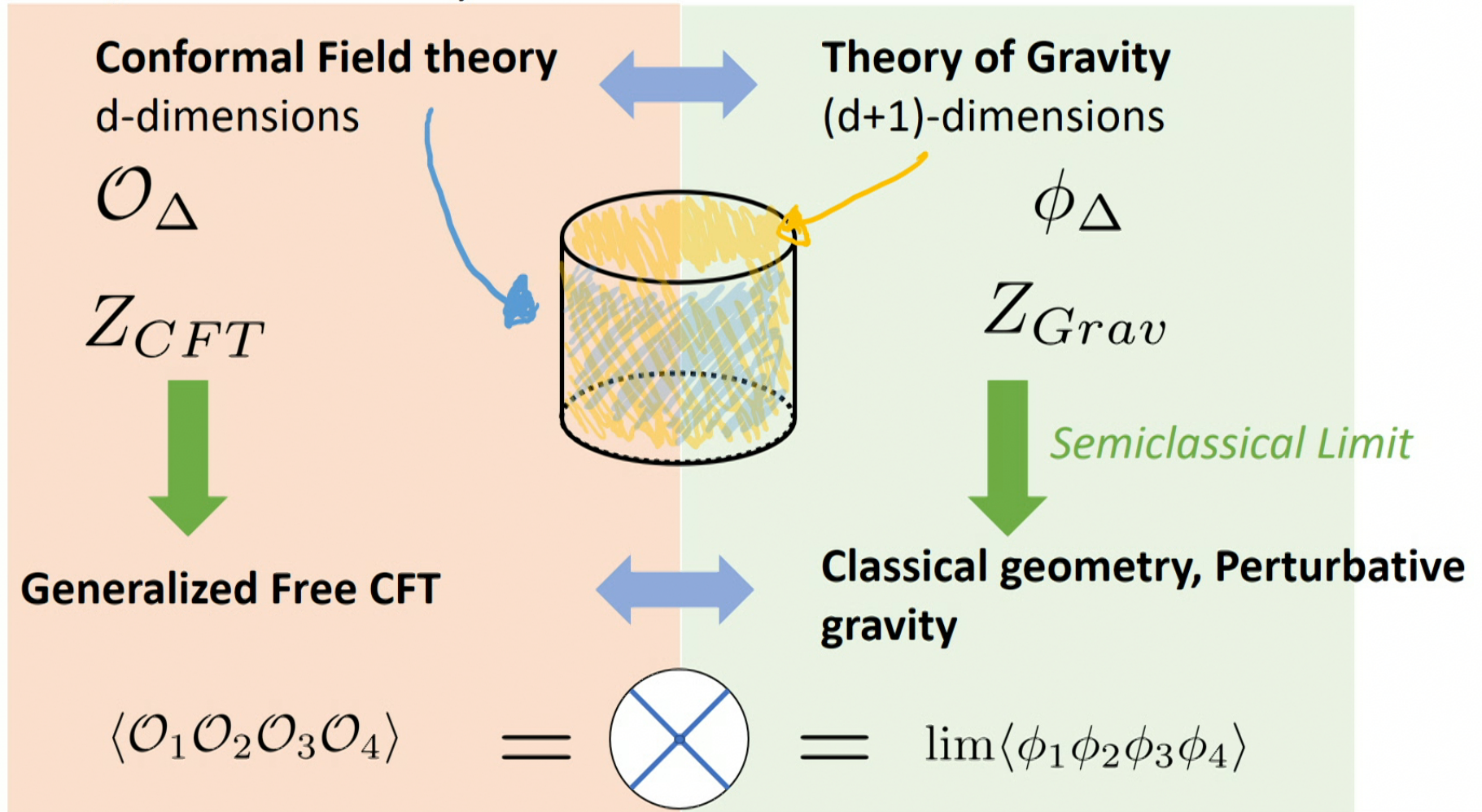


**Theory of Gravity**  
(d+1)-dimensions





# AdS/CFT Duality



# AdS/CFT Duality

[cf. Heemskerck, Penedones, Polchinski, JS]

Only particular class of CFTs expected to have a ‘good’ gravitational dual:

What are these CFTs?

Consider a CFT that has:

1. **Large central charge:**  $c \gg 1 \left( \frac{L}{l_p} \right) \gg 1$

2. Whose **correlators factorize:**

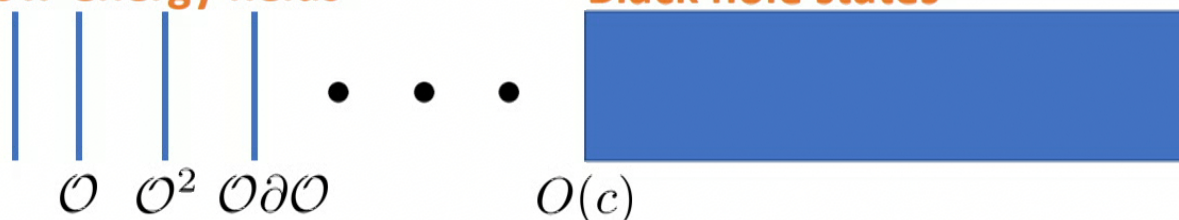
$$\langle \mathcal{O}_1(x_1) \mathcal{O}_1(x_2) \mathcal{O}_2(x_3) \mathcal{O}_2(x_4) \rangle = \langle \mathcal{O}_1(x_1) \mathcal{O}_1(x_2) \rangle \langle \mathcal{O}_2(x_3) \mathcal{O}_2(x_4) \rangle + O(1/c)$$

Perturbative effective fields in bulk

3. Whose **spectrum of conformal dimensions is sparse:**

Low-energy fields

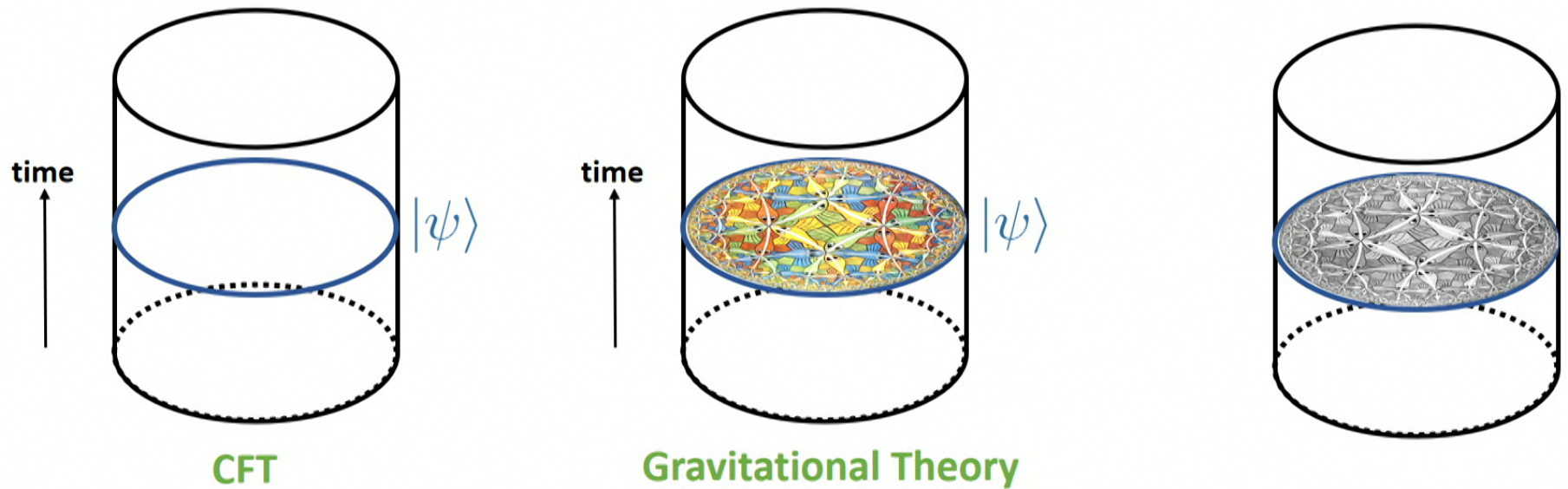
Black hole states



Then this CFT is dual to a theory of gravity whose low-energy energy description is gravity plus effective field theory of field  $\phi$  dual to  $\mathcal{O}$ .

# AdS/CFT Duality

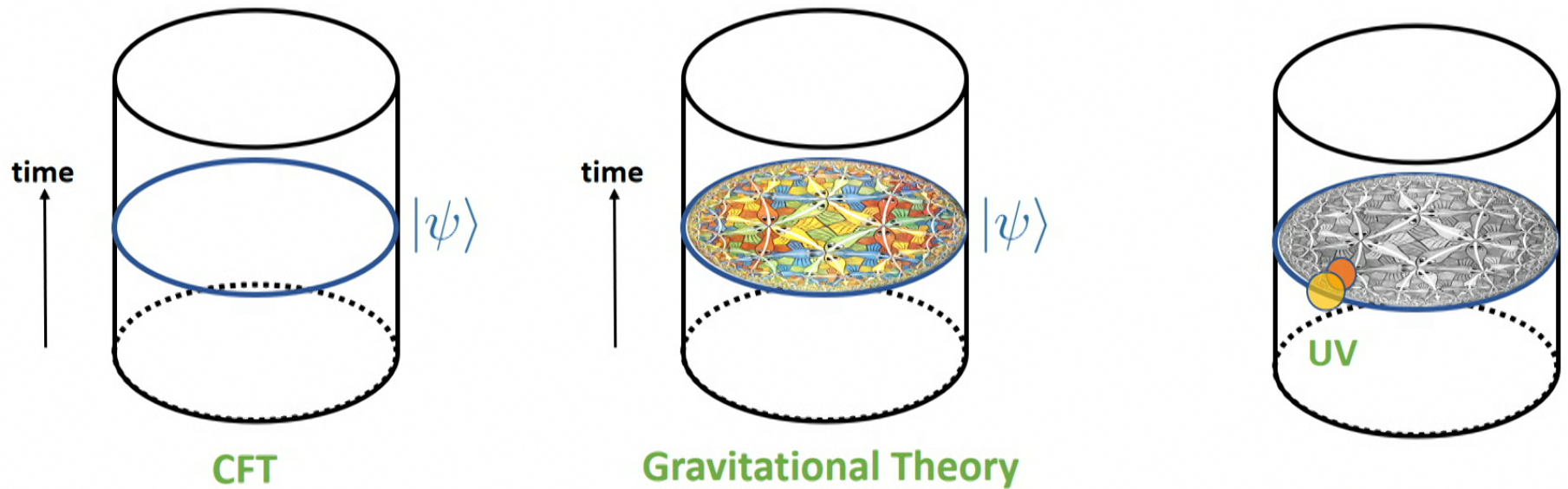
If we consider some state,  $|\psi\rangle$ , in the CFT



it also describes **some geometry in the gravitational theory.**

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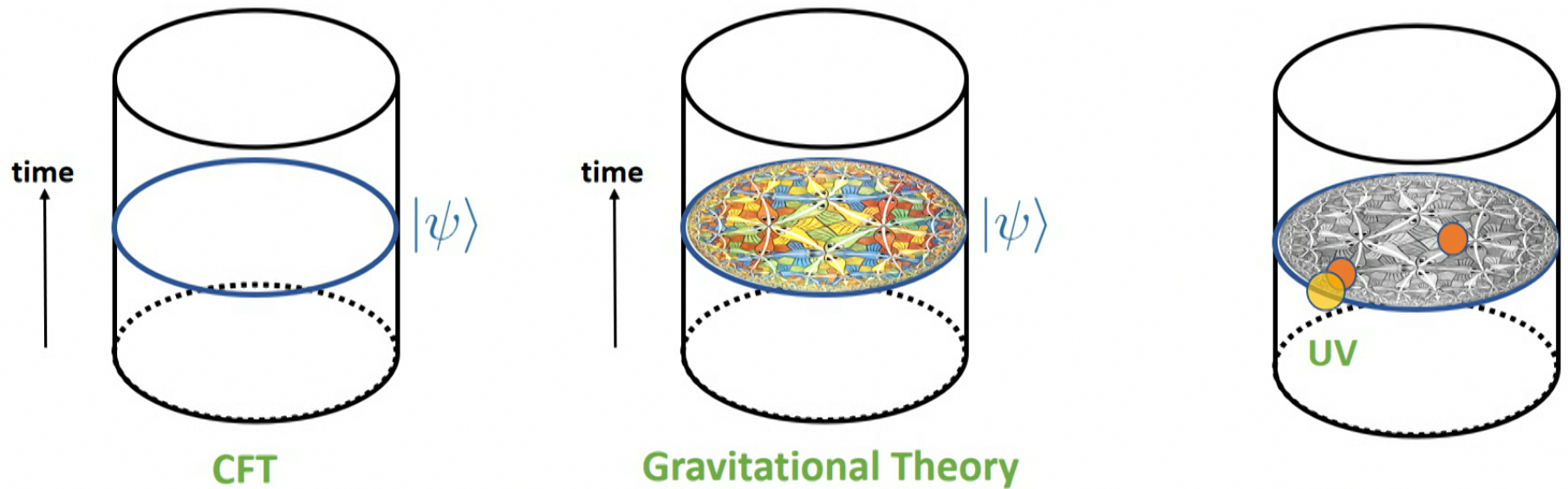
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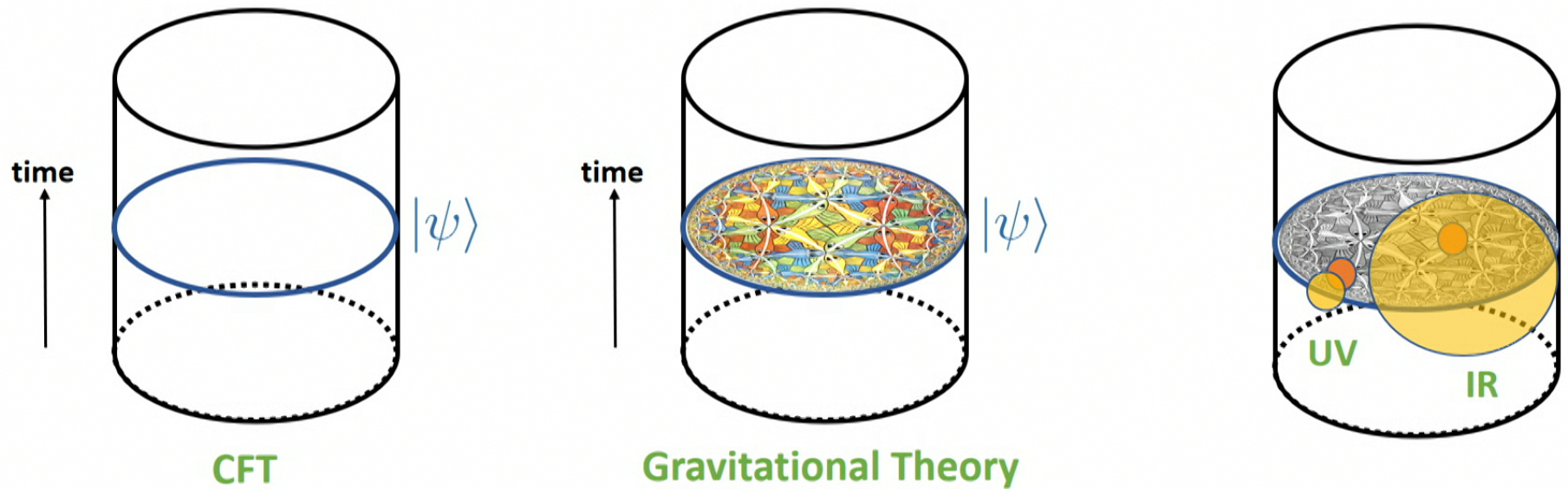
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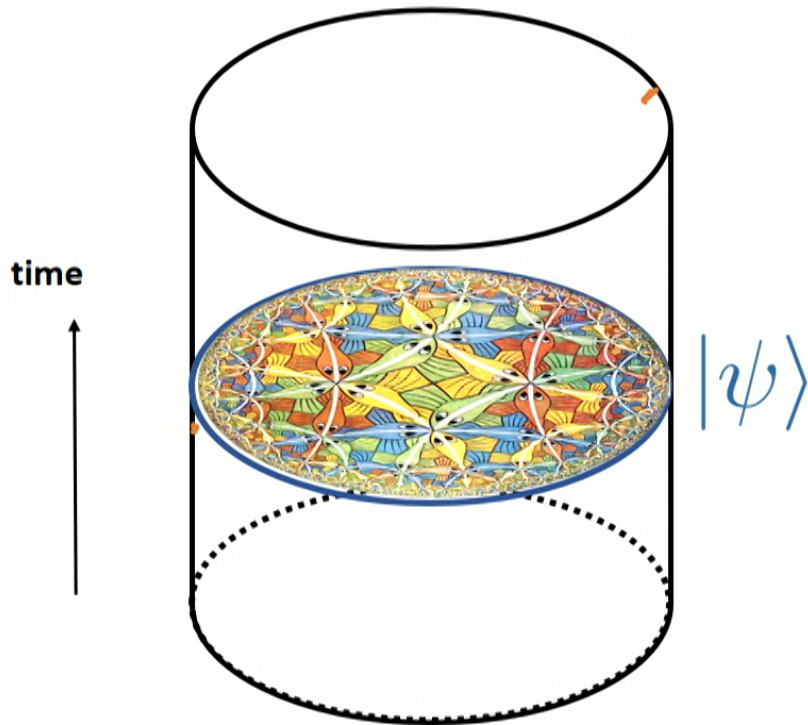
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# AdS/CFT Duality

- In the case where  $|\psi\rangle$  is the **vacuum**, the dual spacetime is maximally symmetric, negative curvature **anti-de Sitter Space (AdS)**

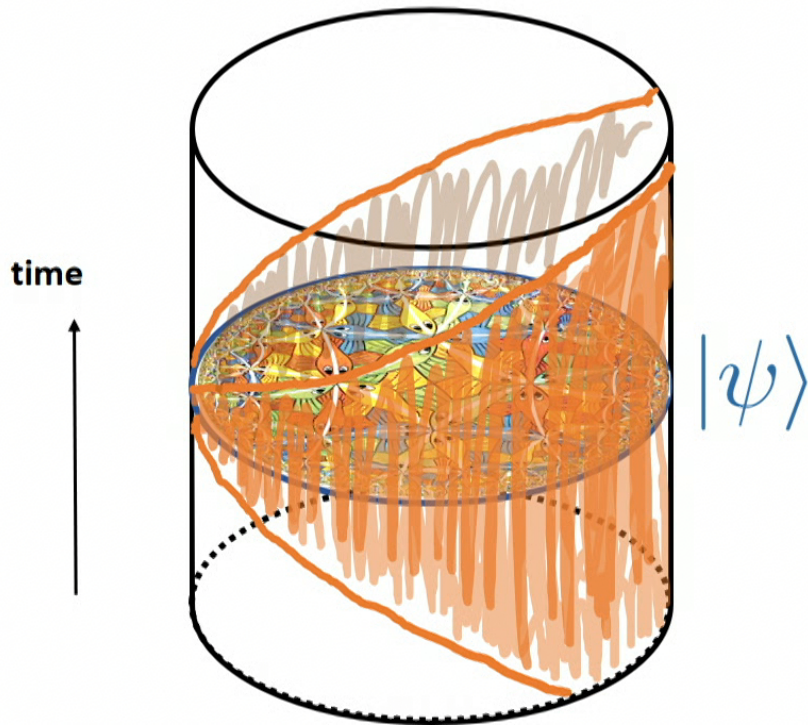


## Global AdS

$$ds^2 = \frac{1}{\cos^2 \rho} (-dt^2 + d\rho^2 + \sin^2 \rho d\Omega_{d-1}^2)$$

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## Poincare Patch

$$ds^2 = \frac{1}{z^2} (dz^2 - dt^2 + d\vec{x}^2)$$

## Spatial Slice: Hyperbolic Disk

$$ds^2 = \frac{1}{z^2} (dz^2 + d\vec{x}^2)$$





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# Reconstructing Bulk Geometry

How do we determine the bulk geometry (and the dynamics of this background) from the CFT?

Given a state  $|\psi\rangle$  want to know:

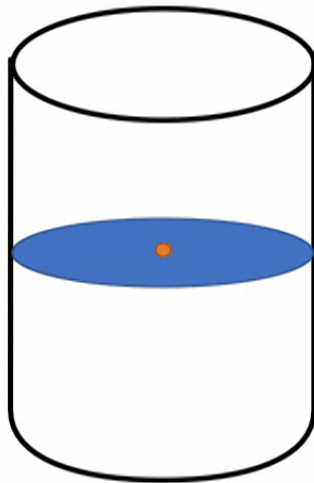
- 1) When is  $|\psi\rangle$  dual to a classical gravitational background?
- 2) What probes can we use to most efficiently determine the classical background?
- 3) How do we describe the local dynamics in these backgrounds?

# Local Bulk Operators

## What does local bulk physics look like in terms of CFT operators?

- Best understood perturbatively about the AdS vacuum:
  - Near the boundary, AdS/CFT dictionary:  $\lim_{z \rightarrow 0} z^{-\Delta} \phi(x, z) = O(x)$
  - Further into the bulk: [HKLL](#)

### Global Reconstruction:



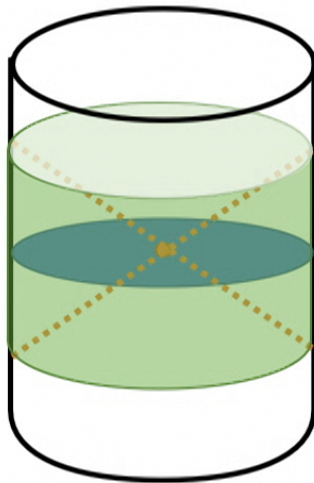
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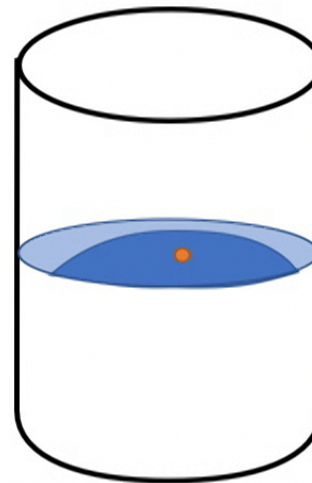
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### Global Reconstruction:

$$\phi(x, z) = \int d^d x' K_g(x, z | x') \mathcal{O}(x')$$



### Rindler Reconstruction:



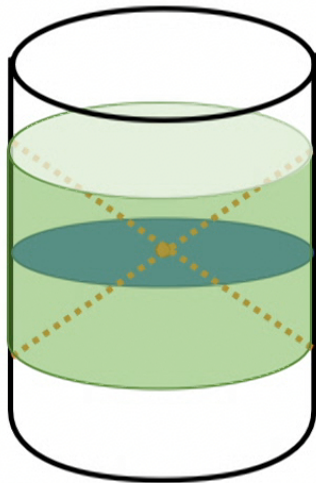
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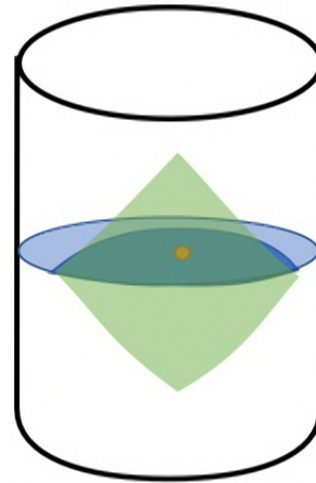
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### Rindler Reconstruction:

$$\phi(x, z) = \int d^d x' K_r(x, z | x') \mathcal{O}(x')$$



- **Local Bulk Operator**  
=  
**Non-local boundary operator**
- Supported on **causal region**

# Geometry and Entropy

## What about non-perturbative probes to study non-trivial backgrounds?

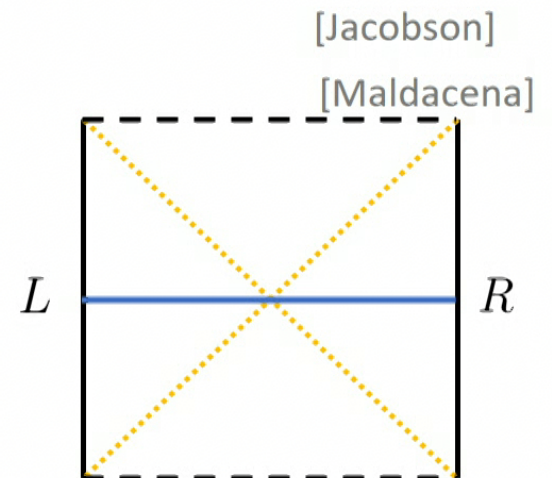
- Long history suggestive that there is a deep connection between spacetime and entropy:

- **Black hole thermodynamics:**  $\delta E = T dS$

$$S_{BH} = \frac{A}{4G_N}$$

- BH thermo can be used to **derive Einstein equation**
- **Thermofield double state** of two CFTs (L and R):

$$|TFD\rangle = N \sum_E e^{-\beta E/2} |E\rangle_L |E\rangle_R$$



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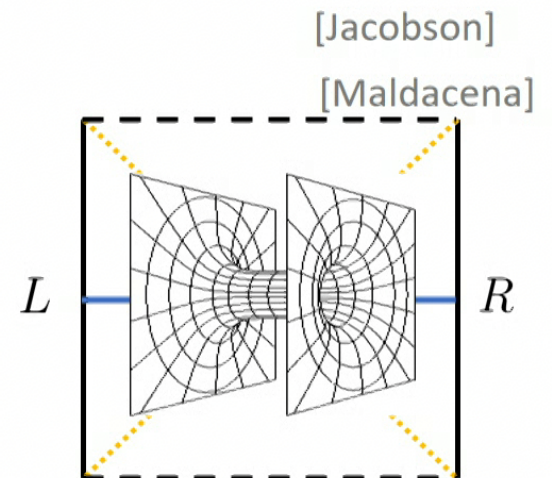
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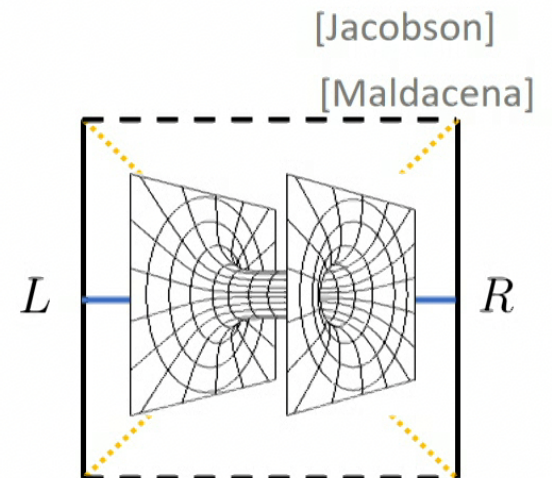
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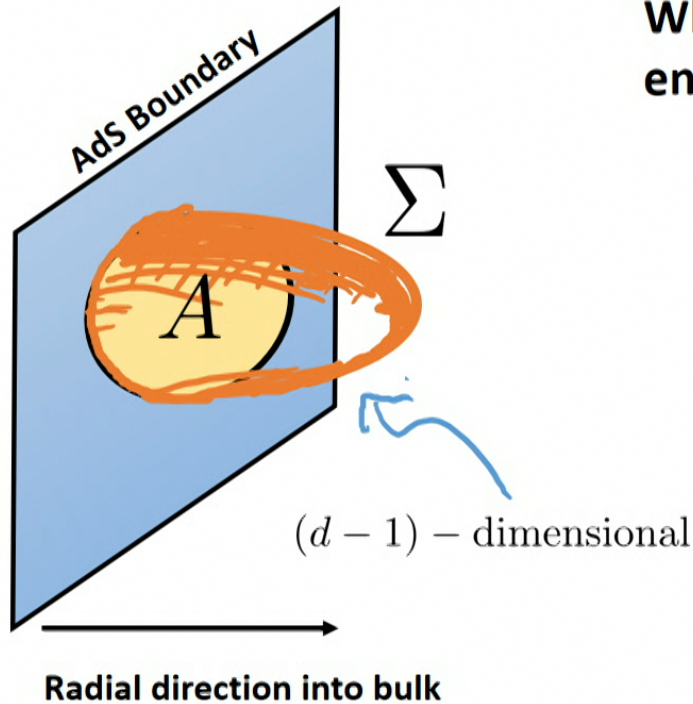
$$S_{BH} = -\text{tr} [\rho_R \log \rho_R]$$

**Black hole entropy is entanglement entropy**



# Ryu-Takayanagi

- Consider a spatial region of a single CFT:



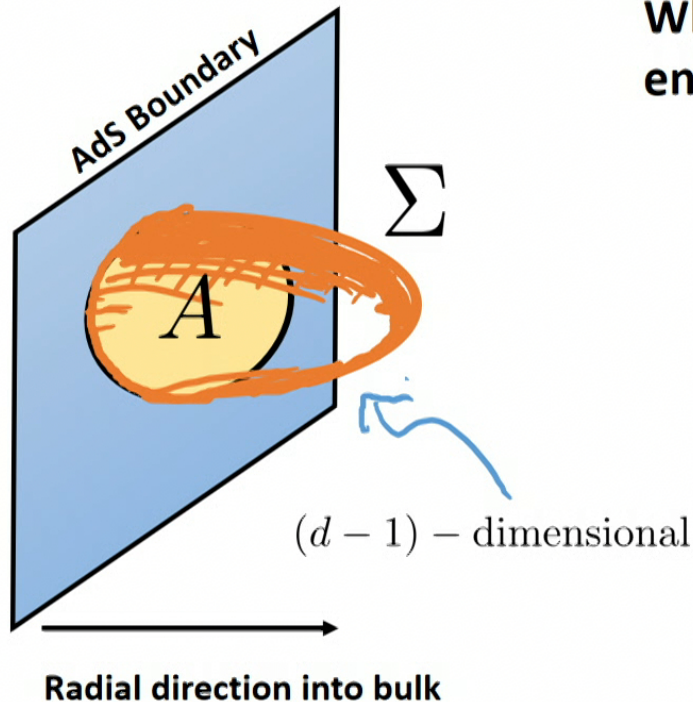
What bulk object computes the entanglement entropy of the region A?

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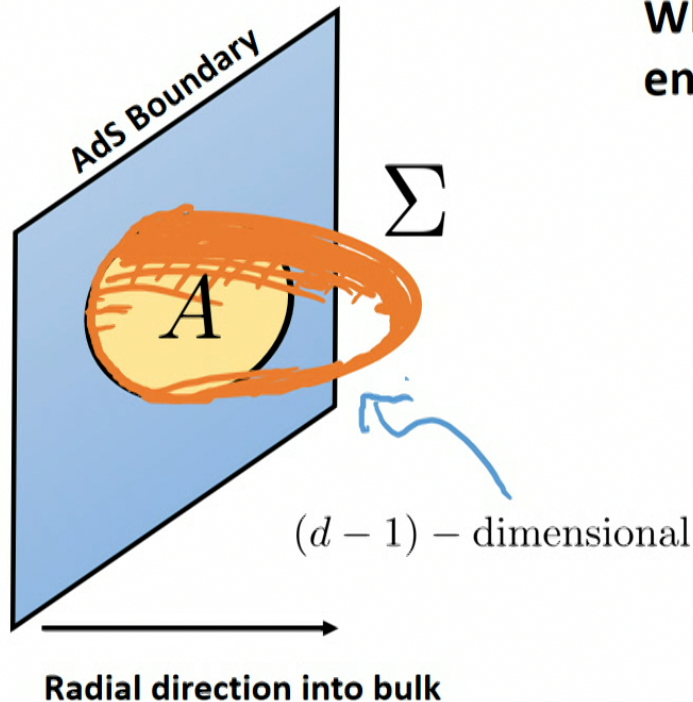
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**Entanglement entropy** gives us **non-perturbative, non-local** probes of the dual **bulk geometry**

---

[Lashkari, McDermott, van Raamsdonk; Faulkner, Guica, Hartman, Myers, van Raamsdonk; Swingle, van Raamsdonk; Lin, Marcollil, Ooguri, Stoica; Czech, Lamprou, McCandlish, Mosk, Sully;]

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Suggestive of picture where CFT **entanglement seems to constitute bulk geometry.**

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- A lot of progress showing how to read off bulk geometry and Einstein equation from this data (**Bulk Tomography**)
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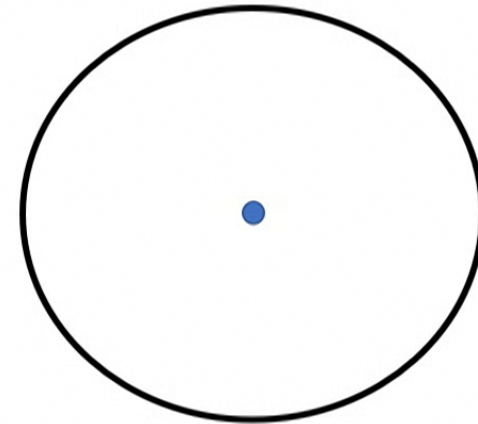
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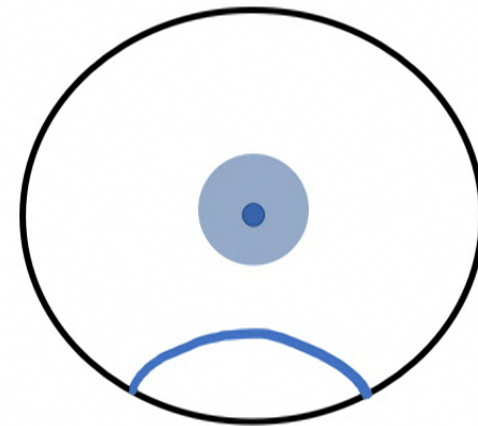
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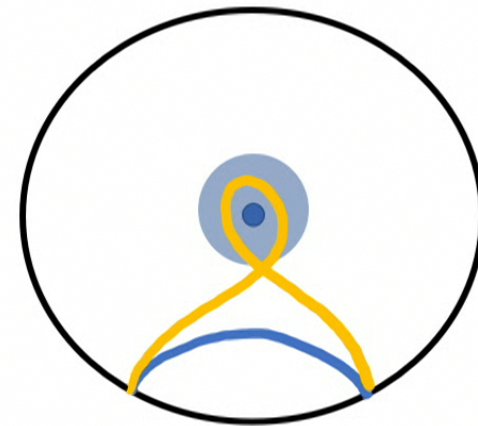
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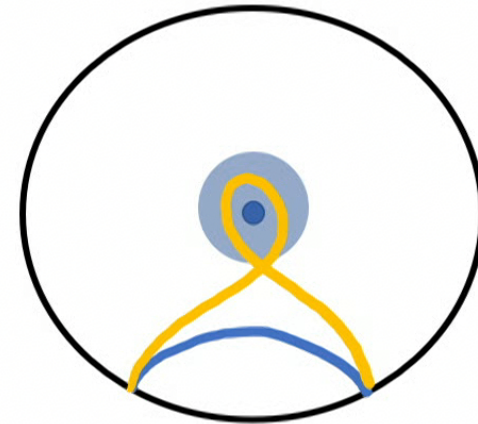
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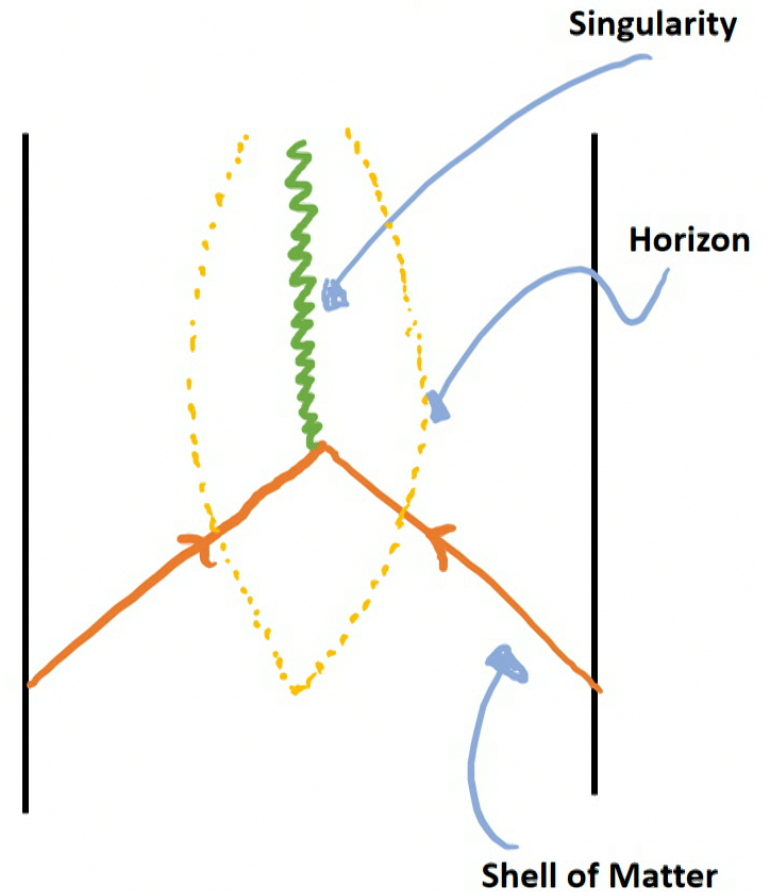
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**What CFT quantities are needed to see the complete bulk geometry?  
Is a different type of CFT correlation responsible for generating these?**

# Black Hole Interior

One of the most important examples of a region not probed by entanglement is the interior of a black hole:

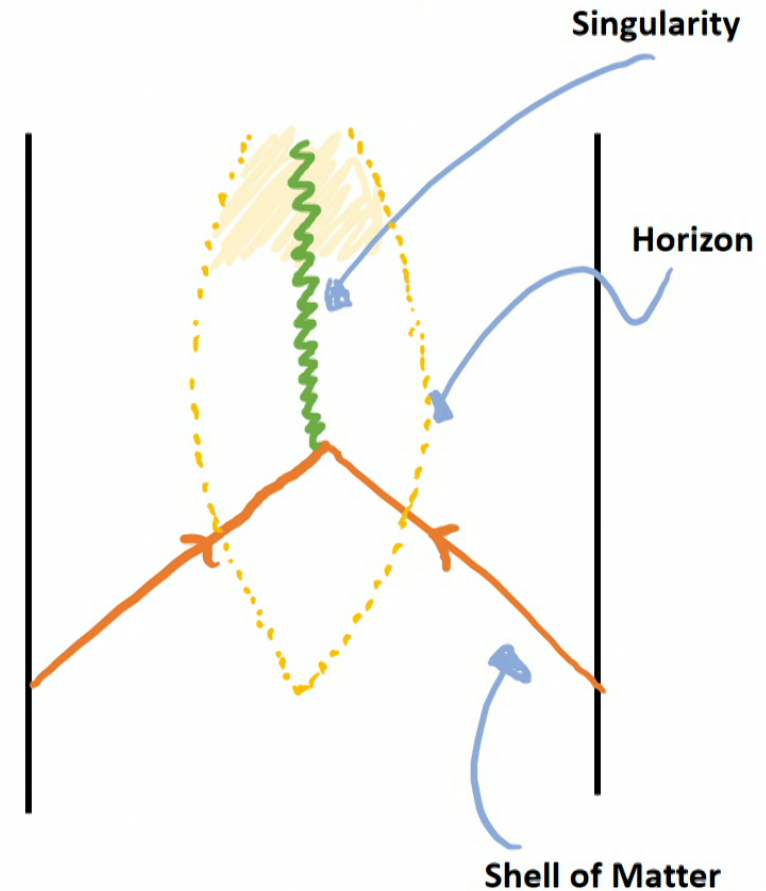




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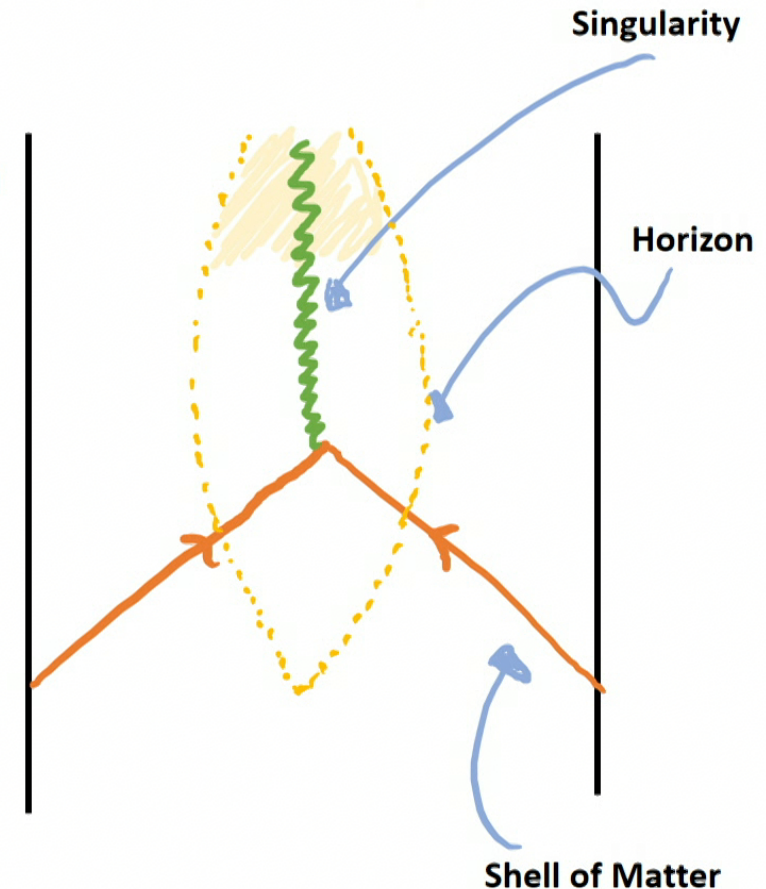


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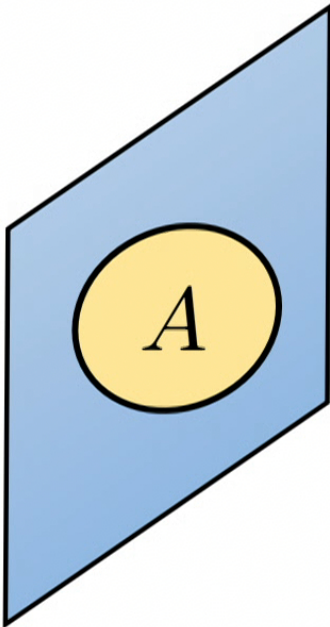
If this region of spacetime exists, how is it encoded by the CFT?



# Sub-region Duality

**How non-local is the mapping between the boundary and the bulk?**

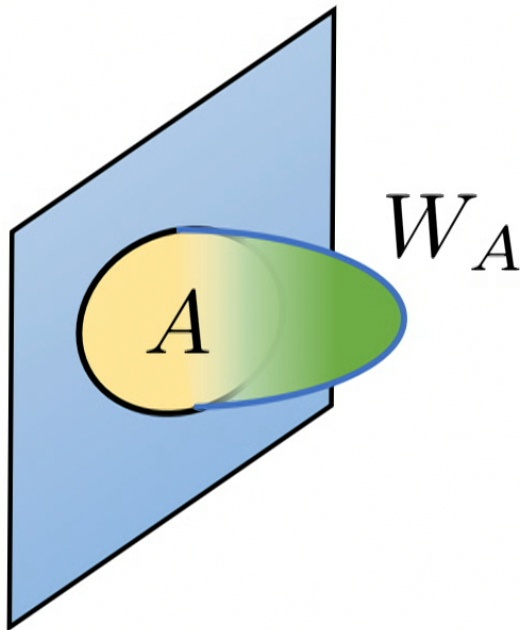
- Given the **density matrix** for a **spatial region** on the boundary, **how much of the bulk** can I reconstruct?
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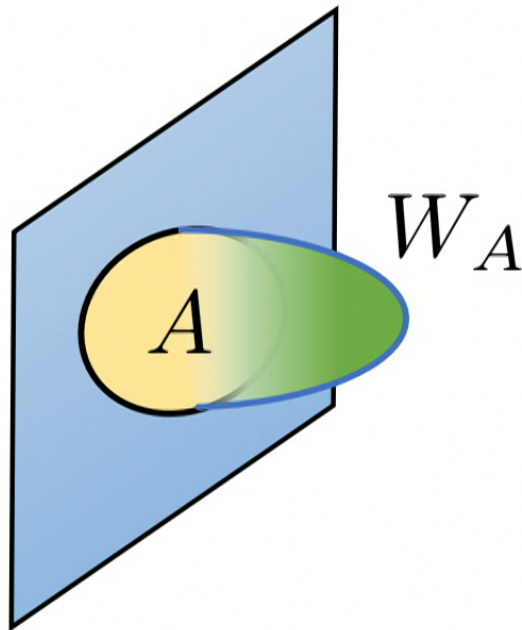


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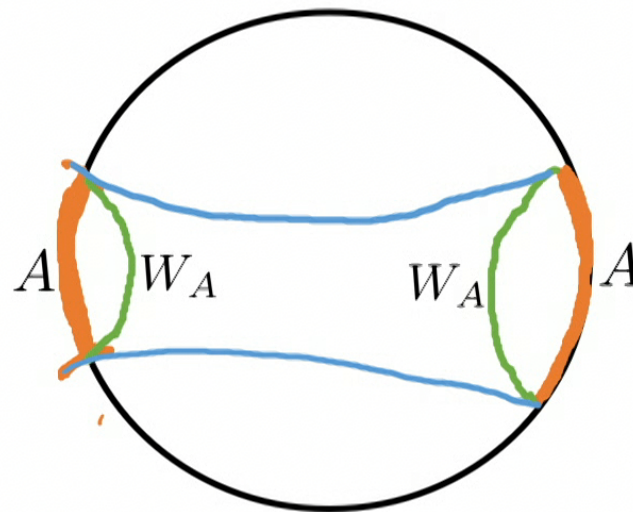
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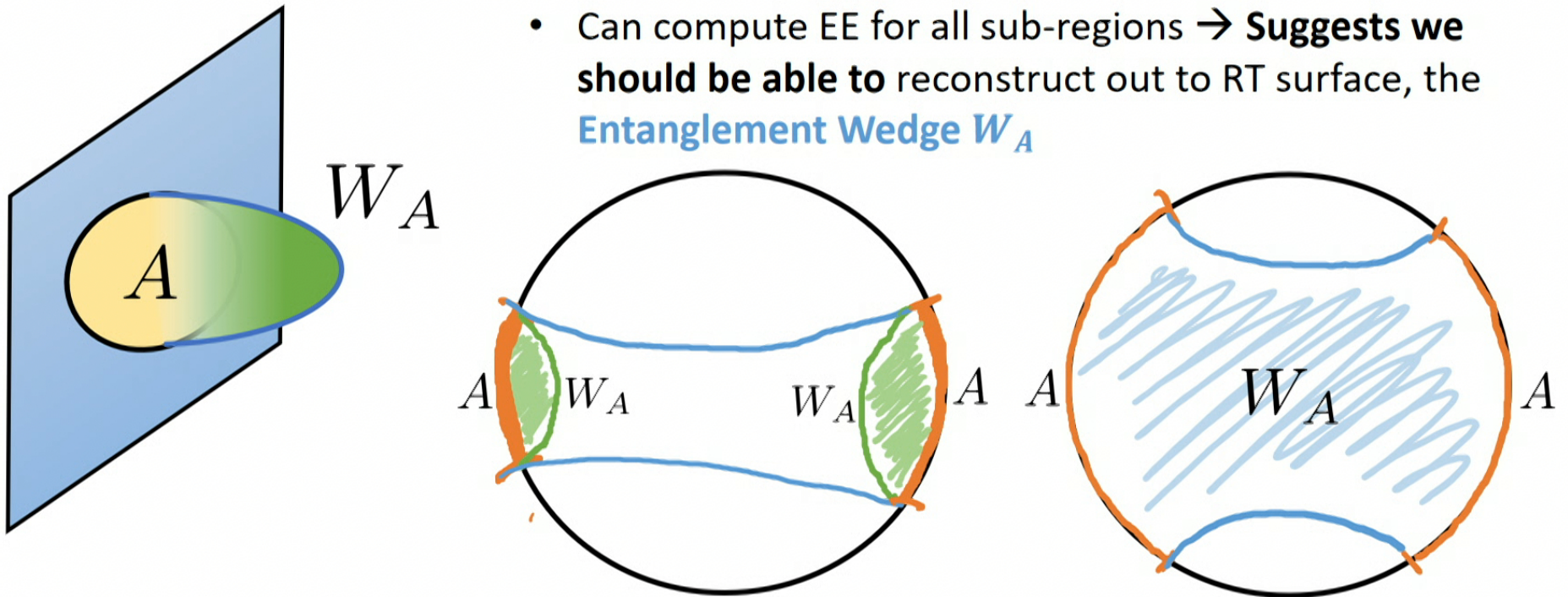


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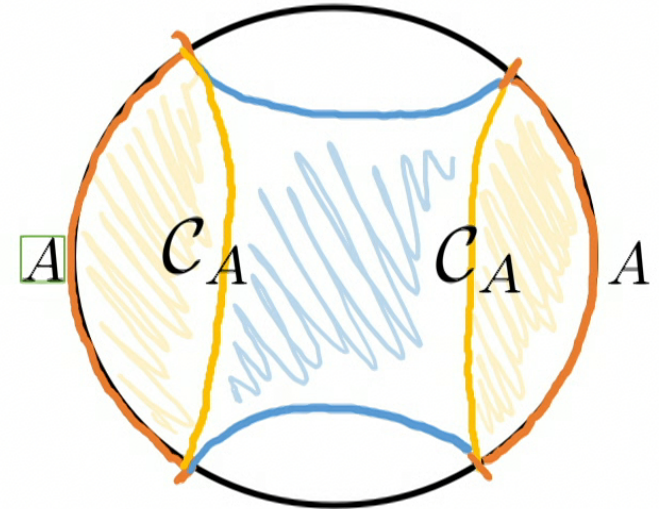


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## What can we actually do?

### Constructive:

- Can build bulk operators in regions **causally connected** to boundary (**HKLL**)



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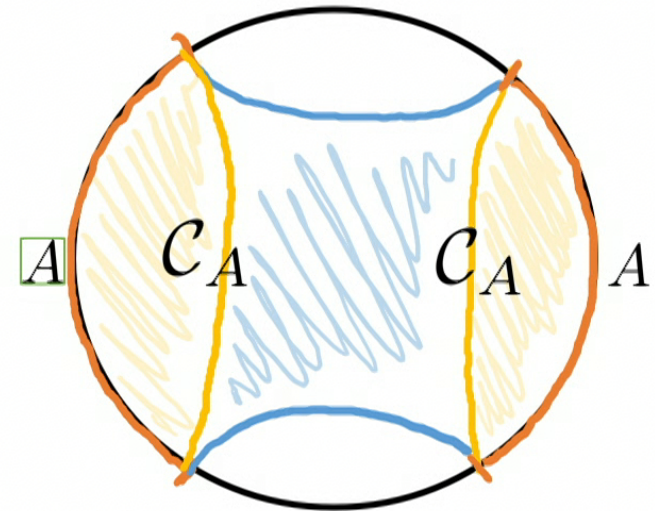
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### Less Constructive:

[Dong, Harlow, Wall]



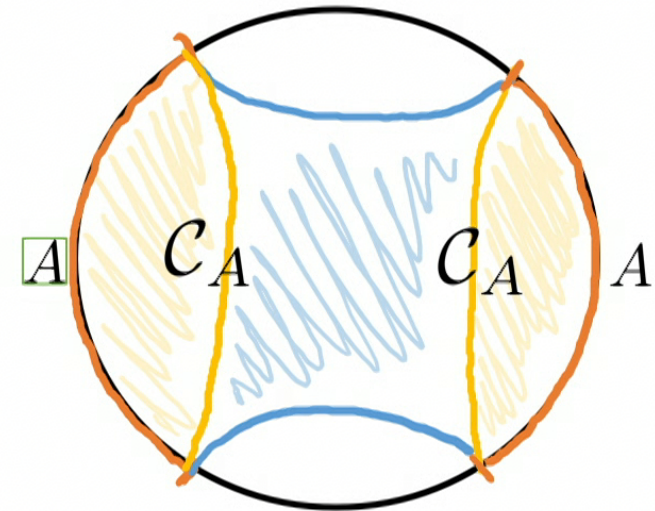


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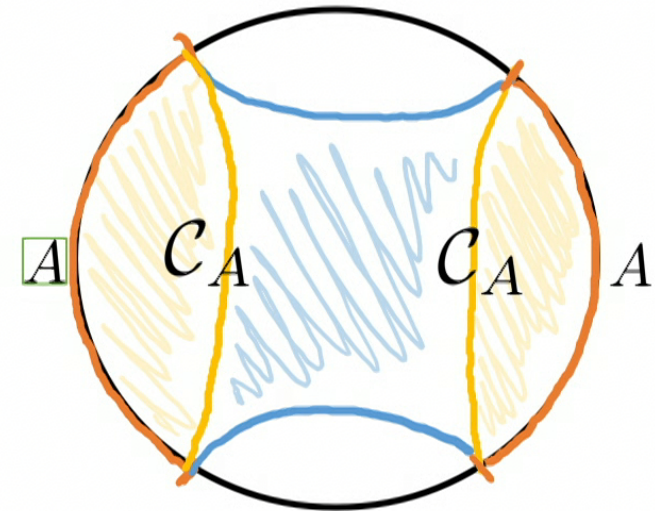
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**Are there concrete scenarios to explore sub-region duality?**

---

# Toy Models for Holography?

To summarize, we are interested in:

1. How do we see the emergence of classical bulk geometry?
2. What are effective non-perturbative probes of this geometry?
3. How do we see into regions shadowed from entanglement?
4. Are there practical approaches to understanding duality for sub-regions?

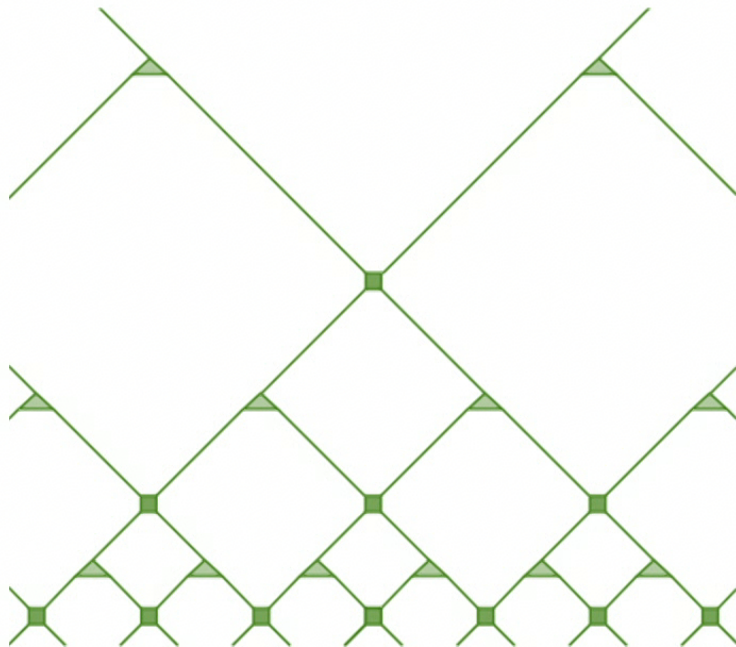
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AdS / MERA

# MERA and Holography

[Swingle]

Sometime circa 2009, Brian Swingle was staring at a MERA network...

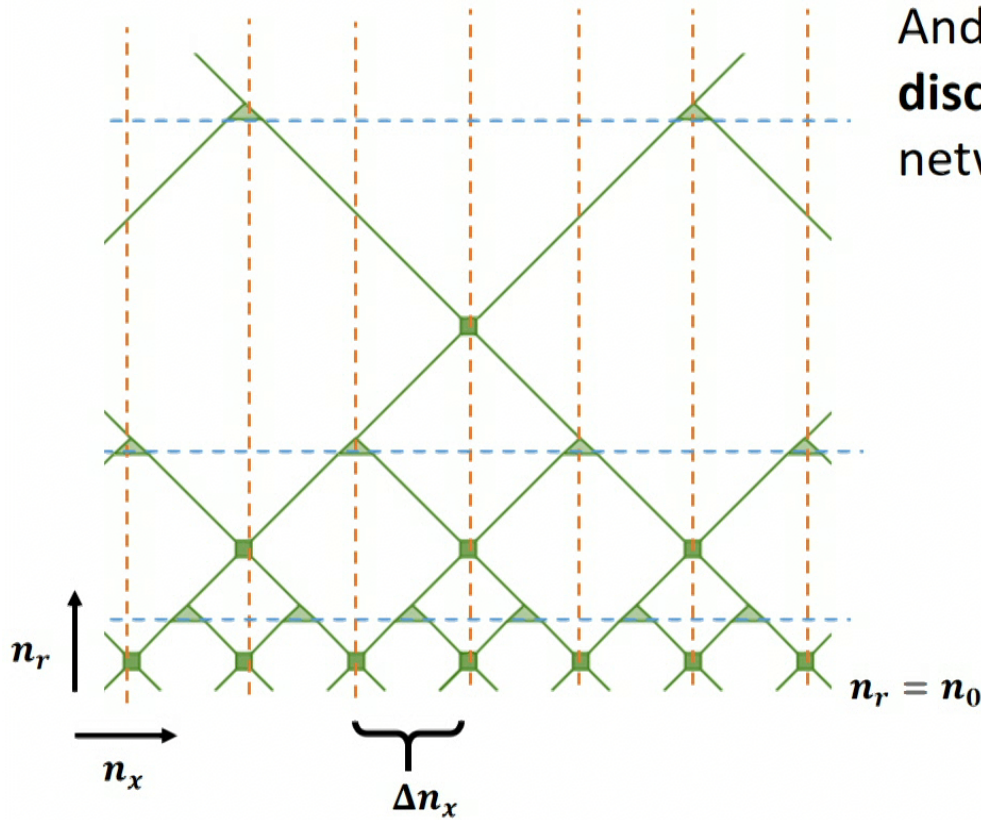


And realized that one can assign a **natural discrete Euclidean metric** to translations in the network:

# MERA and Holography

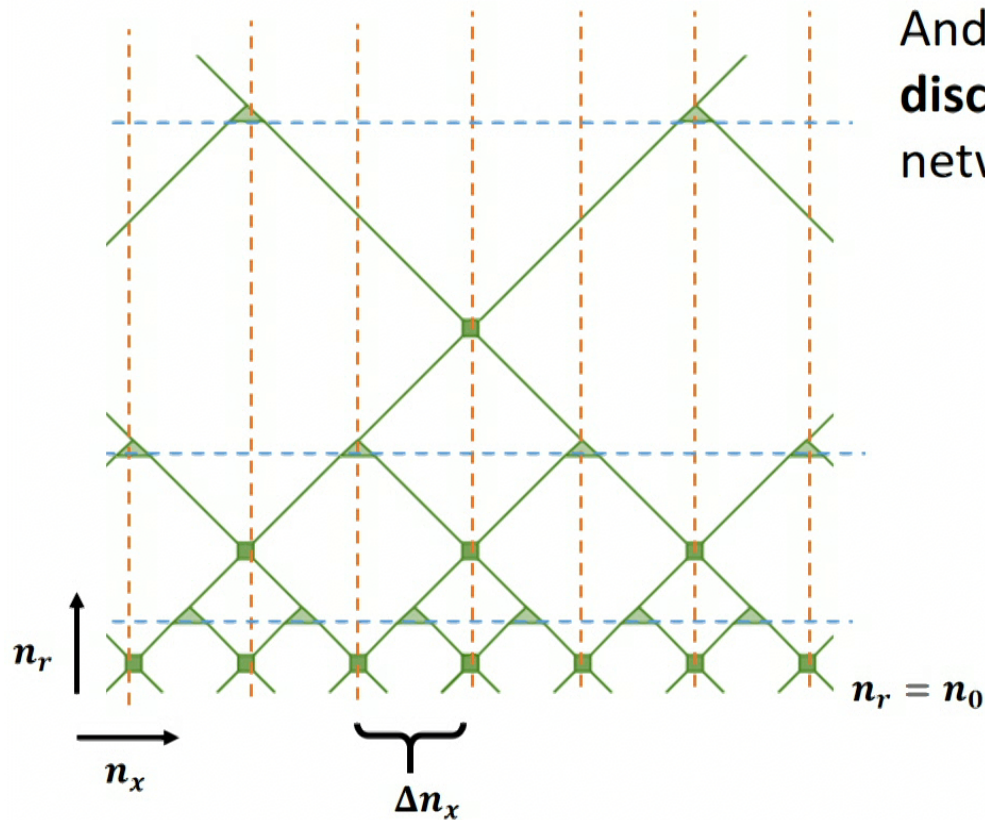
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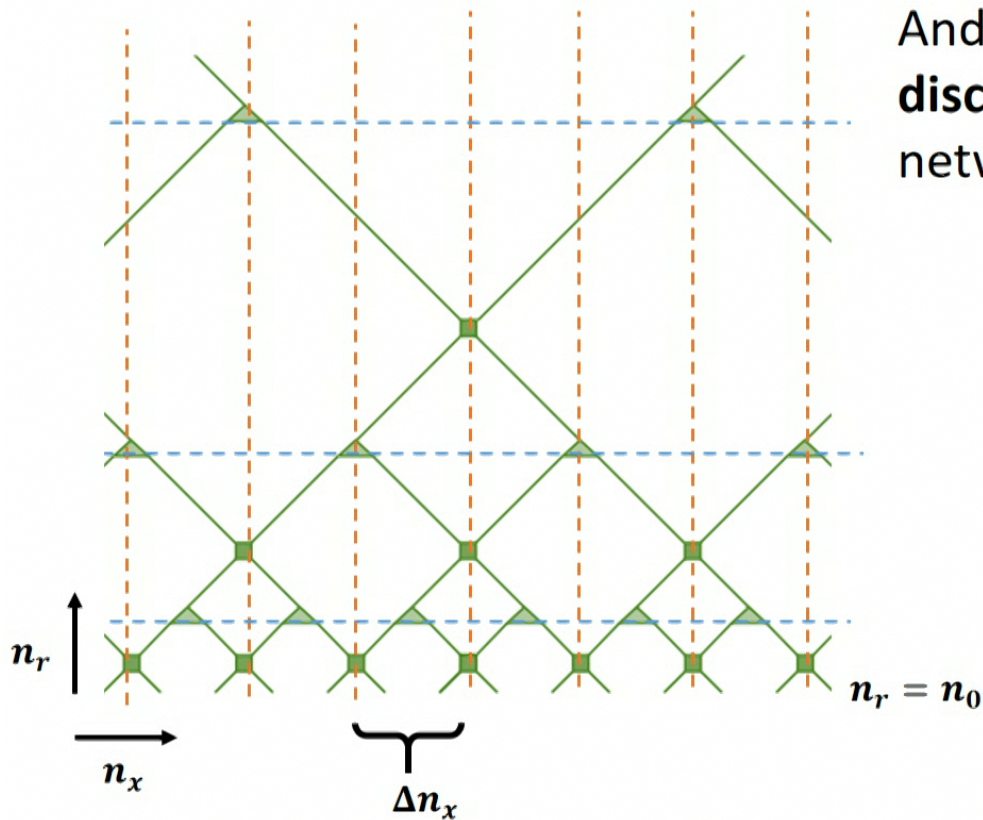


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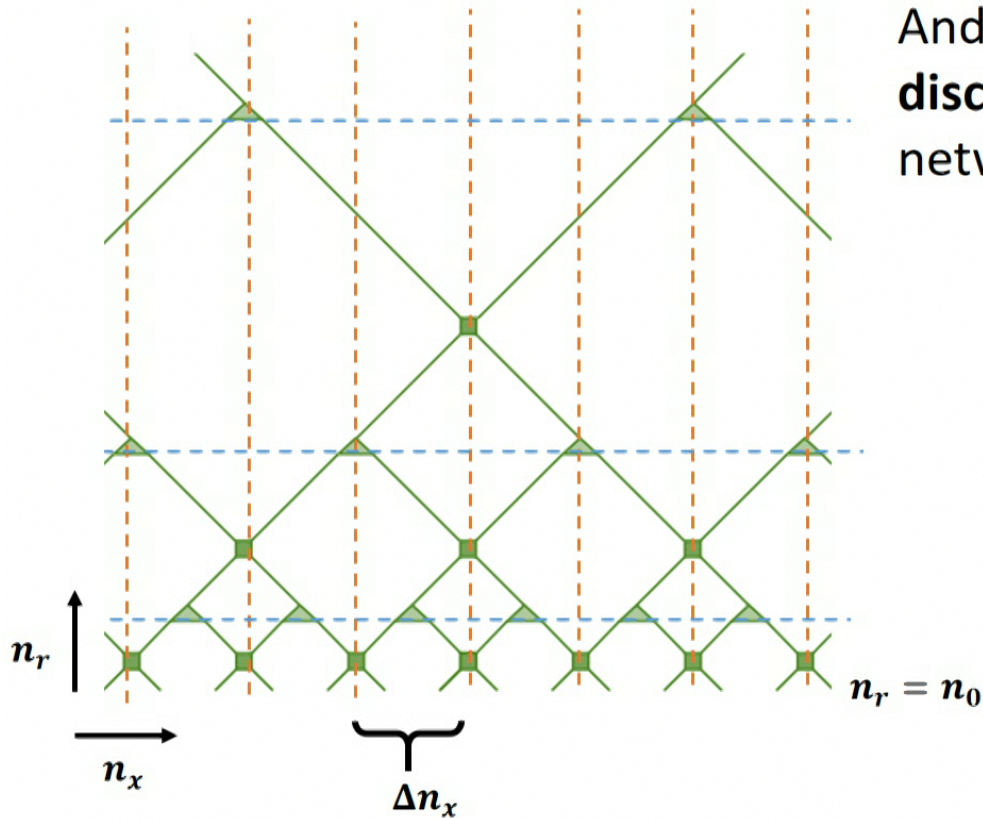
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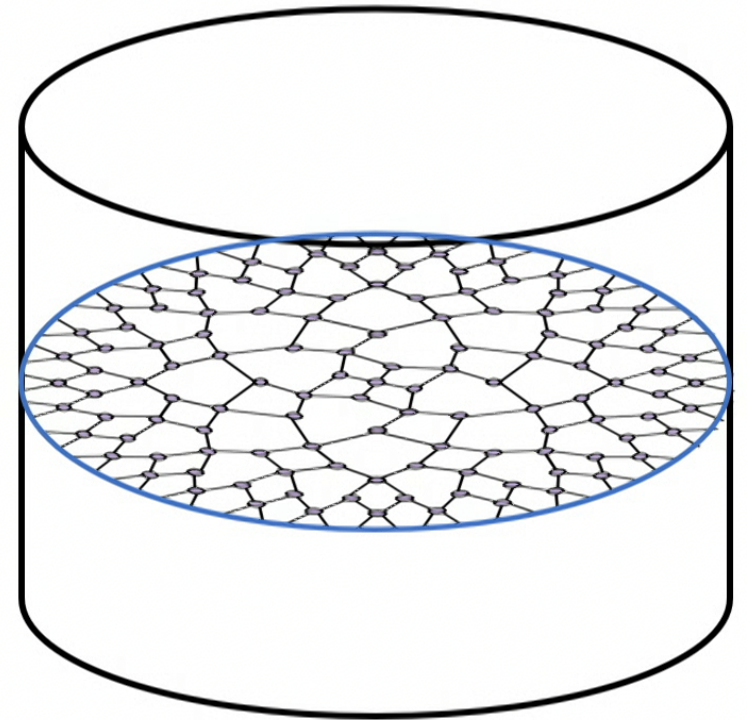
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**Hyperbolic Metric**

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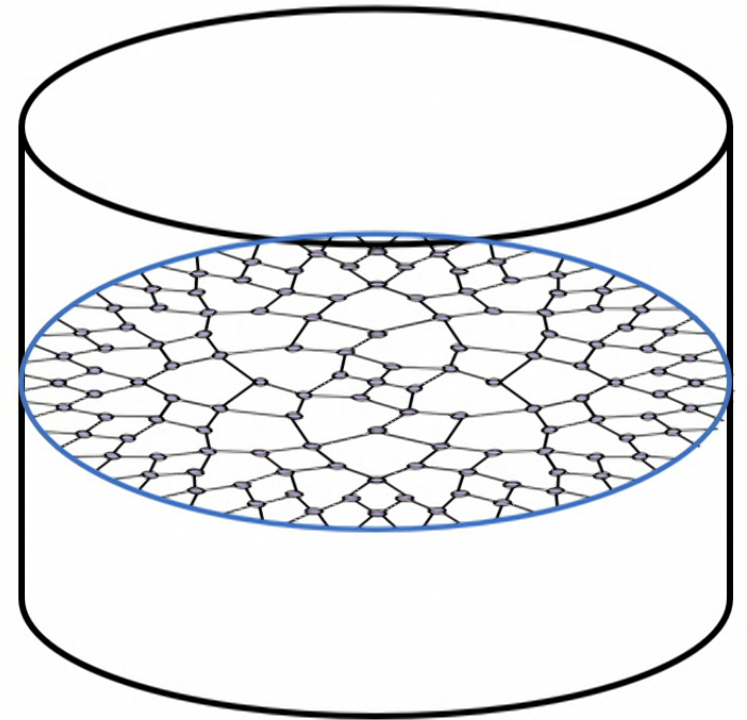
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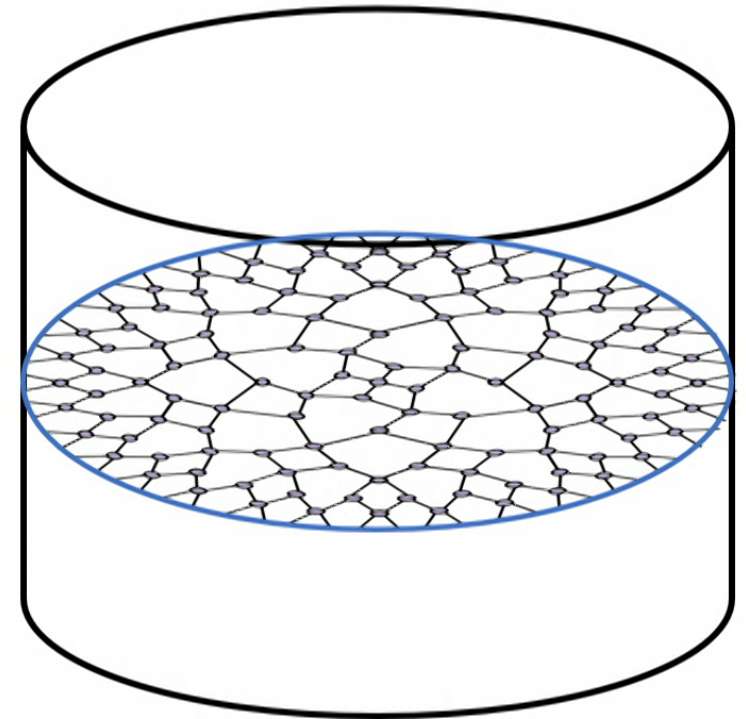
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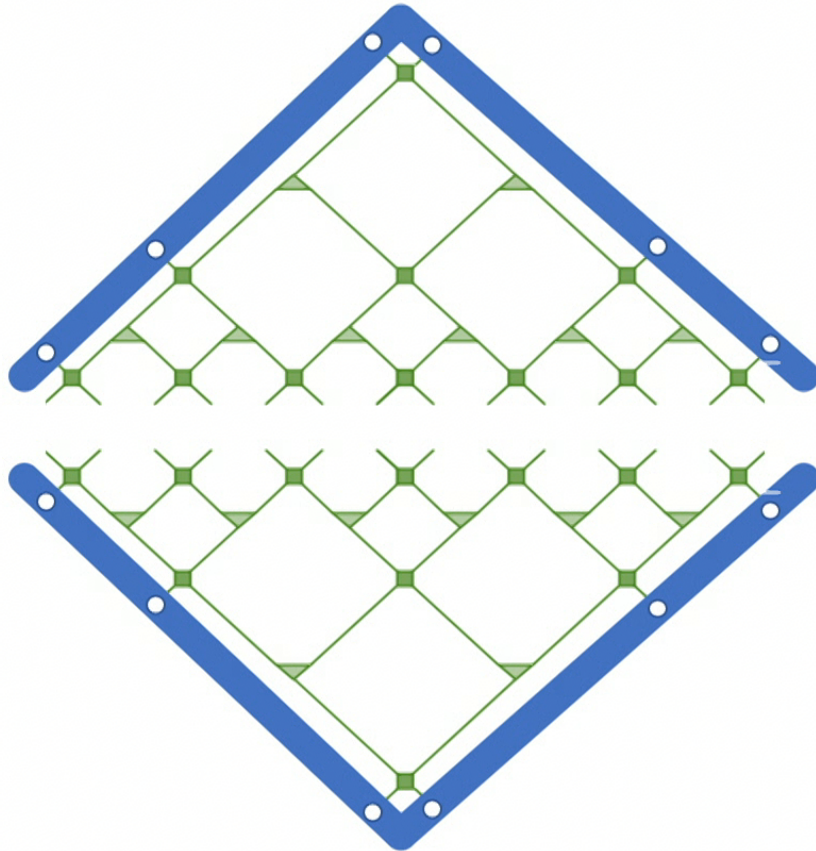
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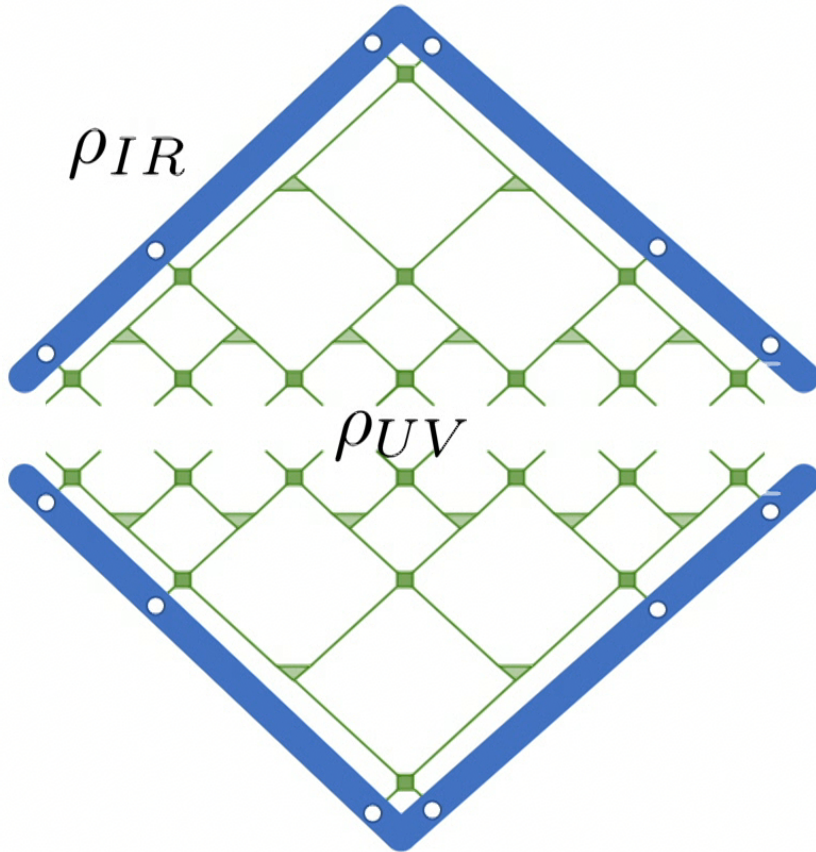
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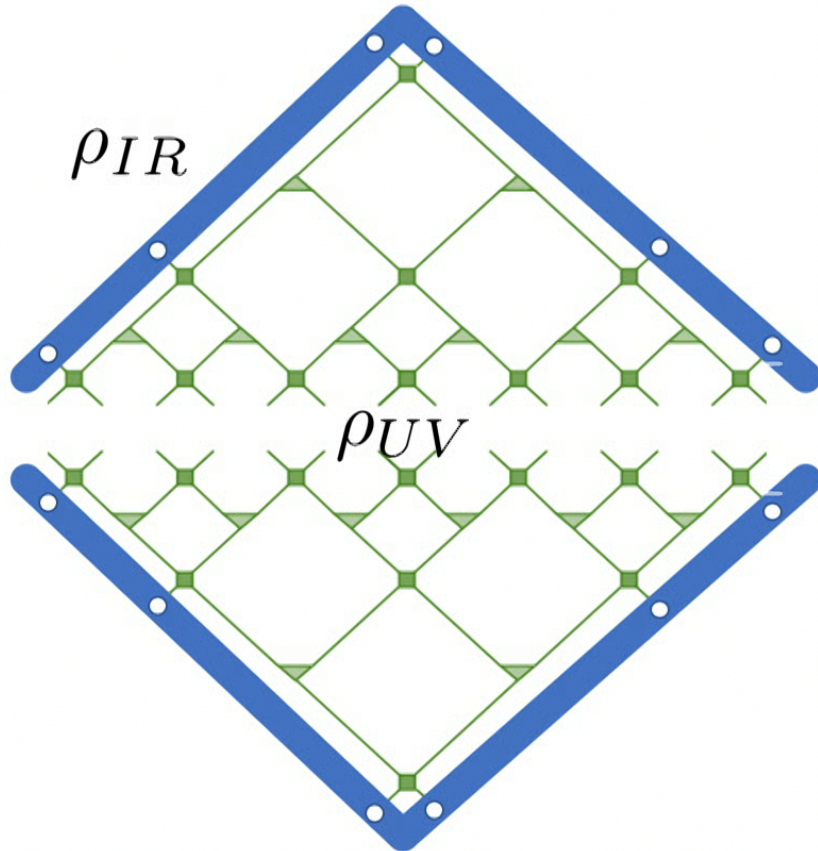
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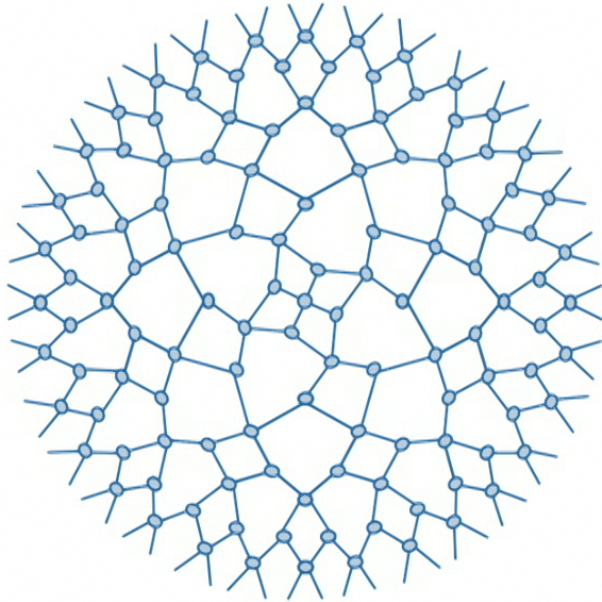
- Better yet, it has been shown that each scale contributes approximately equally to the entropy so that

$$S(A) = n_{cut} \log \chi_{\text{eff}}$$

---

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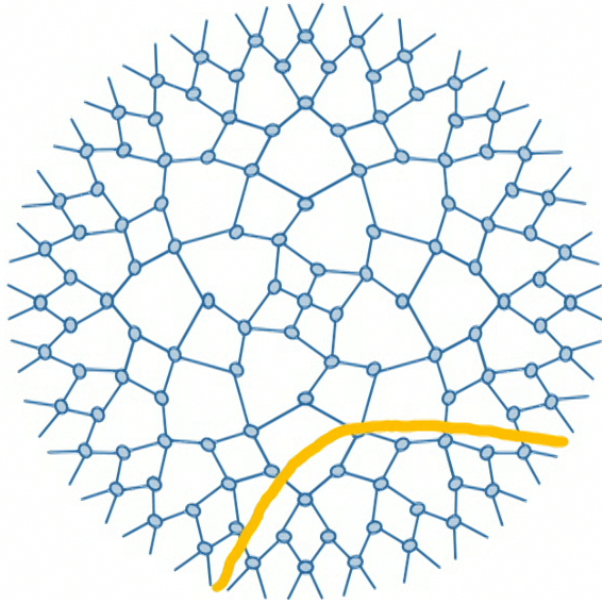
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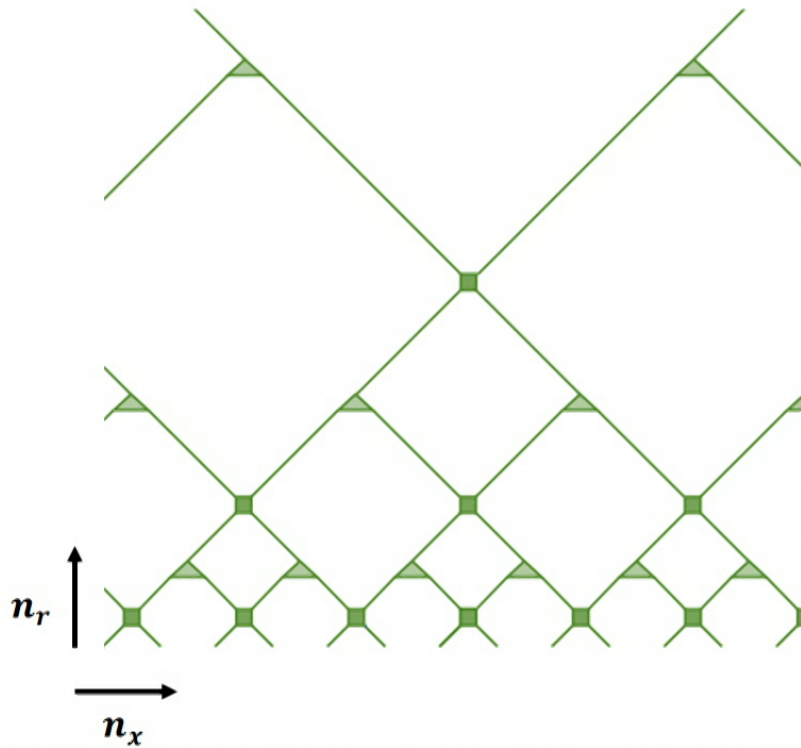


Beautiful correspondence to length of RT surface in spatial slice of AdS!

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[Czech, Lamprou, McCandlish, JS; Beny]

## A short detour:



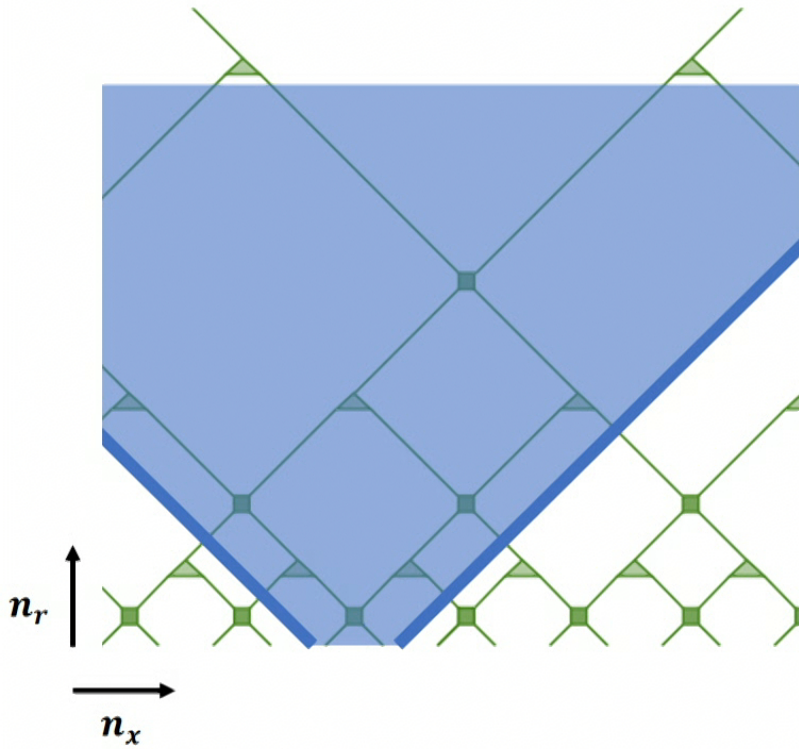
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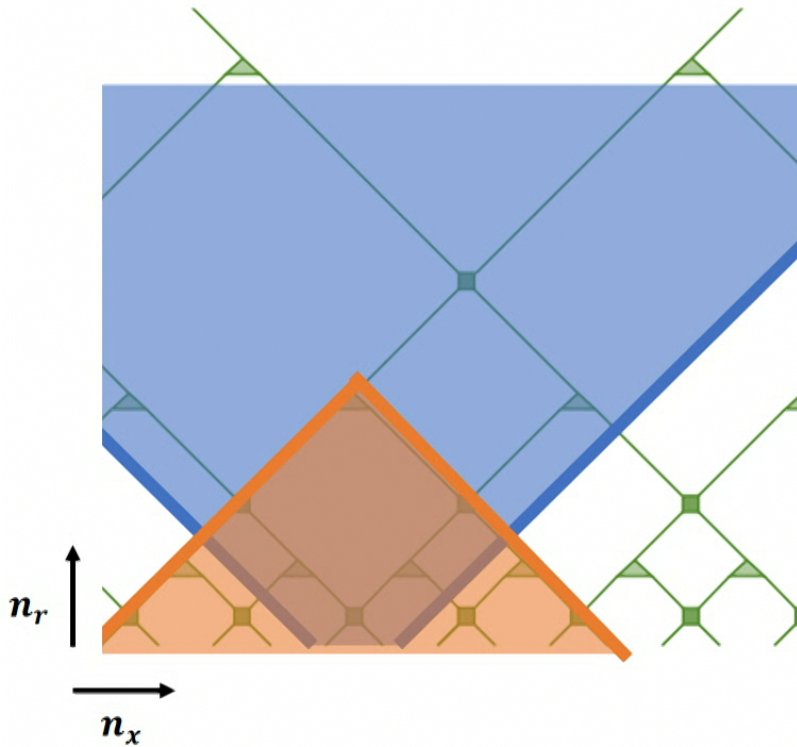


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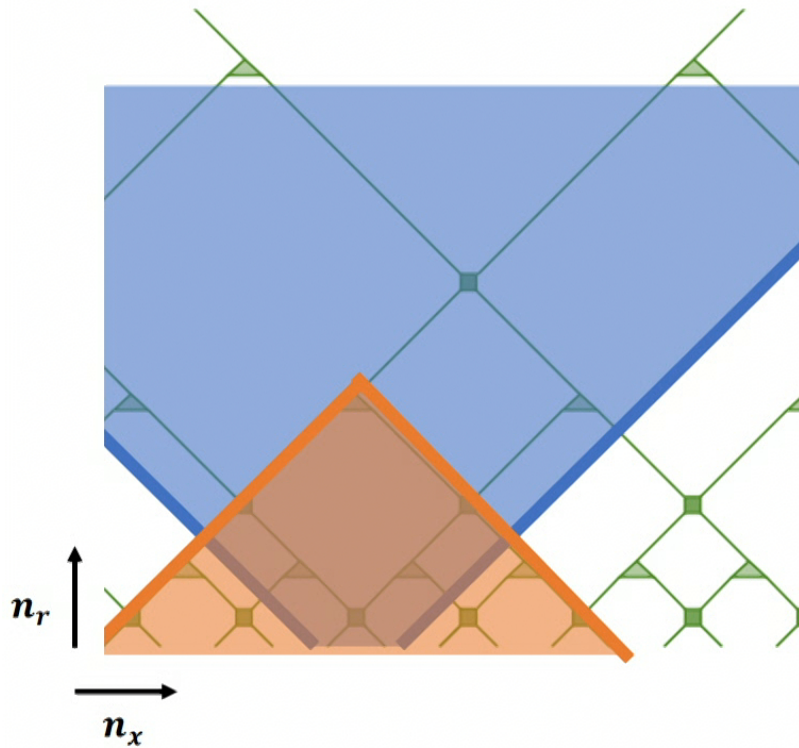
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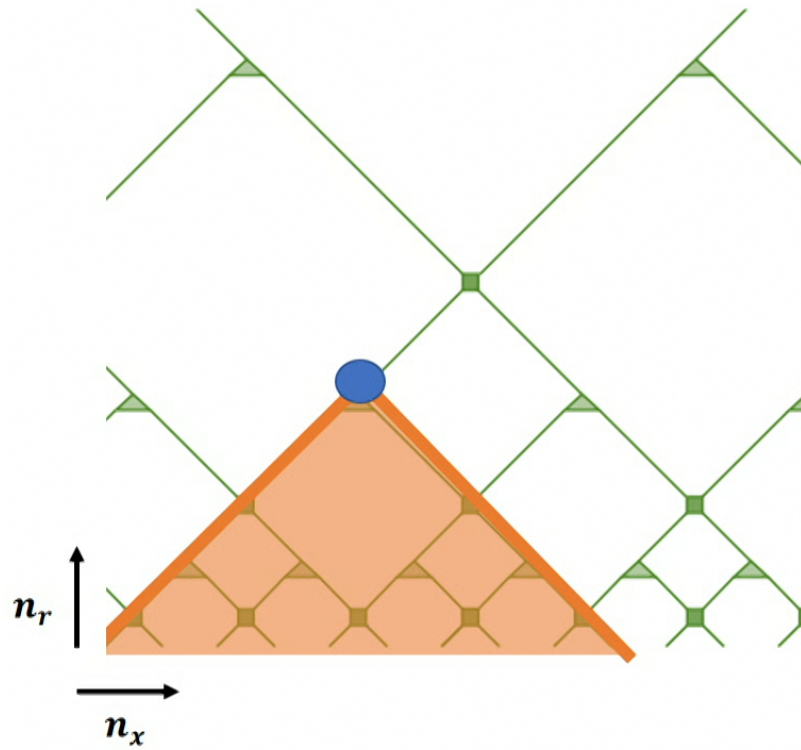
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**de Sitter Metric**

# Kinematic Space

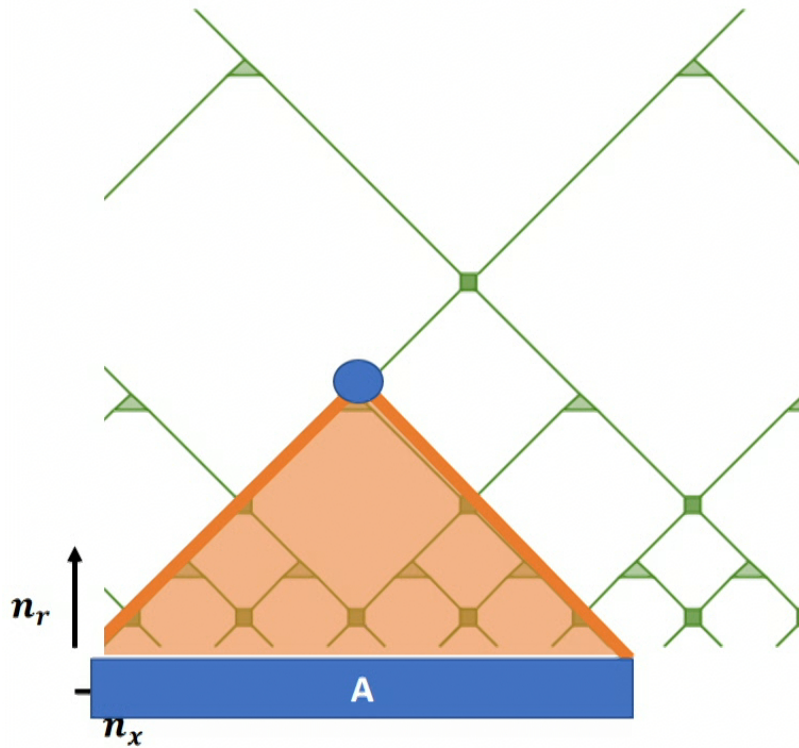
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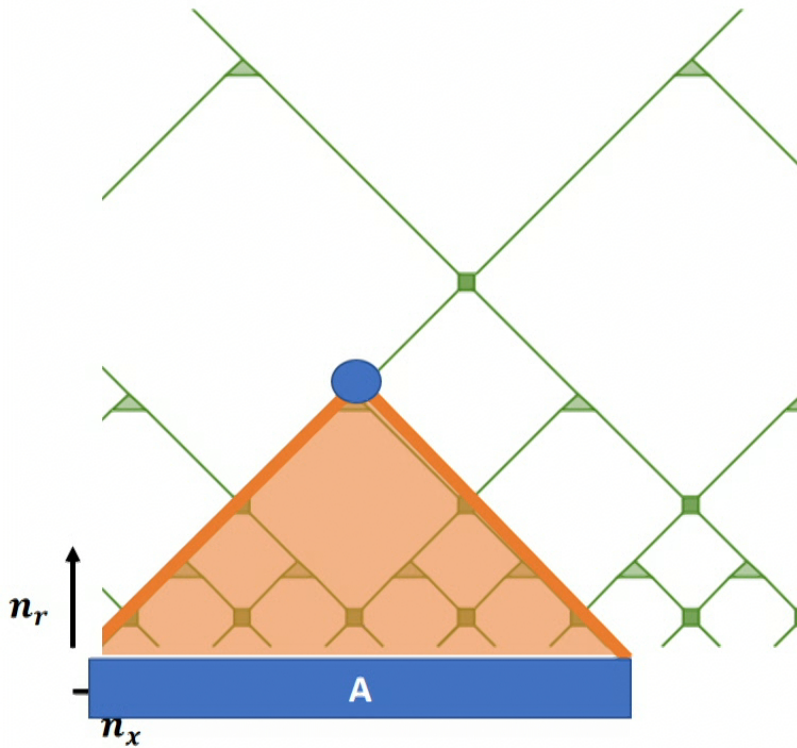


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de Sitter is the space of intervals of the CFT:

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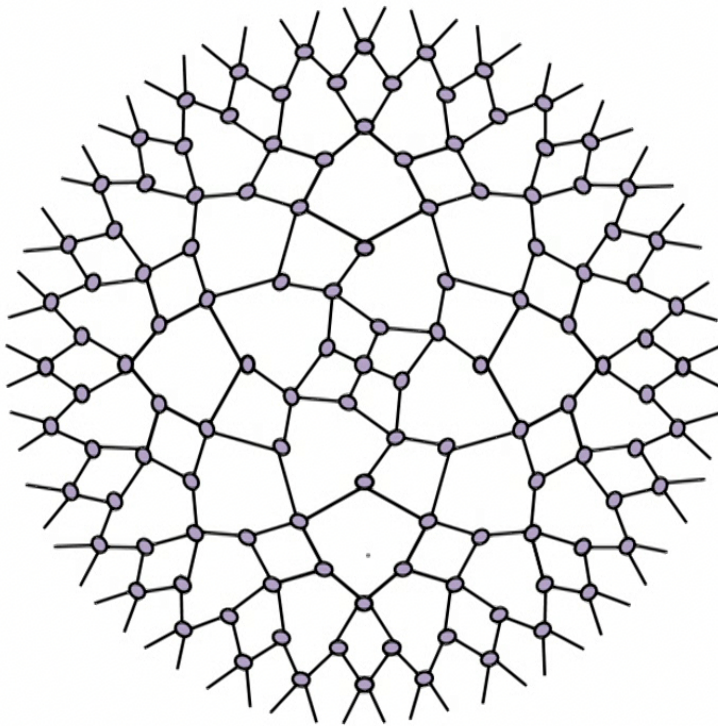
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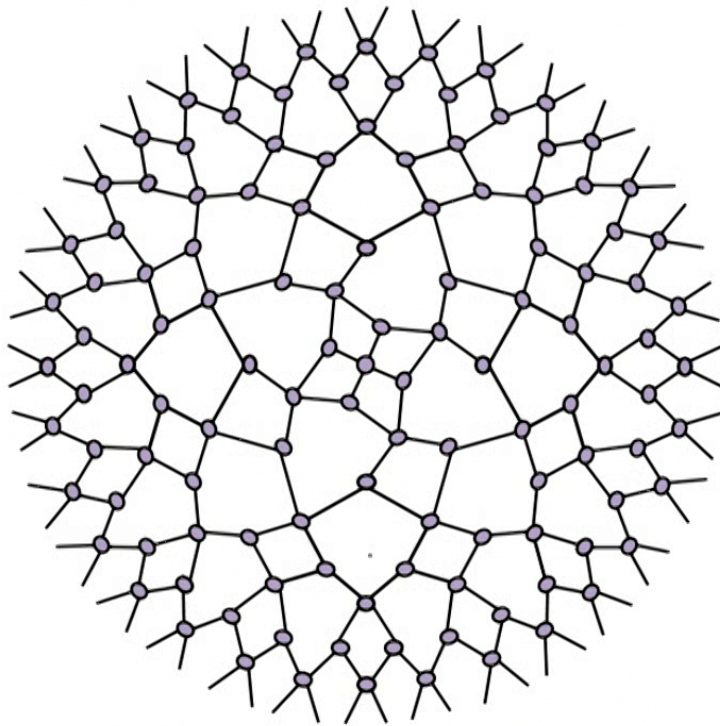


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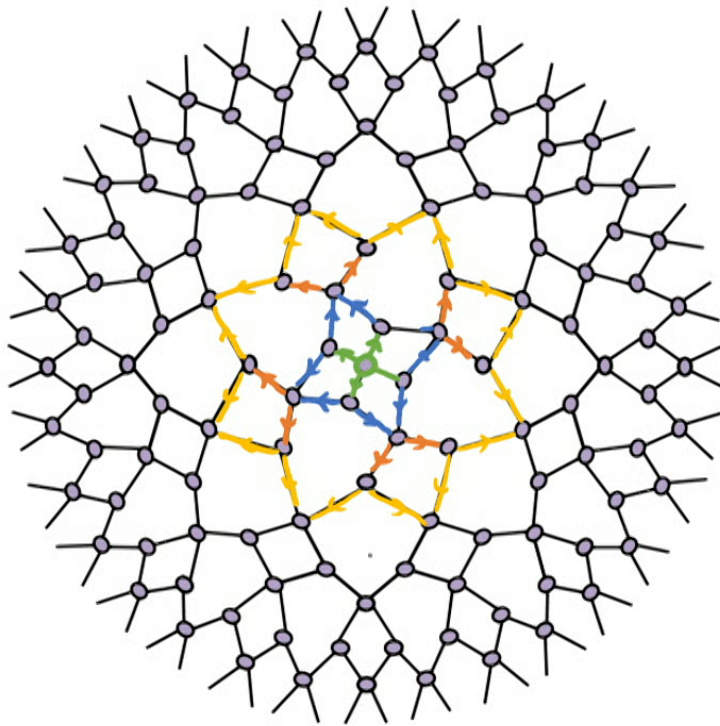


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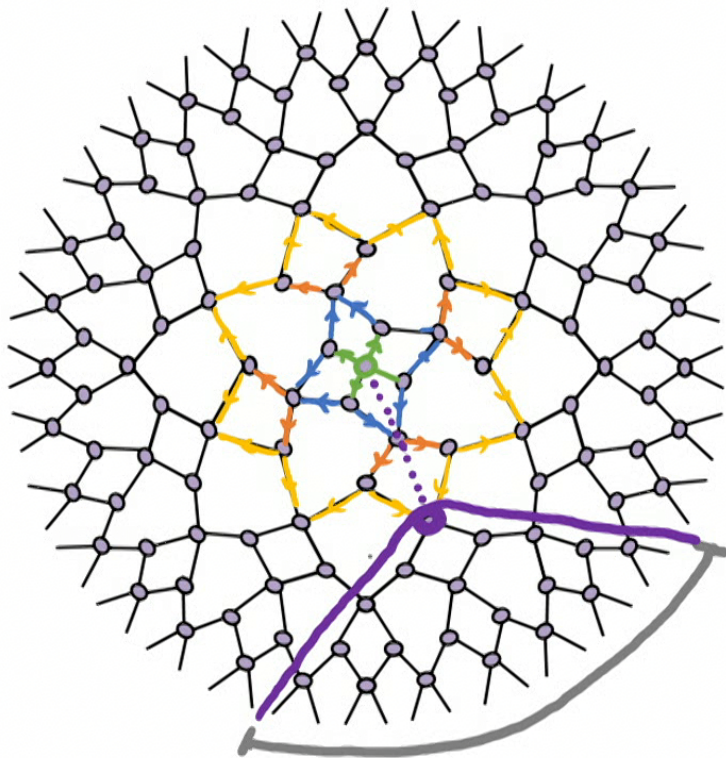


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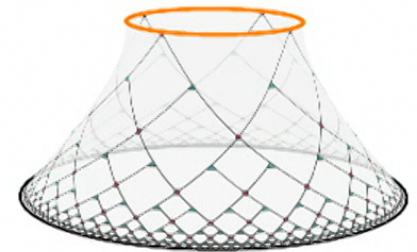
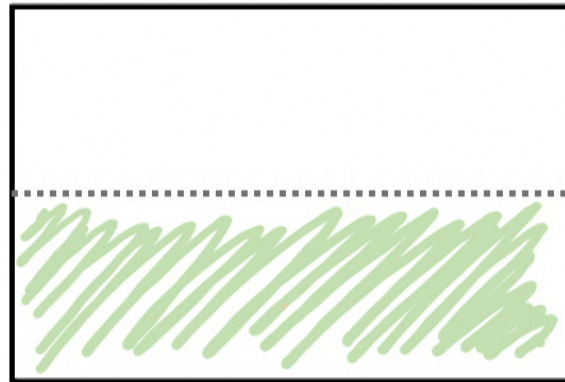
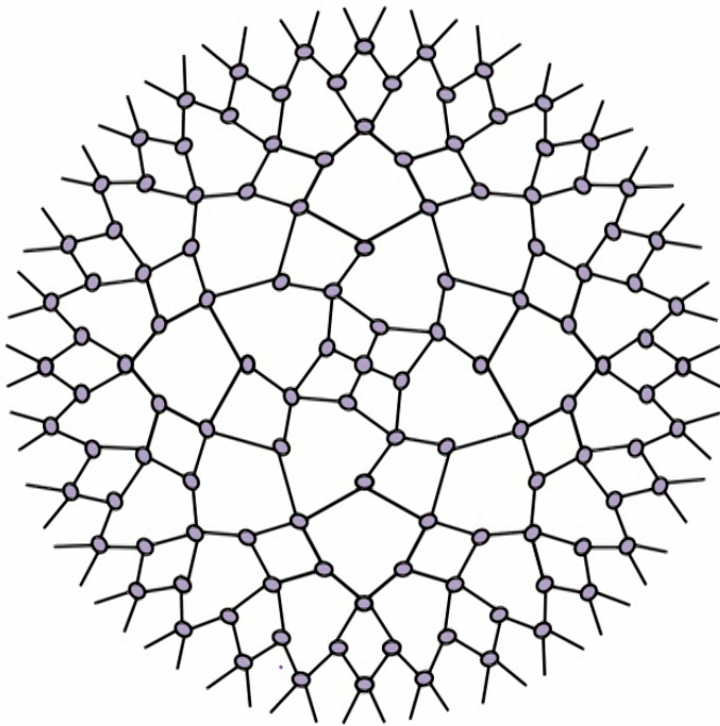


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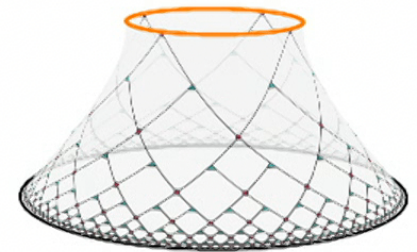
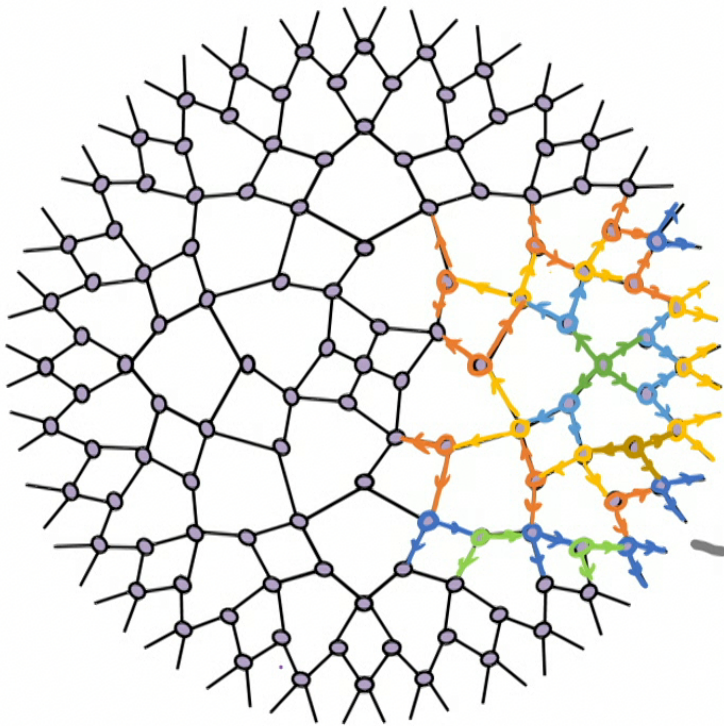
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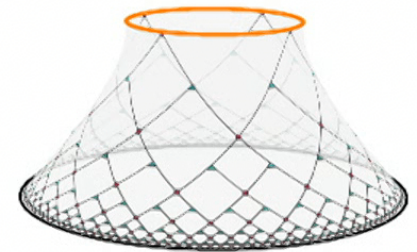
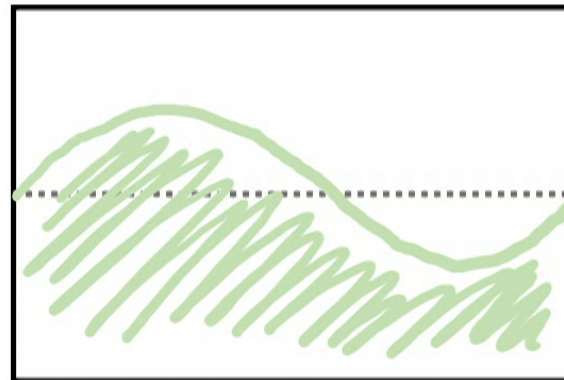
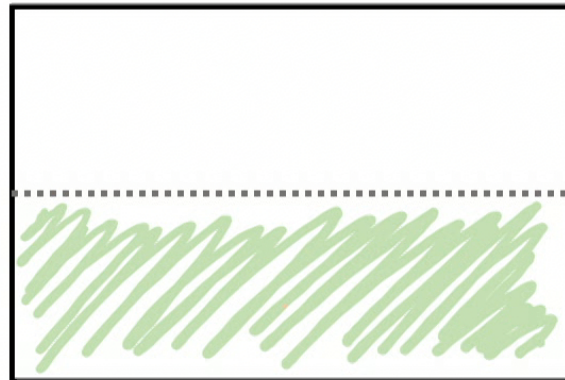
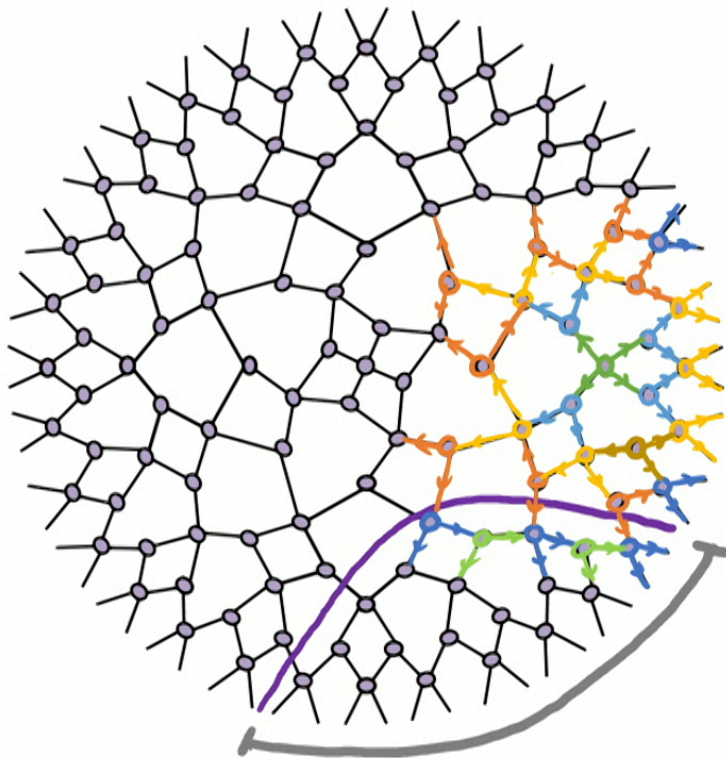




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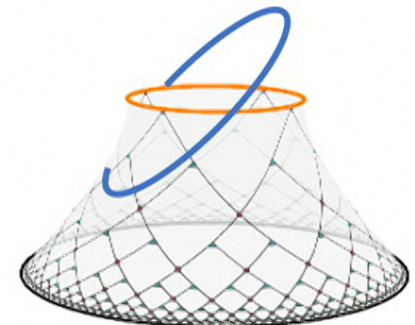
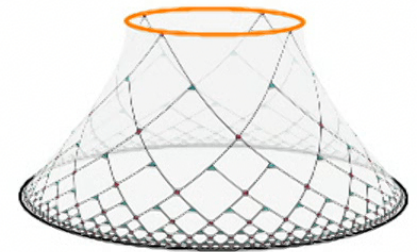
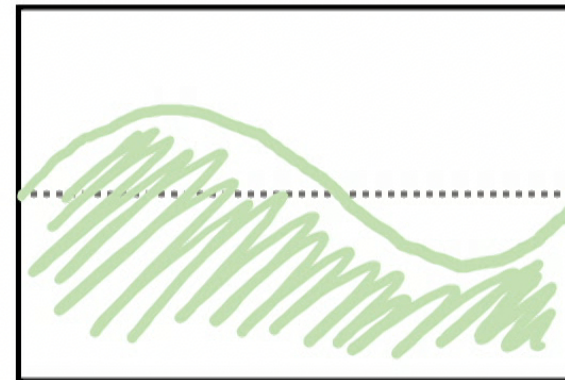
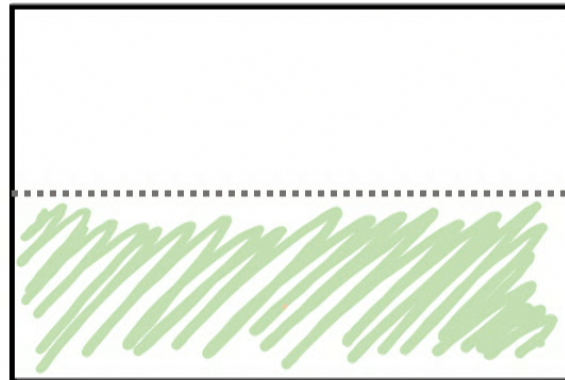
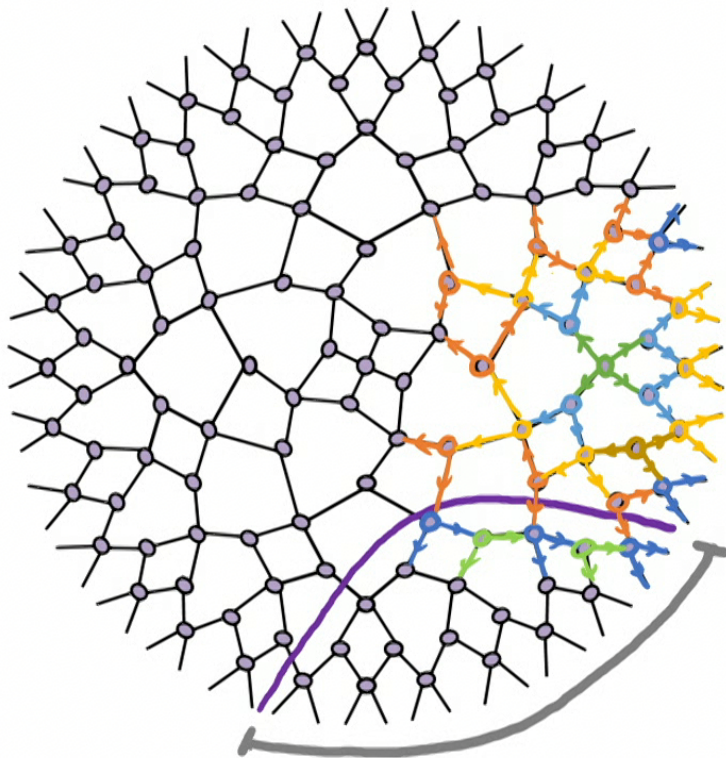
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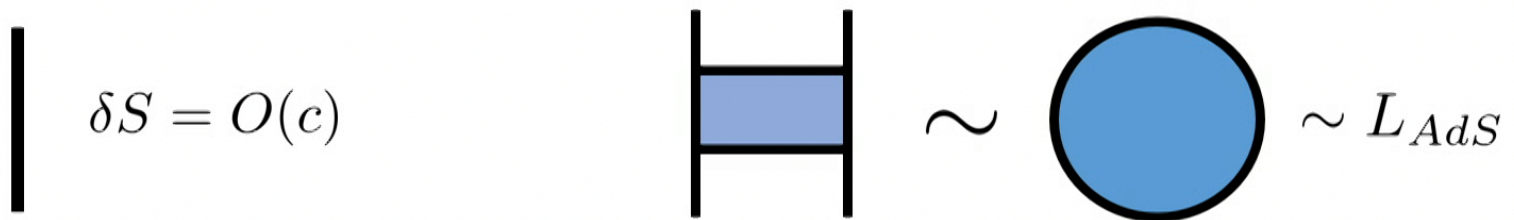
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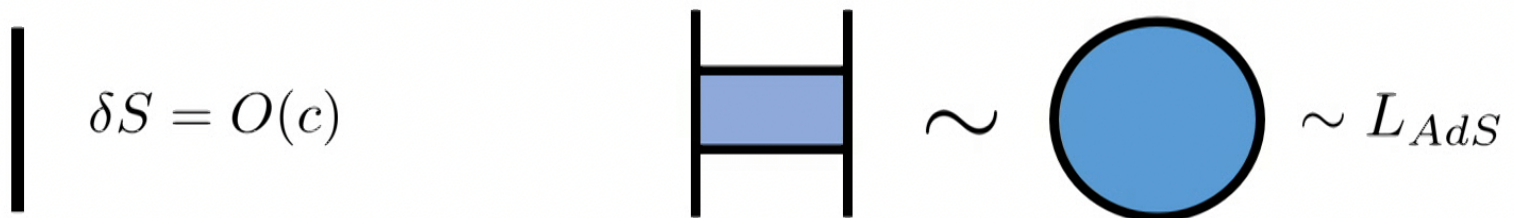
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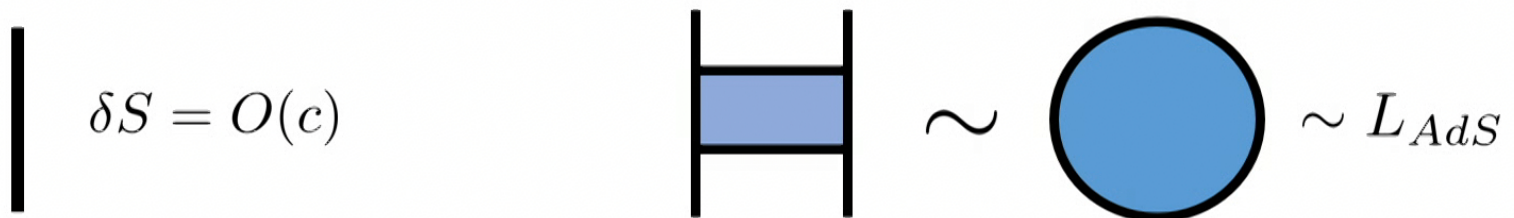


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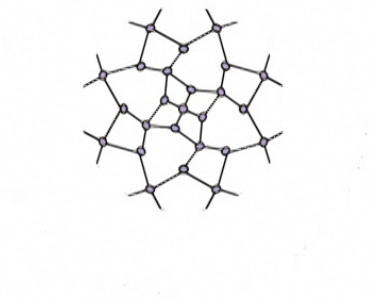
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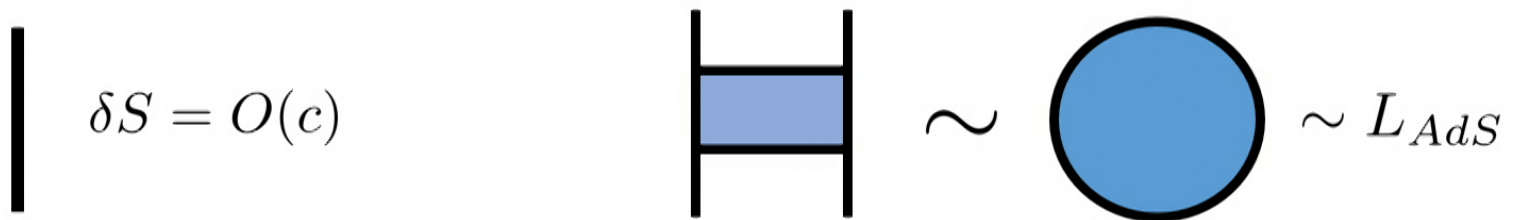
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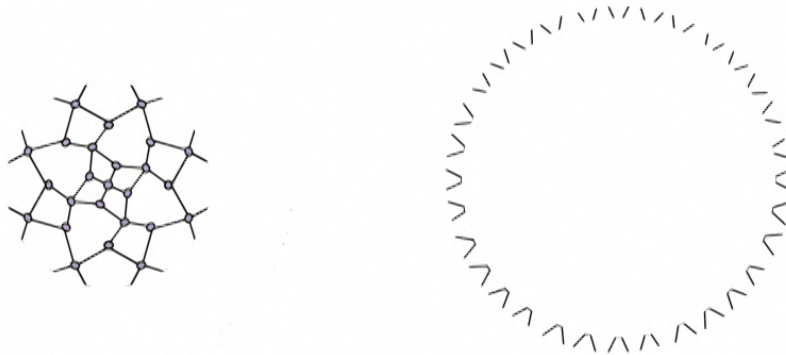
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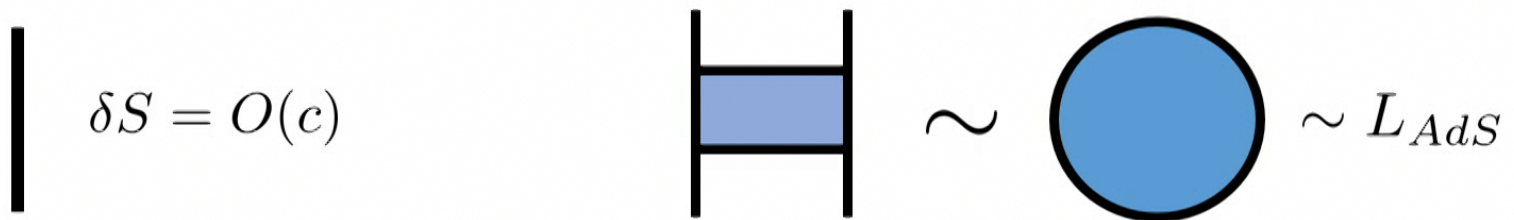
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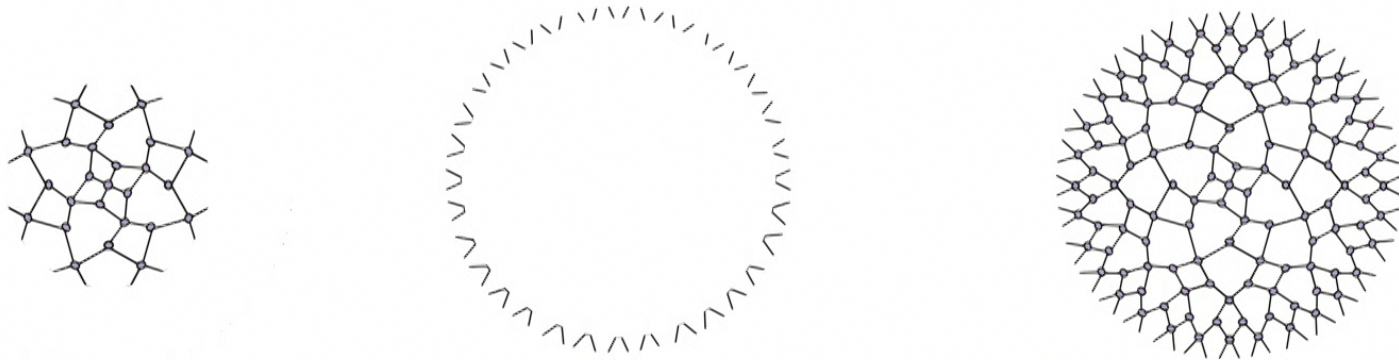
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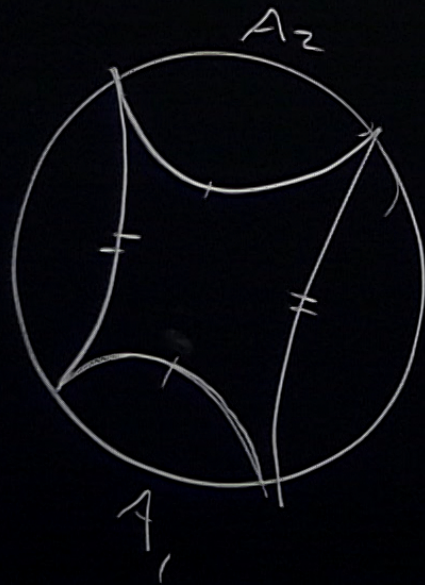
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# Classical Geometry and Random Tensor Networks





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# Beyond the Vacuum

AdS/MERA gives a beautiful description of the ground state. How do we extend this?

Want to understand both directions:

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  - There are many CFT states with same classical geometry, so not unique

---

Geometry  $\rightarrow$  TN  $\rightarrow$  State

[Czech, Hayden, Lashkari, Swingle]

Let's begin with second question:

Finding geometric TN would likely be hard (impossible?) for a *generic* CFT.  
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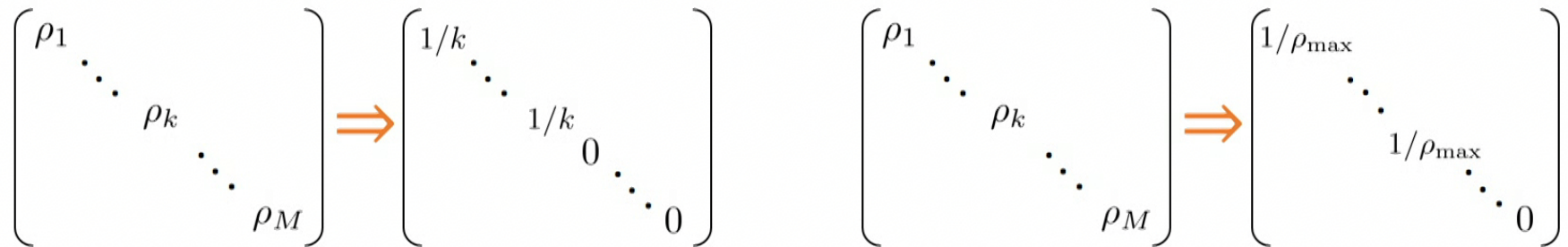
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- **One finds:**  $H_{\min / \max}^{\epsilon}(\rho) = S_{vN}(\rho) \pm O(1/c) \Rightarrow$  **nearly flat spectrum**

---

# Geometry $\rightarrow$ TN

[Hayden, Nezami, Qi, Thomas, Walter, Yang]

## The lesson we should draw:

To good approximation (leading order in  $1/c$ ), we can consider all density matrices to be **maximally mixed**.

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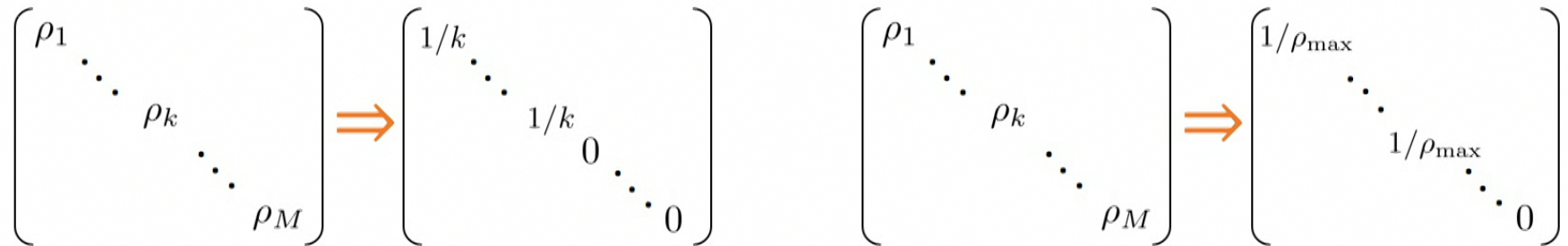
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$$H_{\max}^{\epsilon}(\rho) = \min_{\|\sigma - \rho\| < \epsilon} \log(\text{rank}(\sigma))$$

$$H_{\min}^{\epsilon}(\rho) = \max_{\|\sigma - \rho\| < \epsilon} -\log(\rho_{\max})$$



- **One finds:**  $H_{\min / \max}^{\epsilon}(\rho) = S_{vN}(\rho) \pm O(1/c) \Rightarrow$  **nearly flat spectrum**

---

# Geometry $\rightarrow$ TN

[Hayden, Nezami, Qi, Thomas, Walter, Yang]

## The lesson we should draw:

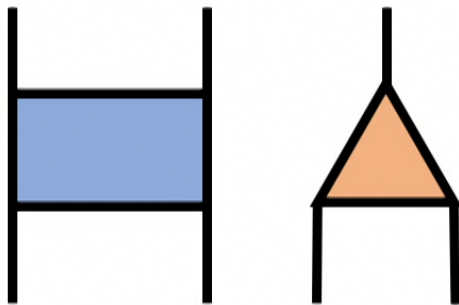
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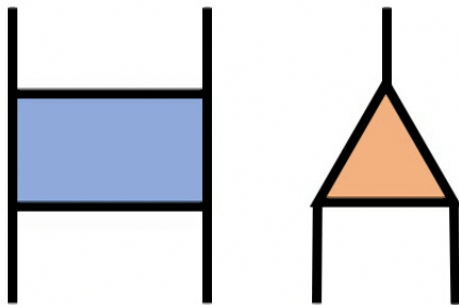
**MERA Tensors**  
Detailed spectrum,  
OPE Coefficients

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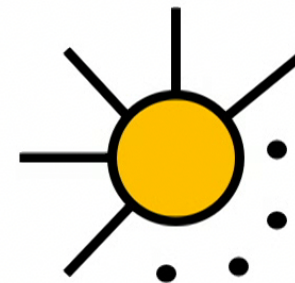
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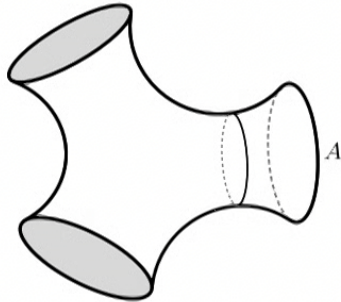
**Random Tensor**  
Flat spectrum  
(large bond dimension)

---

Geometry  $\rightarrow$  TN  $\rightarrow$  State

[Hayden, Nezami, Qi, Thomas, Walter, Yang]

**Consider an arbitrary geometry:**





# Geometry $\rightarrow$ TN $\rightarrow$ State

[Hayden, Nezami, Qi, Thomas, Walter, Yang]

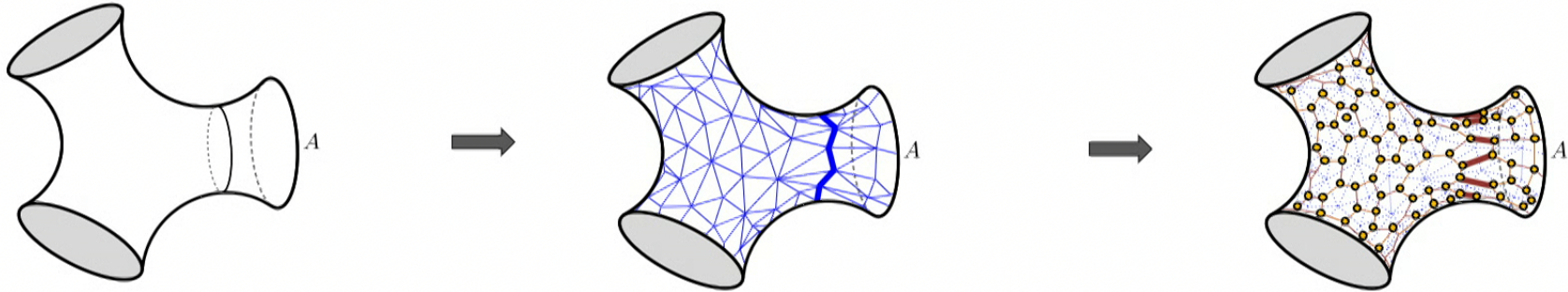
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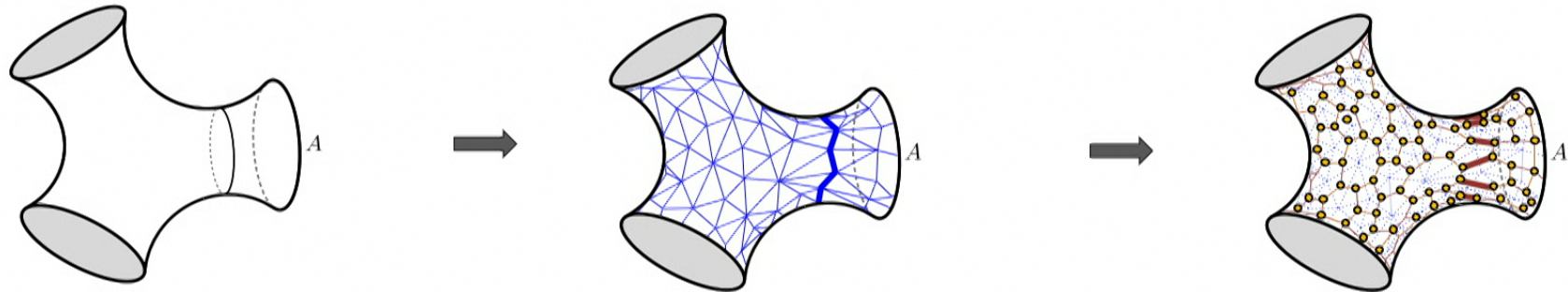
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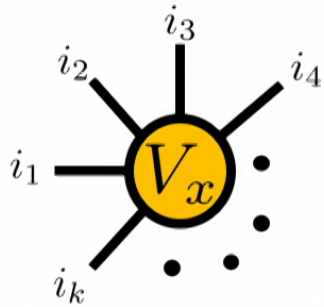


# Geometry $\rightarrow$ TN $\rightarrow$ State

Consider an arbitrary geometry:



- **Number of edges** along cut is (approximately) the **same as bulk area**

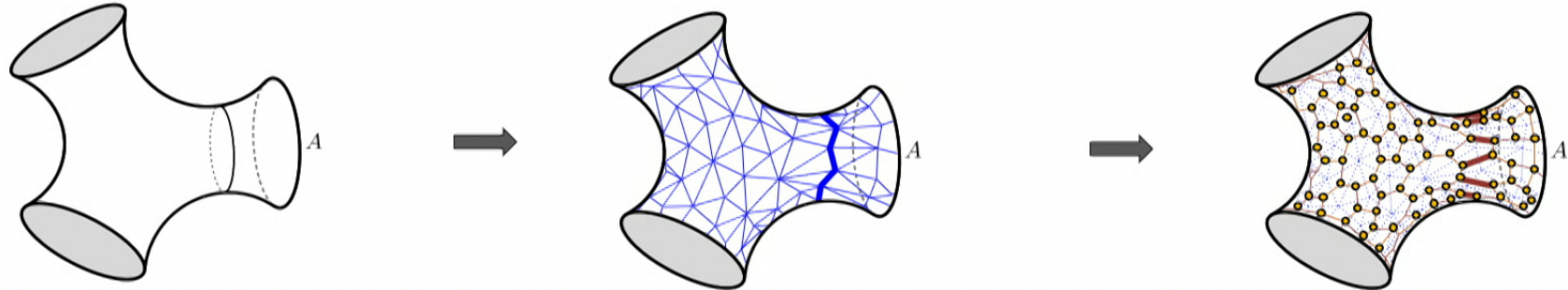


$$|V_x\rangle \in H_D^{\otimes k}$$

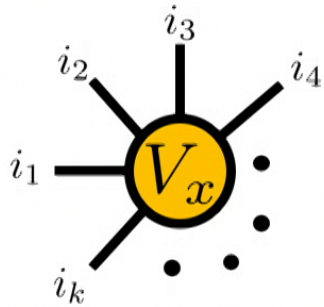
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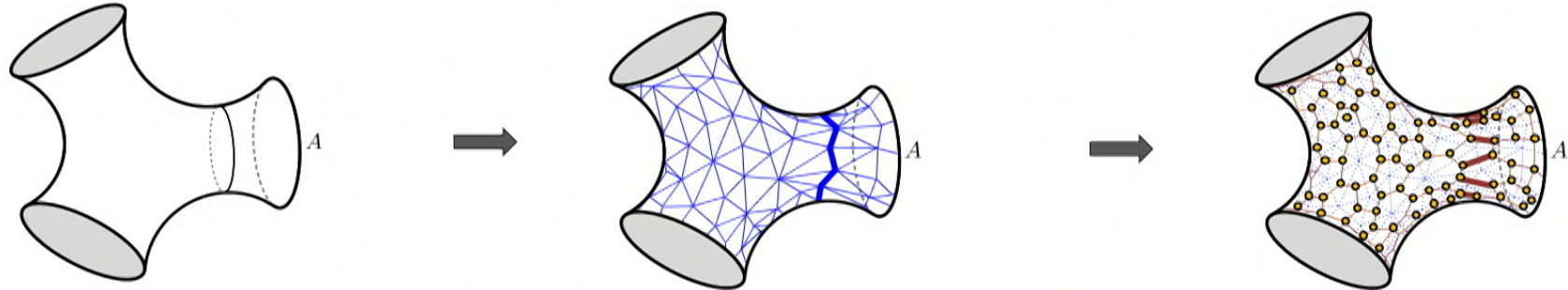
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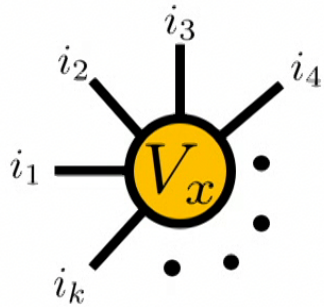
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$$|\psi\rangle = \left( \bigotimes_{\langle xy \rangle \in E} \langle xy| \right) \left( \bigotimes_{x \in V} |V_x\rangle \right)$$

# Geometry $\rightarrow$ TN

## What is the entanglement entropy?

**Bounded from above** by the smallest cut through the tensor network.

**Bounded from below** by the second Renyi entropy:

$$S_2(A) = -\log \text{tr} \rho_A^2$$

**One can show:**

$$\mathbb{E} [\text{tr} \rho_A^2] \propto \sum_{\Gamma_A} e^{-\log(D)|\partial\Gamma_A|}$$

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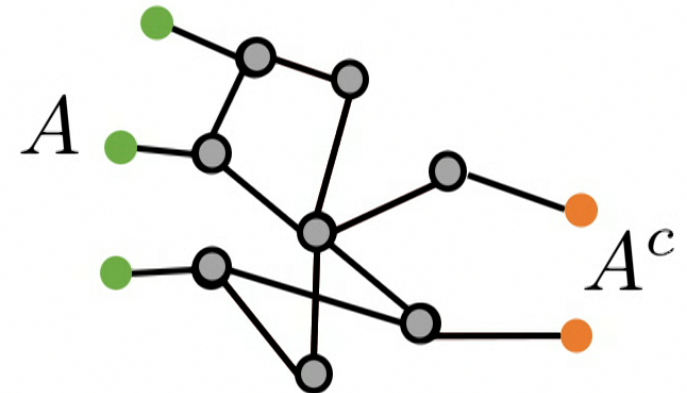
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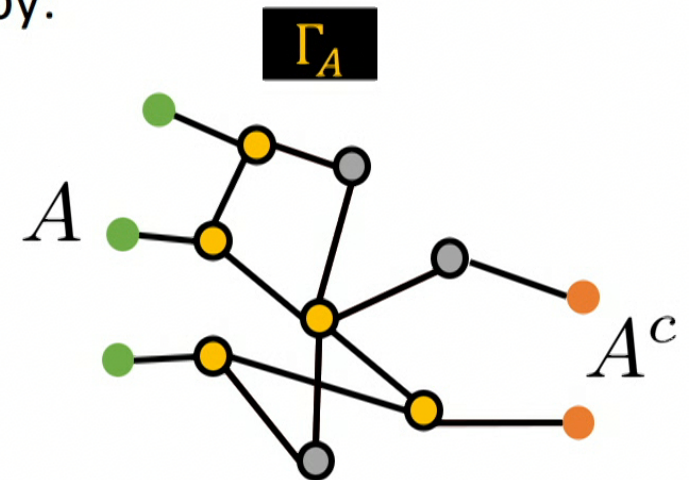
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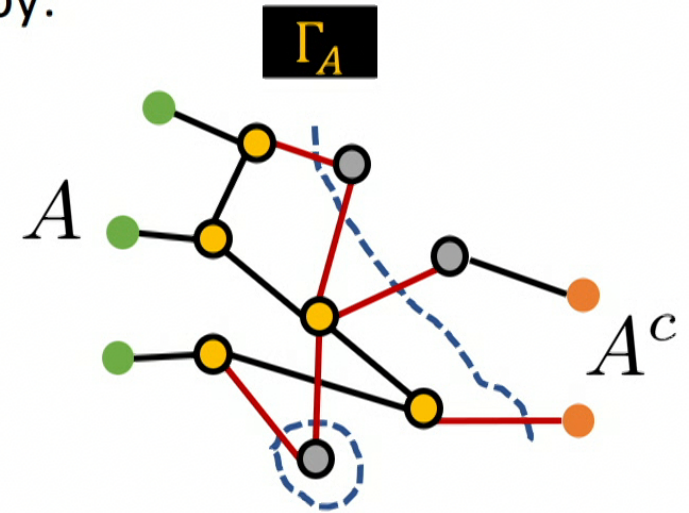
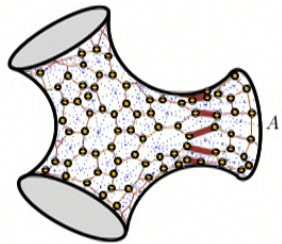
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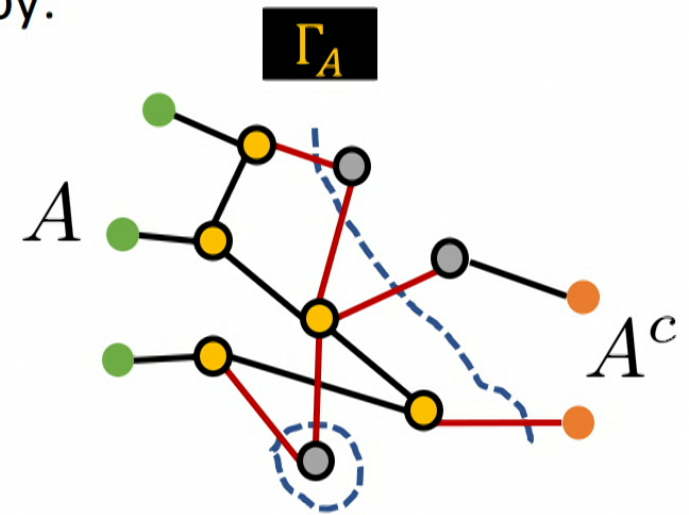
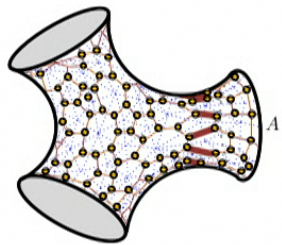
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$$S_2(A) = \min_{\Gamma_A} -\log(D)|\partial\Gamma_A| + \dots$$

---

# Geometry $\rightarrow$ TN $\rightarrow$ State

[Pastawski, Yoshida, Harlow, Preskill]

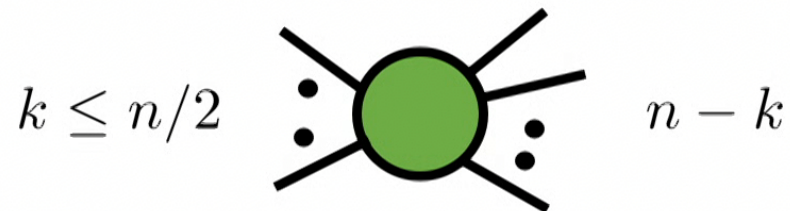
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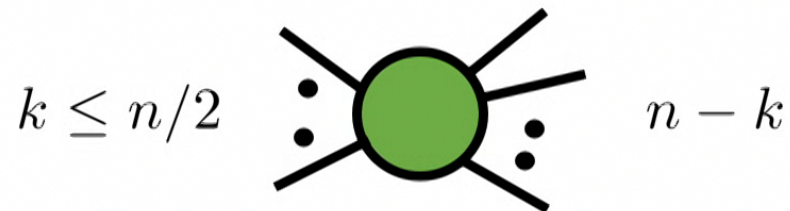
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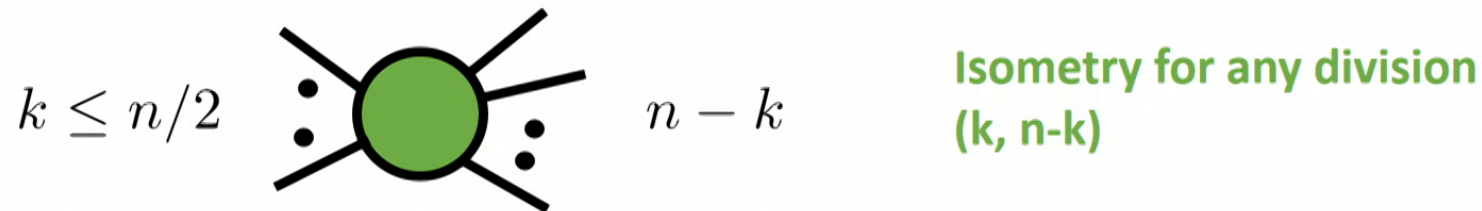


Isometry for any division  
( $k, n-k$ )

# Geometry $\rightarrow$ TN $\rightarrow$ State

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- **Too flat spectrum** in this toy model
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- Idealization of random tensors: **Perfect Tensors**



- Perfect tensor formulation actually predates work with random tensors

---

State  $\rightarrow$  TN  $\rightarrow$  Geometry

### What about the other direction?

How do we take a **CFT state** and **find a tensor network** that will tell us about geometry?

- Not much progress in this direction
- This is an **even more interesting question!**

**Could even ask within scope of ‘maximally entangled’ toy models:**

- Is there an algorithm to find random tensor network description of state?
- Are there constraints needed?

---

# Bulk Physics and Holographic Maps



# Bulk States

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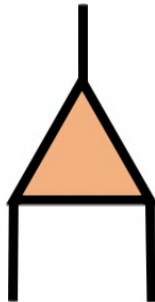
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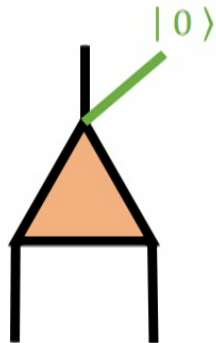


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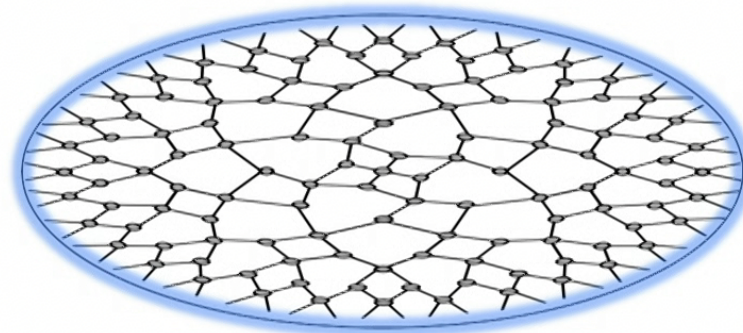
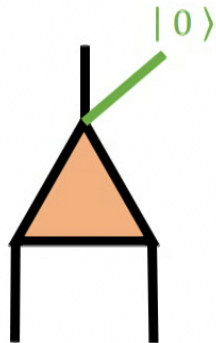


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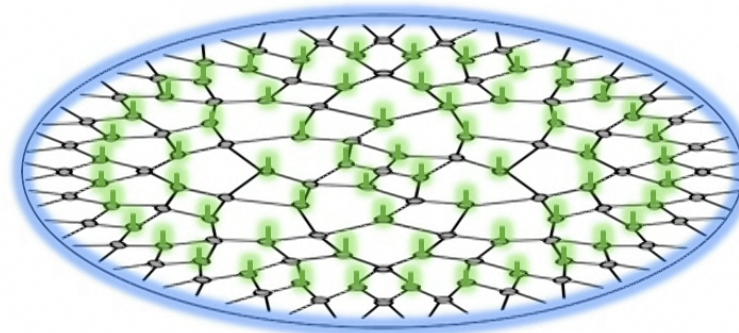
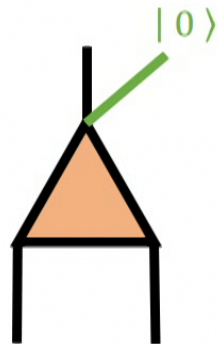
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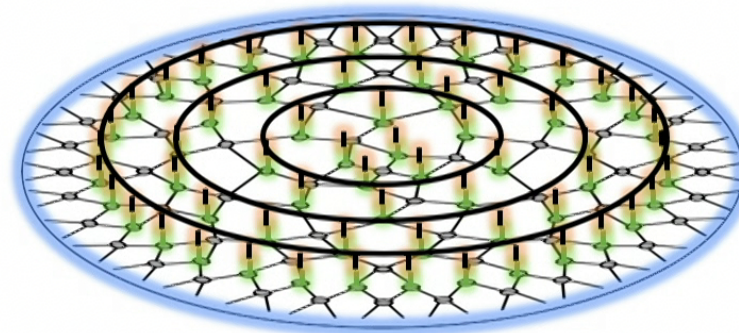
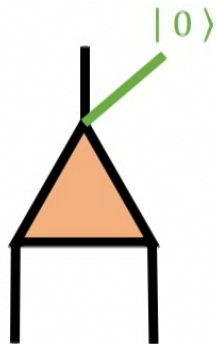


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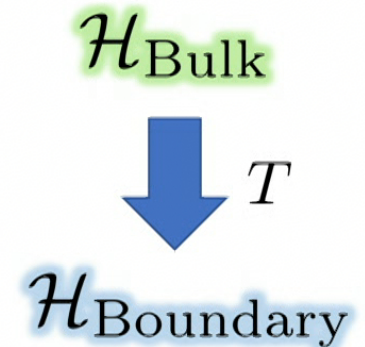
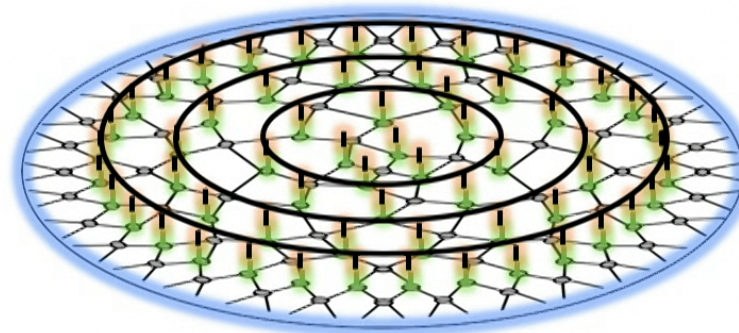
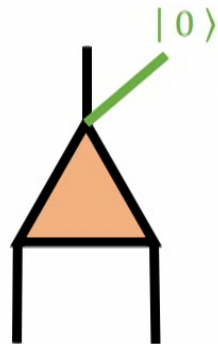


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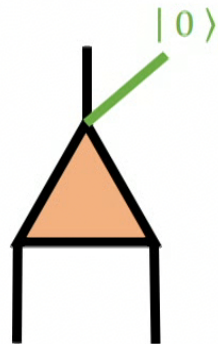


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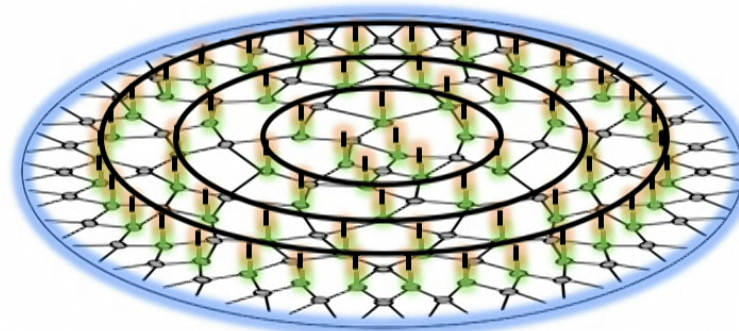
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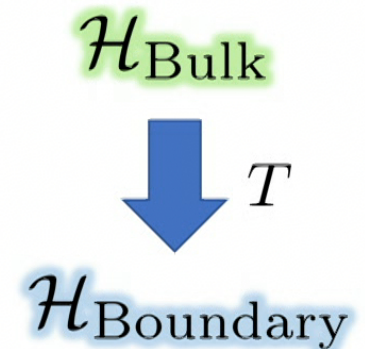
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Holographic Map:



Bulk states  $\leftrightarrow$  Boundary states

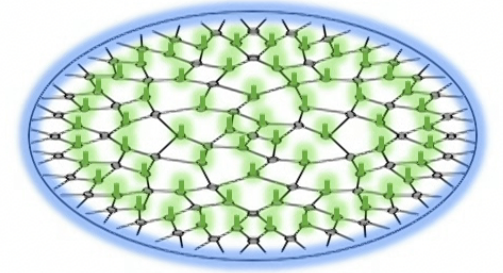


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$$\mathcal{H}_{\text{Bulk}} \cong \mathcal{H}_{\text{Boundary}}$$



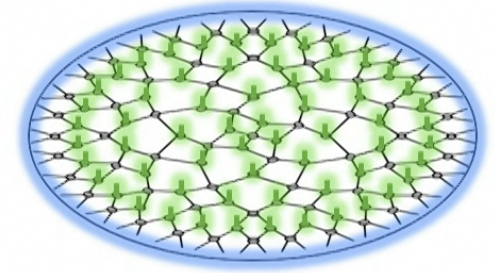
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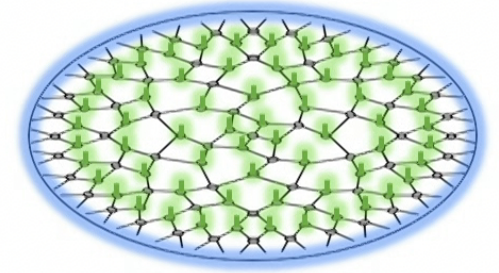
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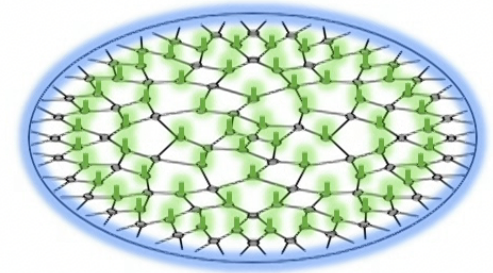
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- Typical bulk state will have **entanglement structure unrelated to the MERA network**.
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A better requirement is to have an **isometry**

$$\mathcal{H}_{\text{Bulk}} \hookrightarrow \mathcal{H}_{\text{Boundary}} \quad \dim \mathcal{H}_{\text{Bulk}} \ll \dim \mathcal{H}_{\text{Boundary}}$$

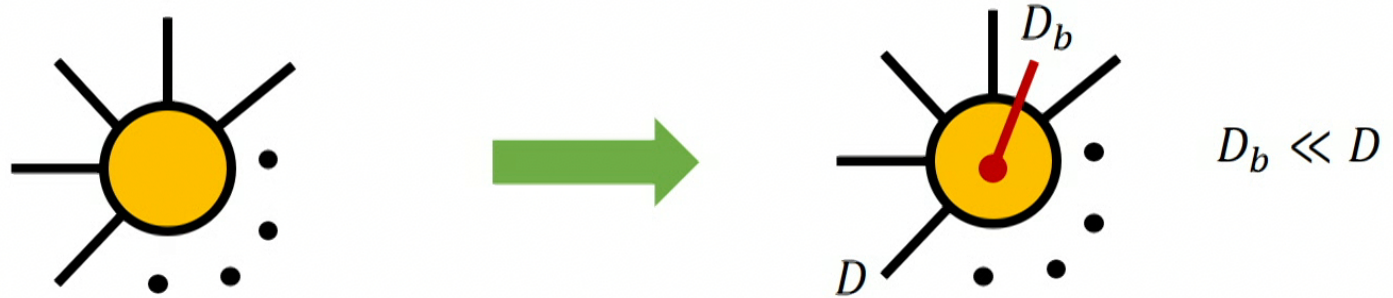
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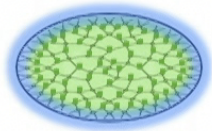


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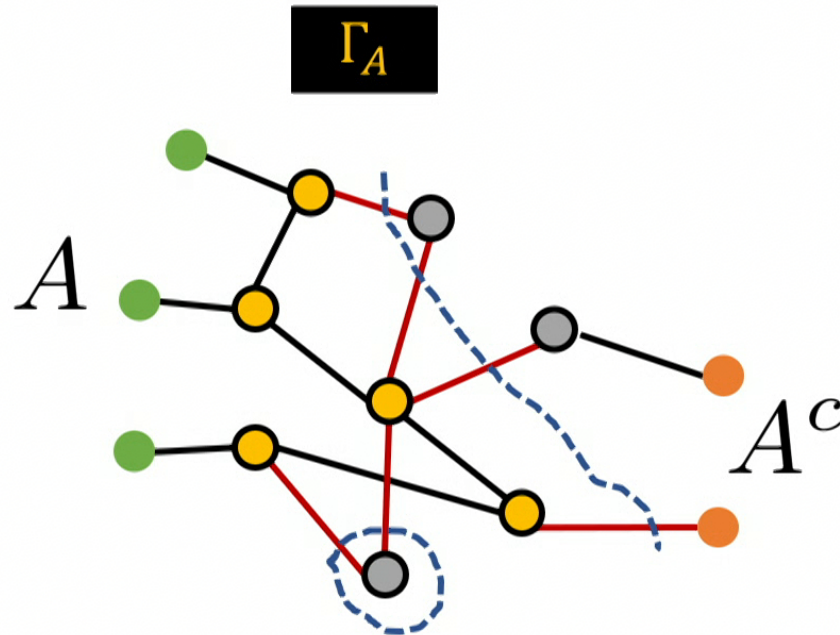


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$$T \quad |\psi_{\text{B}}\rangle \rightarrow |\psi_{\partial}\rangle = T|\psi_{\text{B}}\rangle$$


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Moreover, **reproduce quantum corrections** to the RT formula:



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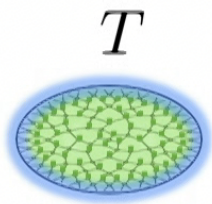


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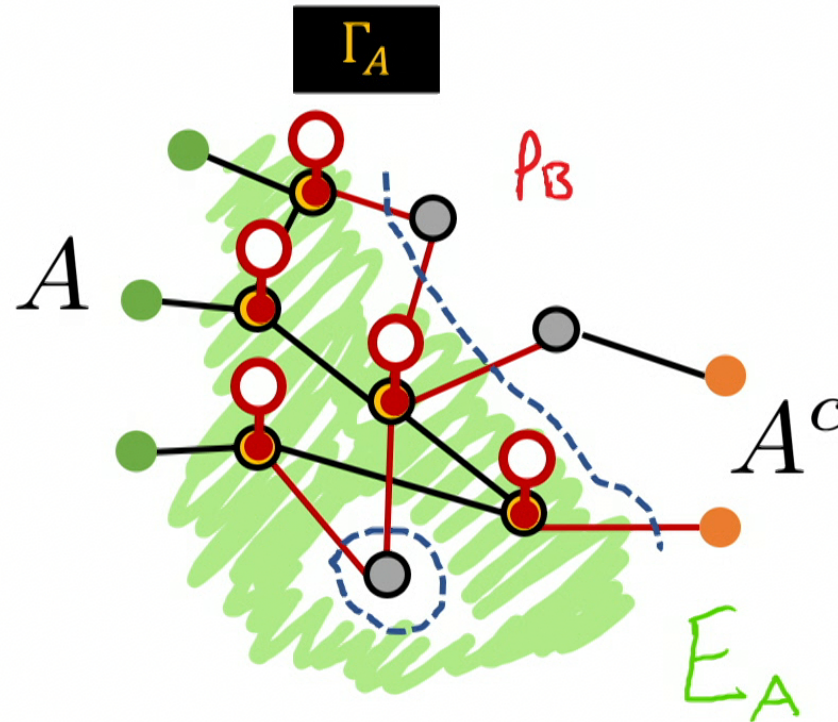
**(Bulk EFT)**

$$\phi_B \rightarrow \mathcal{O}_\partial = T\phi_B T^\dagger$$

**(HKLL for random TN)**

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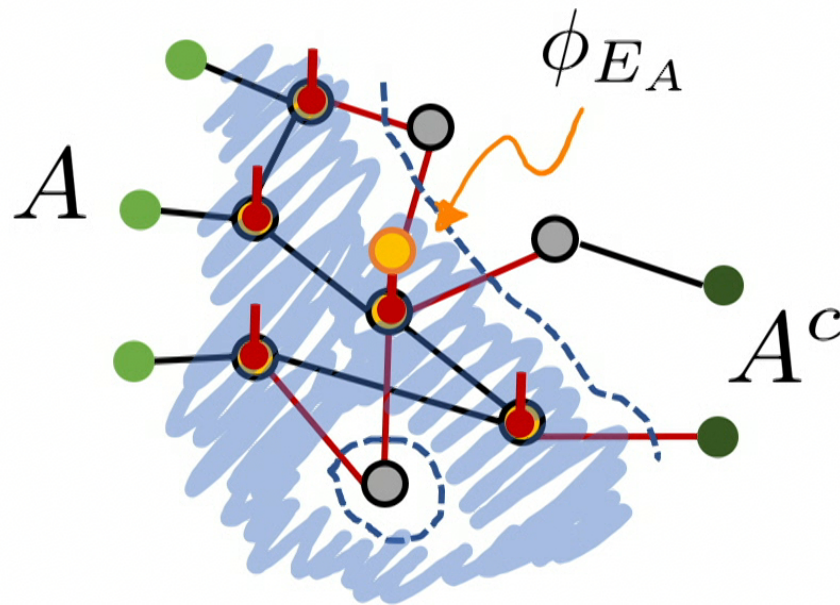
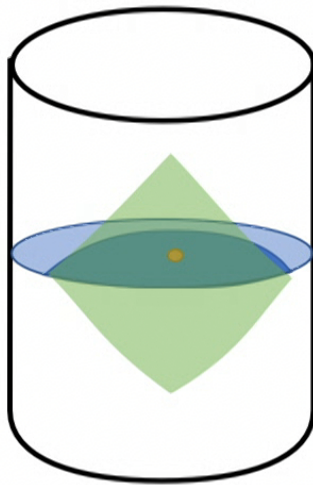
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# Subregion Duality and Error-Correcting Codes

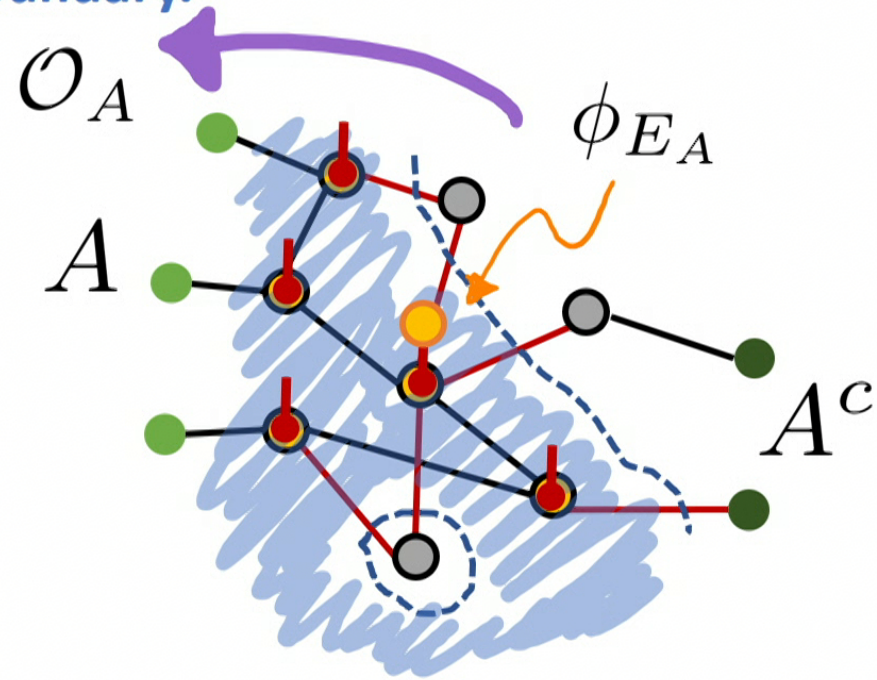
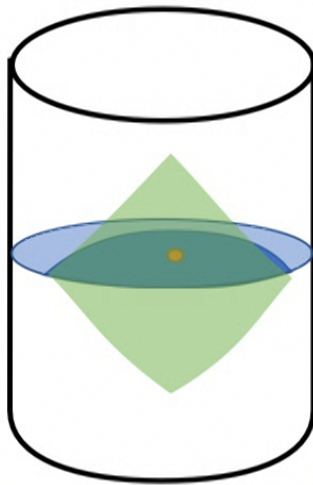
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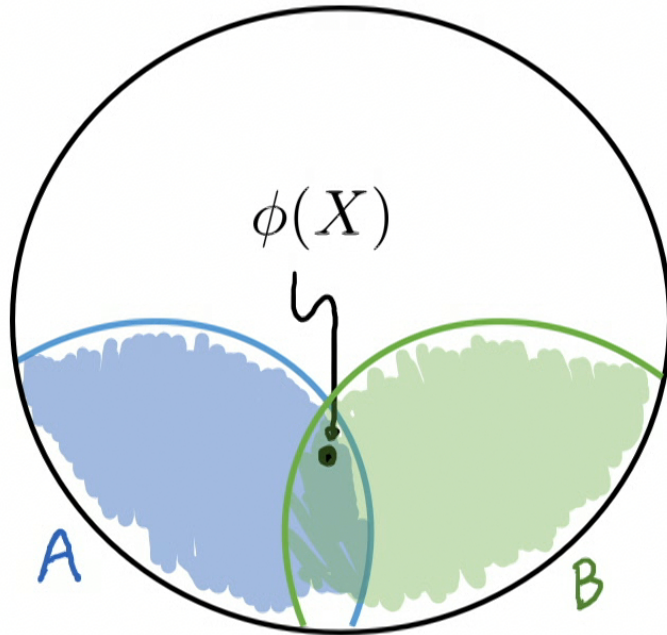


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To understand construction, it's useful to consider a puzzle:

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$\phi(X) \in E_A$  : Dual to operator  $O_{\phi(X)} \in H_A$

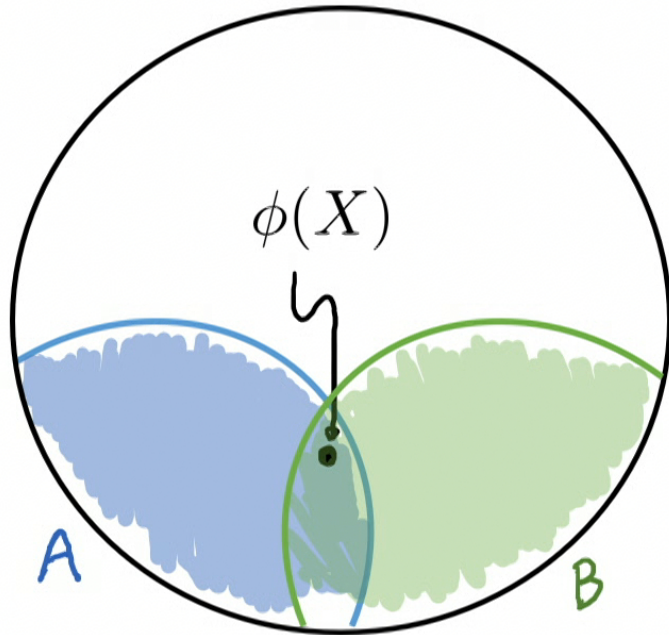


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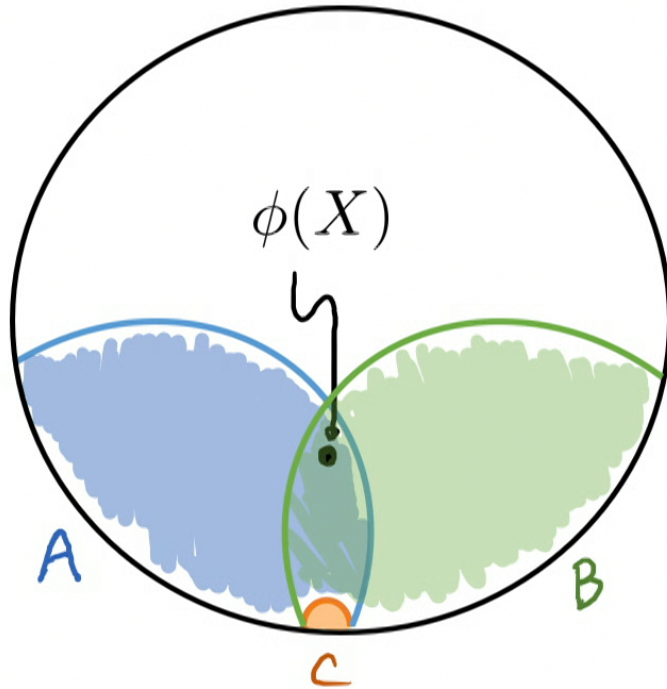
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**If  $\phi(X)$  is a fixed boundary operator**

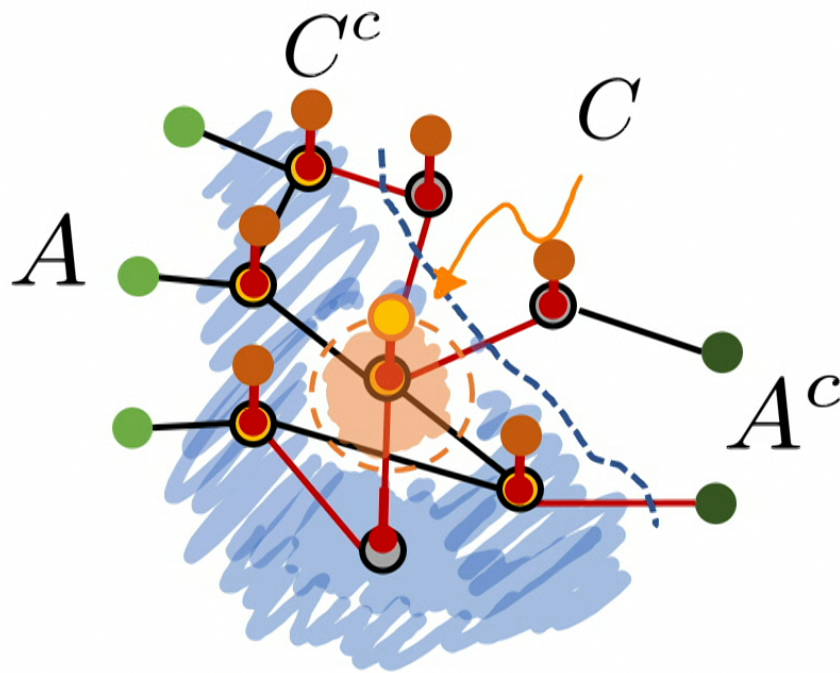
$\Rightarrow O_{\phi(X)} \in H_A \cap H_B = H_C$





# Error-Correction

When can a logical bulk operator be encoded by a physical operator in a region  $A$ ?



Need there to be no mutual information between  $C$  and  $C^c A^c$ :

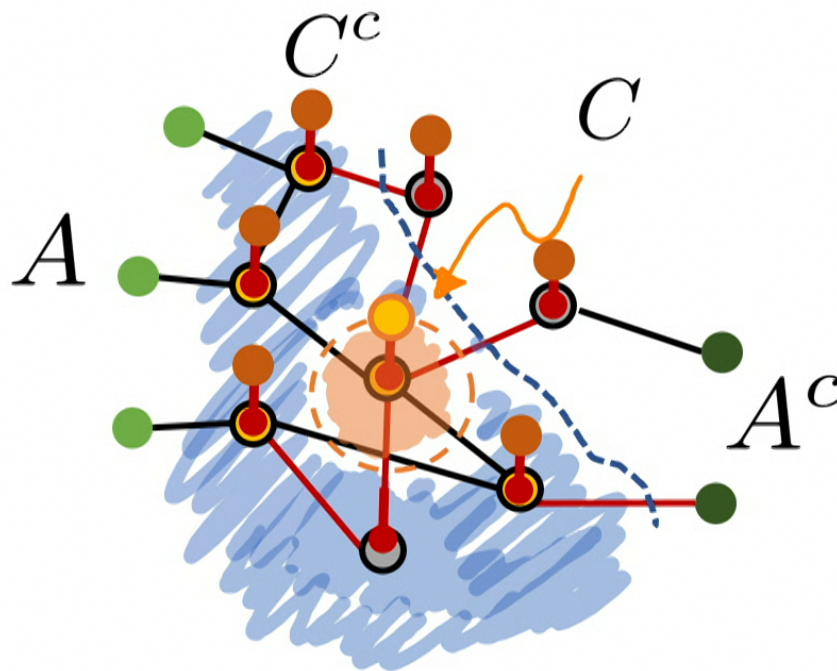
$$S(C) + S(C^c A^c) = S(C C^c A^c)$$

Can calculate by treating bulk legs on equal footing with boundary legs

- Just find min cuts

# Error-Correction

When can a logical bulk operator be encoded by a physical operator in a region  $A$ ?



Need there to be no mutual information between  $C$  and  $C^c A^c$ :

$$S(C) + S(C^c A^c) = S(C C^c A^c)$$

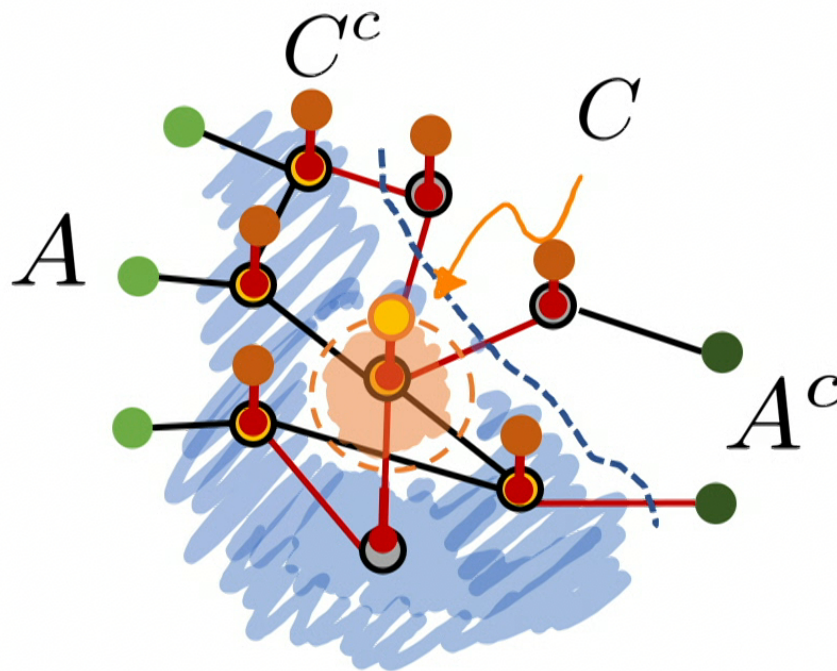
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When is mutual information vanishing?

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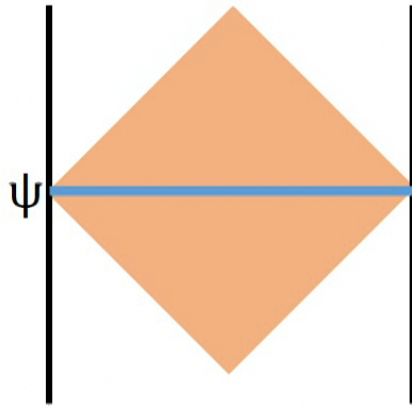
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$$C \subset E_A$$

Reconstruct entanglement wedge using QECC

# Space < Spacetime

So far we have used tensor networks to understand spatial geometry in an emergent extra dimension. What about spacetime?

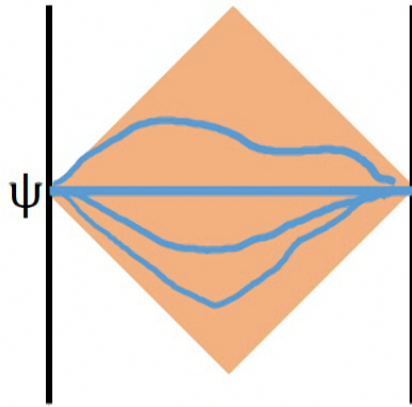


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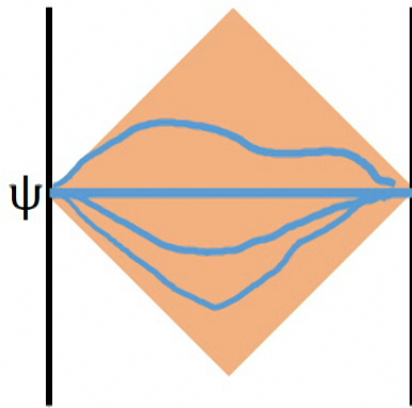
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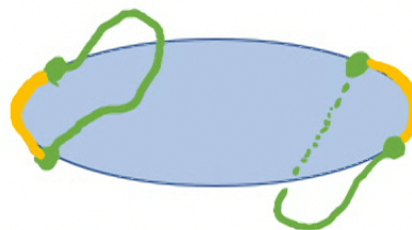
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- Different HRT slices may not lie on same slice

# Real-time Evolution

We would also like to understand real-time evolution:

