

Title:  $W_\infty$  - Higher spins and integrability

Date: Mar 31, 2017 11:00 AM

URL: <http://pirsa.org/17030095>

Abstract: <p>I will discuss the chiral algebra  $W_\infty$  which is obtained from the Virasoro algebra by adding fields of spin 3, 4, .... Via a non-local non-linear map one can show that it is equivalent to Tsybaliuk's Yangian of affine  $u(1)$ . In this way we find an infinite number of commuting conserved charges. Diagonalizing these, the representation theory reduces to combinatorial study of plane partitions, 3-dimensional generalization of the Young diagrams. Tsybaliuk's presentation can be derived from RTT relations using Maulik-Okounkov's free boson R-matrix.</p>

# W algebras

$\mathcal{W}$ -algebras: extensions of the Virasoro algebra by higher spin currents - appear in many different contexts:

- integrable hierarchies of PDE (KdV, KP)
- (old) matrix models
- instanton partition functions and AGT
- holographic dual description of 3d higher spin theories
- quantum Hall effect
- topological strings (toric CY)
- higher spin square (Gaberdiel, Gopakumar)
- combinatorics: plane partitions ( $\equiv$  3d Young diagrams)

## Zamolodchikov $\mathcal{W}_3$ algebra

As in illustration, the  $\mathcal{W}_3$  algebra constructed by Zamolodchikov (1984) has a stress-energy tensor (Virasoro algebra) with OPE

$$T(z)T(w) \sim \frac{c/2}{(z-w)^4} + \frac{2T(w)}{(z-w)^2} + \frac{\partial T(w)}{z-w} + \text{reg.}$$

together with spin 3 primary field  $W(w)$

$$T(z)W(w) \sim \frac{3W(w)}{(z-w)^2} + \frac{\partial W(w)}{z-w} + \text{reg.}$$

We need to find the OPE of  $W$  with itself such that the resulting chiral algebra is 'associative'.

The result:

$$\begin{aligned}
 W(z)W(w) \sim & \frac{c/3}{(z-w)^6} + \frac{2T(w)}{(z-w)^4} + \frac{\partial T(w)}{(z-w)^3} \\
 & + \frac{1}{(z-w)^2} \left( \frac{32}{5c+22} \Lambda(w) + \frac{3}{10} \partial^2 T(w) \right) \\
 & + \frac{1}{z-w} \left( \frac{16}{5c+22} \partial \Lambda(w) + \frac{1}{15} \partial^3 T(w) \right)
 \end{aligned}$$

$\Lambda$  is a quasiprimary 'composite' (spin 4) field,

$$\Lambda(z) = (TT)(z) - \frac{3}{10} \partial^2 T(z).$$

The algebra is non-linear, not a Lie algebra in the usual sense (in fact linearity should not be expected for spins  $\geq 3$ ).

## Construction of $\mathcal{W}_N$ -algebras

We will be interested in specific class of algebras  $\mathcal{W}_N$ , which have generators of spin  $2, 3, \dots, N$ . These algebras are naturally associated to Lie algebras  $\mathfrak{sl}(N)$ . We can construct them in various ways:

- from affine Lie algebras - Casimir algebra, coset construction, Drinfeld-Sokolov reduction ( $\mathfrak{su}(N) \rightarrow \hat{\mathfrak{su}}(N) \rightarrow \mathcal{W}_N$ )
- from free fields - Miura transformation (oper)
- classically ( $c \rightarrow \infty$ ) Poisson brackets in algebra of pseudodifferential operators
- directly - bootstrap (CFT associativity conditions)

## Digression: higher spin algebra $hs(\lambda)$

- so far we defined a family  $\mathcal{W}_N$  with discrete parameter  $N$ ; our main object of study in the following is the algebra  $\mathcal{W}_\infty[\lambda]$  continuously interpolating between these
- there is an analogous construction of interpolating algebra in the case of  $\mathfrak{gl}(N)/\text{End}(V_N)$
- $\mathfrak{sl}(2) \subset \mathfrak{gl}(N)$  (principal) and decompose adjoint  $\rightarrow$  spins  $0, 1, \dots, N-1$

$$T_{m_1}^{l_1} * T_{m_2}^{l_2} = \sum C_{m_1 m_2}^{l_1 l_2 l} (N) T_{m_1 + m_2}^l$$

- the structure constants can be chosen to be rational in  $N$
- we can thus take generators of spin  $0, 1, \dots, \infty$  and define a new associative algebra  $hs(\lambda)$  using these rational functions as its structure constants (analytic continuation,  $N \rightarrow \lambda$ )
- specializing  $\lambda$  to an integer  $N$  (and truncating spins  $\geq N$ ), we get back  $\mathfrak{gl}(N)$



## Digression: higher spin algebra $hs(\lambda)$

- so far we defined a family  $\mathcal{W}_N$  with discrete parameter  $N$ ; our main object of study in the following is the algebra  $\mathcal{W}_\infty[\lambda]$  continuously interpolating between these
- there is an analogous construction of interpolating algebra in the case of  $\mathfrak{gl}(N)/\text{End}(V_N)$
- $\mathfrak{sl}(2) \subset \mathfrak{gl}(N)$  (principal) and decompose adjoint  $\rightarrow$  spins  $0, 1, \dots, N-1$

$$T_{m_1}^{l_1} * T_{m_2}^{l_2} = \sum C_{m_1 m_2}^{l_1 l_2 l} (N) T_{m_1 + m_2}^l$$

- the structure constants can be chosen to be rational in  $N$
- we can thus take generators of spin  $0, 1, \dots, \infty$  and define a new associative algebra  $hs(\lambda)$  using these rational functions as its structure constants (analytic continuation,  $N \rightarrow \lambda$ )
- specializing  $\lambda$  to an integer  $N$  (and truncating spins  $\geq N$ ), we get back  $\mathfrak{gl}(N)$

- explicit formula for multiplication using 3j and 6j symbols (Pope, Shen, Romans)
- there is also a one-line algebraic definition of  $hs(\lambda)$

$$hs(\lambda) = \frac{\mathcal{U}(\mathfrak{sl}(2))}{(X^2 + Y^2 + Z^2 - (\lambda^2 - 1))}$$

- this definition also makes clear the geometric meaning of  $hs(\lambda)$  as the quantum sphere  $S^2$
- expansion at  $\lambda = \infty$ :  $hs(\lambda)$  is deformation quantization ( $\hbar \sim \frac{1}{\lambda}$ ) of algebra of functions on  $S^2$
- truncation to  $gl(N)$  for  $\lambda \in \mathbb{N}$  goes under the name of fuzzy sphere (quantization - integral phase space volume)



$\mathcal{W}_\infty$  from bootstrap

- Gaberdiel-Gopakumar: use the associativity conditions to construct  $\mathcal{W}_\infty$  (spins 2, 3, 4, ...)

$$\sum_m \begin{array}{c} i \quad l \\ \diagdown \quad \diagup \\ \text{---} m \text{---} \\ \diagup \quad \diagdown \\ j \quad k \end{array} = \sum_n \begin{array}{c} i \quad l \\ \diagdown \quad \diagup \\ \text{---} n \text{---} \\ \diagup \quad \diagdown \\ j \quad k \end{array}$$

- equivalently (Jacobi identity for mode operators)

$$[[A_j, B_k], C_l] + [[B_k, C_l], A_j] + [[C_l, A_j], B_k] = 0$$

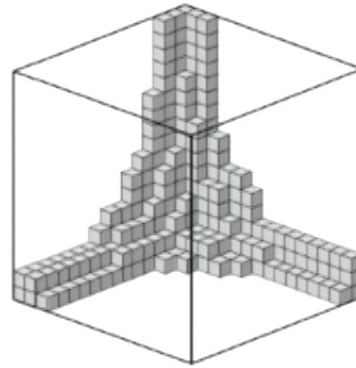
- two-parametric family of solutions: central charge  $c$  and rank parameter  $\lambda$  (Gaberdiel/Gopakumar)
- choosing  $\lambda = N \rightarrow$  truncation of  $\mathcal{W}_\infty$  to  $\mathcal{W}_N$ , so  $\mathcal{W}_\infty$  is interpolating algebra for  $\mathcal{W}_N$  series (cf.  $\mathfrak{gl}(N)$  vs  $hs(\lambda)$ )
- working in basis of fields from Miura transformation, one can find explicit expressions for OPE (quadratic nonlinearity)

The resulting OPEs are very complicated

$$\begin{aligned}
U_3(z)U_5(w) \sim & \frac{1}{(z-w)^7} \left( \frac{1}{2} \alpha(n-3)(n-2)(n-1)n \left( 4\alpha^2 \left( \alpha^2(n(5n-9)+1) - 3n+4 \right) + 1 \right) \right) \\
& + \frac{1}{(z-w)^6} \left( \frac{1}{6} (n-3)(n-2)(n-1) \left( 6\alpha^4 n(2n-3) + \alpha^2(10-9n) + 1 \right) U_1(w) \right) \\
& + \frac{1}{(z-w)^5} \left( -\alpha(n-3)(n-2)(n-1) \left( -4\alpha^2 + 3\alpha^2 n - 1 \right) (U_1 U_1)(w) \right) \\
& + \alpha(n-3)(n-2) \left( 4\alpha^2 n^2 - 4\alpha^2 n - n - 2 \right) U_2(w) \\
& - \frac{1}{2} \alpha^2(n-3)(n-2)(n-1) \left( 4\alpha^2 n(2n-3) - 3n+2 \right) U_1'(w) \\
& + \frac{1}{(z-w)^4} \left( -\alpha(n-3)(n-2)(n-1) \left( \alpha^2(3n-4) - 1 \right) (U_1' U_1)(w) \right) \\
& - \frac{1}{2} (n-3)(n-2) \left( 2\alpha^2(n-1) - 1 \right) (U_1 U_2)(w) \\
& + (n-3) \left( \alpha^2(n^2+2) - 3 \right) U_3(w) \\
& - \frac{1}{4} \alpha^2(n-3)(n-2)(n-1) \left( 4\alpha^2 n(2n-3) - 3n+2 \right) U_1''(w) \\
& + \alpha(n-3)(n-2) \left( \alpha^2(n-1)n - 1 \right) U_2'(w) \\
& + \dots
\end{aligned}$$

# Topological vertex

- characters of other special  $\mathcal{W}_{1+\infty}$  reps are counting plane partitions with specified Young diagram asymptotics



- level  $\leftrightarrow$  number of boxes
- this is exactly the topological vertex of topological strings
- so the topological vertex  $\leftrightarrow$  characters of  $\mathcal{W}_{1+\infty}$  reps
- truncation to integer  $\lambda \leftrightarrow$  restriction the height